# Classical Physics Models

### Yingbo Ma, Chris Rackauckas

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If you're getting some cold feet to jump in to DiffEq land, here are some handcrafted differential equations mini problems to hold your hand along the beginning of your journey.

### 0.1 Radioactive Decay of Carbon-14

#### First order linear ODE

$$f(t, u) = \frac{du}{dt}$$

The Radioactive decay problem is the first order linear ODE problem of an exponential with a negative coefficient, which represents the half-life of the process in question. Should the coefficient be positive, this would represent a population growth equation.

```
using OrdinaryDiffEq, Plots
gr()

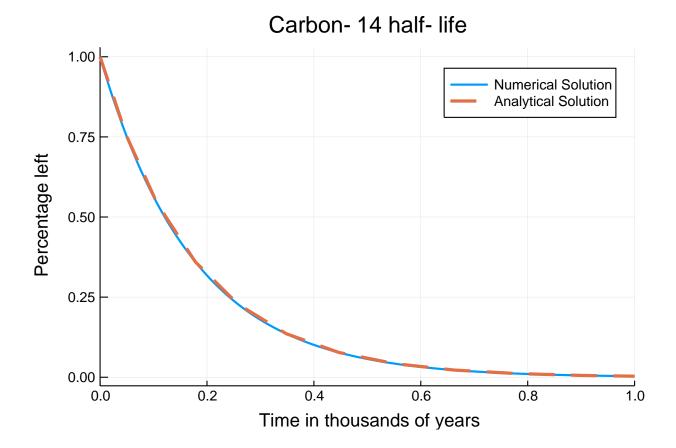
#Half-life of Carbon-14 is 5,730 years.
C_1 = 5.730

#Setup
u_0 = 1.0
tspan = (0.0, 1.0)

#Define the problem
radioactivedecay(u,p,t) = -C_1*u

#Pass to solver
prob = ODEProblem(radioactivedecay,u_0,tspan)
sol = solve(prob,Tsit5())

#Plot
plot(sol,linewidth=2,title ="Carbon-14 half-life", xaxis = "Time in thousands of years", yaxis = "Percentage left", label = "Numerical Solution")
plot!(sol.t, t->exp(-C_1*t),lw=3,ls=:dash,label="Analytical Solution")
```



## 0.2 Simple Pendulum

**Second Order Linear ODE** We will start by solving the pendulum problem. In the physics class, we often solve this problem by small angle approximation, i.e. \$ sin(\theta) \approx \theta\$, because otherwise, we get an elliptic integral which doesn't have an analytic solution. The linearized form is

$$\ddot{\theta} + \frac{g}{L}\theta = 0$$

But we have numerical ODE solvers! Why not solve the *real* pendulum?

$$\ddot{\theta} + \frac{g}{L}\sin(\theta) = 0$$

```
# Simple Pendulum Problem
using OrdinaryDiffEq, Plots

#Constants
const g = 9.81
L = 1.0

#Initial Conditions
u_0 = [0, \pi/2]
tspan = (0.0,6.3)

#Define the problem
function simplependulum(du,u,p,t)
```

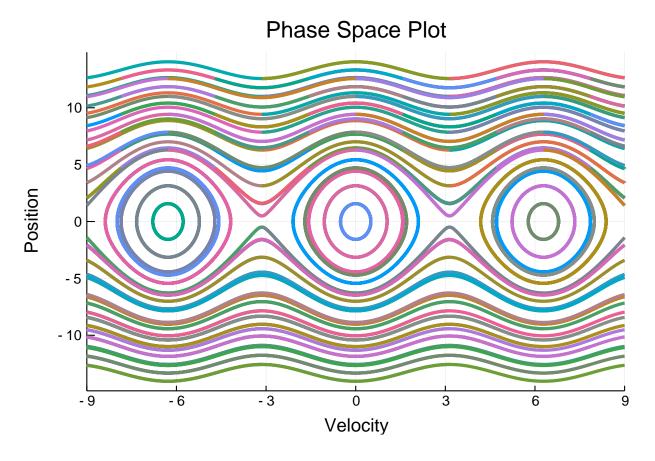
## Simple Pendulum Problem 1.5 Theta dTheta 1.0 0.5 Height 0.0 - 0.5 - 1.0 - 1.5 2 0 3 4 5 6

So now we know that behaviour of the position versus time. However, it will be useful to us to look at the phase space of the pendulum, i.e., and representation of all possible states of the system in question (the pendulum) by looking at its velocity and position. Phase space analysis is ubiquitous in the analysis of dynamical systems, and thus we will provide a few facilities for it.

Time

```
p = plot(sol,vars = (1,2), xlims = (-9,9), title = "Phase Space
   Plot", xaxis = "Velocity", yaxis = "Position", leg=false)
function phase_plot(prob, u0, p, tspan=2pi)
   _prob = ODEProblem(prob.f,u0,(0.0,tspan))
   sol = solve(_prob,Vern9()) # Use Vern9 solver for higher accuracy
   plot!(p,sol,vars = (1,2), xlims = nothing, ylims = nothing)
end
for i in -4pi:pi/2:4π
   for j in -4pi:pi/2:4π
        phase_plot(prob, [j,i], p)
```

```
end
end
plot(p,xlims = (-9,9))
```



## 0.3 Simple Harmonic Oscillator

#### 0.3.1 Double Pendulum

```
#Double Pendulum Problem
using OrdinaryDiffEq, Plots
#Constants and setup
const m_1, m_2, L_1, L_2 = 1, 2, 1, 2
initial = [0, \pi/3, 0, 3pi/5]
tspan = (0.,50.)
#Convenience function for transforming from polar to Cartesian coordinates
function polar2cart(sol;dt=0.02,l1=L_1,l2=L_2,vars=(2,4))
   u = sol.t[1]:dt:sol.t[end]
   p1 = 11*map(x->x[vars[1]], sol.(u))
   p2 = 12*map(y->y[vars[2]], sol.(u))
   x1 = 11*sin.(p1)
   y1 = 11*-cos.(p1)
    (u, (x1 + 12*sin.(p2),
    y1 - 12*cos.(p2)))
end
#Define the Problem
```

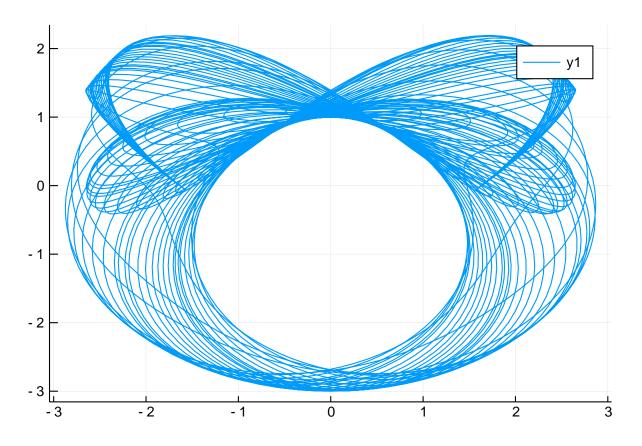
```
function double_pendulum(xdot,x,p,t)
    xdot[1]=x[2]

xdot[2]=-((g*(2*m_1+m_2)*sin(x[1])+m_2*(g*sin(x[1]-2*x[3])+2*(L_2*x[4]^2+L_1*x[2]^2*cos(x[1]-x[3])
    xdot[3]=x[4]

xdot[4]=(((m_1+m_2)*(L_1*x[2]^2+g*cos(x[1]))+L_2*m_2*x[4]^2*cos(x[1]-x[3]))*sin(x[1]-x[3]))/(L_2*(end)

#Pass to Solvers
double_pendulum_problem = ODEProblem(double_pendulum, initial, tspan)
sol = solve(double_pendulum_problem, Vern7(), abs_tol=1e-10, dt=0.05);
```

```
#Obtain coordinates in Cartesian Geometry
ts, ps = polar2cart(sol, 11=L_1, 12=L_2, dt=0.01)
plot(ps...)
```



#### 0.3.2 Poincaré section

The Poincaré section is a contour plot of a higher-dimensional phase space diagram. It helps to understand the dynamic interactions and is wonderfully pretty.

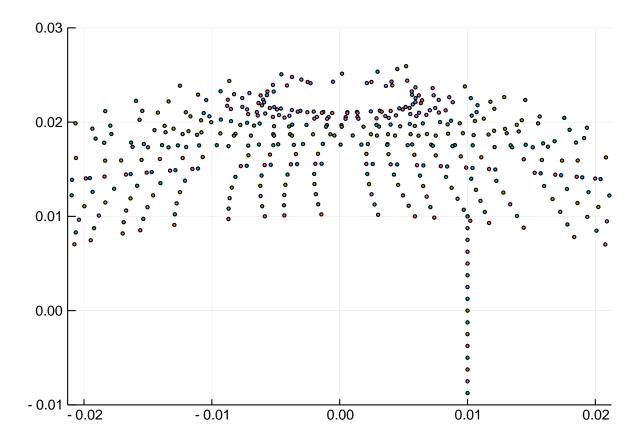
The following equation came from StackOverflow question

$$\frac{d}{dt} \begin{pmatrix} \alpha \\ l_{\alpha} \\ \beta \\ l_{\beta} \end{pmatrix} = \begin{pmatrix} 2\frac{l_{\alpha} - (1+\cos\beta)l_{\beta}}{3-\cos2\beta} \\ -2\sin\alpha - \sin(\alpha+\beta) \\ 2\frac{-(1+\cos\beta)l_{\alpha} + (3+2\cos\beta)l_{\beta}}{3-\cos2\beta} \\ -\sin(\alpha+\beta) - 2\sin(\beta)\frac{(l_{\alpha} - l_{\beta})l_{\beta}}{3-\cos2\beta} + 2\sin(2\beta)\frac{l_{\alpha}^{2} - 2(1+\cos\beta)l_{\alpha}l_{\beta} + (3+2\cos\beta)l_{\beta}^{2}}{(3-\cos2\beta)^{2}} \end{pmatrix}$$

The Poincaré section here is the collection of  $(\beta, l_{\beta})$  when  $\alpha = 0$  and  $\frac{d\alpha}{dt} > 0$ .

**Hamiltonian of a double pendulum** Now we will plot the Hamiltonian of a double pendulum

```
#Constants and setup
 using OrdinaryDiffEq
 initial2 = [0.01, 0.005, 0.01, 0.01]
 tspan2 = (0.,200.)
 #Define the problem
 function double_pendulum_hamiltonian(udot,u,p,t)
                \alpha = u[1]
                1\alpha = u[2]
                \beta = u[3]
                1\beta = u[4]
                udot .=
                 [2(1\alpha - (1+\cos(\beta)))1\beta)/(3-\cos(2\beta)),
                -2\sin(\alpha) - \sin(\alpha+\beta),
                2(-(1+\cos(\beta))1\alpha + (3+2\cos(\beta))1\beta)/(3-\cos(2\beta)),
                -\sin(\alpha+\beta) - 2\sin(\beta)*(((1\alpha-1\beta)1\beta)/(3-\cos(2\beta))) + 2\sin(2\beta)*((1\alpha^2 - 2(1+\cos(\beta))1\alpha*1\beta)) + 2\cos(2\beta)*((1\alpha^2 - 2(1+\cos(\beta))1\alpha*1\beta)) + 2\cos(2\beta)*
                + (3+2\cos(\beta))1\beta^2/(3-\cos(2\beta))^2
 end
 # Construct a ContinuousCallback
 condition(u,t,integrator) = u[1]
 affect!(integrator) = nothing
 cb = ContinuousCallback(condition,affect!,nothing,
                                                                                                 save_positions = (true, false))
 # Construct Problem
 poincare = ODEProblem(double_pendulum_hamiltonian, initial2, tspan2)
 sol2 = solve(poincare, Vern9(), save_everystep = false, callback=cb, abstol=1e-9)
 function poincare_map(prob, u_0, p; callback=cb)
                _prob = ODEProblem(prob.f, [0.01, 0.01, 0.01, u_0], prob.tspan)
                sol = solve(_prob, Vern9(), save_everystep = false, callback=cb, abstol=1e-9)
                scatter!(p, sol, vars=(3,4), markersize = 2)
 end
poincare map (generic function with 1 method)
p = scatter(sol2, vars=(3,4), leg=false, markersize = 2, ylims=(-0.01,0.03))
 for i in -0.01:0.00125:0.01
                poincare_map(poincare, i, p)
plot(p,ylims=(-0.01,0.03))
```



#### Hénon-Heiles System 0.4

The Hénon-Heiles potential occurs when non-linear motion of a star around a galactic center with the motion restricted to a plane.

$$\frac{d^2x}{dt^2} = -\frac{\partial V}{\partial x} \tag{1}$$

$$\frac{d^2x}{dt^2} = -\frac{\partial V}{\partial x}$$

$$\frac{d^2y}{dt^2} = -\frac{\partial V}{\partial y}$$
(1)

where

$$V(x,y) = \frac{1}{2}(x^2 + y^2) + \lambda \left(x^2y - \frac{y^3}{3}\right).$$

We pick  $\lambda = 1$  in this case, so

$$V(x,y) = \frac{1}{2}(x^2 + y^2 + 2x^2y - \frac{2}{3}y^3).$$

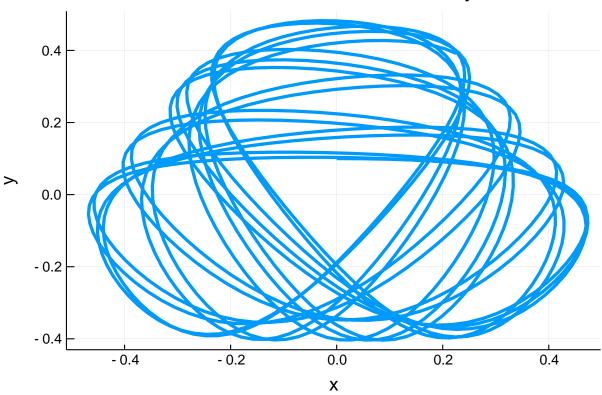
Then the total energy of the system can be expressed by

$$E = T + V = V(x, y) + \frac{1}{2}(\dot{x}^2 + \dot{y}^2).$$

The total energy should conserve as this system evolves.

```
using OrdinaryDiffEq, Plots
#Setup
initial = [0.,0.1,0.5,0]
tspan = (0,100.)
\#Remember, V is the potential of the system and T is the Total Kinetic Energy, thus E
#the total energy of the system.
V(x,y) = 1//2 * (x^2 + y^2 + 2x^2*y - 2//3 * y^3)
E(x,y,dx,dy) = V(x,y) + 1//2 * (dx^2 + dy^2);
#Define the function
function Hénon_Heiles(du,u,p,t)
   x = u[1]
   y = u[2]
   dx = u[3]
   dy = u[4]
   du[1] = dx
   du[2] = dy
   du[3] = -x - 2x*y
   du[4] = y^2 - y - x^2
end
#Pass to solvers
prob = ODEProblem(Hénon_Heiles, initial, tspan)
sol = solve(prob, Vern9(), abs_tol=1e-16, rel_tol=1e-16);
# Plot the orbit
plot(sol, vars=(1,2), title = "The orbit of the Hénon-Heiles
   system", xaxis = "x", yaxis = "y", leg=false)
```

# The orbit of the Hénon-Heiles system

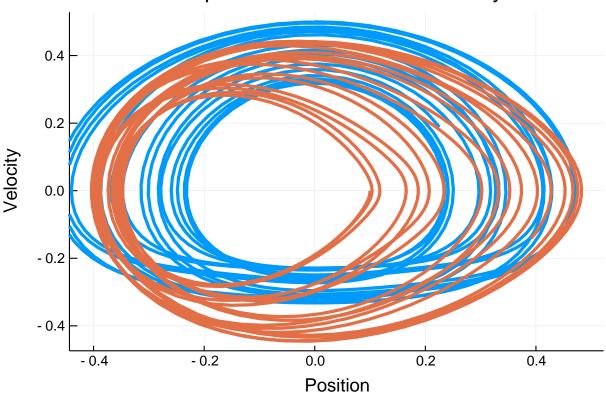


```
#Optional Sanity check - what do you think this returns and why?

@show sol.retcode
```

#### sol.retcode = :Success

# Phase space for the Hénon-Heiles system



#We map the Total energies during the time intervals of the solution (sol.u here) to a new vector

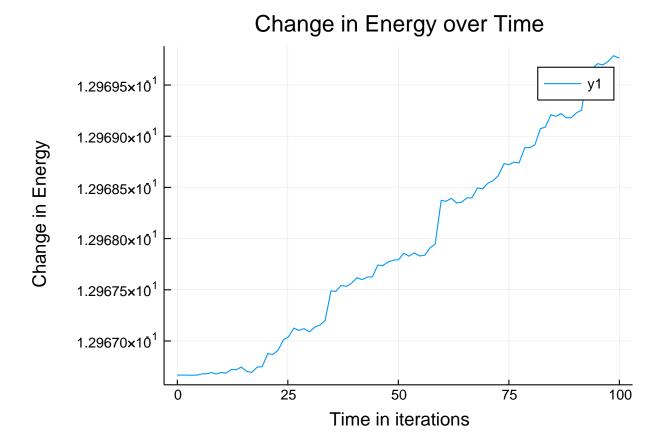
#pass it to the plotter a bit more conveniently energy =  $map(x\rightarrow E(x...), sol.u)$ 

#We use Oshow here to easily spot erratic behaviour in our system by seeing if the loss in energy was too great.

Oshow  $\Delta E = \text{energy}[1] - \text{energy}[\text{end}]$ 

```
\Delta E = \text{energy}[1] - \text{energy}[\text{end}] = -3.099845153070602e-5
```

```
#Plot
plot(sol.t, energy, title = "Change in Energy over Time", xaxis = "Time in
   iterations", yaxis = "Change in Energy")
```



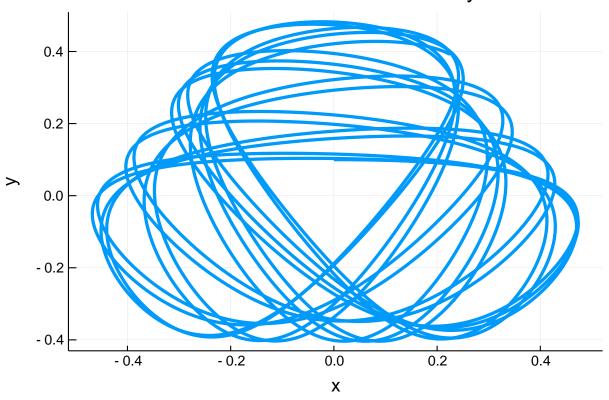
#### 0.4.1 Symplectic Integration

To prevent energy drift, we can instead use a symplectic integrator. We can directly define and solve the SecondOrderODEProblem:

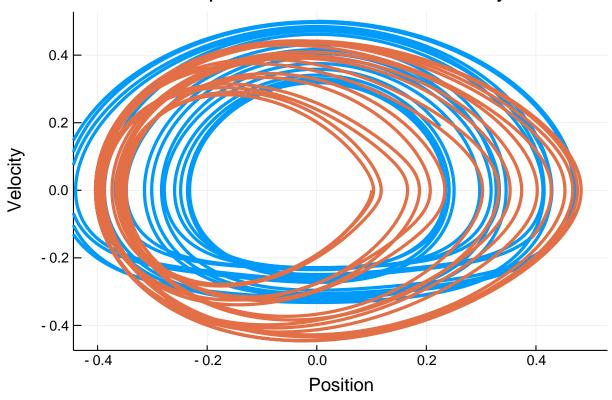
```
function HH_acceleration!(dv,v,u,p,t)
    x,y = u
    dx,dy = dv
    dv[1] = -x - 2x*y
    dv[2] = y^2 - y -x^2
end
initial_positions = [0.0,0.1]
initial_velocities = [0.5,0.0]
prob = SecondOrderODEProblem(HH_acceleration!,initial_velocities,initial_positions,tspan)
sol2 = solve(prob, KahanLi8(), dt=1/10);
```

Notice that we get the same results:

# The orbit of the Hénon-Heiles system



# Phase space for the Hénon-Heiles system



but now the energy change is essentially zero:

```
energy = map(x->E(x[3], x[4], x[1], x[2]), sol2.u)
#We use @show here to easily spot erratic behaviour in our system by seeing if the loss
    in energy was too great.
@show \Delta E = energy[1]-energy[end]

\Delta E = energy[1] - energy[end] = 9.020562075079397e-15

#Plot
plot(sol2.t, energy, title = "Change in Energy over Time", xaxis = "Time in
    iterations", yaxis = "Change in Energy")
```

## Change in Energy over Time

\_\_\_\_ y1

Change in Energy

0

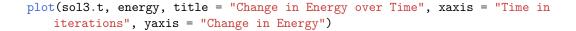
#### Time in iterations

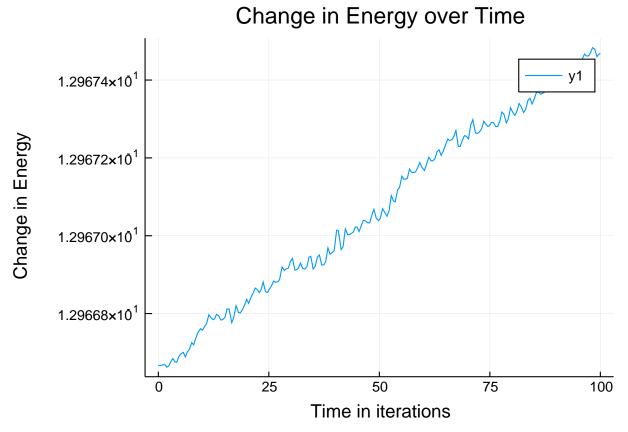
It's so close to zero it breaks GR! And let's try to use a Runge-Kutta-Nyström solver to solve this. Note that Runge-Kutta-Nyström isn't symplectic.

```
sol3 = solve(prob, DPRKN6());
energy = map(x->E(x[3], x[4], x[1], x[2]), sol3.u)
@show ΔE = energy[1]-energy[end]

ΔE = energy[1] - energy[end] = -8.017994408110463e-6

gr()
```





Note that we are using the DPRKN6 sovler at reltol=1e-3 (the default), yet it has a smaller energy variation than Vern9 at abs\_tol=1e-16, rel\_tol=1e-16. Therefore, using specialized solvers to solve its particular problem is very efficient.

```
using DiffEqTutorials
DiffEqTutorials.tutorial_footer(WEAVE_ARGS[:folder],WEAVE_ARGS[:file])
```

## 0.5 Appendix

These benchmarks are part of the DiffEqTutorials.jl repository, found at: https://github.com/JuliaDiffEq/Tutorials.jl repository, found at: https://github.com/Tutorials.jl reposito

```
using DiffEqTutorials
DiffEqTutorials.weave file("models","classical physics.jmd")
```

Computer Information:

```
Julia Version 1.1.0
Commit 80516ca202 (2019-01-21 21:24 UTC)
Platform Info:
    OS: Windows (x86_64-w64-mingw32)
    CPU: Intel(R) Core(TM) i7-8700 CPU @ 3.20GHz
```

WORD\_SIZE: 64 LIBM: libopenlibm

LLVM: libLLVM-6.0.1 (ORCJIT, skylake)

Environment:

JULIA\_EDITOR = "C:\Users\accou\AppData\Local\atom\app-1.34.0\atom.exe" -a

JULIA\_NUM\_THREADS = 6

Package Information:

Status `C:\Users\accou\.julia\external\DiffEqTutorials.jl\src\Project.toml`