

## 1 2-D Special Functions

- Function of two independent spatial variables specifies amplitude (“brightness”) at each spatial coordinate in a plane
  - fulfill usual definition of “image”.
- Three categories:
  1. Cartesian separable functions
    - represented as products of 1-D special functions along orthogonal Cartesian axes
  2. Circularly symmetric functions
    - product of 1-D functions in radial direction and unit constant in orthogonal (“azimuthal”) direction
  3. General functions
    - includes pictorial scenes and 2-D stochastic functions.

## 2 2-D Separable Functions

- “orthogonal multiplication” of two 1-D functions  $f_x[x]$  and  $f_y[y]$ :

$$f[x, y] = f_x[x] \times f_y[y]$$

- General expression for separable function in terms of scaled and translated 1-D functions is:

$$f[x, y] = f_x\left[\frac{x - x_0}{a}\right] \times f_y\left[\frac{y - y_0}{b}\right]$$

- $a$  and  $b$  are real-valued scale factors
  - $x_0$  and  $y_0$  are real-valued translation parameters
- Consider those which have general application to imaging problems.
- Volume of 2-D separable function is product of areas of component 1-D functions:

$$\begin{aligned} \iint_{-\infty}^{+\infty} f[x, y] \, dx \, dy &= \iint_{-\infty}^{+\infty} f_x\left[\frac{x - x_0}{a}\right] f_y\left[\frac{y - y_0}{b}\right] \, dx \, dy \\ &= \left(\int_{-\infty}^{+\infty} f_x\left[\frac{x - x_0}{a}\right] \, dx\right) \left(\int_{-\infty}^{+\infty} f_y\left[\frac{y - y_0}{b}\right] \, dy\right) \end{aligned}$$

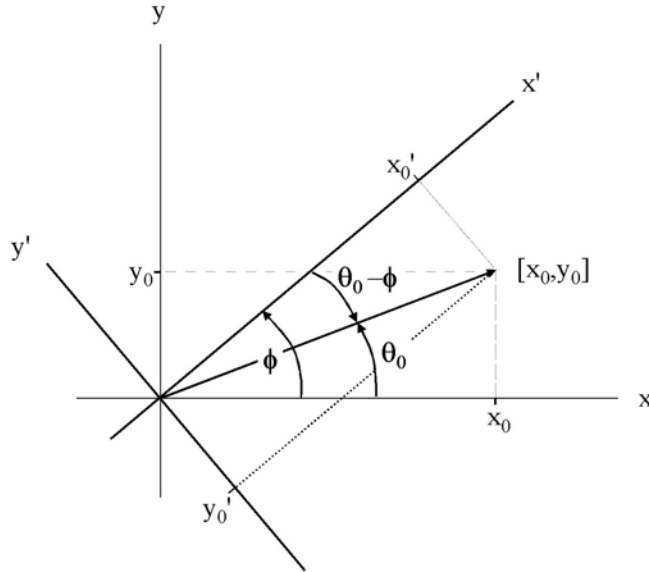
## 2.1 Rotations of 2-D Separable Functions

- Example: square centered at origin with sides parallel to the  $x$ - and  $y$ -axes rotated by  $\pm \frac{\pi}{4}$  to generate “baseball diamond” with vertices on  $x$ - and  $y$ -axes
- Rotation of function about origin is an “imaging system” with 2-D input function  $f[x, y]$  and 2-D output  $g[x, y]$ 
  - Amplitude  $g$  of rotated function at  $[x, y]$  is original amplitude  $f$  at location  $[x', y']$ :

$$\mathcal{O}\{f[x, y]\} = g[x, y] = f[x', y']$$

- Rotation specified completely by mapping that relates coordinates  $[x, y]$  and  $[x', y']$ .
  - Illustrated by considering location with Cartesian coordinates  $[x_0, y_0]$  and polar coordinates  $(r_0, \theta_0)$ 
    - \*  $r_0 = \sqrt{x_0^2 + y_0^2}$
    - \*  $\theta_0 = \tan^{-1} \left[ \frac{y_0}{x_0} \right]$ .
    - \* Radial coordinate  $r_0$  unchanged by rotation
    - \* Azimuthal angle becomes  $\theta' = \theta_0 - \phi$ .
    - \* Cartesian coordinates of original location in rotated coordinates  $[x'_0, y'_0]$  in terms of original coordinates  $[x_0, y_0]$ :

$$\begin{aligned} x'_0 &= |\mathbf{r}_0| \cos[\theta_0 - \phi] = |\mathbf{r}_0| (\cos[\theta_0] \cos[\phi] + \sin[\theta_0] \sin[\phi]) \\ &= x_0 \cos[\phi] + y_0 \sin[\phi] \\ y'_0 &= r_0 \sin[\theta_0 - \phi] = r_0 (\sin[\theta_0] \cos[\phi] - \cos[\theta_0] \sin[\phi]) \\ &= -x_0 \sin[\phi] + y_0 \cos[\phi] \end{aligned}$$



The Cartesian coordinates of a particular location  $[x_0, y_0]$  is evaluated in the rotated coordinate system  $[x', y']$  from the polar coordinates  $(|\mathbf{r}_0|, \theta_0 - \phi)$ :

$$x'_0 = |\mathbf{r}_0| \cos[\theta_0 - \phi] = x_0 \cos[\phi] + y_0 \sin[\phi], \quad y'_0 = |\mathbf{r}_0| \sin[\theta_0 - \phi] = -x_0 \sin[\phi] + y_0 \cos[\phi]$$

## 2.2 Rotated Coordinates as Scalar Products

- Rotated  $x$ -coordinate written as scalar product of position vector  $\underline{\mathbf{r}} \equiv [x, y]$  and unit vector directed along azimuth angle  $\phi$ , which has Cartesian coordinates  $[\cos[\phi], \sin[\phi]]$ , denoted by  $\underline{\hat{\mathbf{p}}}$ :

$$x' = x \cos[\phi] + y \sin[\phi] = \begin{bmatrix} x \\ y \end{bmatrix} \bullet \begin{bmatrix} \cos[\phi] \\ \sin[\phi] \end{bmatrix} \equiv \underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}$$

- Notation for  $x'$  may seem “weird” because rotated 1-D argument  $x'$  is function of both  $x$  and  $y$  through scalar product of two vectors.
  - rotated argument  $x'$  defines a set of points  $\underline{\mathbf{r}} = [x, y]$  that fulfill same conditions as coordinate  $x$  in original function.
  - Rotated coordinate  $x'$  evaluates to 0 for all vectors  $\underline{\mathbf{r}} \perp \underline{\hat{\mathbf{p}}}$ , (equivalent to  $\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}} = 0$ ).
  - Vectors  $\underline{\mathbf{r}}$  specified by condition  $\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}} = 0$  include coordinates on “both” sides of origin
  - Complete set of possible azimuth angles specified by angles  $\phi$  in interval spanning  $\pi$  radians, e.g.,  $-\frac{\pi}{2} \leq \phi < +\frac{\pi}{2}$ . This means that the set of possible rotations is specified by unit vectors  $\underline{\hat{\mathbf{p}}}$  in the first and fourth quadrants.
- Polar form of position vector  $\underline{\mathbf{r}} = (|\underline{\mathbf{r}}|, \theta) = (r, \theta)$  substituted into scalar product in terms of magnitudes of vectors and included angle:

$$\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}} = |\underline{\mathbf{r}}| |\underline{\hat{\mathbf{p}}}| \cos[\theta - \phi] = r \cos[\theta - \phi] \text{ because } |\underline{\hat{\mathbf{p}}}| = 1.$$

- $y$ -coordinate of rotated function written in same way as scalar product of position vector  $\underline{\mathbf{r}}$  and unit vector directed along azimuth angle  $\phi + \frac{\pi}{2}$ ; call it  $\underline{\hat{\mathbf{p}}}^\perp$ :

$$\begin{aligned} y' &= x \cos\left[\phi + \frac{\pi}{2}\right] + y \sin\left[\phi + \frac{\pi}{2}\right] \\ &= -x \sin[\phi] + y \cos[\phi] = \begin{bmatrix} x \\ y \end{bmatrix} \bullet \begin{bmatrix} -\sin[\phi] \\ \cos[\phi] \end{bmatrix} \equiv \underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}^\perp \end{aligned}$$

- Polar form for rotated  $y$ -axis is:

$$\begin{aligned} \underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}^\perp &= |\underline{\mathbf{r}}| |\underline{\hat{\mathbf{p}}}^\perp| \cos\left[\theta - \left(\phi + \frac{\pi}{2}\right)\right] \\ &= r \cos\left[\theta - \phi - \frac{\pi}{2}\right] = r \sin[\theta - \phi] \end{aligned}$$

$$g[x, y] = f_x[x'] f_y[y'] = f_x[\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}] f_y[\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}^\perp]$$

- Formula is applicable to any 2-D separable special function

### 3 Definitions of 2-D Separable Functions

#### 3.1 2-D Constant

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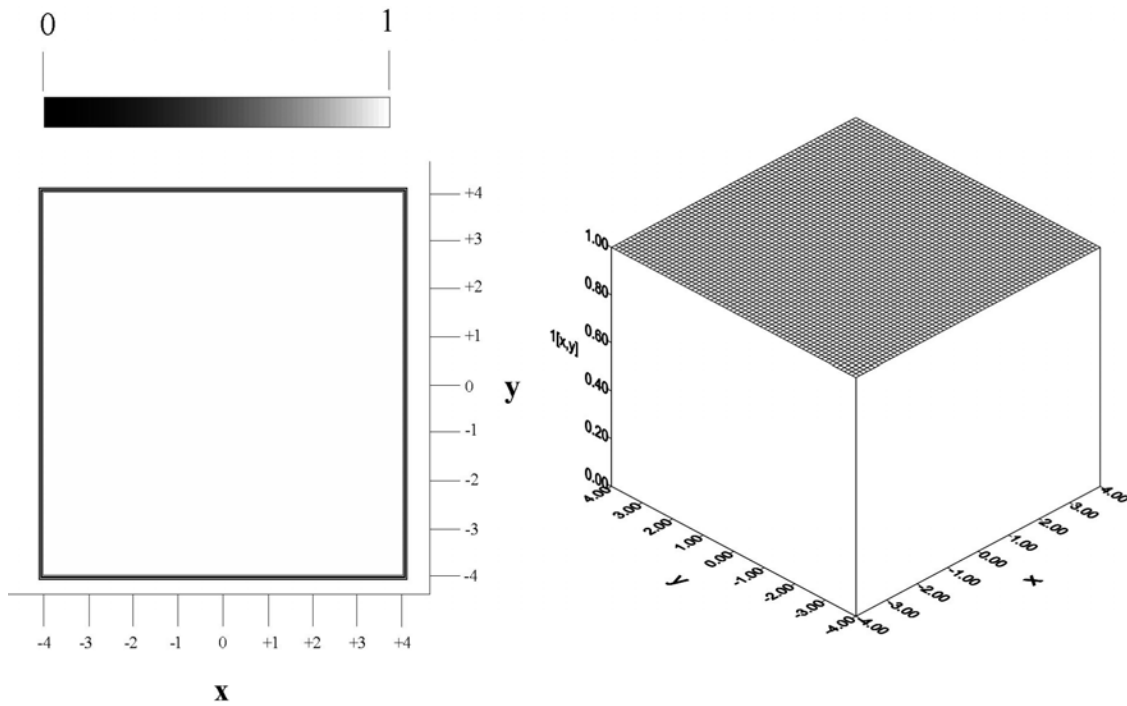
$$1[x, y] = 1[x] \cdot 1[y]$$

$$\iint_{-\infty}^{+\infty} 1[x, y] \, dx \, dy = \infty$$

$$0[x, y] = 0[x] \cdot f_y[y]$$

$$\iint_{-\infty}^{+\infty} 0[x, y] \, dx \, dy = 0$$

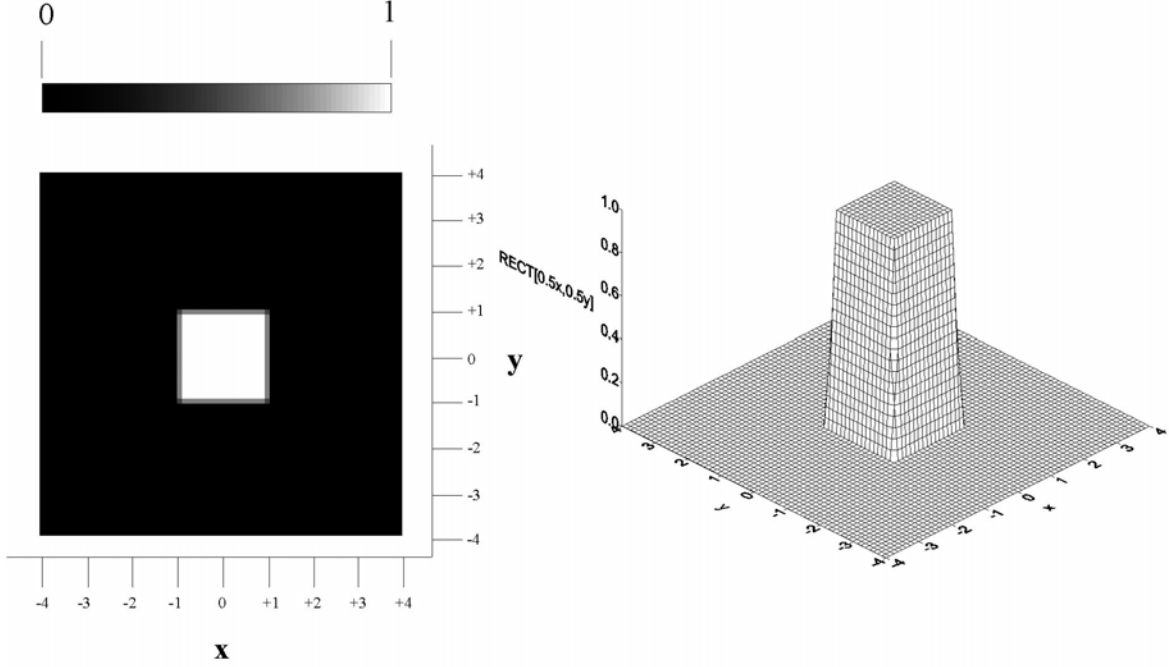
- Neither translations nor rotations affect amplitude of a 2-D constant function at any coordinate.



Representations of the 2-D unit constant function in “image” format, where the amplitude is represented by shades of gray according to the scale, and in “surface” format, where the amplitude is the “height” of a 3-D surface.

### 3.2 2-D Rectangle

$$\begin{aligned} \text{RECT} \left[ \frac{x}{a}, \frac{y}{b} \right] &= \text{RECT} \left[ \frac{x}{a} \right] \text{RECT} \left[ \frac{y}{b} \right] \\ \text{RECT} \left[ \frac{x-x_0}{a}, \frac{y-y_0}{b} \right] &= \text{RECT} \left[ \frac{x-x_0}{a} \right] \text{RECT} \left[ \frac{y-y_0}{b} \right] \end{aligned}$$



Example of the 2-D rectangle function  $f[x, y] = \text{RECT} \left[ \frac{x}{2}, \frac{y}{2} \right]$ .

- Volume obviously is finite when both  $a$  and  $b$  are finite:

$$\begin{aligned} &\iint_{-\infty}^{+\infty} \text{RECT} \left[ \frac{x-x_0}{a}, \frac{y-y_0}{b} \right] dx dy \\ &= \int_{-\infty}^{+\infty} \text{RECT} \left[ \frac{x-x_0}{a} \right] dx \int_{-\infty}^{+\infty} \text{RECT} \left[ \frac{y-y_0}{b} \right] dy = |ab| \end{aligned}$$

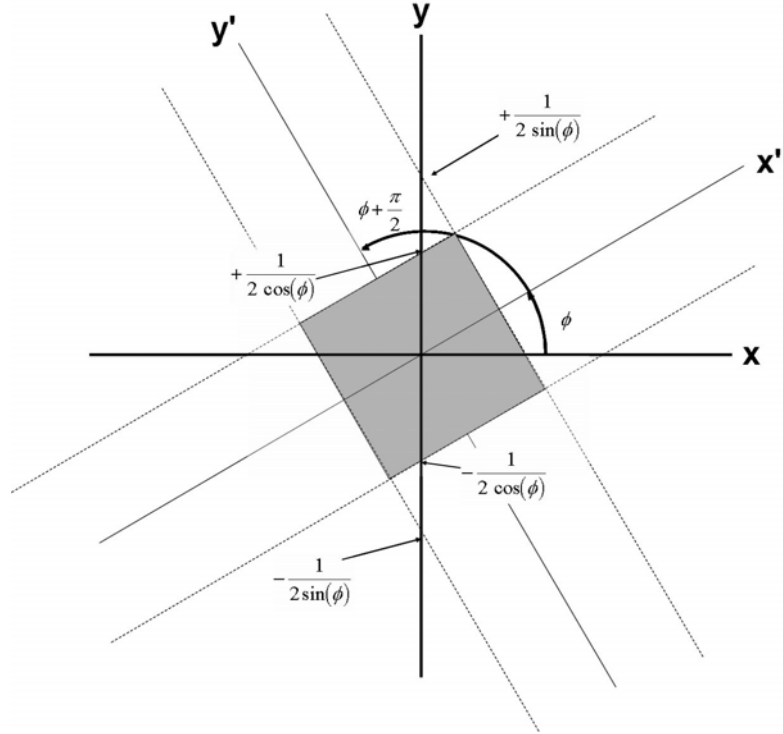
- 2-D  $\text{RECT}$  often modulates functions that have larger domains of support.
- Rotated rectangle by substituting rotated coordinates into definition

$$\begin{aligned} \text{RECT}[x', y'] &= \text{RECT}[\mathbf{r} \bullet \hat{\mathbf{p}}, \mathbf{r} \bullet \hat{\mathbf{p}}^\perp] = \text{RECT}[\mathbf{r} \bullet \hat{\mathbf{p}}] \text{RECT}[\mathbf{r} \bullet \hat{\mathbf{p}}^\perp] \\ &= \text{RECT}[x \cos[\phi] + y \sin[\phi]] \text{RECT}[-x \sin[\phi] + y \cos[\phi]] \end{aligned}$$

- “Transition” locations  $[x, y]$  where  $\text{RECT}[x, y] = \frac{1}{2}$  by substitution:

$$\begin{aligned} x' &= x \cos[\phi] + y \sin[\phi] = \pm \frac{1}{2} \implies y = (-\cot[\phi]) x \pm \frac{1}{2 \sin[\phi]} \\ y' &= -x \sin[\phi] + y \cos[\phi] = \pm \frac{1}{2} \implies y = (\tan[\phi]) x \pm \frac{1}{2 \cos[\phi]} \end{aligned}$$

- Conditions define slopes and  $y$ -intercepts of two pairs of parallel lines in 2-D plane
- slopes  $+\tan[\phi]$  and  $-\cot[\phi] = \tan[\phi + \frac{\pi}{2}]$
- Dashed lines enclose support of rectangle



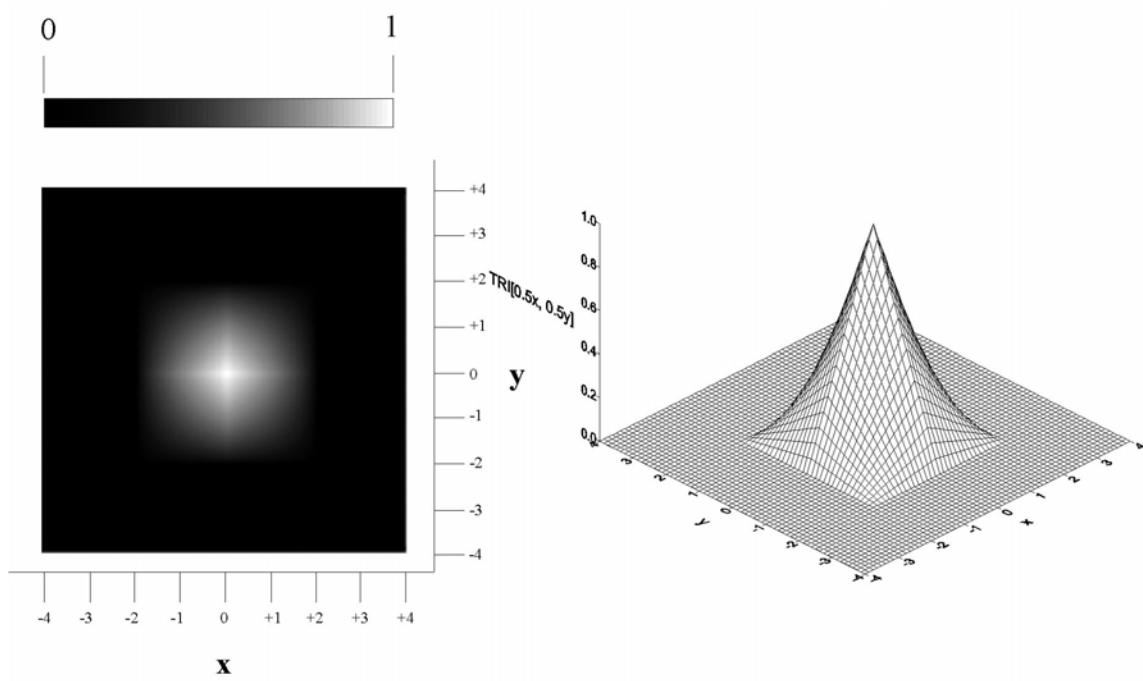
The loci of points that satisfy the conditions  $y = x(-\cot[\phi]) \pm \frac{1}{2\sin[\phi]}$  and  $y = x\tan[\phi] \pm \frac{1}{2\cos[\phi]}$ .

### 3.3 2-D Triangle

$$\begin{aligned}
 TRI \left[ \frac{x}{a}, \frac{y}{b} \right] &= TRI \left[ \frac{x}{a} \right] TRI \left[ \frac{y}{b} \right] \\
 &= \left( 1 - \frac{|x|}{a} \right) \left( 1 - \frac{|y|}{b} \right) RECT \left[ \frac{x}{2a}, \frac{y}{2b} \right] \\
 &= \left( 1 - \frac{|x|}{a} - \frac{|y|}{b} + \frac{|xy|}{ab} \right) RECT \left[ \frac{x}{2a}, \frac{y}{2b} \right]
 \end{aligned}$$

- Profiles of 2-D *TRI* function are straight lines only parallel to  $x$ - or  $y$ -axes.
- Shape of edge profile along other radial line includes quadratic function of radial distance  $\Rightarrow$  parabola.
- Volume of 2-D *TRI* is product of 1-D areas:

$$\iint_{-\infty}^{+\infty} TRI \left[ \frac{x}{a}, \frac{y}{b} \right] dx dy = \left( \int_{-\infty}^{+\infty} TRI \left[ \frac{x}{b} \right] dx \right) \left( \int_{-\infty}^{+\infty} TRI \left[ \frac{y}{b} \right] dy \right) = |ab|$$



2-D triangle function  $TRI[x, y]$ , showing the parabolic character along profiles that do not coincide with the coordinate axes.

### 3.4 2-D Signum and Step

- Encountered rarely
- Possible definition based on “convention”:

$$SGN[x, y] = SGN[x] \cdot SGN[y]$$

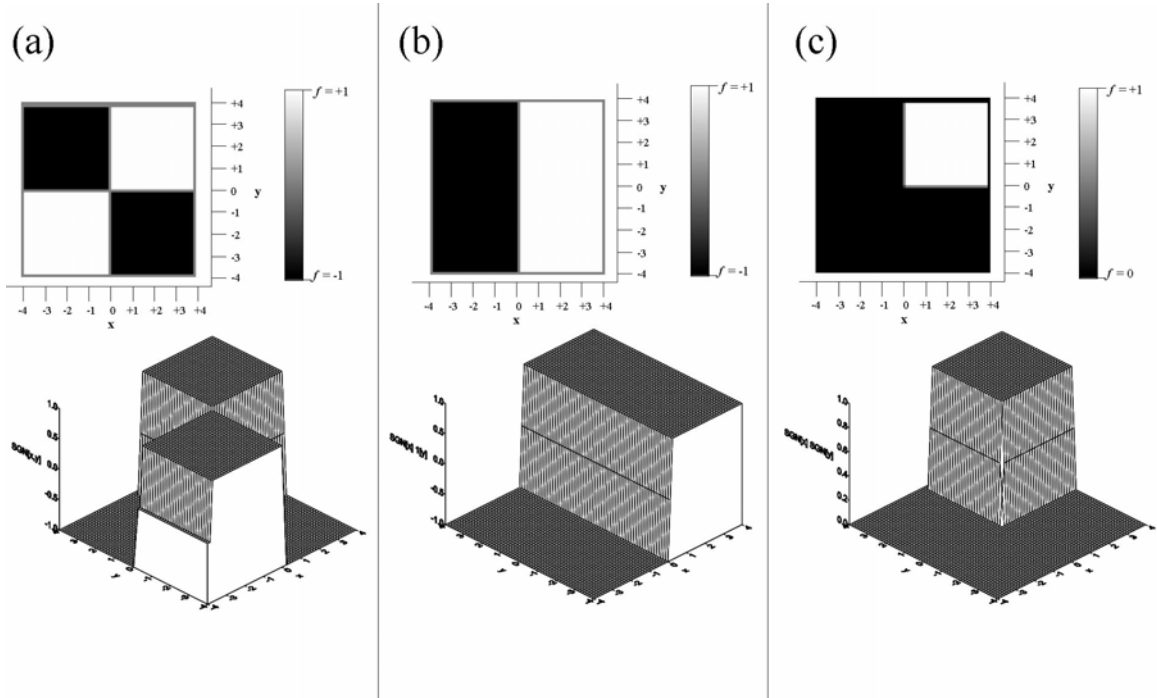
- +1 in first & third quadrants
- -1 in second & fourth quadrants
- 0 along coordinate axes

- Volume is zero due to cancellation of positive and negative regions.
- Another possible definition:

$$f[x, y] = STEP[x] \cdot 1[y]$$

$$f[x, y] = STEP[\mathbf{r} \cdot \hat{\mathbf{p}}] \cdot 1[\mathbf{r} \cdot \hat{\mathbf{p}}^\perp]$$

- Notation for unit constant may seem weird
  - think of 1-D unit constant as having same unit amplitude everywhere in the 1-D domain
  - when applied in 2-D domain, 1-D unit constant takes amplitude of  $STEP[\mathbf{r} \cdot \hat{\mathbf{p}}]$  at  $\mathbf{r}_0$  along direction of  $\hat{\mathbf{p}}$  and assigns it to all points  $\mathbf{r}$  along line perpendicular to  $\hat{\mathbf{p}}$ , i.e., in direction of  $\hat{\mathbf{p}}^\perp$
  - Action of unit constant is to “spread” or “smear” amplitude of 1-D  $STEP$  along perpendicular direction
  - use “rotated” form for 1-D unit constant only to produce 2-D functions by orthogonal
- volume =  $\infty$ .



2-D  $SGN$  and  $STEP$  functions: (a)  $SGN[x] \cdot SGN[y]$ , (b)  $SGN[x] \cdot 1[y]$ , (c)  $STEP[x] \cdot STEP[y]$ .

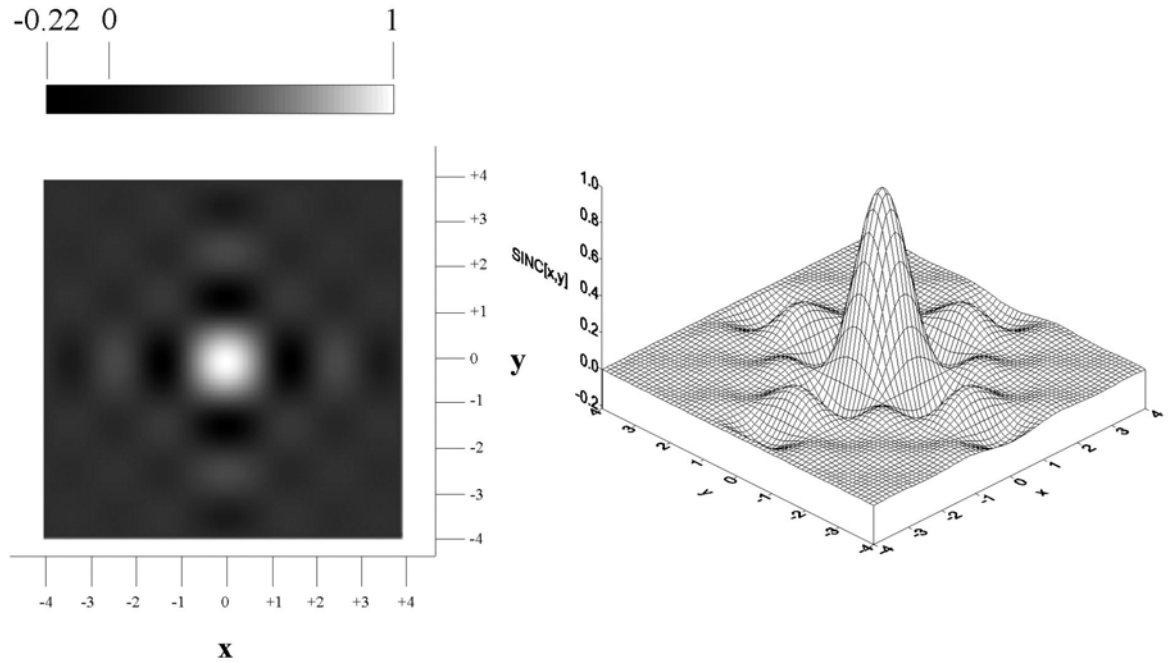


### 3.5 2-D *SINC*

$$SINC\left[\frac{x}{a}, \frac{y}{b}\right] = SINC\left[\frac{x}{a}\right] SINC\left[\frac{y}{b}\right] = \frac{\sin\left[\frac{\pi x}{a}\right]}{\left(\frac{\pi x}{a}\right)} \frac{\sin\left[\frac{\pi y}{b}\right]}{\left(\frac{\pi y}{b}\right)}$$

- Amplitude > 0 in regions where both of 1-D functions are positive or both are negative
- Amplitude < 0 where either (but not both) is negative.
- “checkerboard-like” pattern of positive and negative regions
- volume is product of the areas of the individual functions:

$$\begin{aligned} \iint_{-\infty}^{+\infty} SINC\left[\frac{x}{a}, \frac{y}{b}\right] dx dy &= \left(\int_{-\infty}^{+\infty} SINC\left[\frac{x}{a}\right] dx\right) \left(\int_{-\infty}^{+\infty} SINC\left[\frac{y}{b}\right] dy\right) \\ &= |a| |b| = |ab| \end{aligned}$$



$$SINC[x, y] = SINC[x] \cdot SINC[y]$$

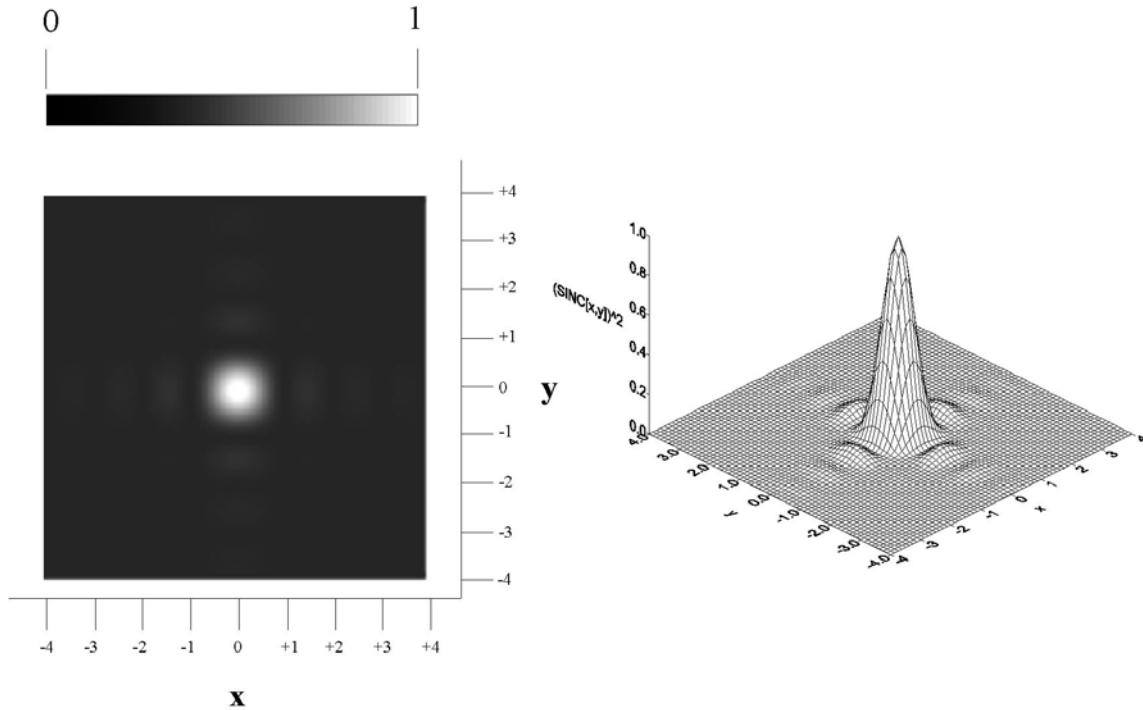
### 3.6 2-D $SINC^2$

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$$\begin{aligned} SINC^2 \left[ \frac{x}{a}, \frac{y}{b} \right] &= SINC^2 \left[ \frac{x}{a} \right] SINC^2 \left[ \frac{y}{b} \right] \\ &= \frac{\sin^2 \left[ \frac{\pi x}{a} \right]}{\left( \frac{\pi x}{a} \right)^2} \frac{\sin^2 \left[ \frac{\pi y}{b} \right]}{\left( \frac{\pi y}{b} \right)^2} \end{aligned}$$

- $SINC^2$  function is everywhere nonnegative
- volume same as that of  $SINC \left[ \frac{x}{a}, \frac{y}{b} \right]$

$$\begin{aligned} \iint_{-\infty}^{+\infty} SINC^2 \left[ \frac{x}{a}, \frac{y}{b} \right] dx dy &= \left( \int_{-\infty}^{+\infty} SINC^2 \left[ \frac{x}{a} \right] dx \right) \left( \int_{-\infty}^{+\infty} SINC^2 \left[ \frac{y}{b} \right] dy \right) \\ &= |a| |b| = |ab| \end{aligned}$$



$$SINC^2 [x, y] = SINC [x] \cdot SINC [y]$$

### 3.7 2-D Gaussian

$$\begin{aligned} GAUS \left[ \frac{x}{a}, \frac{y}{b} \right] &= GAUS \left[ \frac{x}{a} \right] GAUS \left[ \frac{y}{b} \right] = e^{-\frac{\pi x^2}{a^2}} e^{-\frac{\pi y^2}{b^2}} \\ &= e^{-\pi \left( \frac{x^2}{a^2} + \frac{y^2}{b^2} \right)} \end{aligned}$$

- Volume is product of individual areas:

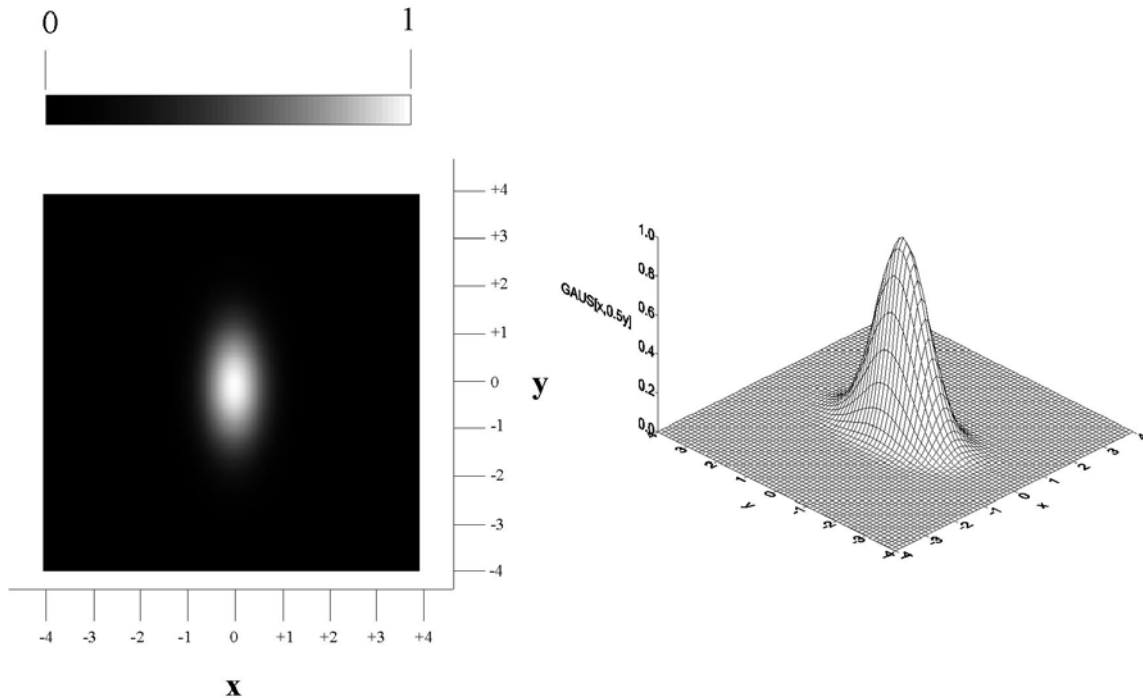
$$\begin{aligned} \iint_{-\infty}^{+\infty} GAUS \left[ \frac{x}{a}, \frac{y}{b} \right] dx dy &= \left( \int_{-\infty}^{+\infty} GAUS \left[ \frac{x}{a} \right] dx \right) \left( \int_{-\infty}^{+\infty} GAUS \left[ \frac{y}{b} \right] dy \right) \\ &= |a| |b| = |ab| \end{aligned}$$

- Rotated version:

$$\begin{aligned} GAUS \left[ \frac{x'}{a}, \frac{y'}{b} \right] &= GAUS \left[ \frac{\mathbf{r} \cdot \hat{\mathbf{p}}}{a}, \frac{\mathbf{r} \cdot \hat{\mathbf{p}}^\perp}{b} \right] \\ &= e^{-\pi \left[ \left( \frac{\mathbf{r} \cdot \hat{\mathbf{p}}}{a} \right)^2 + \left( \frac{\mathbf{r} \cdot \hat{\mathbf{p}}^\perp}{b} \right)^2 \right]} \end{aligned}$$

- Circularly symmetric function generated with equal width parameters:

$$GAUS \left[ \frac{x}{d}, \frac{y}{d} \right] = e^{-\pi \frac{x^2+y^2}{d^2}} = e^{-\pi \left( \frac{r}{d} \right)^2}$$



Example of 2-D separable Gaussian function  $GAUS \left[ x, \frac{y}{2} \right] = GAUS[x] \cdot GAUS \left[ \frac{y}{2} \right]$

### 3.8 2-D Sinusoid

- “meaningful” definition varies along one direction but is constant in the orthogonal direction

$$\begin{aligned} f[x, y] &= A_0 \cos[2\pi\xi_0 x + \phi_0] \cdot 1[y] \\ &= A_0 \cos[2\pi\xi_0 x + \phi_0] \cos[2\pi \times 0 \times y] \end{aligned}$$

- Gaskill suggests image of field plowed with sinusoidal furrows
- Rotated:

$$\begin{aligned} f[x', y'] &= A_0 \cos(2\pi\xi_0 x' + \phi_0) \cdot 1[y'] = A_0 \cos[2\pi\xi_0 (\mathbf{r} \bullet \hat{\mathbf{p}}) + \phi_0] \cdot 1[\mathbf{r} \bullet \hat{\mathbf{p}}^\perp] \\ &= A_0 \cos[2\pi(\xi_0 \cos[\theta] x + (\xi_0 \sin[\theta]) y) + \phi_0] \cdot 1[-x \sin[\theta] + y \cos[\theta]] \\ &\equiv A_0 \cos[2\pi(\xi_1 x + \eta_1 y) + \phi_0] \cdot 1[-x \sin[\theta] + y \cos[\theta]] \end{aligned}$$

- Spatial frequencies of rotated function along “new”  $\xi$ - and  $\eta$ -axes are  $\xi_1$  and  $\eta_1$

$$\begin{aligned} \xi_1 &\equiv +\xi_0 \cos[\theta] \\ \eta_1 &\equiv +\xi_0 \sin[\theta] \end{aligned}$$

- specify rates of sinusoidal variation along  $x$ - and  $y$ -axes
- specify periods of 2-D sinusoid along  $x$ - and  $y$ -directions, respectively:

$$\begin{aligned} f[x, y] &= A_0 \cos[2\pi(\xi_1 x + \eta_1 y) + \phi_0] \\ &= A_0 \cos\left[2\pi\left(\frac{x}{X_1} + \frac{y}{Y_1}\right) + \phi_0\right] \end{aligned}$$

$$\begin{aligned} X_1 &= \frac{1}{\xi_0 \cos[\theta]} \\ Y_1 &= \frac{1}{\xi_0 \sin[\theta]} \end{aligned}$$

- 2-D sinusoid oscillating along arbitrary direction recast into form where 2-D spatial coordinate is 2-D radius vector  $\mathbf{r} \equiv [x, y]$ .
- phase  $[\xi_0 x + \eta_0 y]$  is scalar product of  $\mathbf{r}$  with “polar spatial frequency vector”  $\underline{\rho}_0 \equiv \begin{bmatrix} \xi_0 \\ \eta_0 \end{bmatrix}$
- Polar representation of vector is  $\underline{\rho}_0 = (\rho_0, \psi_0)$

$$|\underline{\rho}_0| = \sqrt{\xi_0^2 + \eta_0^2}$$

$$\psi_0 = \tan^{-1} \left[ \frac{\eta_0}{\xi_0} \right]$$

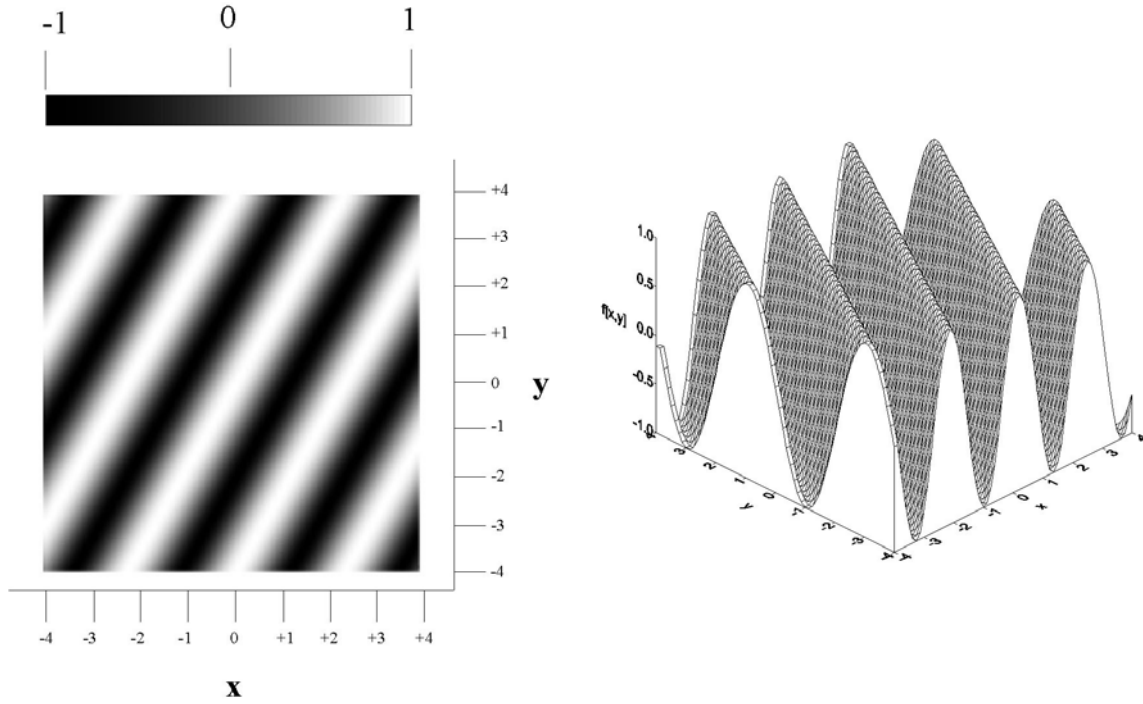
$$\begin{aligned} f(\mathbf{r}) &= A_0 \cos\left[2\pi(\mathbf{r} \bullet \underline{\rho}_0) + \phi_0\right] \\ &= A_0 \cos\left[2\pi\left(|\mathbf{r}| |\underline{\rho}_0| \cos[\theta - \psi_0]\right) + \phi_0\right] \end{aligned}$$

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$$\begin{aligned} \xi_0 &= |\underline{\rho}_0| \cos[\psi_0] \\ \eta_0 &= |\underline{\rho}_0| \sin[\psi_0] \end{aligned}$$

- spatial period is reciprocal of magnitude of spatial frequency vector:

$$R_0 = \frac{1}{|\underline{\rho}_0|} = \frac{1}{\sqrt{\xi_0^2 + \eta_0^2}}$$



2-D sinusoid function rotated in azimuth about the origin by  $\theta = -\frac{\pi}{6}$  radians.

## 4 2-D Dirac Delta Function and its Relatives

- Several flavors of 2-D Dirac delta function possible
  - different separable versions of 2-D Dirac delta function defined in Cartesian and polar coordinates
- Properties:

$$\delta [x - x_0, y - y_0] = 0 \text{ for } x - x_0 \neq 0 \text{ or } y - y_0 \neq 0$$

$$\iint_{-\infty}^{+\infty} \delta [x - x_0, y - y_0] \, dx \, dy = 1$$

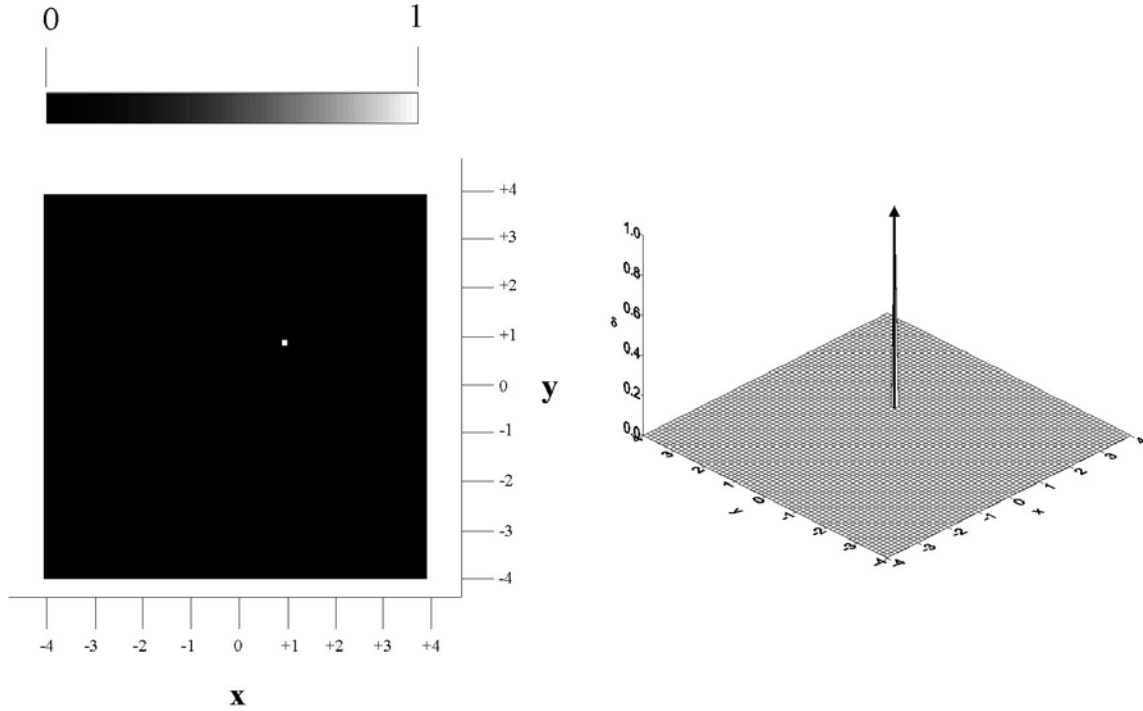
- Expressed as limit of sequence of 2-D functions, each with unit volume
- May be generated from 2-D *RECT*, *TRI*, *GAUS*, *SINC*, and *SINC*<sup>2</sup> in Cartesian coordinates
- Separability of functions ensures that corresponding representation of  $\delta [x, y]$  is separable
- Most common representation of  $\delta [x, y]$  based on 2-D *RECT* function of unit volume in the limit of infinitesimal area:

$$\delta [x, y] = \lim_{b \rightarrow 0} \lim_{d \rightarrow 0} \left\{ \frac{1}{bd} \, \text{RECT} \left[ \frac{x}{b}, \frac{y}{d} \right] \right\}$$

$$= \lim_{b \rightarrow 0} \left\{ \frac{1}{b} \text{RECT} \left[ \frac{x}{b} \right] \right\} \times \lim_{d \rightarrow 0} \left\{ \frac{1}{d} \text{RECT} \left[ \frac{y}{d} \right] \right\}$$

### 4.1 Properties of Separable 2-D Dirac Delta Function, Cartesian Coordinates

$$\delta [x, y] = \delta [x] \, \delta [y]$$



$$\text{2-D separable Dirac delta function } \delta [x - 1, y - 1] = \delta [x - 1] \cdot \delta [y - 1].$$

- 2-D Dirac delta function may be synthesized by summing unit-amplitude 2-D complex linear-phase exponentials with all spatial frequencies.

$$\begin{aligned}
\delta[x] \times \delta[y] &= \left( \int_{-\infty}^{+\infty} e^{+2\pi i \xi x} d\xi \right) \left( \int_{-\infty}^{+\infty} e^{+2\pi i \eta y} d\eta \right) \\
&= \iint_{-\infty}^{+\infty} e^{+2\pi i \xi x} e^{+2\pi i \eta y} d\xi d\eta \\
&= \iint_{-\infty}^{+\infty} e^{+2\pi i (\xi x + \eta y)} d\xi d\eta = \delta[x, y]
\end{aligned}$$

- Because odd imaginary parts cancel:

$$\delta[x, y] = \iint_{-\infty}^{+\infty} \cos[2\pi(\xi x + \eta y)] d\xi d\eta$$

- Scaling property of  $\delta[x] \Rightarrow$

$$\begin{aligned}
\delta\left[\frac{x}{b}, \frac{y}{d}\right] &= \delta\left[\frac{x}{b}\right] \delta\left[\frac{y}{d}\right] \\
&= |b| \delta[x] |d| \delta[y] \\
&= |bd| \delta[x] \delta[y] = |bd| \delta[x, y]
\end{aligned}$$

- Translat by shifting separable components:

$$\delta\left[\frac{x-x_0}{b}, \frac{y-y_0}{d}\right] = |bd| \delta[x-x_0] \delta[y-y_0]$$

- Sifting property evaluates amplitude at a specific location

$$\iint_{-\infty}^{+\infty} f[x, y] \delta[x-x_0, y-y_0] dx dy = f[x_0, y_0]$$

## 4.2 Properties of 2-D Dirac Delta Function, Polar Coordinates

- 2-D Dirac delta function located on  $x$ -axis at distance  $\alpha > 0$  from origin
  - polar coordinates are  $r_0 = \alpha$  and  $\theta_0 = 0$ .

$$\delta[x - \alpha, y] = \delta[x - \alpha] \delta[y]$$

- Polar representation easy to derive because radial coordinate directed along  $x$ -axis
- Azimuthal displacement due to angle parallel to  $y$ -axis.
- Identify  $x \rightarrow r$ ,  $\alpha \rightarrow r_0$ , and  $y \rightarrow r_0 \theta$  and substitute directly, use 1-D scaling property

$$\begin{aligned} \delta[x - \alpha] \delta[y] &= \delta[r - r_0] \delta[r_0 \theta] \\ &= \delta[r - r_0] \left( \frac{1}{|r_0|} \delta[\theta] \right) = \frac{\delta[r - r_0]}{r_0} \delta[\theta] \end{aligned}$$

- last step follows from observation that polar radial coordinate  $r_0 \geq 0$
- Domain of azimuthal coordinate  $\theta$  is a continuous interval of  $2\pi$  radians, e.g.,  $-\pi \leq \theta < +\pi$
- Generalize for a 2-D Dirac delta function located at same radial distance  $r_0$  from the origin but at different azimuth  $\theta_0$ 
  - Cartesian coordinates  $x_0 = r_0 \cos[\theta_0]$ ,  $y_0 = r_0 \sin[\theta_0]$

$$\begin{aligned} \delta[x - x_0, y - y_0] &= \delta[x - r_0 \cos[\theta_0]] \delta[y - r_0 \sin[\theta_0]] \\ &= \frac{\delta[r - r_0]}{r_0} \delta[\theta - \theta_0] \equiv \delta(\underline{\mathbf{r}} - \underline{\mathbf{r}}_0) \end{aligned}$$

- Polar form is product of 1-D Dirac delta functions in the radial and azimuthal directions
- Amplitude scaled by reciprocal of the radial distance.
- Confirm that expression satisfies criteria for 2-D Dirac delta function.
- Easy to show that support is infinitesimal
- Volume is evaluated easily:

$$\begin{aligned} \iint_{-\infty}^{+\infty} \delta(\underline{\mathbf{r}} - \underline{\mathbf{r}}_0) d\underline{\mathbf{r}} &= \int_{-\pi}^{+\pi} d\theta \int_0^{+\infty} \left( \frac{\delta[r - r_0]}{r_0} \delta[\theta - \theta_0] \right) r dr \\ &= \left( \int_{-\pi}^{+\pi} \delta[\theta - \theta_0] d\theta \right) \cdot \left( \int_0^{+\infty} \frac{\delta[r - r_0]}{r_0} r dr \right) \text{ where } r_0 > 0 \\ &= 1 \cdot \int_0^{+\infty} \frac{\delta[r - r_0]}{r_0} r dr = 1 \cdot \int_0^{+\infty} \delta[r - r_0] dr = 1 \end{aligned}$$

- Extend derivation of polar form of 2-D Dirac delta function to  $r_0 = 0$

$$\delta(\underline{\mathbf{r}} - \underline{\mathbf{0}})$$

- Domain of  $r$  is semiclosed single-sided interval  $[0, +\infty)$ 
  - 1-D radial part  $\delta(r - r_0)$  at origin cannot be symmetric
  - $\theta_0$  is indeterminate at origin



- not valid for  $\mathbf{r}_0 = \mathbf{0}$
- 2-D Dirac delta function at the origin must be circularly symmetric and therefore a function of  $r$  only
- No dependence on azimuth angle  $\theta$ :

$$\begin{aligned}\delta(\mathbf{r} - \mathbf{0}) &= \alpha \delta[|\mathbf{r}|] \times 1[\theta] \\ &= \alpha \delta[r] \times 1[\theta]\end{aligned}$$

- \*  $\alpha$  is scaling parameter ensures that  $\delta(\mathbf{r})$  has unit volume
- Domains of polar arguments are  $0 \leq r < +\infty$  and  $-\pi \leq \theta < +\pi$ .
- Modify domains
  - \* Radial variable over domain  $(-\infty, +\infty)$
  - \* Azimuthal domain constrained to interval of  $\pi$  radians, e.g.,  $[0, +\pi)$  or  $[-\frac{\pi}{2}, +\frac{\pi}{2})$ .

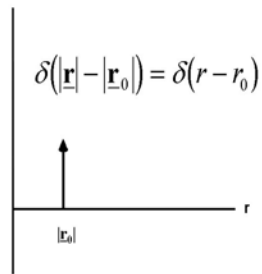
$$\begin{aligned}1 &= \int_{-\pi}^{+\pi} \int_0^{+\infty} \alpha \delta[r] 1[\theta] r dr d\theta = \int_{-\frac{\pi}{2}}^{+\frac{\pi}{2}} \int_{-\infty}^{+\infty} \alpha \delta[r] 1[\theta] r dr d\theta \\ &= \int_{-\frac{\pi}{2}}^{+\frac{\pi}{2}} 1[\theta] d\theta \int_{-\infty}^{+\infty} \alpha \delta[r] r dr = \int_{-\infty}^{+\infty} \alpha \delta[r] \pi r dr\end{aligned}$$

- Area of  $\delta(r)$  over  $(-\infty, +\infty)$  is unity
  - \*  $\alpha = (\pi r)^{-1}$ .

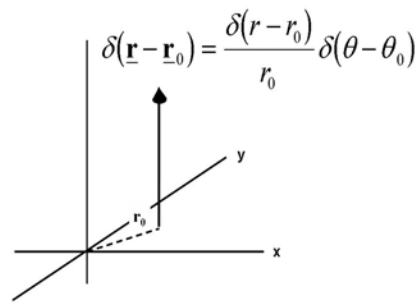
$$\begin{aligned}\delta(\mathbf{r} - \mathbf{0}) &= \delta(\mathbf{r}) \\ &= \left(\frac{1}{\pi r}\right) \delta(r) 1[\theta] \\ &= \frac{\delta(r)}{\pi r}\end{aligned}$$

(a)

1-D Radial Function

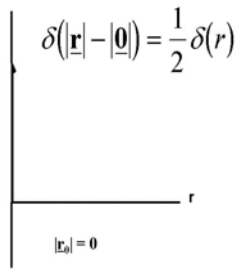


2-D Function

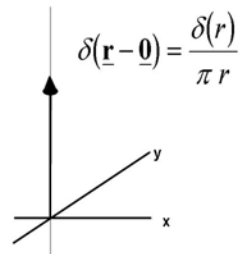


(b)

1-D Radial Function



2-D Function

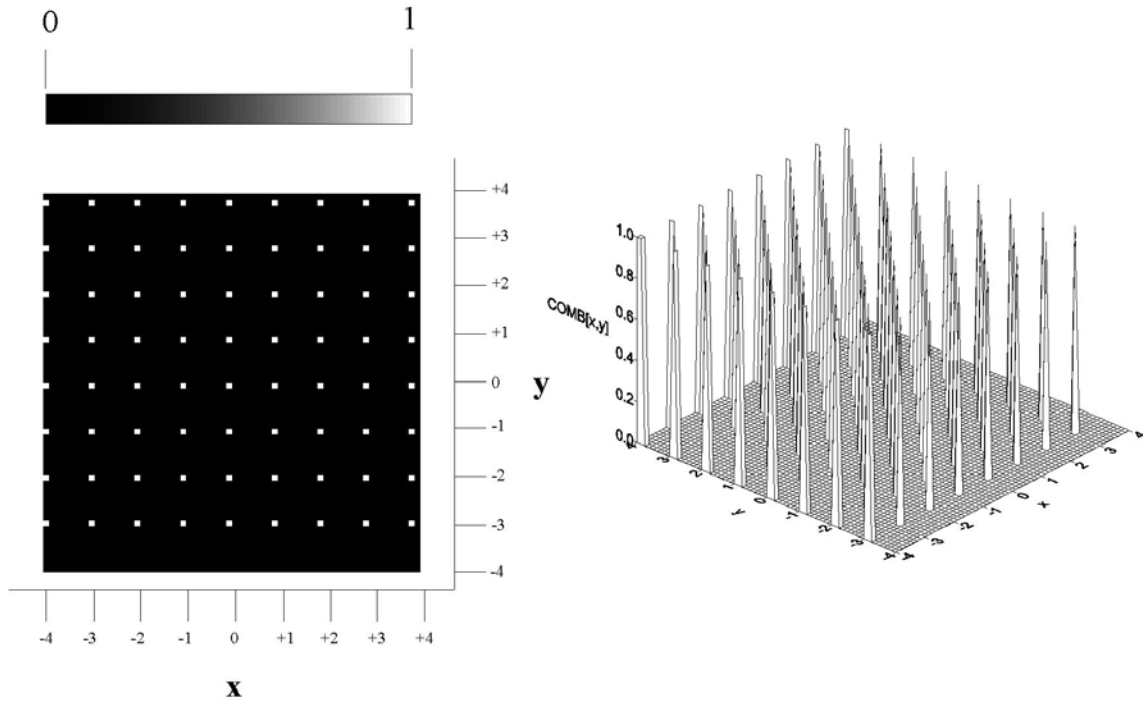


2-D “radial” representation of Dirac delta function.

### 4.3 2-D Separable Comb Function

$$\begin{aligned} \text{COMB}[x, y] &= \text{COMB}[x] \text{COMB}[y] \\ &= \left( \sum_{n=-\infty}^{+\infty} \delta[x - n] \right) \left( \sum_{\ell=-\infty}^{+\infty} \delta[y - \ell] \right) \end{aligned}$$

- Gaskill's "bed of nails"
- Volume is infinite.
- Most important application is to model 2-D sampled functions.



2-D separable COMB function  $\text{COMB}[x, y] = \text{COMB}[x] \cdot \text{COMB}[y]$ .

## 4.4 2-D “Line Delta” Function

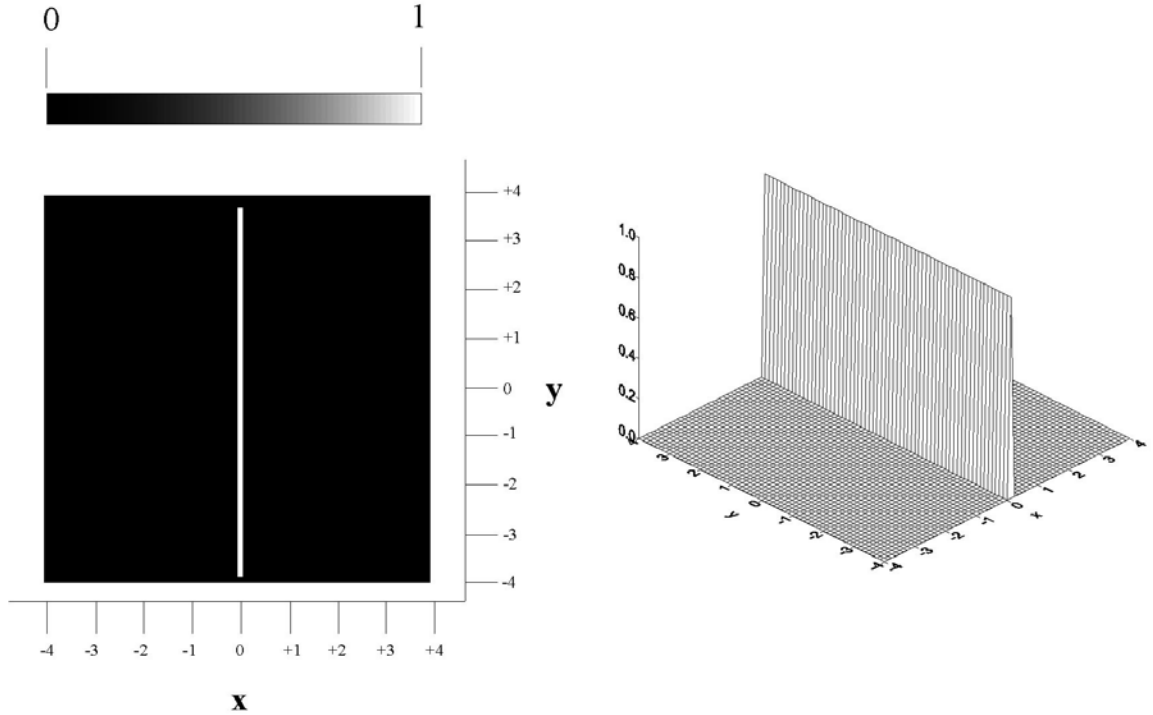
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$$m_1[x, y] = \delta[x] \cdot 1[y]$$

- a “line” or “wall” of 1-D Dirac delta functions along  $y$ -axis
- “line delta function”, “line mass”, or “straight-line impulse”
- Most authors delete explicit unit constant  $1[y]$ 
  - do not distinguish 1-D Dirac delta function and 2-D line delta function along the  $y$ -axis.
- Volume is infinite:

$$\iint_{-\infty}^{+\infty} \delta[x] \cdot 1[y] \, dx \, dy = \int_{-\infty}^{+\infty} \delta[x] \, dx \int_{-\infty}^{+\infty} 1[y] \, dy = 1 \times \infty$$

- Useful to define Radon transform (mathematical basis for medical computed tomography and magnetic resonance imaging)



2-D line Dirac delta function  $\delta[x] \cdot 1[y]$  produces a “wall” of Dirac delta functions along the  $y$ -axis.

- Product of arbitrary function  $f[x, y]$  with line delta function
  - apply sifting property:

$$\begin{aligned} f[x, y] (\delta[x] \cdot 1[y]) &= (f[x, y] \delta[x]) \cdot 1[y] \\ &= (f[0, y] \delta[x]) \cdot 1[y] \\ &= f[0, y] (\delta[x] \cdot 1[y]) \end{aligned}$$

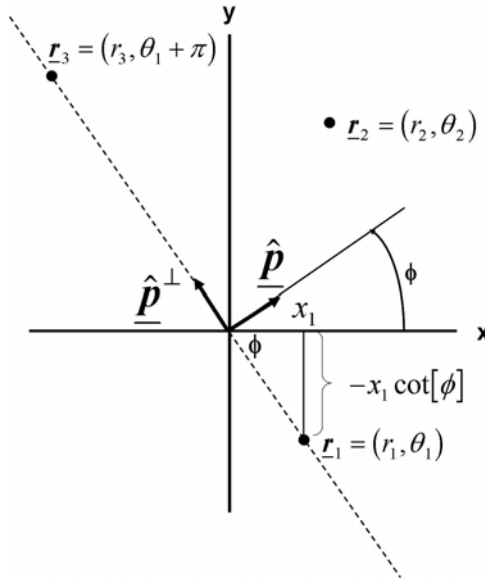
- Volume of product of functions is:

$$\begin{aligned}
\iint_{-\infty}^{+\infty} f[x, y] \delta[x] 1[y] dx dy &= \iint_{-\infty}^{+\infty} f[0, y] \delta[x] 1[y] dx dy \\
&= \int_{-\infty}^{+\infty} f[0, y] 1[y] dy \int_{-\infty}^{+\infty} \delta[x] dx \\
&= \int_{-\infty}^{+\infty} f[0, y] dy
\end{aligned}$$

- Line delta function  $\delta[x] 1[y]$  “sifts out” area of  $f[x, y]$  evaluated along the  $y$ -axis
- Line delta functions oriented along arbitrary azimuth angle:

$$\begin{aligned}
m_2[x, y] &= \delta[\mathbf{r} \cdot \hat{\mathbf{p}}] 1[\mathbf{r} \cdot \hat{\mathbf{p}}^\perp] \\
&= \delta[x \cos[\phi] + y \sin[\phi]] 1[-x \sin[\phi] + y \cos[\phi]]
\end{aligned}$$

- Line delta function oriented along the  $y$ -axis obtained by setting  $\phi = 0$
- Notation for 1-D Dirac delta function may seem “weird” because the argument is the scalar product of two vectors.
- Argument of Dirac delta function is zero for all vectors  $\mathbf{r}$  in 2-D plane for which  $\mathbf{r} \cdot \hat{\mathbf{p}} = 0$ , for all  $\mathbf{r} \perp \hat{\mathbf{p}}$ .
- Complete set of line delta functions at all possible azimuth angles by selecting all angles  $\phi$  in interval  $-\frac{\pi}{2} \leq \phi < +\frac{\pi}{2}$
- Defining unit vector  $\hat{\mathbf{p}}$  lies in first or fourth quadrants.
- Action of unit constant is to “spread” or “smear” amplitude of 1-D Dirac delta function along  $\hat{\mathbf{p}}$ .



The rotated form of the line delta function  $m_2[x, y] = \delta[\mathbf{r} \cdot \hat{\mathbf{p}}] \cdot 1[\mathbf{r} \cdot \hat{\mathbf{p}}^\perp]$ , where  $\hat{\mathbf{p}} = (\cos[\phi], \sin[\phi])$ . Three cases of  $\mathbf{r}$  are shown:  $\mathbf{r}_1$  and  $\mathbf{r}_3$  lie upon the line delta function and  $\mathbf{r}_2$  does not.

- Use polar form of radius vector  $\underline{\mathbf{r}} = (r, \phi)$ , where  $0 \leq r < +\infty$  and  $-\pi \leq \phi < +\pi$

$$\begin{aligned}
m_2[x, y] &= \delta[\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}] \, 1[\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}^\perp] \\
&= \delta[|\underline{\mathbf{r}}| |\underline{\hat{\mathbf{p}}}| \cos[\theta - \phi]] \, 1[\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}^\perp] \\
&= \delta[|\underline{\mathbf{r}}| \cos[\theta - \phi]] \, 1[\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}^\perp] \\
&= \delta[r \cos[\phi - \theta]]
\end{aligned}$$

- Expression determines the set of values of  $(r, \theta)$  on radial line through origin perpendicular to azimuthal angle  $\phi$ .
- Write as function of azimuthal angle  $\phi$  by recasting into more convenient form
  - Apply expression for Dirac delta function with a functional argument

$$\begin{aligned}
\delta[\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}] \, 1[\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}^\perp] &= \delta[r \cos[\theta - \phi]] \, 1[\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}^\perp] \\
&= \delta[g[\phi]] \, 1[\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}^\perp] \\
&= \frac{1}{|g'(\phi_0)|} \delta[\phi - \phi_0] \, 1[\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}^\perp] \\
&= \frac{1}{r |\sin[\theta - \phi_0]|} \delta[\phi - \phi_0] \, 1[\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}^\perp]
\end{aligned}$$

- $\phi_0$  is angle that satisfies condition  $\cos[\theta - \phi_0] = 0 \implies \phi_0 = \theta \pm \frac{\pi}{2}$

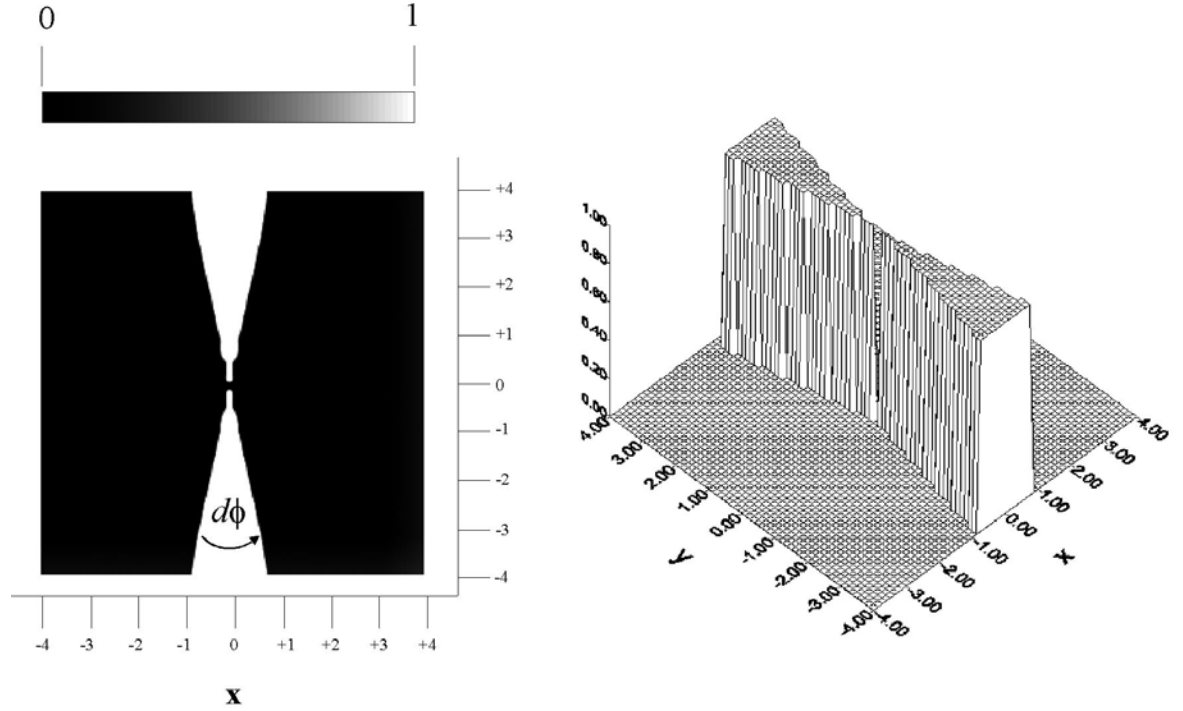
$$\begin{aligned}
\delta[\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}] \, 1[\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}^\perp] &= \frac{1}{r |\sin[\mp \frac{\pi}{2}]|} \delta\left[\phi - \left(\theta \pm \frac{\pi}{2}\right)\right] \\
&= \frac{1}{r} \delta\left[\phi - \left(\theta \pm \frac{\pi}{2}\right)\right]
\end{aligned}$$

- Line delta function through origin lying along radial line perpendicular to azimuthal angle  $\phi$  is equivalent to amplitude-weighted 1-D Dirac delta function of angle  $\phi$  that is nonzero only for  $\phi = \theta \pm \frac{\pi}{2}$
- Sign selected to ensure that  $\phi$  lies within usual domain of polar coordinates:  $-\pi \leq \phi < +\pi$ .
- Each unit increment of radial distance along  $y$ -axis contributes same unit volume.
- This expression seems to indicate otherwise for rotated function
- Contribution to volume affected by factor  $r^{-1}$ .
- o resolve apparent comundrum, substitute expression for Dirac delta function as limit of rectangle function

$$\delta\left[\phi - \left(\theta \pm \frac{\pi}{2}\right)\right] = \lim_{b \rightarrow 0} \left\{ \frac{1}{b} \, RECT\left[\frac{\phi - \left(\theta \pm \frac{\pi}{2}\right)}{b}\right] \right\}$$

- \*  $b$  measured in radians
- \* Function consists of two symmetric “wedges” of fixed amplitude  $b^{-1}$ .
  - Second “wedge” about the angle  $\phi = -\frac{\pi}{2}$  results from convention that domain of azimuth angle is  $-\frac{\pi}{2} \leq \phi < +\frac{\pi}{2}$ .
- \* In limit  $b \rightarrow 0$ , the angular “spread” of wedges decreases while amplitude increases.

- \* Contribution of segments of wedges with unit radial extent increases in proportion to radial distance  $r$  from origin.
- \* Factor of  $r^{-1}$  compensates for increase in volume to ensure that contributions to volume of segments with equal radial extent remain constant.

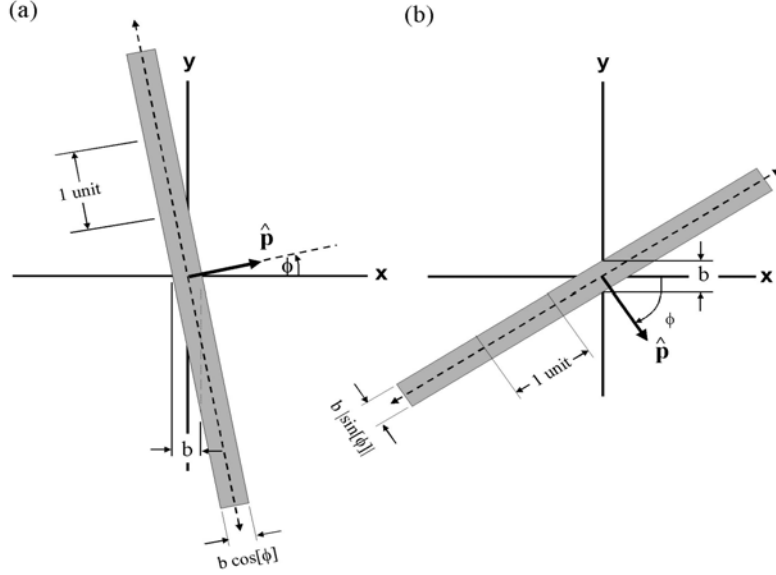


“Angular” delta function as limit of rectangle,  $\lim_{d\phi \rightarrow 0} \left\{ \frac{1}{d\phi} \text{RECT} \left[ \frac{\phi - \frac{\pi}{2}}{d\phi} \right] \right\}$ .

- Other forms may be derived by manipulating argument
- Slope-intercept form of a line in the 2-D plane

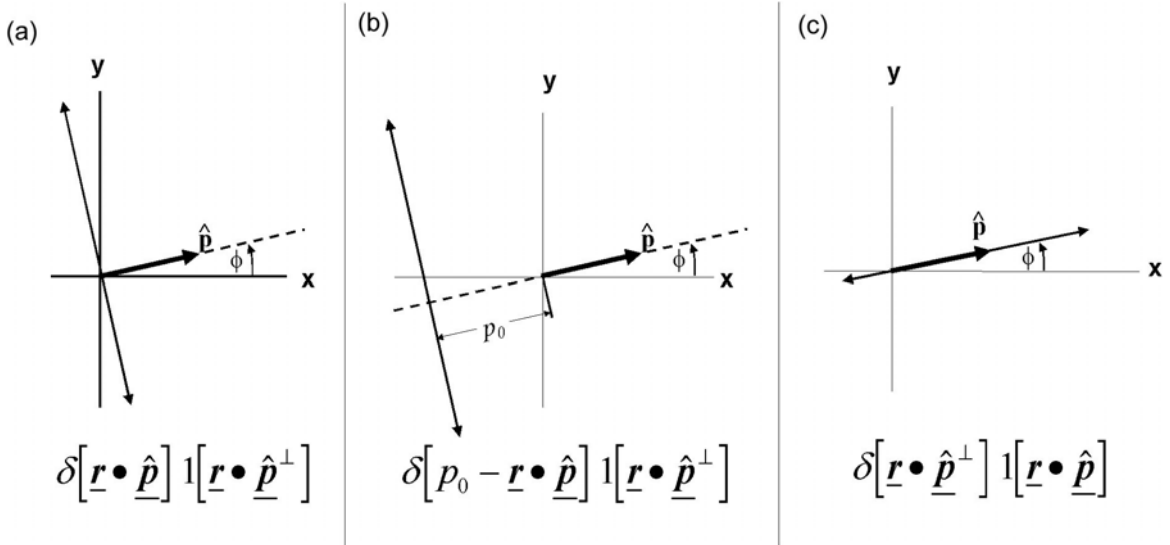
$$\begin{aligned} \delta [\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}] \, 1 [\underline{\mathbf{r}} \bullet \underline{\hat{\mathbf{p}}}^\perp] &= \delta \left[ +\sin[\phi] \left( y + \frac{x \cos[\phi]}{\sin[\phi]} \right) \right] \, 1 [x \sin[\phi] - y \cos[\phi]] \\ &= \left| \frac{1}{\sin[\phi]} \right| \delta [(y - (\cot[-\phi] x + 0))] \, 1 [x \sin[\phi] - y \cos[\phi]] \end{aligned}$$

- Dirac delta function evaluates to zero except at  $[x, y]$  that satisfy slope-intercept form of straight line
- $y$ -intercept is zero
- Slope is  $s = \cot[-\phi] = -\cot[\phi]$



Two examples of rotated line delta functions as 1-D rectangle functions of width  $b$ . In both cases, the amplitude of the gray area is  $b^{-1}$ . (a) The rectangle is a function of  $x$  with width  $b$  measured on the  $x$ -axis and “perpendicular width”  $b \cos(\phi)$ ; (b) the rectangle is a function of  $y$  with “perpendicular width”  $b \sin[\phi]$ .

Three examples of line delta functions:



Different “flavors” of line delta functions: (a)  $\delta[\underline{\mathbf{r}} \cdot \underline{\hat{\mathbf{p}}}] \, 1[\underline{\mathbf{r}} \cdot \underline{\hat{\mathbf{p}}^\perp}]$  through the origin perpendicular to  $\underline{\hat{\mathbf{p}}}$ ; (b)  $\delta[p_0 - \underline{\mathbf{r}} \cdot \underline{\hat{\mathbf{p}}}] \, 1[\underline{\mathbf{r}} \cdot \underline{\hat{\mathbf{p}}^\perp}]$  perpendicular to  $\underline{\hat{\mathbf{p}}}$  at a distance  $p_0 < 0$  from the origin; (c)  $\delta[\underline{\mathbf{r}} \cdot \underline{\hat{\mathbf{p}}^\perp}] \, 1[\underline{\mathbf{r}} \cdot \underline{\hat{\mathbf{p}}}]$  through the origin parallel to  $\underline{\hat{\mathbf{p}}}$ .

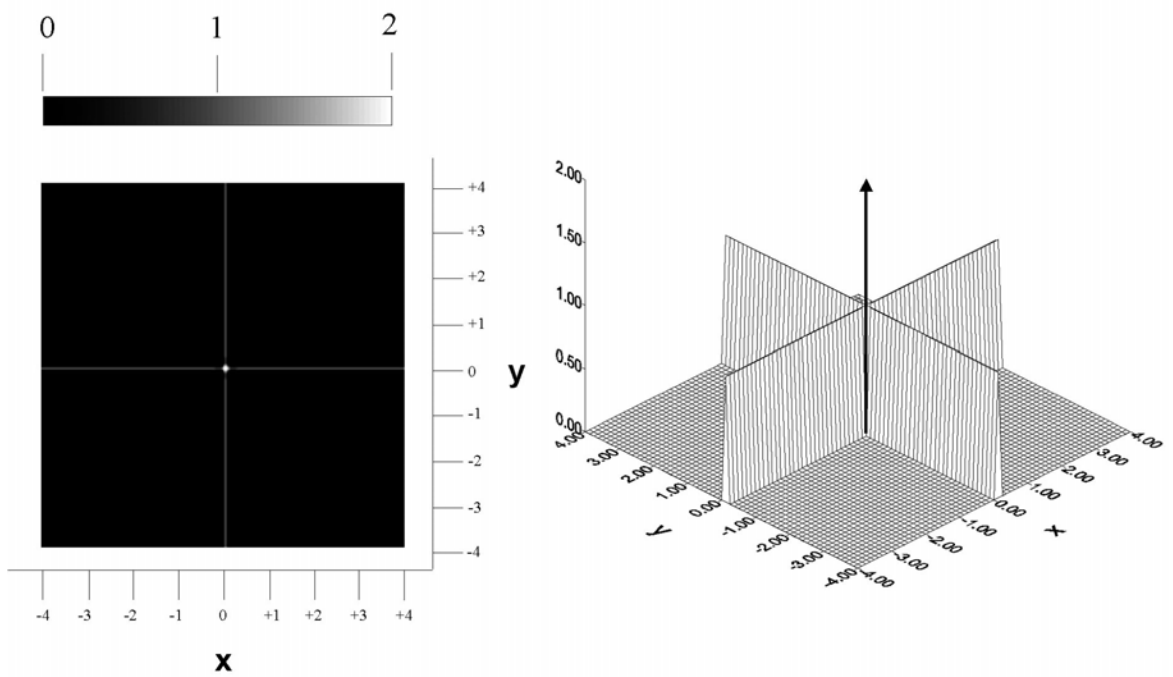


- Sum of two orthogonal line delta functions.
- Simplest version defined with “lines” coincident with Cartesian axes:

$$CROSS[x, y] = \delta[x] \cdot 1[y] + 1[x] \cdot \delta[y]$$

- $CROSS[0, 0] = 2 \cdot \delta[x, y]$
- Cross function with “arms” oriented along the vectors  $\hat{\underline{\mathbf{p}}}$  and  $\hat{\underline{\mathbf{p}}}^\perp$  by substituting  $\underline{\mathbf{r}} \bullet \hat{\underline{\mathbf{p}}}$  for  $x$  and  $\underline{\mathbf{r}} \bullet \hat{\underline{\mathbf{p}}}^\perp$  for  $y$ :

$$CROSS[\underline{\mathbf{r}} \bullet \hat{\underline{\mathbf{p}}}, \underline{\mathbf{r}} \bullet \hat{\underline{\mathbf{p}}}^\perp] = \delta[\underline{\mathbf{r}} \bullet \hat{\underline{\mathbf{p}}}] \cdot 1[\underline{\mathbf{r}} \bullet \hat{\underline{\mathbf{p}}}^\perp] + 1[\underline{\mathbf{r}} \bullet \hat{\underline{\mathbf{p}}}] \cdot \delta[\underline{\mathbf{r}} \bullet \hat{\underline{\mathbf{p}}}^\perp]$$



$CROSS[x, y] = \delta[x] \cdot 1[y] + 1[x] \cdot \delta[y]$ . Note that  $CROSS[0, 0] = 2 \cdot \delta[x, y]$ .

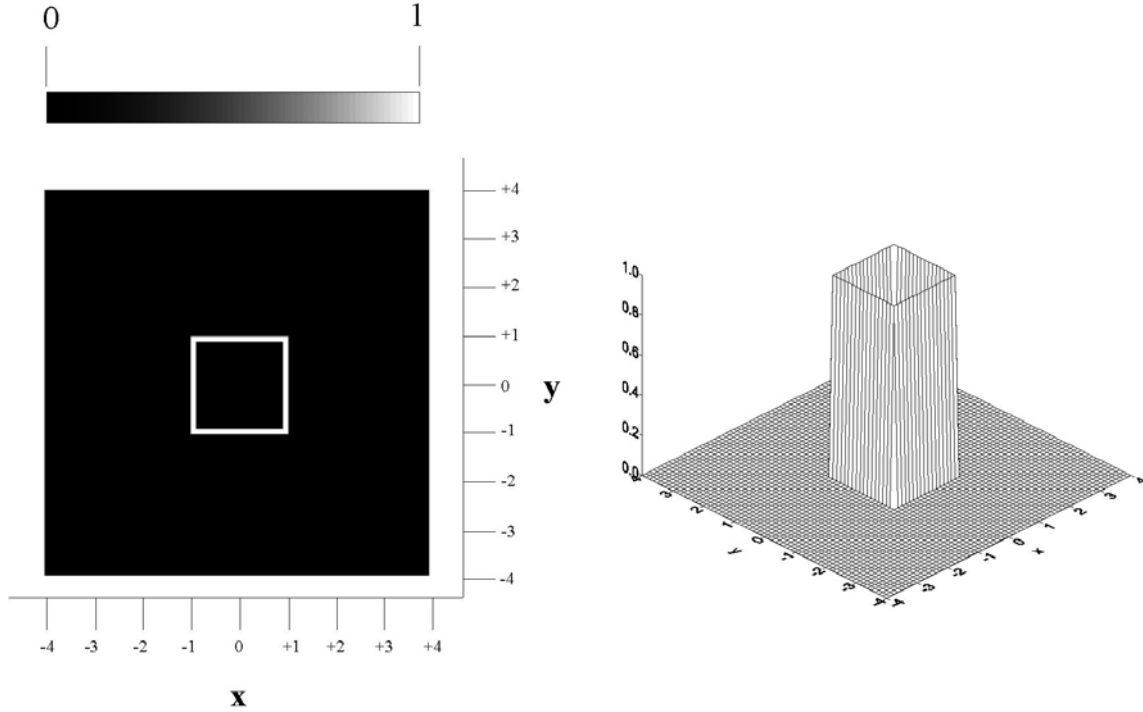
## 4.5 2-D “Corral” Function

- Superpose truncated line delta functions to construct rectangular “stockade” of Dirac delta functions
- Sum of four sections:

$$\begin{aligned}
 COR \left[ \frac{x}{b}, \frac{y}{d} \right] &\equiv \delta \left[ x + \frac{b}{2} \right] RECT \left[ \frac{y}{d} \right] + RECT \left[ \frac{x}{b} \right] \delta \left[ y + \frac{d}{2} \right] \\
 &\quad + \delta \left[ x - \frac{b}{2} \right] RECT \left[ \frac{y}{d} \right] + RECT \left[ \frac{x}{b} \right] \delta \left[ y - \left( -\frac{d}{2} \right) \right] \\
 &= \left( \delta \left[ x + \frac{b}{2} \right] + \delta \left[ x - \left( -\frac{b}{2} \right) \right] \right) RECT \left[ \frac{y}{d} \right] + RECT \left[ \frac{x}{b} \right] \left( \delta \left[ y - \frac{d}{2} \right] + \delta \left[ y - \frac{d}{2} \right] \right)
 \end{aligned}$$

- Volume computed from separable parts:

$$\begin{aligned}
 \iint_{-\infty}^{+\infty} COR \left[ \frac{x}{b}, \frac{y}{d} \right] dx dy &= \int_{-\infty}^{+\infty} \left( \delta \left[ x + \frac{b}{2} \right] + \delta \left[ x - \frac{b}{2} \right] \right) dx \int_{-\infty}^{+\infty} RECT \left[ \frac{y}{d} \right] dy \\
 &\quad + \int_{-\infty}^{+\infty} RECT \left[ \frac{x}{b} \right] dx \int_{-\infty}^{+\infty} \left( \delta \left[ y - \frac{d}{2} \right] + \delta \left[ y - \frac{d}{2} \right] \right) dy \\
 &= 2d + 2b = 2(b + d)
 \end{aligned}$$



## 5 2-D Functions with Circular Symmetry

- Vary along radial direction
- Constant in azimuthal direction
- All profiles along radial lines are identical
- Important in optics
- Optical systems are constructed from lenses with circular cross sections.
- Same amplitude at all points on circle of radius  $r_0$  centered at origin
- Amplitude of radial function  $f_r(r_0)$  replicated for  $[x, y]$  that satisfy  $x^2 + y^2 = r_0^2$ .
- Circularly symmetric function expressed as orthogonal product of 1-D radial profile and unit constant in azimuthal direction:

$$f[x, y] \implies f(\mathbf{r}) = f_r(r) \cdot 1[\theta], \quad 0 \leq r < +\infty, \quad -\pi \leq \theta < +\pi$$

- Rotation about origin has no effect on amplitude
- Symmetric with respect to the origin
  - Domains of radial and azimuthal variables may be recast with symmetric radial interval  $-\infty < r < +\infty$
  - Azimuthal domain  $-\frac{\pi}{2} \leq \theta < +\frac{\pi}{2}$ .
- Volume calculated in polar coordinates
- area element  $dx dy$  replaced by area element in polar coordinates  $r dr d\theta$ :

$$\begin{aligned} \iint_{-\infty}^{+\infty} f[x, y] dx dy &= \int_{\theta=-\pi}^{\theta=+\pi} \int_{r=0}^{r=+\infty} f(\mathbf{r}) r dr \\ &= 2\pi \int_0^{+\infty} f_r(r) r dr \end{aligned}$$

- Center of symmetry relocated to  $[x_0, y_0]$  by adding vector to argument:

$$\begin{aligned} f(\mathbf{r} - \mathbf{r}_0) &= f[x, y] \\ &= f[|\mathbf{r}| \cos[\theta] - |\mathbf{r}_0| \cos[\theta_0], |\mathbf{r}|] \end{aligned}$$

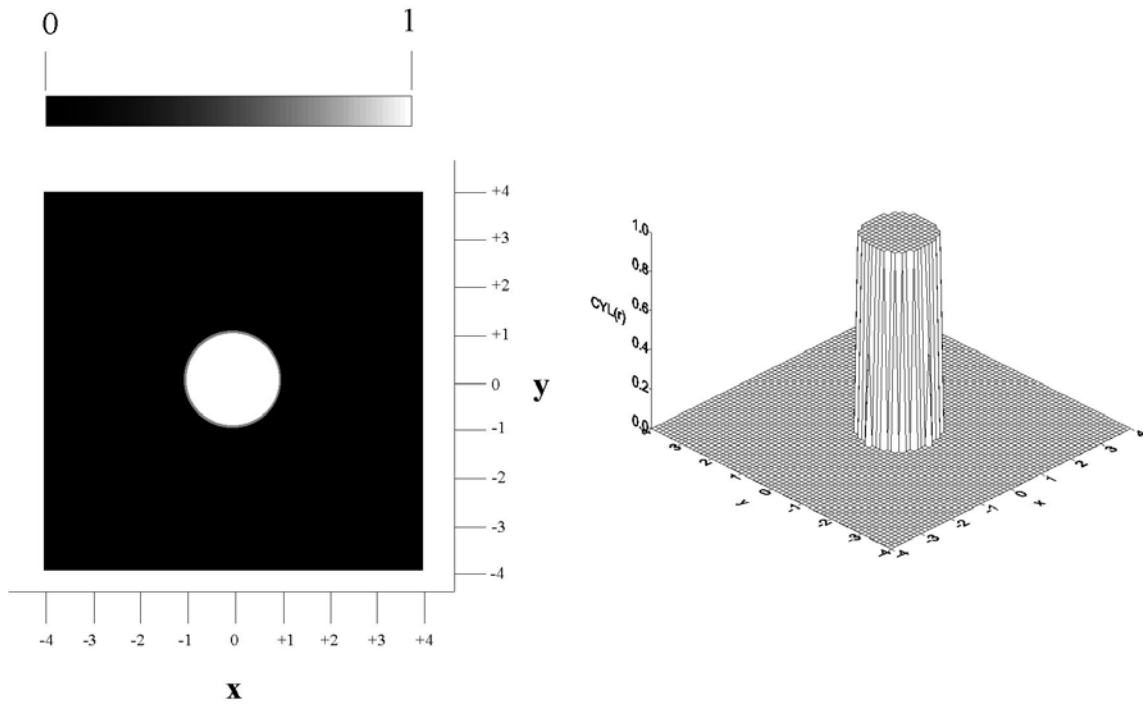
## 5.1 Cylinder (Circle) Function

- Unit amplitude inside radius  $r_0$
- Null amplitude outside
- Circularly symmetric version of 2-D rectangle

$$CYL\left(\frac{r}{d_0}\right) = \begin{cases} 1 & \text{for } r < \frac{d_0}{2} \\ \frac{1}{2} & \text{for } r = \frac{d_0}{2} \\ 0 & \text{for } r > \frac{d_0}{2} \end{cases}$$

- Area of enclosed circle of unit diameter is  $\frac{\pi}{4} \simeq 0.7854 < \text{unit area of } RECT[x, y]$ .
- Volume of  $CYL\left(\frac{r}{d_0}\right)$ :

$$\int_{-\pi}^{+\pi} d\theta \int_0^{+\infty} CYL\left(\frac{r}{d_0}\right) r d_0 r = \frac{\pi d_0^2}{4}$$



Example of 2-D cylinder function  $CYL\left(\frac{r}{2}\right)$ .

## 5.2 Circularly Symmetric Gaussian

- Circularly symmetric Gaussian function identical to separable Gaussian function if scale parameters  $a = b \equiv d$

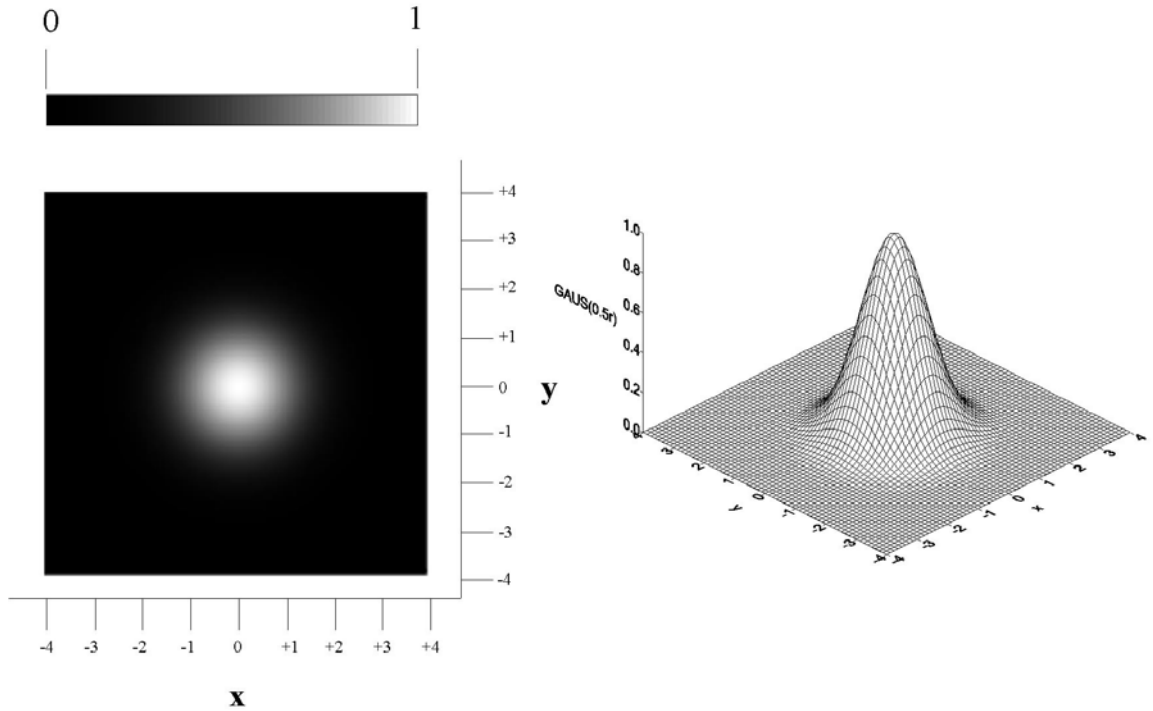
$$\begin{aligned} e^{-\frac{\pi x^2}{d^2}} e^{-\frac{\pi y^2}{d^2}} &= e^{-\pi \frac{x^2+y^2}{d^2}} \\ &= e^{-\pi \left(\frac{r}{d}\right)^2} \\ &\equiv GAUS\left(\frac{r}{d}\right) \end{aligned}$$

- Volume by integrating over polar coordinates after making appropriate change of variable:

$$\begin{aligned} \iint_{-\infty}^{+\infty} GAUS\left(\frac{r}{d}\right) dx dy &= \int_{-\pi}^{+\pi} \int_0^{+\infty} GAUS\left(\frac{r}{d}\right) r dr d\theta \\ &= 2\pi \int_0^{+\infty} r e^{-\pi \left(\frac{r}{d}\right)^2} dr \end{aligned}$$

- Define new integration variable  $\alpha \equiv \pi \left(\frac{r}{d}\right)^2$

$$\iint_{-\infty}^{+\infty} GAUS\left(\frac{r}{d}\right) dx dy = 2\pi \int_0^{+\infty} \frac{d^2}{2\pi} e^{-\alpha} d\alpha = d^2$$



Example of 2-D circularly symmetric Gaussian  $f(r) = e^{-\pi \left(\frac{r}{d}\right)^2}$ .

### 5.3 Circularly Symmetric Bessel Function, Zero Order

$$J_0 [2\pi r \rho_0] \ 1 [\theta] = J_0 \left( 2\pi \sqrt{x^2 + y^2} \ \rho_0 \right)$$

- Selectable parameter  $\rho_0$  analogous to spatial frequency of sinusoid
  - Larger  $\rho_0 \implies$  shorter interval between successive maxima of Bessel function
- Appears in several imaging applications.
- $J_0 [2\pi r \rho_0] \ 1 [\theta]$  generated by summing 2-D cosine functions with the same period “directed” along all azimuthal directions.
- Constituent functions have form  $\cos [2\pi (\eta x + \xi y)]$ , where  $\sqrt{\xi^2 + \eta^2} = \rho_0$ .

$$J_0 [2\pi r \rho_0] \ 1 [\theta] = \int_{-\pi/2}^{+\pi/2} \cos [2\pi (\xi x + \eta y)] \ d\phi \text{ where } \xi^2 + \eta^2 = \rho_0^2$$

- Symmetry of integrand used to evaluate integral over domain of  $2\pi$  radians:

$$\begin{aligned} J_0 [2\pi r \rho_0] &= \frac{1}{2} \int_{-\pi}^{+\pi} \cos [2\pi (\xi x + \eta y)] \ d\phi, \text{ where } \xi^2 + \eta^2 = \rho_0^2 \\ &= \frac{1}{2} \int_{-\pi}^{+\pi} \cos [2\pi (r \xi \cos [\phi] + r \eta \sin [\phi])] \ d\phi \end{aligned}$$

- Rewrite integrals by substituting corresponding linear-phase complex exponential for cosine function

$$\begin{aligned} J_0 [2\pi r \rho_0] \ 1 [\theta] &= \frac{1}{\pi} \int_{-\pi/2}^{+\pi/2} e^{\pm 2\pi i (\xi x + \eta y)} d\phi \\ &= \frac{1}{2\pi} \int_{-\pi}^{+\pi} e^{\pm 2\pi i (\xi x + \eta y)} d\phi \\ &= \frac{1}{2\pi} \int_{-\pi}^{+\pi} \cos [2\pi r (\xi \cos [\phi] + \eta \sin [\phi])] \ d\phi, \text{ where } \xi^2 + \eta^2 = \rho_0^2 \end{aligned}$$

- Profile along the x-axis by setting  $\eta = 0$ . Spatial frequency  $\xi$  along  $x$ -axis becomes  $\xi = \rho_0$ :

$$\begin{aligned} J_0 [2\pi r \rho_0] \ 1 [\theta] |_{\eta=0} &= J_0 [2\pi x \rho_0] \\ &= \frac{1}{\pi} \int_{-\pi/2}^{+\pi/2} \cos [2\pi r \rho_0 \cos [\phi]] \ d\phi \\ &= \frac{1}{2\pi} \int_{-\pi}^{+\pi} \cos [2\pi r \rho_0 \cos [\phi]] \ d\phi \\ &= J_0 [2\pi x \xi] |_{r=x, \xi=\rho_0} = J_0 [2\pi r \rho_0] \end{aligned}$$

- $\rho_0 \cos [\phi]$  is spatial frequency of constituent 1-D cosines of Bessel function
  - Suggests alternate interpretation that 1-D  $J_0$  Bessel function is sum of 1-D cosines with spatial frequencies in interval  $-\rho_0 \leq \xi \leq +\rho_0$  but weighted in “density” by  $\cos [\phi]$
  - Largest spatial frequency exists when  $\phi = 0$ , while the cosine is the unit constant when  $\phi = \pm \frac{\pi}{2}$ .

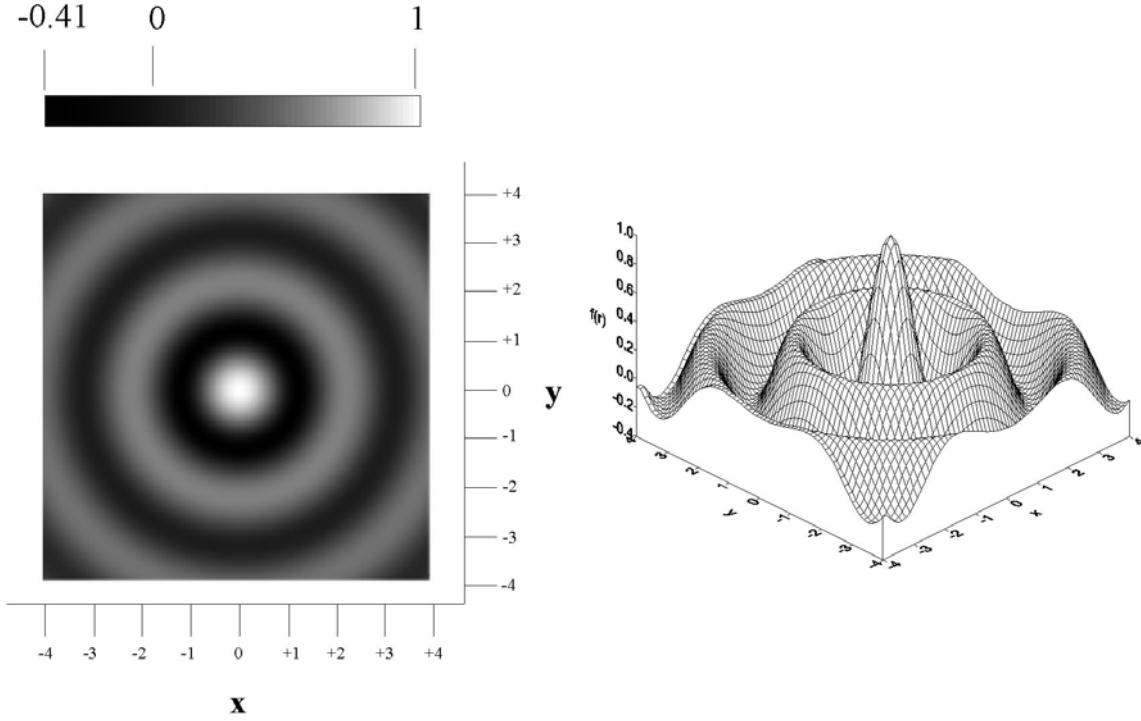
- Equivalent expressions for 1-D Bessel function obtained by setting  $\xi = 0$  so that  $\eta = \rho_0$ 
  - equivalent to projecting argument onto  $y$ -axis:

$$\begin{aligned}
J_0 [2\pi y \eta] \big|_{\xi=0} &= \frac{1}{\pi} \int_{-\pi/2}^{+\pi/2} \cos [2\pi r (\rho_0 \sin [\phi])] d\phi \\
&= \frac{1}{2\pi} \int_0^{+\frac{\pi}{2}} \cos [2\pi r (\rho_0 \sin [\phi])] d\phi \\
&= J_0 [2\pi y \eta] \big|_{r=y, \eta=\rho_0} = J_0 [2\pi r \rho_0]
\end{aligned}$$

- Projection of complex-valued formulations onto  $x$ -axis is:

$$(J_0 [2\pi y \rho_0] \big|_{\xi=0}) = \frac{1}{\pi} \int_{-\pi/2}^{+\pi/2} e^{\pm 2\pi i r \rho_0 \cos[\phi]} d\phi \text{ where } \xi^2 + \eta^2 = \rho_0^2$$

- May be used as an equivalent definition of the Bessel function.

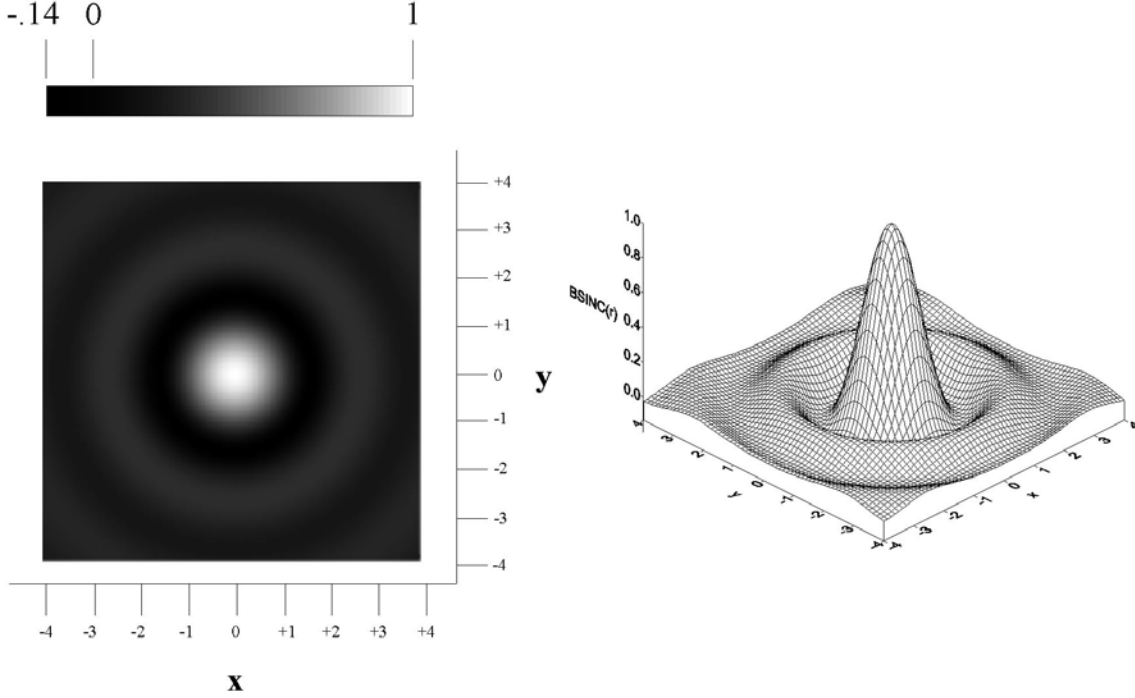


Circularly symmetric function  $J_0 (2\pi \rho_0 r)$  for  $\rho_0 = \frac{1}{2}$ .

## 5.4 “Sombrero” (Besinc) Function

- Circularly symmetric analogue to *SINC* function
- Ratio of two functions with null amplitude at origin
  - numerator is Bessel function of the first kind of order unity  $J_1(\pi r)$
  - denominator is factor proportional to  $r$ .

$$SOMB(r) = \frac{2 J_1(\pi r)}{\pi r}$$



$$SOMB(r) = 2 \frac{J_1(\pi r)}{\pi r}$$

- Amplitude of  $SOMB(r)$  at origin obtained via l'Hôpital's rule:

$$\begin{aligned} SOMB(0) &= 2 \lim_{r \rightarrow 0} \left\{ \frac{\frac{d}{dr}(J_1[\pi r])}{\frac{d}{dr}(\pi r)} \right\} \\ &= 2 \frac{\frac{d}{dr} \left( \frac{\pi r}{2} - \frac{(\pi r)^3}{16} + \frac{(\pi r)^5}{384} - \frac{(\pi r)^7}{18,432} + \dots \right) \Big|_{r=0}}{\frac{d}{dr}(\pi r) \Big|_{r=0}} \\ &= 2 \frac{\left( \frac{\pi}{2} \right)}{\pi} = 1 \end{aligned}$$

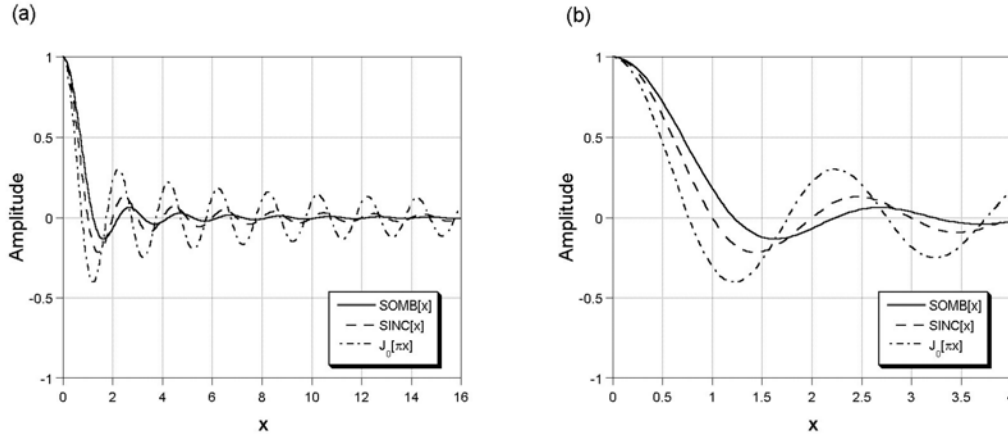
- Asymptotic form of  $J_1[x]$  for large  $x$  used to derive expression for  $SOMB(r)$  in limit of large  $r$ :

$$\begin{aligned} \lim_{r \rightarrow +\infty} \{SOMB(r)\} &= \frac{\sqrt{\frac{2}{\pi r}} \sin \left[ \pi r - \frac{\pi}{4} \right]}{\pi r} \\ &\propto r^{-\frac{3}{2}} \sin \left[ \pi r - \frac{\pi}{4} \right] \end{aligned}$$



- Compare  $x$ -axis profiles of  $J_0[\pi x]$ ,  $SINC[x]$ , and  $SOMB[x]$ 
  - All have unit amplitude at origin
  - Product of periodic or pseudoperiodic oscillation and decaying function of  $x$
  - Peak amplitudes of  $J_0$  decay most slowly with increasing  $x$  ( $|J_0[\pi x]| \propto x^{-\frac{1}{2}}$ );
  - $SINC[x]$  decreases as  $x^{-1}$
  - $SOMB[x]$  as  $x^{-\frac{3}{2}}$
  - Differences in locations of zeros
    - \* zeros of  $SINC[x]$  located at integer values of  $x$
    - \* First two zeros of  $J_0[\pi x]$  located at  $x_1 \simeq \frac{2.4048}{\pi} \simeq 0.7654$  and  $x_2 \simeq \frac{5.5201}{\pi} \simeq 1.7571$ 
      - interval slightly less than unity
    - \* Interval between successive pairs of zeros of  $J_0[\pi x]$  decreases and asymptotically approaches unity as  $x \rightarrow \infty$
    - \* First two zeros of  $J_1[\pi x]$  (and therefore of  $SOMB(r)$ ) located at  $x \simeq 1.2197$  and  $x \simeq 2.2331$ 
      - Interval between zeros of  $SOMB(r)$  decreases with increasing  $r$
      - asymptotically approaches unity as  $r \rightarrow \infty$ .

$$\int_{-\pi}^{+\pi} d\theta \int_0^{+\infty} SOMB\left(\frac{r}{d}\right) r dr = \frac{4d^2}{\pi}$$



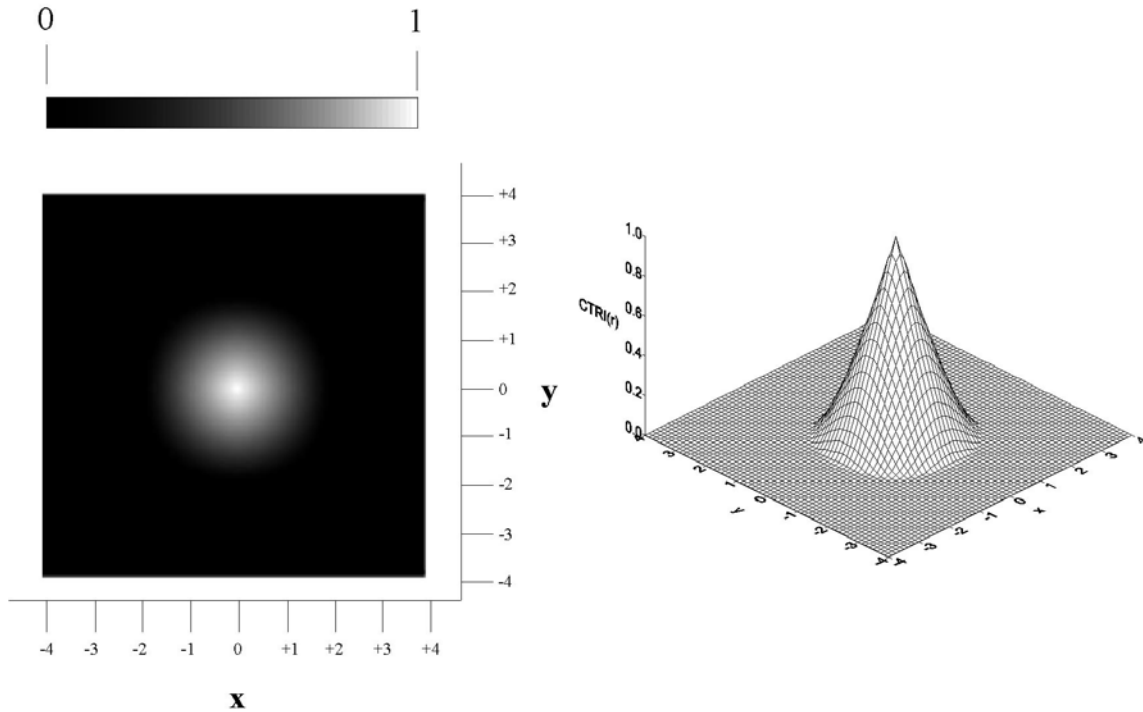
Comparison of profiles of  $BSINC[x]$ ,  $SINC[x]$ , and  $J_0[\pi x]$  (a) for  $0 \leq x \leq 8$  and (b) magnified view for  $0 \leq x \leq 2$ .  $BSINC[x]$  and  $J_0[\pi x]$  “fall off” at the most rapid and slowest rates, respectively.

## 5.5 Circular Triangle Function

- Circularly symmetric analogue of 2-D triangle function
- Most easily constructed as a 2-D autocorrelation
- Describes spatial response of optical imaging systems constructed from elements with circular cross sections.

$$CTRI(r) \equiv \frac{2}{\pi} \left( \cos^{-1}[r] - r\sqrt{1-r^2} \right) CYL\left(\frac{r}{2}\right)$$

- Amplitude decreases approximately linearly until  $r \simeq 0.8$  where slope “flattens out” slightly
- Profiles of both 2-D  $CTRI$  and 2-D  $TRI$  are not straight lines.



Circularly symmetric “triangle” function  $CTRI(r)$ .

## 5.6 Ring Delta Function

- Circularly symmetric analogue of rectangular “corral”:

$$\begin{aligned} f(\mathbf{r}) &= \delta(|\mathbf{r}| - r_0) \\ &\equiv \delta(r - r_0) \\ &= \delta(r - r_0) 1[\theta] \end{aligned}$$

- Resembles polar representation of 2-D Dirac delta function
- Radial arguments  $r$  and  $r_0$  are scalars
- Volume evaluated in polar coordinates by applying sifting property of 1-D Dirac delta function:

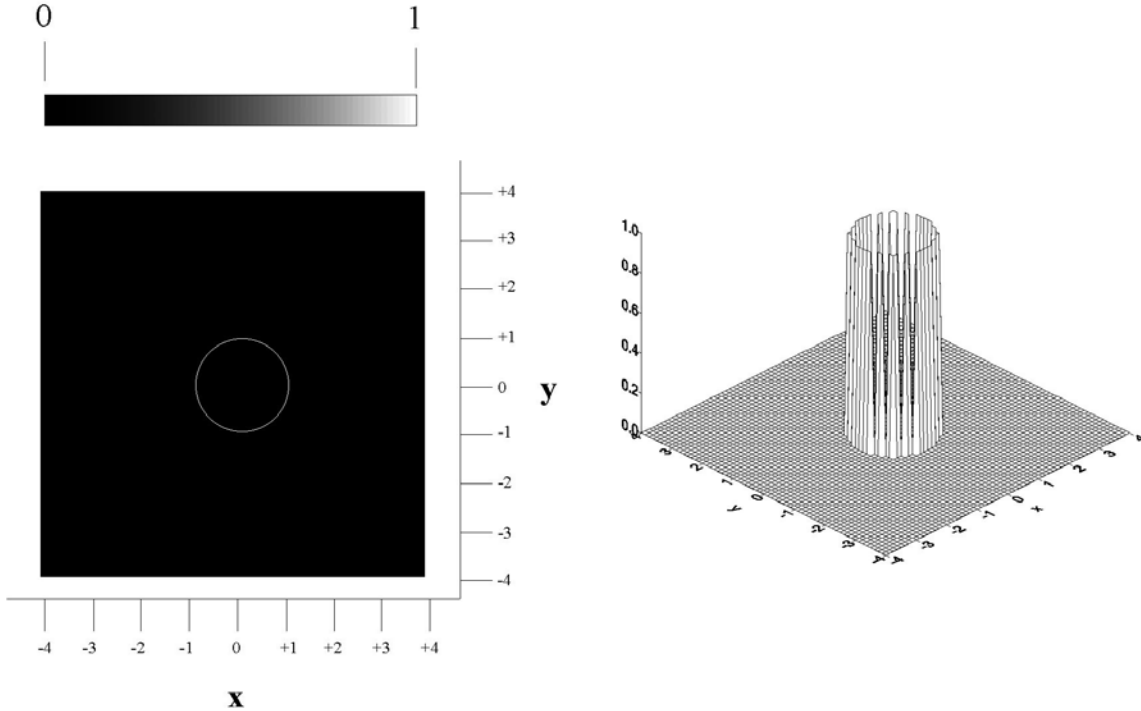
$$\begin{aligned} \int_{-\pi}^{+\pi} \int_0^{+\infty} \delta(r - r_0) 1[\theta] r dr d\theta &= \int_{-\frac{\pi}{2}}^{+\frac{\pi}{2}} \int_{-\infty}^{+\infty} \delta(r - r_0) 1[\theta] r dr d\theta \\ &= \pi \int_{-\infty}^{+\infty} [\delta(r + r_0) + \delta(r - r_0)] r dr = \pi(2r_0) = 2\pi r_0 \end{aligned}$$

- Scaled difference of two cylinder functions with diameters  $2r_0 + \Delta$  and  $2r_0 - \Delta$ :

$$f(r) = \lim_{\Delta \rightarrow 0} \left\{ \frac{1}{\Delta} \left[ CYL\left(\frac{r}{2r_0 + \Delta}\right) - CYL\left(\frac{r}{2r_0 - \Delta}\right) \right] \right\}$$

- Volume:

$$\begin{aligned} \int_{-\pi}^{+\pi} \int_0^{+\infty} f(r) r dr d\theta &= \lim_{\Delta \rightarrow 0} \left\{ \frac{1}{\Delta} \left[ \frac{\pi(2r_0 + \Delta)^2}{4} - \frac{\pi(2r_0 - \Delta)^2}{4} \right] \right\} \\ &= \lim_{\Delta \rightarrow 0} \left\{ \frac{\pi}{4\Delta} [(2r_0 + \Delta)^2 - (2r_0 - \Delta)^2] \right\} \\ &= \lim_{\Delta \rightarrow 0} \left\{ \frac{\pi}{4\Delta} [8r_0\Delta] \right\} = \lim_{\Delta \rightarrow 0} \{2\pi r_0\} = 2\pi r_0 \end{aligned}$$



2-D “ring” Dirac delta function  $f(r) = \delta(r - r_0)$ .

## 6 Complex-Valued 2-D Functions

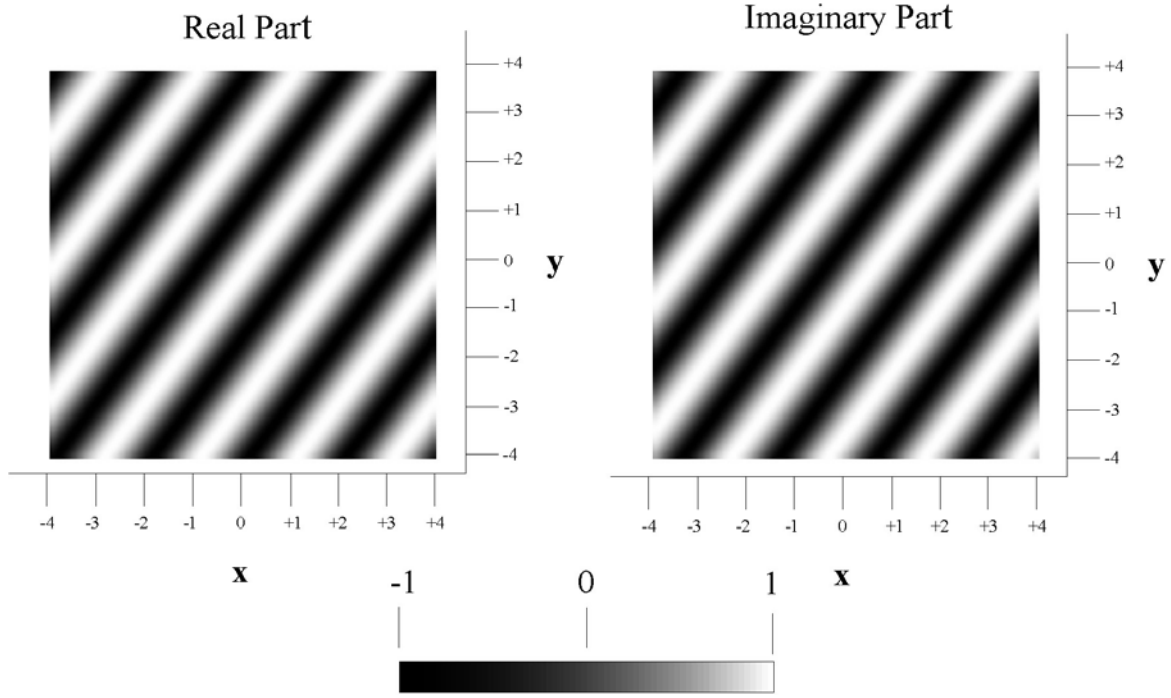
### 6.1 Complex 2-D Sinusoid

$$\begin{aligned} f[x, y] &= A_0 (\cos [2\pi\xi_0 x + \phi_0] + i \sin [2\pi\xi_0 x + \phi_0]) \times 1[y] \\ &= A_0 \left( \cos [2\pi\xi_0 x + \phi_0] + i \cos \left[ 2\pi\xi_0 x + \phi_0 - \frac{\pi}{2} \right] \right) \times 1[y] \end{aligned}$$

$$f[x, y] = A_0 e^{i(2\pi\xi_0 x + \phi_0)} 1[y]$$

Azimuth may be oriented in an arbitrary direction:

$$\begin{aligned} f[x, y] &= \cos [2\pi\xi_0 (\mathbf{r} \bullet \hat{\mathbf{p}})] + i \cos \left[ 2\pi\xi_0 (\mathbf{r} \bullet \hat{\mathbf{p}}) - \frac{\pi}{2} \right] \\ &= \cos [2\pi (\xi_0 x + \eta_0 y)] + i \sin [2\pi (\xi_0 x + \eta_0 y)] = e^{+2\pi i(\xi_0 x + \eta_0 y)} \end{aligned}$$



Real and imaginary parts of 2-D sinusoid.

## 6.2 2-D Complex Quadratic-Phase Sinusoid (“Chirp”)

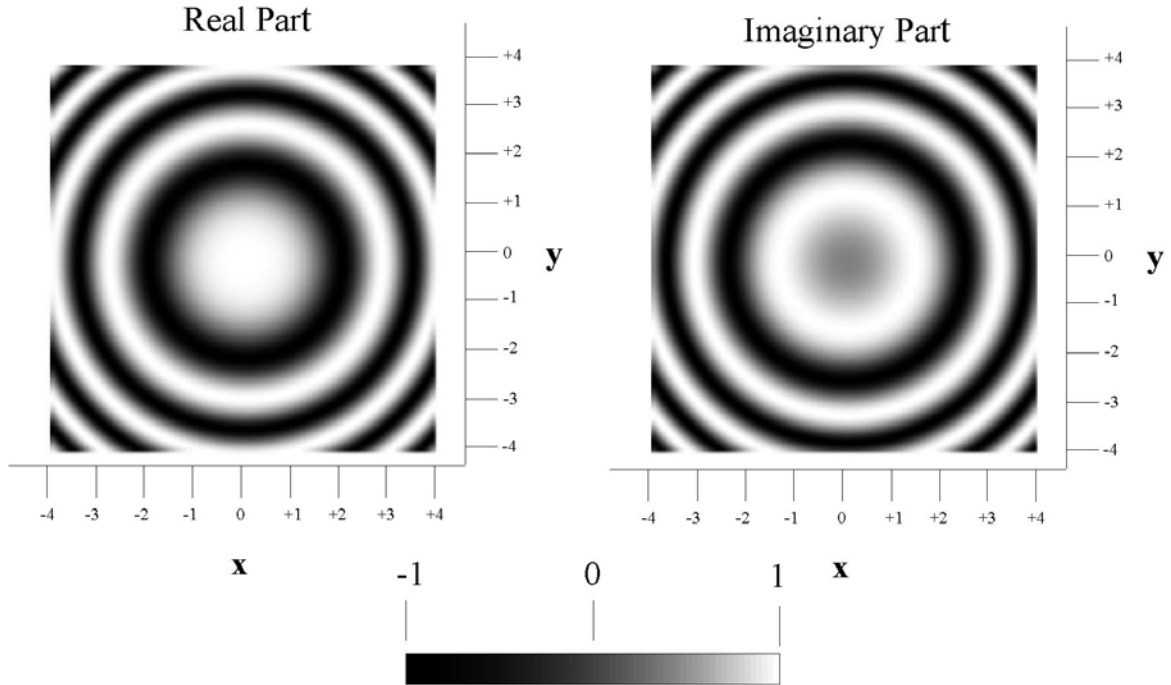
- Important in describing Fresnel diffraction
- 2-D complex “chirp” function:

$$\begin{aligned} f[x, y] &= A e^{\pm i\pi \left(\frac{x}{a}\right)^2} 1[y] \\ &= A_0 \left( \cos \left[ \frac{\pi x^2}{a^2} \right] \pm i \sin \left[ \frac{\pi x^2}{a^2} \right] \right) \times 1[y] \end{aligned}$$

$$\begin{aligned} f[x, y] &= A_0 e^{\pm i\pi \left(\frac{x}{a}\right)^2} B_0 e^{\pm i\pi \left(\frac{y}{b}\right)^2} \\ &= A_0 B_0 e^{\pm i\pi \left( \left(\frac{x}{a}\right)^2 \pm \left(\frac{y}{b}\right)^2 \right)} \end{aligned}$$

- Magnitude of second function is not constant
  - Contours of constant magnitude are elliptical
- Set  $a$  and  $b$  to  $d$  to obtain circularly symmetric complex chirp function:

$$\begin{aligned} f[x, y] &= A_0 e^{\pm i\pi \left(\frac{x}{d}\right)^2} B_0 e^{\pm i\pi \left(\frac{y}{d}\right)^2} \\ &= A_0 B_0 e^{\pm i\pi \frac{(x^2 + y^2)}{d^2}} \end{aligned}$$



Real and imaginary parts of 2-D circularly symmetric chirp function  $f(r) = e^{+i\pi \left(\frac{r}{2}\right)^2}$ .