Math 4990 - Final Exam

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Problem 1: Consider the integral function

$$J[y] = \int_1^2 \frac{y'(x)^2}{x} dx$$

The integrand here is strongly convex (on an appropriately defined set). Find the unique $y \in D$ that minimizes J[y] over D, for the following cases. In each case, is the minimizer (if it exists) unique?

Part (a)
$$D = \{y \in C^1[1,2] : y(1) = 0, y(2) = 3\}$$

Solution

If J is a minimum at $y \in D$ then y must satisfy the Euler-Lagrange equation $F_y(x, y, y') - \frac{d}{dx} F_{y'}(x, y, y') = 0$ where $F(x, y, y') = \frac{y'(x)^2}{x}$.

$$\frac{\partial F}{\partial y} - \frac{d}{dx} \frac{\partial F}{\partial y'} = 0 \Longrightarrow \frac{\partial F}{\partial y} = \frac{d}{dx} \frac{\partial F}{\partial y'}$$

$$\Longrightarrow 0 = \frac{d}{dx} \frac{2y'}{x} \Longrightarrow \int 0 dx = \int \frac{d}{dx} \frac{2y'}{x} dx$$

$$\Longrightarrow \frac{2y'}{x} = C \quad (C \in \mathbb{R})$$

$$\Longrightarrow y' = Bx \quad (B \in \mathbb{R})$$

$$\Longrightarrow \int y' dx = \int Bx dx$$

$$\Longrightarrow y = c_1 x^2 + c_2 \quad (c_1, c_2 \in \mathbb{R})$$

y must also satisfy the boundary conditions:

$$\begin{cases} y(1) = c_1 + c_2 = 0 \\ y(2) = 4c_1 + c_2 = 3 \end{cases} \implies c_1 = \frac{3}{5}, c_2 = -\frac{3}{5} \implies \boxed{y = \frac{3}{5}x^2 - \frac{3}{5}}$$

We have shown with Euler-Lagrange equation the necessary conditions for y to minimize J[y] and because we are given the integrand is strongly convex on the defined set D, we have the sufficient conditions to guarantee y is the unique minimizer.

Part (b)
$$D = \{y \in C^1[1,2] : y(2) = 3\}$$

Solution

From the notes given in class, we have the following:

Theorem. If F(x, y, y') is strongly convex in (y, y') then a solution y(x) of the Differential Euler-Lagrange equation uniquely minimizes

$$J[y] = \int_{a}^{b} F(x, y, y') dx$$

over $\{y \in C'[a, b] : y(b) = B\}$ if $F_{y'}(a, y(a), y'(a)) = 0$.

We find the solution to the Euler-Lagrange equation $F_y(x, y, y') - \frac{d}{dx}F_{y'}(x, y, y') = 0$ where $F(x, y, y') = \frac{y'(x)^2}{x}$.

$$\frac{\partial F}{\partial y} - \frac{d}{dx} \frac{\partial F}{\partial y'} = 0 \Longrightarrow \frac{\partial F}{\partial y} = \frac{d}{dx} \frac{\partial F}{\partial y'}$$

$$\Longrightarrow 0 = \frac{d}{dx} \frac{2y'}{x} \Longrightarrow \int 0 dx = \int \frac{d}{dx} \frac{2y'}{x} dx$$

$$\Longrightarrow \frac{2y'}{x} = C \quad (C \in \mathbb{R})$$

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$$\Longrightarrow \int y' dx = \int Bx dx$$

$$\Longrightarrow y = c_1 x^2 + c_2 \quad (c_1, c_2 \in \mathbb{R})$$

y must also satisfy the boundary condition:

$$y(2) = 4c_1 + c_2 = 3 \Longrightarrow c_2 = 3 - 4c_1$$
$$\Longrightarrow y = c_1 x^2 + (3 - 4c_1)$$
$$\Longrightarrow y = c_1 (x^2 - 4) + 3$$

thus

$$F_{y'}(1, y(1), y'(1)) = 0 \Longrightarrow \frac{2y'(1)}{1} = 0$$
$$\Longrightarrow y'(1) = 0$$
$$\Longrightarrow 2c_1(1)$$
$$\Longrightarrow c_1 = 0$$

which implies

$$y(x) = 3$$

Given the integrand is strongly convex on the defined set D along with the theorem mentioned above, y uniquely minimizes J[y] over D.

Part (c) $D = C^1[1, 2]$

Solution

If J is a minimum at $y \in D$ then y must satisfy the Euler-Lagrange equation $F_y(x, y, y') - \frac{d}{dx}F_{y'}(x, y, y') = 0$ where $F(x, y, y') = \frac{y'(x)^2}{x}$.

$$\frac{\partial F}{\partial y} - \frac{d}{dx} \frac{\partial F}{\partial y'} = 0 \Longrightarrow \frac{\partial F}{\partial y} = \frac{d}{dx} \frac{\partial F}{\partial y'}$$

$$\Longrightarrow 0 = \frac{d}{dx} \frac{2y'}{x} \Longrightarrow \int 0 dx = \int \frac{d}{dx} \frac{2y'}{x} dx$$

$$\Longrightarrow \frac{2y'}{x} = C \quad (C \in \mathbb{R})$$

$$\Longrightarrow y' = Bx \quad (B \in \mathbb{R})$$

$$\Longrightarrow \int y' dx = \int Bx dx$$

$$\Longrightarrow y = c_1 x^2 + c_2 \quad (c_1, c_2 \in \mathbb{R})$$

From strong convexity and observation we see if $y = c_2$, we achieve a minimum for $J[y] = \int_1^2 \frac{y'(x)^2}{x} dx$ since $y'(x)^2 \ge 0$ and $x \in [1, 2]$ which implies $J[y] \ge 0$. However y is not a unique minimizer since $c^2 \in \mathbb{R}$.

Problem 2: Let $\delta(x)$ be the Dirac delta function. Justify the identity

$$\delta(1-x^2) = \frac{\delta(x-1) + \delta(x+1)}{2}$$

Hint: consider the integral $\int_{-\infty}^{\infty} \delta(1-x^2)dx$

Solution For any continuous function f(x), we have

$$\int_{-\infty}^{\infty} f(x)\delta(1-x^2)dx = \underbrace{\int_{0}^{\infty} f(x)\delta(1-x^2)dx}_{(1)} + \underbrace{\int_{-\infty}^{0} f(x)\delta(1-x^2)dx}_{(2)}$$

If we let $\sqrt{u} = x$ and $\frac{1}{2\sqrt{u}}du = dx$

$$(1) = \int_0^\infty \frac{f(\sqrt{u})\delta(1-u)}{2\sqrt{u}} du = \frac{f(1)}{2}$$

If we let $-\sqrt{u} = x$ and $-\frac{1}{2\sqrt{u}}du = dx$

$$(2) = -\int_{-\infty}^{0} \frac{f(-\sqrt{u})\delta(1-u)}{2\sqrt{u}} du = \int_{0}^{\infty} \frac{f(-\sqrt{u})\delta(1-u)}{2\sqrt{u}} du = \frac{f(-1)}{2}$$

thus

$$\int_{\infty}^{\infty} f(x)\delta(1-x^2)dx = \frac{1}{2}[f(1)+f(-1)]$$

$$= \frac{1}{2}\left[\int_{-\infty}^{\infty} f(x)\delta(x-1)dx + \int_{-\infty}^{\infty} f(x)\delta(x+1)dx\right]$$

$$= \frac{1}{2}\int_{-\infty}^{\infty} f(x)(\delta(x-1)+\delta(x+1))dx$$

$$= \int_{-\infty}^{\infty} f(x)\frac{\delta(x-1)+\delta(x+1)}{2}dx$$

Since δ functions are defined only by how they behave in integrals we conclude

$$\delta(1-x^2) = \frac{\delta(x-1) + \delta(x+1)}{2}$$

Problem 3: Let $f:[0,1] \to \mathbb{R}$ be a given continuous function and consider the following differential equation for the function y:

$$\begin{cases} y'''(x) = f(x) & 0 \le x \le 1\\ y(0) = y'(0) = y''(0) = 0 \end{cases}$$
 (1)

Intuitively, finding y(x) requires three integrations. This problem illustrates that by employing a Green's function we can actually find y via a single integral.

Part (a) Show that if the Green's function u = G(x; s) satisfies

$$\begin{cases} u'''(x) = \delta(x - s) & 0 \le x \le 1 \\ u(0) = u'(0) = u''(0) = 0 \end{cases}$$

then the function

$$y(x) = \int_0^1 f(s)G(x;s)ds \tag{2}$$

satisfies (1).

Solution

From the notes given in class, we have the following:

Theorem.

The function $y(x) = \int_0^1 f(s)G(x;s)ds$ satisfies the differential equation

$$\begin{cases} L[y] = y'''(x) = f(x) \\ y(0) = y'(0) = y''(0) = 0 \end{cases} 0 \le x \le 1$$

Let us first observe

$$L[y] = y''' = \frac{d^3}{dx^3} \left(\int_0^1 f(s)G(x;s)ds \right)$$

$$= \int_0^1 \frac{d^3}{dx^3} \left(f(s)G(x;s) \right) ds$$

$$= \int_0^1 f(s) \frac{d^3}{dx^3} \left(G(x;s) \right) ds$$

$$= \int_0^1 f(s)G'''(x;s) ds$$

$$= \int_0^1 f(s)L[G] ds$$

and

$$L[G] = G'''(x; s) = \delta(x - s), \qquad 0 \le x \le 1$$

thus

$$L[y] = L\left[\int_0^1 f(s)G(x;s)ds\right]$$
$$= \int_0^1 L\left[f(s)G(x;s)\right]ds$$
$$= \int_0^1 f(s)L[G]ds$$
$$= \int_0^1 f(s)\delta(x-s)ds$$
$$= f(x)$$

along with the boundary conditions given G(0; s) = 0

$$y(0) = \int_0^1 f(s) \underbrace{G(0; s)}_0 ds = 0$$
$$y'(0) = \int_0^1 f'(s) \underbrace{G(0; s)}_0 ds = 0$$
$$y''(0) = \int_0^1 f''(s) \underbrace{G(0; s)}_0 ds = 0$$

We have showed (2) satisfies (1).

Part (b) Find the Green's function G(x:s) for this problem

$$\begin{cases} y'''(x) = f(x) & 0 \le x \le 1 \\ y(0) = y'(0) = y''(0) = 0 \end{cases}$$

Solution To find Green's function we solve for

$$\begin{cases} u'''(x) = \delta(x - s) & 0 \le x \le 1 \\ u(0) = u'(0) = u''(0) = 0 \end{cases}$$

where u(x) = G(x; s).

For $x \neq s$ the differential equation is

$$u'''(x) = 0 \qquad 0 \le x \le 1$$

whose general solution is found to be

$$u(x) = Ax^2 + Bx + C$$
 $(A, B, C \in \mathbb{R})$

thus

$$u(x) = \begin{cases} Ax^2 + Bx + C & x < s \\ Ex^2 + Dx + F & x > s \end{cases}$$

The boundary conditions require

$$u(0) = 0 \implies Ax^2 + Bx + C = 0 \implies C = 0$$

 $u'(0) = 0 \implies 2Ax + B = 0 \implies B = 0$
 $u''(0) = 0 \implies 2A = 0 \implies A = 0.$

The Green function now has the form

$$u(x) = \begin{cases} 0 & x < s \\ Ex^2 + Dx + F & x > s. \end{cases}$$

We also need to satisfy the second derivative discontinuity condition since $u''' = \delta(x - s) \Rightarrow u'' = H(x - s)$

$$u''(s^+) - u''(s^-) = 1 \Longrightarrow 2E - 0 = 1$$

 $\Longrightarrow 2E = 1$
 $\Longrightarrow E = \frac{1}{2}$

The Green function now has the form

$$u(x) = \begin{cases} 0 & x < s \\ \frac{1}{2}x^2 + Dx + F & x > s. \end{cases}$$

Together with the need to satisfy the first derivative continuity condition since $u' = \int H(x-s)dx$ is continuous.

$$u'(s^+) = u''(s^-) \Longrightarrow s + D = 0$$

 $\Longrightarrow D = -s$

The Green function now has the form

$$u(x) = \begin{cases} 0 & x < s \\ \frac{1}{2}x^2 - sx + F & x > s. \end{cases}$$

Along with the continuity condition of u since u' is continuous which implies u is continuous.

$$u(s^{+}) = u(s^{-}) \Longrightarrow \frac{1}{2}s^{2} - s^{2} + F = 0$$
$$\Longrightarrow -\frac{1}{2}s^{2} + F = 0$$
$$\Longrightarrow F = \frac{1}{2}s^{2}.$$

Thus the Green function is as follows

$$u(x) = \begin{cases} 0 & x < s \\ \frac{1}{2}(x^2 + s^2) - sx & x > s \end{cases}$$

Part (c) For the case f(x) = x, solve (1) directly by integrating thrice. Verify that (2) gives the same result.

Solution We are given

$$\begin{cases} y'''(x) = x & 0 \le x \le 1\\ y(0) = y'(0) = y''(0) = 0 \end{cases}$$

By integrating f(x) = x thrice,

$$\iiint xdx \Longrightarrow \iint \frac{x^2}{2} + c_1 dx$$

$$\Longrightarrow \int \frac{x^3}{6} + c_1 x + c_2 dx$$

$$\Longrightarrow \frac{x^4}{24} + \frac{c_1 x^2}{2} + c_2 x + c_3 \qquad (c_1, c_2, c_3 \in R)$$

and in order to satisfy the boundary conditions, $c_1=c_2=c_3=0$ since

$$\begin{cases} y(0) &= \frac{0^4}{24} = 0\\ y'(0) &= \frac{0^3}{6} = 0\\ y''(0) &= \frac{0^2}{3} = 0. \end{cases}$$

Thus

$$y = \frac{x^4}{24}.$$

Now to verify green's function with $y(x) = \int_0^1 f(s)G(x;s)ds$, we have

$$G(x;s) = \begin{cases} 0 & x < s \\ \frac{1}{2}(x^2 + s^2) - sx & x > s \end{cases}$$

therefore

$$y(x) = \int_0^x s(\frac{1}{2}(x^2 + s^2) - sx)ds + \int_x^1 0ds$$

$$= \int_0^x \frac{sx^2}{2} + \frac{s^3}{2} - s^2xds$$

$$= \left[\frac{s^2x^2}{4} + \frac{s^4}{8} - \frac{s^3x}{3}\right]_0^x$$

$$= \frac{x^4}{4} + \frac{x^4}{8} - \frac{x^4}{3}$$

$$= \frac{x^4}{24}.$$

We have showed both methods give the same results.

Problem 4: The Fourier transform is often useful for finding Green's functions. Consider the following partial differential equation for the function u(x,t):

$$\begin{cases} u_{tt} - u_{xx} + u = 0 & -\infty < x < \infty \\ u(x,0) = \delta(x-s) & s \in \mathbb{R} \\ u_t(x,0) = 0. \end{cases}$$

Part (a) Apply a Fourier transform with respect to x to obtain a 2nd-order ordinary differential equation for the function $\hat{u}(\omega, t)$.

Solution

By applying the Fourier transform with respect to x such that $\mathcal{F}\{u(x,t)\}=\hat{u}(\omega,t)$ we get

$$\begin{cases} u_{tt} - u_{xx} + u &= 0 \\ u(x,0) &= \delta(x-s) \\ u_t(x,0) &= 0 \end{cases} \xrightarrow{\mathcal{F}} \begin{cases} \hat{u}_{tt} - (i\omega)^2 \hat{u} + \hat{u} &= 0 \\ \hat{u}(\omega,0) &= \frac{1}{2\pi} e^{-i\omega s} \\ \hat{u}_t(\omega,0) &= 0. \end{cases}$$
(1)

which is a 2nd order ordinary differential equation. Second order due to \hat{u}_{tt} and ordinary due to derivatives only involving t.

Part (b) Find $\hat{u}(\omega, t)$ by solving the differential equation from part (a).

Solution From part a we have the following:

$$\begin{cases} \hat{u}_{tt} - (i\omega)^2 \hat{u} + \hat{u} &= 0\\ \hat{u}(\omega, 0) &= \frac{1}{2\pi} e^{-i\omega s}\\ \hat{u}_t(\omega, 0) &= 0. \end{cases}$$

For $\hat{u}_{tt} - (i\omega)^2 \hat{u} + \hat{u} = 0$

$$\hat{u}_{tt} = (i\omega)^2 \hat{u} - \hat{u} \Longrightarrow \hat{u}_{tt} = (\omega^2 - 1)\hat{u}$$

$$\Longrightarrow \hat{u} = c_1 e^{\sqrt{\omega^2 - 1}t} + c_2 e^{-\sqrt{\omega^2 - 1}t} \quad (c_1, c_2 \in \mathbb{R})$$
(1)

⁽¹⁾By solving for second-order linearly ordinary differential equation on Wolfram Alpha

together with $\hat{u}_t(\omega,0)=0$

$$\hat{u}_t(\omega,0) = 0 \Longrightarrow \frac{d}{dt} (c_1 e^{\sqrt{\omega^2 - 1}t} + c_2 e^{-\sqrt{\omega^2 - 1}t}) = 0$$

$$\Longrightarrow c_1 \sqrt{\omega^2 - 1} e^{\sqrt{\omega^2 - 1}(0)} - c_2 \sqrt{\omega^2 - 1} e^{-\sqrt{\omega^2 - 1}(0)} = 0$$

$$\Longrightarrow c_1 \sqrt{\omega^2 - 1} = c_2 \sqrt{\omega^2 - 1}$$

$$\Longrightarrow c_1 = c_2$$

⁽¹⁾By derivative property of Fourier Transforms.

⁽²⁾ By shift property and Fourier Transform of the Dirac Delta function.

along with $\hat{u}(\omega, 0) = \frac{1}{2\pi} e^{-i\omega s}$

$$\implies \frac{1}{2\pi} e^{-i\omega s} \underbrace{\left(c_1 e^{\sqrt{\omega^2 - 1}(0)} + c_1 e^{-\sqrt{\omega^2 - 1}(0)}\right)}_{=1}$$

$$\implies c_1 = \frac{1}{2}$$

therefore

$$\hat{u}(\omega, t) = \frac{1}{4\pi} e^{-i\omega s} \left(e^{\sqrt{\omega^2 - 1}t} + e^{-\sqrt{\omega^2 - 1}t} \right)$$
$$= \left[\frac{1}{2\pi} e^{-i\omega s} \cosh\left(t\sqrt{\omega^2 - 1}\right) \right] \quad \text{since } e^x + e^- x = 2\cosh(x)$$

Part (c) Invert the Fourier transform to obtain u(x,t)

Solution

$$u(x,t) = \int_{-\infty}^{\infty} \hat{u}(\omega,t)e^{i\omega t}d\omega$$

$$= \int_{-\infty}^{\infty} \left[\frac{1}{4\pi} e^{-i\omega s} \left(e^{\sqrt{\omega^2 - 1}t} + e^{-\sqrt{\omega^2 - 1}t} \right) \right] e^{i\omega t}d\omega$$

$$= \frac{1}{4\pi} \int_{-\infty}^{\infty} \left[e^{-i\omega s} \left(e^{\sqrt{\omega^2 - 1}t} + e^{-\sqrt{\omega^2 - 1}t} \right) \right] e^{i\omega t}d\omega$$

$$= \frac{1}{4\pi} \int_{-\infty}^{\infty} \left(e^{\sqrt{\omega^2 - 1}t} + e^{-\sqrt{\omega^2 - 1}t} \right) e^{i\omega(t-s)}d\omega$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} \cosh(t\sqrt{\omega^2 - 1}) e^{i\omega(t-s)}d\omega$$