

Ash
01/3/19

ELECTROMAGNETIC THEORY

Vector Analysis

Any physical quantity may be represented either as a scalar or as a vector.

→ Scalar may be defined as a quantity which is completely characterised by its magnitude.

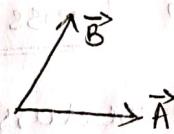
ex: mass, density, volume, area, time, temp etc...

→ Vector is defined as a quantity which is completely characterised by both magnitude & direction.

ex: force, velocity, displacement, acceleration, weight etc...

* vector quantities can be represented by a directed line segment (\vec{A}).

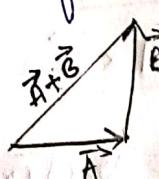
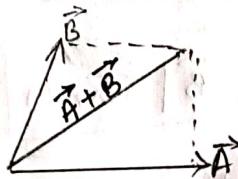
• The orientation depicts direction of the vector & length depicts the magnitude.



Addition of vectors:

→ The sum of 2 vectors results into a vector. if \vec{A} & \vec{B} are two vectors then;

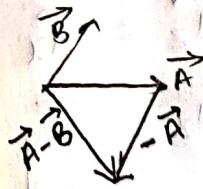
$$\vec{A} + \vec{B} = \vec{B} + \vec{A} \Rightarrow \vec{C}$$



Subtraction of vectors:

→ The difference b/w 2 vectors also results into a vector i.e

$$\vec{A} - \vec{B} = \vec{C} \quad \text{but} \quad \vec{A} - \vec{B} \neq \vec{B} - \vec{A}$$



Multiplication of vectors

→ Vector multiplication & scalar Multiplication

i) Scalar product or Dot product (\cdot):

By this multiplication of 2 vectors we get scalar.
Value. If \vec{A} & \vec{B} are two vectors, then their dot product is given by:

$$\vec{A} \cdot \vec{B} = |\vec{A}| |\vec{B}| \cos \theta$$

④ Dot product of unit vectors follow this rule:

$$\vec{a}_i \cdot \vec{a}_j = 1 \Rightarrow \text{if } i=j$$

$$\vec{a}_i \cdot \vec{a}_j = 0 \Rightarrow \text{if } i \neq j$$

$$\begin{aligned} \rightarrow \vec{A} \cdot \vec{B} &= (A_x \vec{a}_x + A_y \vec{a}_y + A_z \vec{a}_z) \cdot (B_x \vec{a}_x + B_y \vec{a}_y + B_z \vec{a}_z) \\ &= A_x B_x + A_y B_y + A_z B_z \end{aligned}$$

ii) Vector product or Cross product (\times):-

→ cross product of 2 vectors always results into a vector quantity in the direction normal to the plane of vectors. It is given by:

$$\vec{A} \times \vec{B} = |\vec{A}| |\vec{B}| \sin \theta$$

⑤ Unit vectors follow this rule for \times product:

$$\vec{a}_x \times \vec{a}_y = \vec{a}_z$$

$$\vec{a}_y \times \vec{a}_z = \vec{a}_x$$

$$\vec{a}_z \times \vec{a}_x = \vec{a}_y$$

$$\vec{a}_y \times \vec{a}_x = -\vec{a}_z$$

$$\vec{a}_z \times \vec{a}_y = -\vec{a}_x$$

$$\vec{a}_x \times \vec{a}_z = -\vec{a}_y$$

Coordinate Systems

- To describe a vector accurately & to express a vector in terms of its components, it is necessary to have some reference directions. Such directions are represented by coordinate systems.
- Coordinate systems are used to represent in three dimensional space.

There are 3 coordinate systems:

- a) Cartesian system
- b) Circular cylindrical system
- c) Spherical coordinate system

a) Cartesian System:-

- It has 3 axes, x, y, z mutually \perp to each other
- Orientation of x, y, z follow right hand thumb rule

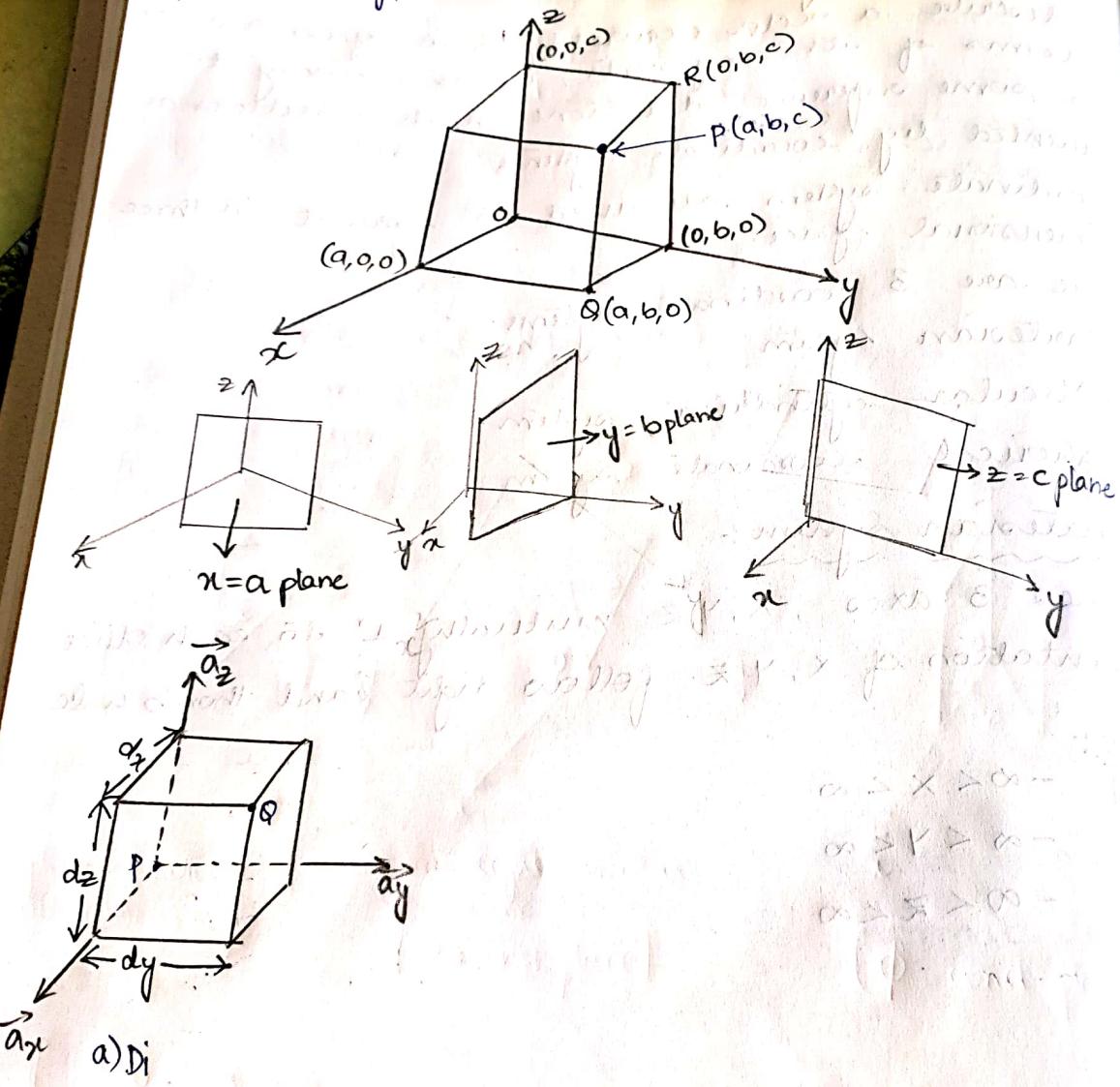
Range:-

$$-\infty < x < \infty$$

$$-\infty < y < \infty$$

$$-\infty < z < \infty$$

④ point $P(x, y, z)$ is the P.I. of 3 planes.



Differential length: $dl = dx \vec{a}_x + dy \vec{a}_y + dz \vec{a}_z$

Differential surfaces:

$$dS_x = dy dz \vec{a}_x$$

$$dS_y = dx dz \vec{a}_y$$

$$dS_z = dx dy \vec{a}_z$$

Diff. Volume: $dV = dx \cdot dy \cdot dz$

⑤ cylindrical
This is const
& $z = \text{const}$

⑥ A point P

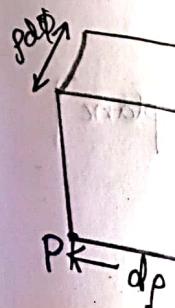
Range:

$$0 \leq x \leq a$$

$$0 \leq y \leq b$$

$$-\infty < z < \infty$$

A vector



Differential

② cylindrical coordinate system :-
 this is constructed by $\sigma = \text{const plane}$; $\phi = \text{const plane}$
 & $z = \text{const plane}$.

③ A point $P(\rho, \phi, z)$ is the intersection of three 3 planes

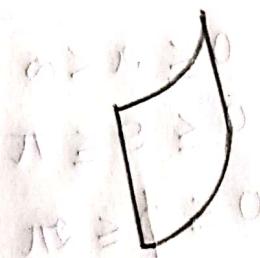
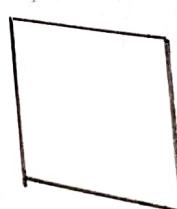
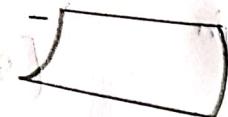
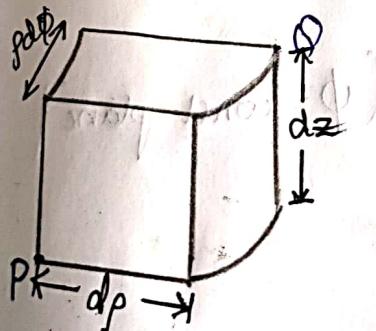
Range:-

$$0 \leq \rho < \infty$$

$$0 \leq \phi \leq 2\pi$$

$$-\infty < z < \infty$$

A vector is represented by: $OP = A\rho a_\rho + A\phi a_\phi + Az a_z$



Differential length: $dl = d\rho a_\rho + \rho d\phi a_\phi + dz a_z$

$$\rho dA + \rho dA + dA = \rho dA$$

Differential surfaces: $d\sigma = f d\phi dz \cdot a_\phi$ in ϕ -direction

$d\phi = df dz a_\phi$ in ϕ -direction

$dz = f df d\phi \cdot a_z$ in z -direction

Differential volume: $dV = f \cdot df d\phi dz$.

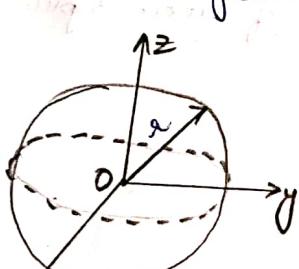
③ Spherical coordinate system :-

It is constructed by foll. 3 planes:

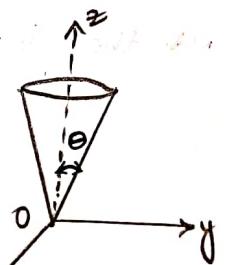
$r = \text{const}$ plane

$\theta = \text{const}$ plane

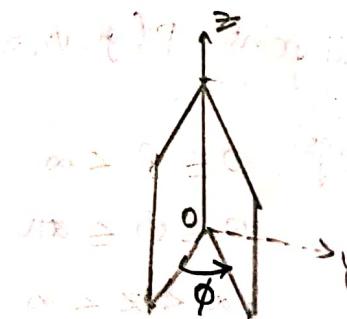
$\phi = \text{const}$ plane.



sphere of radius 'r'
with centre as origin



Right circular
cone
 $(\theta = \text{const. plane})$



$(\phi = \text{const. plane})$

Range:

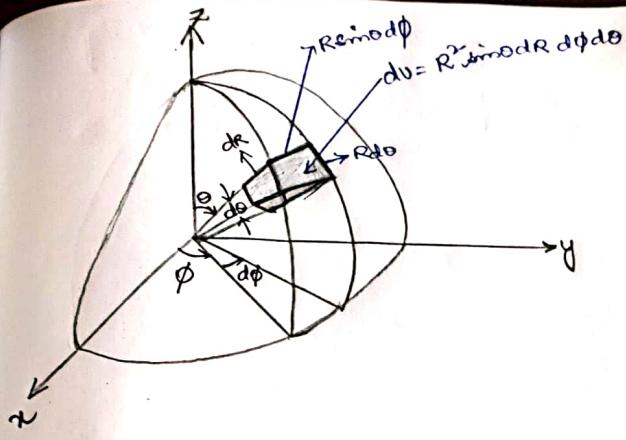
$$0 \leq r < \infty$$

$$0 \leq \theta \leq \pi$$

$$0 \leq \phi \leq 2\pi$$

→ Vector is represented by: $OP = A_r a_r + A_\theta a_\theta + A_\phi a_\phi$

in ρ -direction
in θ -direction
in ϕ -direction



$d\rho$ = differential change in ρ -direction

$d\theta$ = " " " " θ -direction

$d\sin\theta d\phi$ = " " " " ϕ -direction

Differential length: $dl = d\rho_a + \rho d\theta a_\theta + \rho \sin\theta d\phi a_\phi$

Differential surface:

$$ds_r = \rho^2 \sin\theta d\phi a_\theta$$

$$ds_\theta = \rho \sin\theta d\theta a_\theta$$

$$ds_\phi = \rho d\phi a_\phi$$

Dif. volume: $dv = \rho^2 \sin\theta d\rho d\theta d\phi$

[et-plane]

* ELECTROSTATICS

- Electromagnetism deals with study of charges at rest & in motion
- * principles of electromagnetism find its applications in disciplines such as microwaves, antenna, satellite comm., power generation, bio-electromagnetics, radar systems, EMI/EMC [electro-magnetic Interference] electronic warfare, fibre optics etc...
- A field produced by static charge is called "static electric field". It is an interesting phenomena with application in areas like electric power transmission, CRT, X-Ray machines, electric printers, lightning protection, solid state electronics, touchpads, capacitance, keyboards, LCD display, sorting machines in agriculture & mining etc...
- Static electric field is completely governed by two laws:
 - ⇒ Coulomb's law
 - ⇒ Gauss' law.

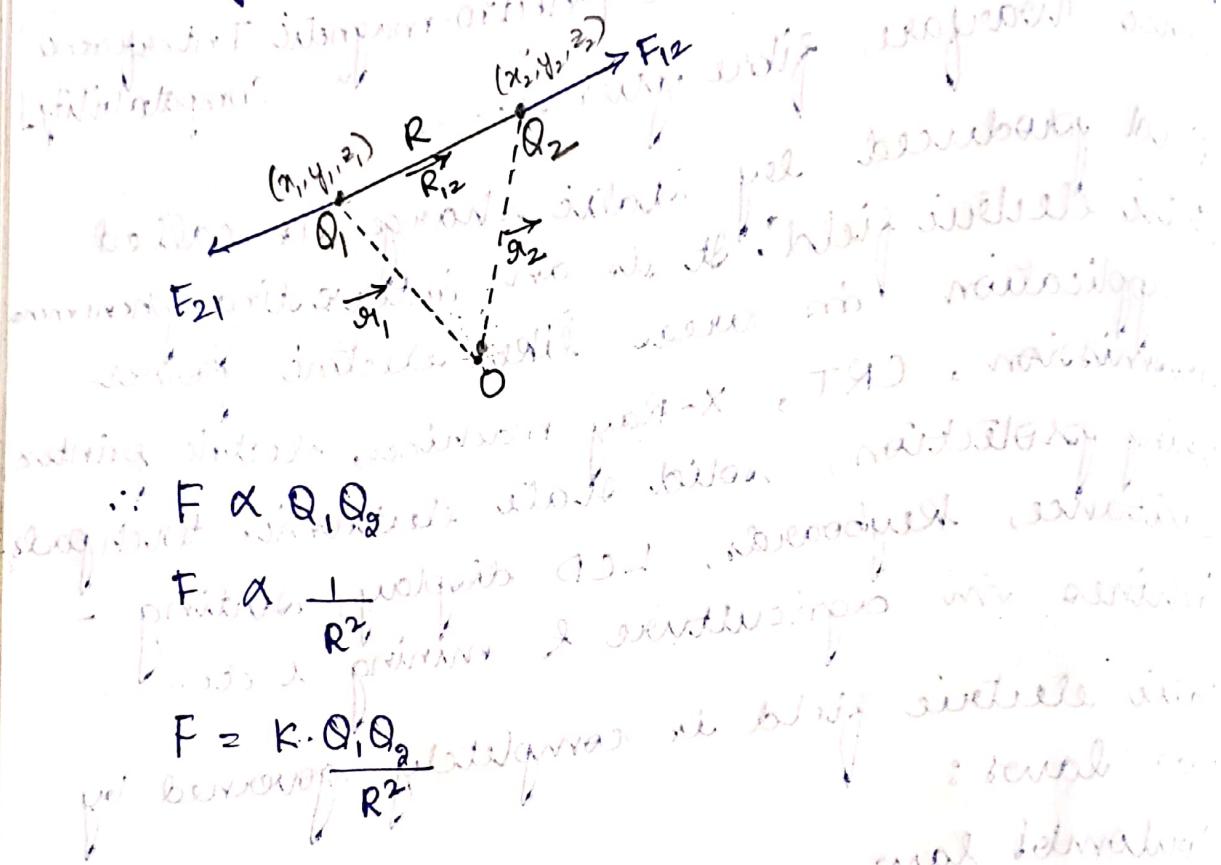
① COULOMB'S LAW:

It was formulated by Charles Augustin De Coulomb in 1785.

- It describes the force b/w two point charges
- Coulomb's law states that: "If two point charges Q_1 & Q_2 are separated by large distance 'R' in a vacuum or free space; then the force b/w

two charges is:

- 1) Directly proportional to the product of the charges Q_1 & Q_2 .
- 2) Inversely proportional to the square of distance between them.
- 3) Along the line joining them. [Direction of Force]



$$\text{where } K = \frac{1}{4\pi\epsilon}$$

$$\therefore E = E_0 \epsilon_0$$

for free space $\Rightarrow \epsilon_0 = 1/\epsilon_0$

$$\Rightarrow K = \frac{1}{4\pi\epsilon_0} \text{ coulombs per unit area of each charged sphere}$$
$$\therefore F = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q_1 Q_2}{R^2}$$

$$\vec{F}_{12} = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q_1 Q_2}{R^3} \vec{a}_{12}$$

$$\vec{a}_{12} = \frac{\vec{R}_{12}}{|R_{12}|} \Rightarrow \frac{\vec{R}_{12}}{R}$$

$$\vec{F}_{12} = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q_1 Q_2}{R^3} \cdot \vec{R}_{12} \Rightarrow \frac{1}{4\pi\epsilon_0} \frac{Q_1 Q_2}{(R_{12})^3} (\vec{r}_2 - \vec{r}_1)$$

∴ $\vec{F}_{21} = \frac{1}{4\pi\epsilon_0} \frac{Q_1 Q_2}{R^3} \cdot \vec{R}_{21} \Rightarrow \frac{1}{4\pi\epsilon_0} \frac{Q_1 Q_2}{(R_{21})^3} (\vec{r}_1 - \vec{r}_2)$

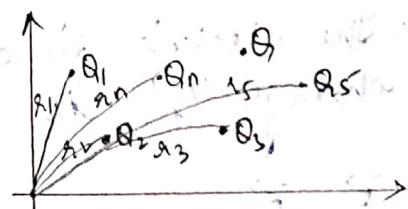
$$\therefore \boxed{\vec{F}_{21} = -\vec{F}_{12}}$$

If a charge 'Q' is placed in 'N' no. of charges, then the total force experienced by charge 'Q' due to $Q_1, Q_2, Q_3, \dots, Q_N$ is the vector sum of individual forces by each charge i.e.

$$F = F_1 + F_2 + F_3 + \dots + F_N$$

$$= \frac{1}{4\pi\epsilon_0} \frac{Q_1 Q}{R_1^3} \cdot \vec{R}_1 + \frac{1}{4\pi\epsilon_0} \frac{Q_2 Q}{R_2^3} \cdot \vec{R}_2 + \dots + \frac{1}{4\pi\epsilon_0} \frac{Q_N Q}{R_N^3} \cdot \vec{R}_N$$

$$\boxed{F = \frac{1}{4\pi\epsilon_0} \sum_{k=1}^N \frac{Q Q_k}{R_k^3} \cdot \vec{R}_k}$$



NOTE:

- ① As $\vec{F}_{21} = -\vec{F}_{12}$, the force of Q_2 on Q_1 is the force of Q_1 on Q_2 in magnitude but opposite in direction.
- ② Like charges repel each other and opposite charges attract.
- ③ The distance b/w the charges must be very large.
- ④ $Q_1, Q_2, Q_3, \dots, Q_N$ must be static.

- ⑤ The polarity of charges must be taken into consideration while calculating force
- ⑥ Force is measured in "Newton".

Electric Field Intensity :

- It is defined as the force per unit charge when placed in the electric field.
- It is also called as 'Electric field strength' denoted by \vec{E} & measured in "Newton/coulomb" or "Volts/meter". Hence $\vec{E} = \frac{\vec{F}}{Q}$
- Let 'Q' be any charge moved into the field of charge Q_1 , then the force on Q due to Q_1 is given by:

$$\vec{F}_r = \frac{1}{4\pi\epsilon_0} \cdot \frac{QQ_1}{R^3} \cdot \vec{R}$$

∴ The electric field intensity of Q_1 experienced by charge 'Q' is given by:

$$\vec{E} = \frac{\vec{F}_r}{Q} = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q_1}{R^3} \cdot \vec{R}$$

- Electric field intensity is a vector with radially outward direction from the source
- As the force on the test charge 'Q' is dependent on the magnitude of test charge, if the test charge becomes zero then, the force on it will also become 'zero'. But the force per unit charge remains same.

If a unit charge 'Q' is placed among the 'N' no. of charges, then the total electric field Intensity on the test charge is equal to the electric field intensity due to each charge.

$$\vec{E} = \vec{E}_1 + \vec{E}_2 + \vec{E}_3 + \dots + \vec{E}_N$$

$$= \frac{1}{4\pi\epsilon_0} \frac{Q_1}{R_1^3} \cdot \vec{R}_1 + \frac{1}{4\pi\epsilon_0} \frac{Q_2}{R_2^3} \cdot \vec{R}_2 + \dots + \frac{1}{4\pi\epsilon_0} \frac{Q_N}{R_N^3} \cdot \vec{R}_N$$

$$\therefore \vec{E} = \frac{1}{4\pi\epsilon_0} \sum_{k=1}^N \frac{Q_k}{R_k^3} \cdot \vec{R}_k$$

Electric field Intensity due to various charge distribution

There are basically 4 types of charge distribution

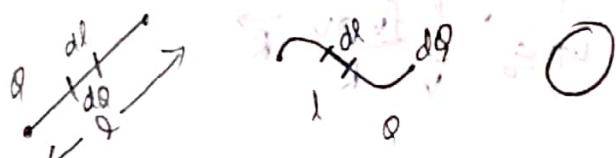
- 1) point charge
- 2) Line charge
- 3) Surface charge
- 4) Volume charge.

D) point charge:

A point charge distribution with no physical dimension is called point charge.

D) Line charge:

A charge distribution with single dimension is called a line charge distribution.



$$\text{Line charge density } \rho_L = \frac{Q}{l} = \frac{dq}{dl}$$

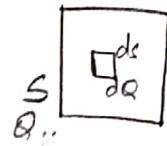
$$\Rightarrow l dq = \rho_L dl$$

$$\text{Total charge } Q = \int dQ = \int \rho_L dl$$

$$\text{Electric field intensity; } E = \frac{1}{4\pi\epsilon_0} \int \frac{\rho_L dl}{R^3} \cdot \vec{R}$$

③ Surface charge distribution: A 2D charge distribution that has some surface area is called surface charge distribution.

$$\rightarrow \text{surface charge density } \rho_s = \frac{Q}{S} = \frac{dQ}{ds}$$



$$dQ = \rho_s ds$$

$$\text{Total surface charge } Q = \int dQ = \iint \rho_s ds$$

$$\text{Electric Field Intensity; } E = \frac{1}{4\pi\epsilon_0} \int \frac{\rho_s ds}{R^3} \cdot \vec{R}$$

④ Volume charge distribution: A 3D charge distribution is called volume charge distribution.

$$\rightarrow \text{volume charge density } \rho_v = \frac{Q}{V} = \frac{dQ}{dv}$$

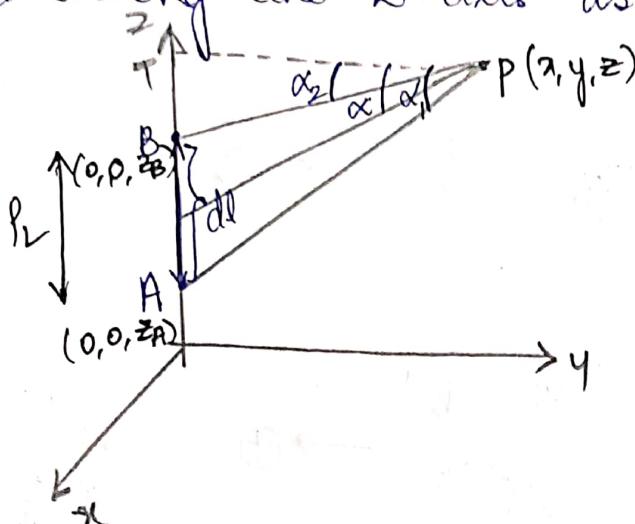


$$\text{Total volume charge } Q = \int dQ = \iiint \rho_v dv$$

$$\text{Electric field Intensity; } E = \frac{1}{4\pi\epsilon_0} \cdot \int \frac{\rho_v dv}{R^3} \cdot \vec{R}$$

Electric field Intensity due to a finite line charge:

consider a line charge of length 'l' & density ρ_L placed along the z-axis as shown in the fig.



To obtain the electric field intensity due to this line charge at a distant 'P(x,y,z)'. Consider an infinitesimal length 'dl' with charge 'dQ'

$$\rho_L = \frac{dQ}{dl}$$

$$dQ = \rho_L dl = \rho_L dz'$$

$$dE = \frac{1}{4\pi\epsilon_0} \cdot \frac{dQ}{R^3} \cdot \vec{R} \quad \rightarrow ①$$

$$\text{&} \quad \vec{R} = x\vec{a}_x + y\vec{a}_y + (z-z')\vec{a}_z$$

$$R = |\vec{R}| = \sqrt{x^2 + y^2 + (z-z')^2}$$

$$\text{also } \vec{TP} = x\vec{a}_x + y\vec{a}_y$$

$$\vec{TP} = \rho \vec{a}_p = \vec{p}$$

$$|\vec{p}| = \sqrt{p^2} = p$$

$$\therefore \vec{R} = f\vec{\alpha}_p + (z - z')\vec{\alpha}_z \quad \rightarrow \textcircled{2}$$

$$R = \sqrt{f^2 + (z - z')^2} \quad \rightarrow \textcircled{3}$$

$$\therefore \vec{dE} = \frac{1}{4\pi\epsilon_0} \cdot \frac{\beta_L dz}{(f^2 + (z - z')^2)^{3/2}} \cdot [f\vec{\alpha}_p + (z - z')\vec{\alpha}_z]$$

$$\vec{E} = \int_{(0,0,z_A)}^{(0,0,z_B)} \frac{f_L (f\vec{\alpha}_p + (z - z')\vec{\alpha}_z) \cdot dz'}{4\pi\epsilon_0 (f^2 + (z - z')^2)^{3/2}} \quad \rightarrow \textcircled{4}$$

from $\Delta \text{lie} = T P Q$

$$\sin \alpha = \frac{z - z'}{R}$$

$$\Rightarrow z - z' = R \sin \alpha \quad \rightarrow \textcircled{5}$$

$$\text{Also } \cos \alpha = \frac{f}{R} \Rightarrow R = \frac{f}{\cos \alpha}$$

sub 'R' in $\textcircled{5}$

$$z - z' = f \tan \alpha \quad \rightarrow \textcircled{5}$$

diff w.r.t z'

$$\Rightarrow 0 - 1 = f \sec^2 \alpha \cdot \frac{dx}{dz'}$$

$$\Rightarrow dz' = -f \sec^2 \alpha dx \quad \rightarrow \textcircled{6}$$

$\textcircled{7}$ when $z' = z_A$; $\boxed{\alpha = \alpha_1}$

$$z - z_A = f \tan \alpha_1$$

$$\therefore \alpha_1 = \tan^{-1} \left[\frac{z - z_A}{f} \right]$$

$\textcircled{8}$ when $z' = z_B$; $\boxed{\alpha = \alpha_2}$

$$\alpha_2 = \tan^{-1} \left[\frac{z - z_B}{f} \right]$$

sub $\textcircled{5}, \textcircled{6}, \textcircled{7}$

$$\vec{E} = \frac{\beta_L}{4\pi\epsilon_0} \int_{\alpha_1}^{\alpha_2} \frac{f_L (f\vec{\alpha}_p + (z - z')\vec{\alpha}_z)}{(f^2 + (z - z')^2)^{3/2}} dz$$

$$= \frac{-\beta_L}{4\pi\epsilon_0} \int_{\alpha_1}^{\alpha_2} f \vec{\alpha}_z dz$$

$$= \frac{-\beta_L}{4\pi\epsilon_0} \int_{x_1}^{x_2} f \vec{\alpha}_z dz$$

$$= \frac{-\beta_L}{4\pi\epsilon_0} \int_{x_1}^{x_2} f \vec{\alpha}_z dz$$

$$= \frac{-\beta_L}{4\pi\epsilon_0 f}$$

$$= \frac{-\beta_L}{4\pi\epsilon_0 f}$$

if
the

$$\therefore \alpha_1 = \tan^{-1} \left[\frac{z - z_A}{r} \right] \rightarrow \textcircled{4}$$

$\textcircled{5}$ when $z' = z_B$; $\alpha = \alpha_2$

$$\alpha_2 = \tan^{-1} \left[\frac{z - z_B}{r} \right] \rightarrow \textcircled{6}$$

subⁿ $\textcircled{5}, \textcircled{6}, \textcircled{4}$ & $\textcircled{6}$ in $\textcircled{4}$

$$\begin{aligned} \vec{E} &= \frac{\rho_L}{4\pi\epsilon_0} \int_{\alpha_1}^{\alpha_2} \frac{[\rho \vec{a}_p + r \tan \alpha \vec{a}_z]}{(r^2 + r^2 \tan^2 \alpha)^{3/2}} \cdot (-r \sec \alpha d\alpha) \\ &= \frac{-\rho_L}{4\pi\epsilon_0} \int_{\alpha_1}^{\alpha_2} \frac{r [\vec{a}_p + \tan \alpha \vec{a}_z]}{r^3 \sec^3 \alpha} \cdot r \sec \alpha d\alpha \\ &= \frac{-\rho_L}{4\pi\epsilon_0} \int_{\alpha_1}^{\alpha_2} \frac{(\vec{a}_p + \tan \alpha \vec{a}_z)}{\sec \alpha} \cdot d\alpha \\ &= \frac{-\rho_L}{4\pi\epsilon_0} \int_{\alpha_1}^{\alpha_2} (\cos \alpha \vec{a}_p + \sin \alpha \vec{a}_z) d\alpha \\ &= \frac{-\rho_L}{4\pi\epsilon_0 r} \left[[\sin \alpha \vec{a}_p]_{\alpha_1}^{\alpha_2} - [\cos \alpha \vec{a}_z]_{\alpha_1}^{\alpha_2} \right] \\ &= \frac{-\rho_L}{4\pi\epsilon_0 r} \left[(\sin \alpha_2 - \sin \alpha_1) \vec{a}_p - (\cos \alpha_2 - \cos \alpha_1) \vec{a}_z \right] \rightarrow \textcircled{7} \end{aligned}$$

if line charge is infinite
then A(0, 0, ∞) & B(0, 0, ∞)

$$\therefore z_A = -\infty \quad z_B = \infty$$

$$\Rightarrow \alpha_1 = \tan^{-1} \left[\frac{z - z_A}{r} \right] = \tan^{-1}(\infty) = \pi/2$$

$$\Rightarrow \alpha_2 = \tan^{-1} \left[\frac{z - z_B}{r} \right] = \tan^{-1}(-\infty) = -\pi/2$$

sub^n in ①

$$\therefore \vec{E} = \frac{-\rho_L}{4\pi\epsilon_0 r} \left[(-1-1) \vec{a}_p - (0) \vec{a}_z \right]$$

$$= \frac{\rho_L}{4\pi\epsilon_0 r} \cdot \vec{a}_p$$

$$\boxed{\vec{E} = \frac{\rho_L}{4\pi\epsilon_0 r} \cdot \vec{a}_p}$$

\therefore The electric field intensity will always be towards radial direction.

Electric field Intensity due to Infinite Surface charge:

Consider an infinite sheet of charge on X-Y plane with uniform charge density ' s '

- To obtain the electric field intensity due to this arrangement, consider infinitesimal incremental surface ' ds ' with charge ' dQ ' as shown in the fig.

$$ds = \rho d\rho d\phi \rightarrow ①$$

$$\vec{R} = (0-\rho)\vec{a}_\rho + (h-0)\vec{a}_z$$

$$= -\rho \vec{a}_\rho + h \vec{a}_z \rightarrow ②$$

$$R = |\vec{R}| = \sqrt{\rho^2 + h^2} \rightarrow ③$$

$$\vec{dE} = \frac{1}{4\pi\epsilon_0} \cdot \frac{dQ}{R^3} \cdot \vec{R} \rightarrow ④$$

$$f_s = \frac{dQ}{ds} \Rightarrow dQ = f_s ds \rightarrow ⑤$$

subⁿ ①, ②, ③ & ⑤ in ④

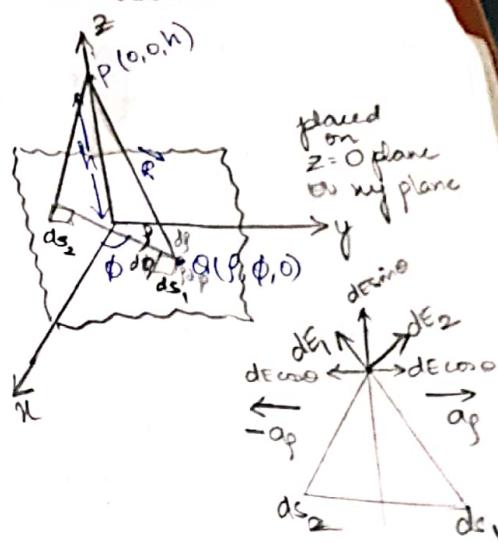
$$\Rightarrow \vec{dE} = \frac{1}{4\pi\epsilon_0} \cdot \frac{f_s \rho d\rho d\phi}{(\rho^2 + h^2)^{3/2}} (-\rho \vec{a}_\rho + h \vec{a}_z)$$

Total field due to infinite surface charge is

$$\vec{E} = \iint \vec{dE}$$

$$\vec{E} = \int_{\rho=0}^{\infty} \int_{\phi=0}^{2\pi} \frac{1}{4\pi\epsilon_0} \cdot \frac{f_s \rho d\rho d\phi}{(\rho^2 + h^2)^{3/2}} (-\rho \vec{a}_\rho + h \vec{a}_z)$$

The total electric field intensity due to infinite surface charge is the integration of the above eqn over the limits: $\rho \rightarrow 0 \text{ to } \infty$ & $\phi \rightarrow 0 \text{ to } 2\pi$



Due to a symmetry of charge distribution,
 "ap" component due to charge on surface '1' is
 equal to "ap" component due to charge on
 surface '2'. But, in opposite direction resulting
 in "zero" ap component.

$$\therefore \vec{E} = \frac{\rho_s}{4\pi\epsilon_0} \int_{r=0}^{\infty} \int_{\phi=0}^{2\pi} \frac{\rho dr d\phi h \hat{a}_z}{(r+h)^{3/2}}$$

$$= \frac{\rho_s h}{4\pi\epsilon_0} \int_{r=0}^{\infty} \frac{\rho dr \hat{a}_z}{(r+h)^{3/2}} \int_{\phi=0}^{2\pi} d\phi \xrightarrow{(2\pi-0) = 2\pi}$$

$$= \frac{\rho_s h}{4\pi\epsilon_0} \cdot 2\pi \int_{r=0}^{\infty} \frac{\rho dr \hat{a}_z}{(r+h)^{3/2}}$$

$$= \frac{\rho_s h}{4\pi\epsilon_0} \hat{a}_z \int_{r=0}^{\infty} \frac{\rho dr}{(r+h)^{3/2}}$$

$$\text{let } r+h = t^2$$

diff w.r.t 't'

$$dr = dt \frac{dt}{dp}$$

$$\rho dr = t dt$$

$$\text{when } r \rightarrow 0 ; t = h \\ r \rightarrow \infty ; t = \infty$$

tribution,
urface '1' is
charge on
resisting

$$\begin{aligned}\vec{E} &= \frac{\rho_s h \hat{a}_z}{2\epsilon_0} \int_h^\infty \frac{t dt}{t^3} \\&= \frac{\rho_s h \hat{a}_z}{2\epsilon_0} \int_h^\infty t^{-2} dt \\&= \frac{\rho_s h \hat{a}_z}{2\epsilon_0} \left[\frac{-1}{t} \right]_h^\infty \\&= -\frac{\rho_s h \hat{a}_z}{2\epsilon_0} \left[\frac{1}{t} \right]_h^\infty \\&= -\frac{\rho_s h \hat{a}_z}{2\epsilon_0} (0 - \frac{1}{h}) \\&\therefore \vec{E} = \frac{\rho_s \hat{a}_z}{2\epsilon_0}\end{aligned}$$

problem:

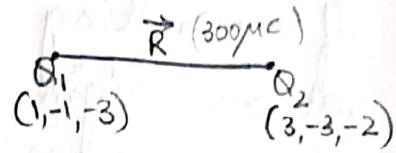
- Q1) Two point charges of magnitude '2 millicoulomb & -3 millicoulomb' are located at pt $(4, 4, -5)$ & $(-3, 2, -9)$ respectively.
Evaluate
- Distance b/w two points
 - Vector force on charge at point 2
 - Vector force on charge at point 1

negative & composite. It is nothing but a composite diapositive
of the original (1 - part) & it is not the true photograph.

It is a composite of the three photographs taken by the
original camera. It is a composite of the three photographs taken by the
original camera. It is a composite of the three photographs taken by the

A point charge Q_1 is equal to $300 \mu\text{coulomb}$ is located at $(1, -1, -3)$ experiences a force $\vec{F} = 8\vec{a}_x - 8\vec{a}_y + 4\vec{a}_z$ at $(3, -3, -2)$. Determine Q_2

$$F = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q_1 Q_2}{R^3} \cdot \vec{R} \quad \text{--- (1)}$$



$$\text{where } \vec{R} = (1-3)\vec{a}_x + (-1+3)\vec{a}_y + (-3+2)\vec{a}_z \\ = -2\vec{a}_x + 2\vec{a}_y - 1\vec{a}_z$$

$$R = |\vec{R}| = \sqrt{4+4+1} = 3$$

subⁿ in (1)

$$(8\vec{a}_x - 8\vec{a}_y + 4\vec{a}_z) = \frac{1}{4\pi\epsilon_0} \cdot \frac{300 \times 10^{-6} \times Q_2}{3^3} [-2\vec{a}_x + 2\vec{a}_y - 1\vec{a}_z]$$

$$\epsilon_0 = 8.854 \times 10^{-12}$$

$$-4(-2\vec{a}_x + 2\vec{a}_y - \vec{a}_z) = \frac{1}{4\pi \times 8.854 \times 10^{-12}} \cdot \frac{300 \times 10^{-6} \times Q_2}{(24)} [-2\vec{a}_x + 2\vec{a}_y - \vec{a}_z]$$

$$-4 = \frac{9 \times 10^9 \times 300 \times 10^{-6} \times Q_2}{8 \times 3}$$

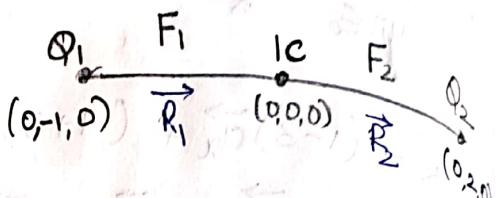
$$-4 = 100 Q_2 \times 10^3$$

$$Q_2 = \frac{-4}{10^5} = -4 \times 10^{-5}$$

$$\therefore Q_2 = -40 \mu\text{coulomb}$$

Q.) A charge Q_1 at $(0, -1, 0)$ & another charge Q_2 at $(0, 2, 0)$ are located. Find the ratio if resulting force on a $1C$ test charge at an origin is zero. All the charges are placed of same polarity.

$$\therefore F = F_1 + F_2 = 0 \quad \& \boxed{Q = 1C}$$



$$\& F_1 = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q_1 Q}{R_1^3} \cdot \vec{R}_1 ; \quad F_2 = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q_2 Q}{R_2^3} \cdot \vec{R}_2$$

$$\text{let } \vec{R}_1 = \vec{a}_y$$

$$R_1 = \sqrt{1} = 1$$

$$\vec{R}_2 = -2\vec{a}_y$$

$$R_2 = \sqrt{4} = 2$$

$$F_1 = 9 \times 10^9 \cdot \frac{Q_1}{1^3} \cdot \vec{a}_y ; \quad F_2 = 9 \times 10^9 \cdot \frac{Q_2}{2^3} (-2\vec{a}_y)$$

$$= -9 \times 10^9 \frac{Q_2}{4} \vec{a}_y$$

$$\therefore \boxed{F_1 + F_2 = 0}$$

$$\Rightarrow 9 \times 10^9 Q_1 \vec{a}_y - 9 \times 10^9 Q_2 \vec{a}_y$$

$$9 \times 10^9 \left[Q_1 - \frac{Q_2}{4} \right] = 0$$

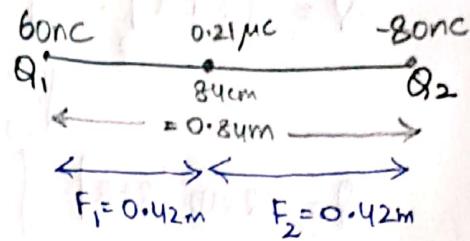
$$Q_1 - \frac{Q_2}{4} = 0$$

$$Q_1 = \frac{Q_2}{4}$$

$$\therefore \boxed{\frac{Q_2}{Q_1} = \frac{1}{4}}$$

A 1) A charge of 60nC is 84cm apart from a charge of -80nC . What force will these charges of $0.21\mu\text{C}$ placed half way b/w them.
 (Take air as medium)

$$\therefore F_1 = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q_1 Q_2}{R^3} \cdot \vec{R}$$



$$F_1 = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q_1 Q_3}{R^2}$$

$$\therefore F_2 = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q_2 Q_3}{R_2^2}$$

$$F_1 = \frac{9 \times 10^9 \times 60 \times 10^{-9} \times 0.21 \times 10^{-6}}{(0.42)^2}$$

$$F_2 = \frac{9 \times 10^9 \times (-80 \times 10^{-9}) \times 0.21 \times 10^{-6}}{(0.42)^2}$$

$$F_1 = 0.642 \times 10^{-3} \text{ N}$$

$$F_2 = -0.857 \times 10^{-3} \text{ N}$$

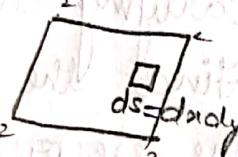
$$\therefore |F| = |F_1| + |F_2|$$

$$= 0.642 \times 10^{-3} + 0.857 \times 10^{-3}$$

$$F = 1.5 \text{ mN away from } Q_1$$

2) A square plane described by $-2 \leq x \leq 2$, $-2 \leq y \leq 2$ & $z=0$ carries a charge density $12|y| \text{ nC/m}^2$. Find the total charge on the plane.

$$\text{W.K.T } Q = \iint f_s \, ds \text{ where } f_s = 12|y| \text{ nC/m}^2$$



$$\therefore Q = \int_{x=-2}^2 \int_{y=-2}^2 12|y| \times 10^{-9} \, dx \cdot dy$$

$$Q = 12 \times 10^{-3} \int_{x=-2}^2 dx \int_{y=-2}^2 |y| dy$$

$$Q = 12 \times 10^{-3} [x]_{-2}^2 \int_{y=-2}^2 |y| dy$$

$$Q = 12 \times 10^{-3} [x+2]_{-2}^{+2} \int_{y=-2}^2 |y| dy$$

$$Q = 48 \times 10^{-3} \int_{-2}^{+2} |y| dy$$

$$Q = 48 \times 10^{-3} \left[- \int_{-2}^0 y dy + \int_0^2 y dy \right]$$

$$Q = 48 \times 10^{-3} \left[- \left(\frac{y^2}{2} \right) \Big|_{-2}^0 + \left(\frac{y^2}{2} \right) \Big|_0^2 \right]$$

$$Q = 48 \times 10^{-3} \left[- \left(\frac{0 - (-2)^2}{2} \right) + \frac{2^2}{2} \right]$$

$$= 48 \times 10^{-3} \left[\frac{4}{2} + \frac{4}{2} \right]$$

$$= 48 \times 10^{-3} \left[4 \left(\frac{1}{2} \right) \right]$$

$$\boxed{Q = 19.2 \text{ mC}}$$

- Q) A finite sheet $0 \leq x \leq 1, 0 \leq y \leq 1$ lie on $[z=0]$ plane carrying charge density $\rho_s = xy(x^2 + y^2 + 25)^{3/2}$ nc/m².
- i) Find the total charge on the sheet.
 - ii) The electric field at $(0, 0, 5)$
 - iii) The force experienced by a -1mc charge located at $(0, 0, 5)$

$$\text{Q} = \iint_S p_0 \, ds \quad \text{where } ds = dx \, dy$$

$$= \int_{x=0}^1 \int_{y=0}^1 xy (x^2 + y^2 + 25)^{3/2} \times 10^{-9} \, dx \, dy$$

$$\text{let } x^2 = t, \quad y^2 = K$$

$$2x = \frac{dt}{dx}, \quad 2y = \frac{dK}{dy}$$

$$x \, dx = \frac{dt}{2}, \quad y \, dy = \frac{dK}{2}$$

$$\left. \begin{array}{l} \text{if } x=0; \, t=0 \\ x=1; \, t=1 \end{array} \right| \quad \left. \begin{array}{l} \text{if } y=0; \, K=0 \\ y=1; \, K=1 \end{array} \right|$$

$$\Rightarrow Q = \int_{t=0}^1 \int_{K=0}^1 \frac{(t+K+25)^{3/2}}{4} dt \cdot dK \times 10^{-9}$$

$$= \frac{10^{-9}}{4} \int_{K=0}^1 \left[\frac{(t+K+25)^{5/2}}{5/2} \right]_0^1 dK$$

$$= \frac{10^{-9}}{10} \int_{K=0}^1 \left[(K+26)^{5/2} - (K+25)^{5/2} \right] dK$$

$$= \frac{10^{-9}}{10} \left[\left[\frac{(K+26)^{7/2}}{7/2} \right]_0^1 - \left[\frac{(K+25)^{7/2}}{7/2} \right]_0^1 \right]$$

$$= \frac{10^{-9}}{35} \left[(27)^{7/2} - (26)^{7/2} - (26)^{7/2} + (25)^{7/2} \right]$$

$$\boxed{Q = 33.14 \text{ nC}}$$

iii.

$$E = \iint \frac{1}{4\pi\epsilon_0} \frac{\rho_s ds}{R^3} \cdot \vec{R}$$

$$\vec{R} = -x\vec{a}_x - y\vec{a}_y + 5\vec{a}_z$$

$$R = \sqrt{x^2 + y^2 + 25}$$

$$\Rightarrow E = \int_{x=0}^1 \int_{y=0}^1 \frac{xy(x^2+y^2+25)^{-1/2} (-x\vec{a}_x - y\vec{a}_y + 5\vec{a}_z)}{4\pi\epsilon_0 (x^2+y^2+25)^{3/2}} dx dy \times 10^9$$

$$E = \frac{10^{-9}}{4\pi\epsilon_0} \int_{x=0}^1 \int_{y=0}^1 -x^2 y \vec{a}_x - xy^2 \vec{a}_y + 5xy \vec{a}_z dy dx$$

$$= \frac{10^{-9}}{4\pi\epsilon_0} \left[- \int_{x=0}^1 \int_{y=0}^1 x^2 y \vec{a}_x dy dx - \int_{x=0}^1 \int_{y=0}^1 xy^2 \vec{a}_y dy dx + 5 \int_{x=0}^1 \int_{y=0}^1 xy \vec{a}_z dy dx \right]$$

$$= \frac{10^{-9}}{4\pi\epsilon_0} \left[-\frac{1}{2} \int_{x=0}^1 x^3 \vec{a}_x - \frac{1}{2} \int_{y=0}^1 y^3 \vec{a}_y + \frac{5}{2} \int_{x=0}^1 x^2 \vec{a}_z \right]$$

$$= \frac{10^{-9}}{4\pi\epsilon_0} \left[-\frac{1}{6} \vec{a}_x - \frac{1}{6} \vec{a}_y + \frac{5}{4} \vec{a}_z \right]$$

$$= 9 \times 10^9 \times 10^{-9} \left[-0.16 \vec{a}_x - 0.16 \vec{a}_y + 1.25 \vec{a}_z \right]$$

| | | | | |
|--|---|---|----|-----|
| $\vec{E} = -1.5 \vec{a}_x - 1.5 \vec{a}_y + 11.25 \vec{a}_z$ | N | C | on | V/m |
|--|---|---|----|-----|

$$\text{iii) } E = \frac{F}{Q}$$

$$F = EQ$$

$$F = (1.54\vec{a}_x + 1.54\vec{a}_y - 11.25\vec{a}_z) \times 10^3 \text{ N}$$

A uniform charge distribution infinite in extent lies along the x-axis with $\rho_L = 20 \text{ n coulomb/m}$. Find the electric field at (6, 8, 3) m

$$\text{Sol: } \vec{E} = \frac{\rho_L}{2\pi\epsilon_0 r} \cdot \vec{a}_r$$

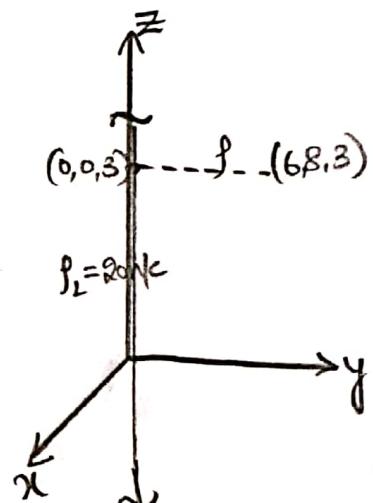
$$\vec{r} = 6\vec{a}_x + 8\vec{a}_y$$

$$r = \sqrt{36+64} = 10$$

$$\begin{aligned} \Rightarrow \vec{a}_r &= \frac{\vec{r}}{r} = \frac{6\vec{a}_x + 8\vec{a}_y}{10} \\ &= 0.6\vec{a}_x + 0.8\vec{a}_y \end{aligned}$$

$$\therefore \vec{E} = \frac{20 \times 10^{-9}}{2 \times \pi \times 8.854 \times 10^{-12} \times 10} (0.6\vec{a}_x + 0.8\vec{a}_y)$$

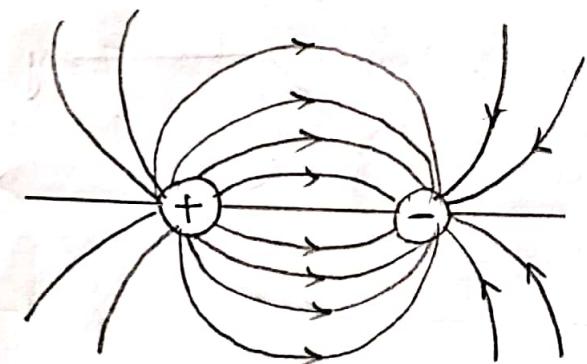
$$= 35.96 (0.6\vec{a}_x + 0.8\vec{a}_y)$$



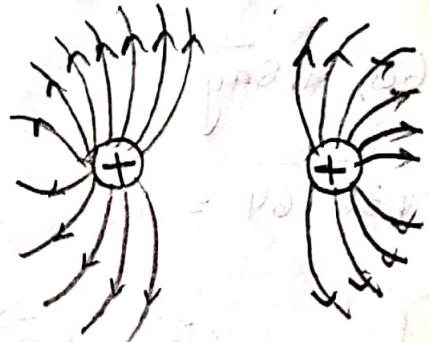
$$\boxed{\vec{E} = 21.57\vec{a}_x + 28.76\vec{a}_y \text{ N/C or V/m}}$$

Electric flux: [electric displacement] :-

- Field lines emanate from a point charge symmetrically in all directions
- Field lines originate on positive charges and terminate on negative ones
- They cannot simply stop in midair, though they may extend out to infinity
- Field lines can never cross.



equal but opposite charges



equal charges

→ These lines of forces constituting an electric field are collectively termed as "electric flux". It is denoted by ' Ψ '.

→ Electric flux is found to be directly proportional to charge Q in S.I units, the value of proportionality constant is '1'.

$$\therefore \boxed{\Psi = Q}$$

Q is measured in coulombs

Electric flux density:

As there are definite no. of lines of force i.e. flux, passed through unit surface area which is normal to the direction of flux, Electric-flux density is defined as flux per unit normal area. denoted by " \vec{D} " measured in "coulombs per m²".

$$\text{i.e. } \vec{D} = \frac{\Psi}{S} \quad \text{or} \quad \vec{D} = \frac{d\Psi}{ds}$$



$$\therefore \Psi = \oint d\Psi = \oint \vec{D} \cdot d\vec{s}$$

GAUSS LAW:

"Gauss law states that total electric flux " Ψ " through any closed surface is equal to the charge enclosed by that surface".

$$\Psi = Q_{\text{enclosed}} \quad \text{--- (1)}$$

$$\text{i.e. } \Psi = \oint d\Psi = \oint \vec{D} \cdot d\vec{s}$$

$$= \text{total charge enclosed } Q = \int p_v dv \quad \text{--- (2)}$$

$$\text{or } Q = \oint \vec{D} \cdot d\vec{s} = \int p_v dv \quad \text{--- (3)}$$

From the divergence theorem:

$$\oint \vec{D} \cdot d\vec{s} = \int \nabla \cdot \vec{D} dv \quad \text{--- (4)}$$

$$p_v = \nabla \cdot \vec{D} \rightarrow \text{Maxwell's 1st eqn}$$

∇ = diff operator

∴ GAUSS LAW is an alternative statement of Coulomb's law

- procedure for Applying GAUSS's Law's
- To calculate the electric field using Gauss law
 - 1st know the symmetric axis
 - If the symmetric charge distribution exists, consider a mathematical closed surface known as Gaussian surface consisting of field point
 - Choose the surface such that displacement density is normal or tangential to Gaussian surface.

→ If 'D' is normal to surface : $\vec{D} \cdot d\vec{s} = D ds$

→ If 'D' is tangent to surface : $\vec{D} \cdot d\vec{s} = 0$

E [electric field Intensity] due to point charge:

w.k.t

$$\Psi = Q_{\text{enclosed}} \rightarrow ①$$

$$Q_{\text{enclosed}} = Q \rightarrow ②$$

$$\Psi = \oint \vec{D} \cdot d\vec{s}$$

where $\vec{D} = D_r \hat{a}_r$ & $d\vec{s} = r \sin\theta d\theta d\phi \hat{a}_\phi$

$$\Rightarrow \Psi = \oint D_r \hat{a}_r \cdot r \sin\theta d\theta d\phi \hat{a}_\phi$$

$$\Psi = \int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} D_r r^2 \sin\theta d\theta d\phi$$

$$= -2\pi D_k r^2 [\cos \theta]_0^\pi$$

$$= -2\pi D_k r^2 [-1 - 1]$$

$$\therefore \psi = 4\pi D_k r^2 \rightarrow \textcircled{3}$$

but from ①

$$\Rightarrow U\pi D_k r^2 = Q$$

$$D_k = \frac{Q}{U\pi r^2}$$

$$\vec{D} = \frac{Q}{U\pi r^2} \cdot \vec{r}$$

w.r.t

$$\vec{E} = \frac{Q}{U\pi \epsilon_0 r^2} \vec{r}$$

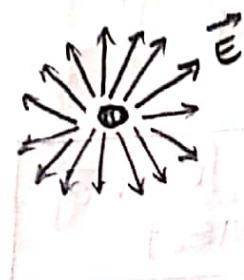
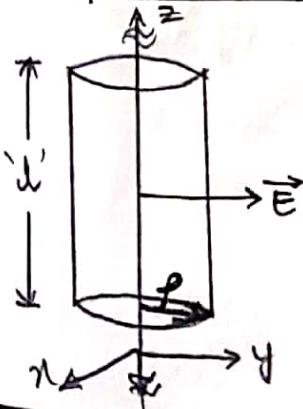
$$\vec{E} = \frac{\vec{D}}{\epsilon_0}$$

$$\boxed{\vec{D} = \epsilon_0 \vec{E}}$$

Applications of Gauss law:

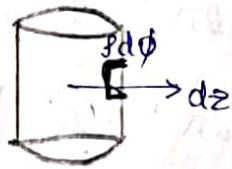
Gauss law is particularly useful in computing $E \propto D$, where the charge distribution has some symmetry

AN INFINITE LINE CHARGE:



W.K.T

$$\Psi = Q_{\text{enclosed}} \rightarrow ①$$



$$Q_{\text{enclosed}} = f_l \cdot l$$

$$\Psi = \oint \vec{D} \cdot d\vec{s} \rightarrow ②$$

$$l \vec{D} = D_p \cdot \vec{a}_p \quad \& \quad d\vec{s} = \int d\phi dz \vec{a}_p$$

sub in ②

$$\begin{aligned} \Psi &= \oint D_p \vec{a}_p \cdot \int d\phi dz \cdot \vec{a}_p \\ &= \int_{z=0}^l \int_{\phi=0}^{2\pi} D_p \cdot p d\phi dz \\ &= 2\pi D_p \cdot p [z]_0^l \end{aligned}$$

$$\Psi = 2\pi D_p p \cdot l$$

From ①:

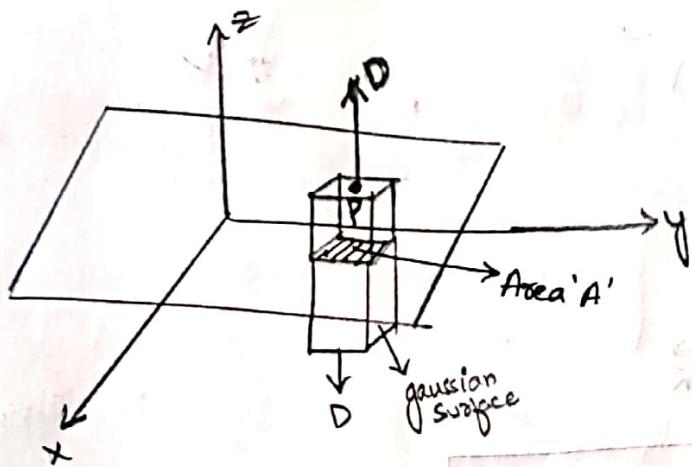
$$2\pi D_p \cdot p l = P_e l$$

$$D_p = \frac{P_e}{2\pi p}$$

$$\vec{D} = \frac{P_e}{2\pi p} \vec{a}_p$$

$$\therefore \vec{E} = \frac{\vec{D}}{\epsilon_0} = \frac{P_e}{2\pi \epsilon_0 p} \cdot \vec{a}_p$$

Electric field intensity due to Infinite Surface Charge:



w.k.t

$$\Psi = Q_{\text{enclosed}} \rightarrow ①$$

$$Q_{\text{enclosed}} = \rho_s \cdot A \rightarrow ②$$

$$\Psi = \oint \vec{D} \cdot d\vec{s} \rightarrow ③$$

$$\text{where } \vec{D} = D_z \vec{a}_z \text{ & } d\vec{s} = dx dy \vec{a}_z \Rightarrow ds \vec{a}_z$$

sub in ③

$$\Psi = \iint D_z \vec{a}_z \cdot ds \vec{a}_z$$

$$= \iint D_z ds$$

$$= D_z \iint ds$$

$$= D_z \left[\iint_{\text{top}} ds + \iint_{\text{bottom}} ds \right]$$

$$\therefore \Psi = D_z [A + A]$$

$$\Psi = 2A D_z \rightarrow ④$$

From ①

$$2AD_z = P_s A$$

$$D_z = \frac{P_s}{2}$$

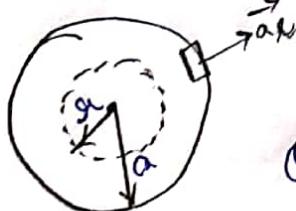
$$\vec{D} = \frac{P_s}{2} \cdot \vec{a}_z$$

$$\therefore \vec{E} = \frac{\vec{D}}{\epsilon_0} = \frac{P_s}{2\epsilon_0} \cdot \vec{a}_z$$

Electric field Intensity due to Uniformly charged sphere:

consider a volume charge density ρ_v enclosed in a gaussian surface of radius ' a ', then its electric field intensity is given by the following cases:

case i: $0 \leq r < a$:



$$\Psi = Q_{\text{enclosed}} \rightarrow ①$$

$$Q_{\text{enc}} = \iiint \rho_v \, dv$$

$$= \rho_v \int_{r=0}^a \int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} r^2 \sin\theta \, dr \, d\theta \, d\phi$$

$$= 2\pi \rho_v \int_{r=0}^a [-\cos\theta]_{0}^{\pi} r^2 \, dr$$



$$\Rightarrow Q_{\text{enc}} = 4\pi \rho_v \int_{r=0}^a r^2 dr$$

$$Q_{\text{enc}} = \frac{4}{3} \pi \rho_v a^3 \rightarrow \textcircled{2}$$

$$\therefore \Psi = \iint \vec{D} d\vec{a}$$

$$\text{Now, } \vec{D} = D_r \hat{a}_r$$

$$d\vec{a} = r^2 \sin\theta d\theta d\phi \hat{a}_r$$

$$\text{sub^n in } \Psi = \int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} D_r \cdot r^2 \sin\theta d\phi$$

$$= 2\pi D_r \cdot r^2 \int_{\theta=0}^{\pi} \sin\theta d\theta$$

$$\Psi = 4\pi D_r r^2 \rightarrow \textcircled{3}$$

$$\therefore \textcircled{2} = \textcircled{3}$$

$$\frac{4\pi \rho_v a^3}{3} = 4\pi D_r r^2$$

$$D_r = \frac{\rho_v r}{3}$$

$$\vec{D} = \frac{\rho_v r}{3} \hat{a}_r$$

$$\boxed{\vec{E} = \frac{\vec{D}}{\epsilon_0} = \frac{\rho_v r}{3\epsilon_0} \hat{a}_r}$$

case-iii: $r \geq a$



$$Q_{\text{enc}} = \iiint f_v dv$$

$$= \int_{a=0}^a \int_{\theta=0}^{\pi} \int_{\phi=0}^{2\pi} f_v r^2 \sin\theta d\theta d\phi dr$$

$$\Rightarrow \therefore Q_{\text{enc}} = \frac{4}{3}\pi f_v a^3$$

$$\Rightarrow \Psi = \iint \vec{D} \cdot \vec{ds}$$

$$= 4\pi D_a r^2$$

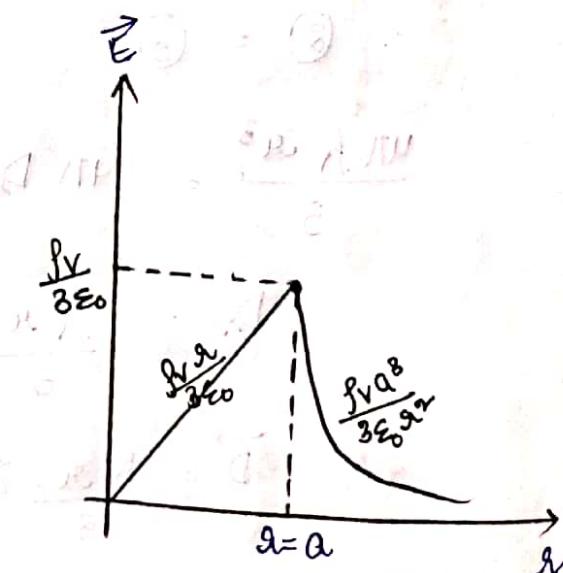
Now, $\Psi = Q_{\text{enclosed}}$

$$4\pi D_a r^2 = \frac{4}{3}\pi f_v a^3$$

$$D_a = \frac{f_v}{3} \cdot \frac{a^3}{r^2}$$

$$\vec{D} = \frac{f_v a^3}{3r^2} \cdot \vec{a_r}$$

$$\boxed{\vec{E} = \frac{\vec{D}}{\epsilon_0} = \frac{f_v a^3}{3\epsilon_0 r^2} \cdot \vec{a_r}}$$



$$\vec{D} = (2y^2 + z) \vec{a}_x + 4xy \vec{a}_y + x \vec{a}_z \text{ C/m}^2. \text{ Find}$$

i) volume charge density at $(-1, 0, 3)$
ii) the flux through the cube defined by $0 \leq y \leq 1; 0 \leq z \leq 1$
iii) the total charge enclosed by 'Q'

$$i) f_V = ?$$

from Gauss law, we have $f_V = \nabla \cdot D$

$$\text{where } \nabla = \frac{\partial}{\partial x} \vec{a}_x + \frac{\partial}{\partial y} \vec{a}_y + \frac{\partial}{\partial z} \vec{a}_z$$

$$\therefore \nabla \cdot D = \frac{\partial}{\partial x} (2y^2 + z) + \frac{\partial}{\partial y} (4xy) + \frac{\partial}{\partial z} (x)$$

$$\boxed{\nabla \cdot D = f_V = 4x \text{ C/m}^3} \quad \text{at } (-1, 0, 3) \Rightarrow \boxed{f_V = -4 \text{ C/m}^3}$$

$$ii) \Psi = \iint \vec{D} \cdot d\vec{s}$$

$$\begin{aligned} \Psi &= Q_{\text{enc}} = \iiint f_V \cdot dV \\ &= \int_{x=0}^1 \int_{y=0}^1 \int_{z=0}^1 4x \, dx \, dy \, dz \\ &= 4 \left[\frac{x^2}{2} \right]_0^1 \\ &= 4 \left[\frac{1}{2} \right] \end{aligned}$$

$$\therefore \boxed{\Psi = 2 \text{ coulombs}}$$

$$iii) \text{The total charge enclosed} \Rightarrow \boxed{Q_{\text{enc}} = 2C}$$

ELECTRIC POTENTIAL:

It is defined as the work done in moving a charge ' Q ' from one point to another point in an electric field \vec{E} .

- Suppose ' Q_t ' is the test charge moved in electric field of point charge ' Q ' from point A to B. Then,

$$V = \frac{W}{Q_t}$$

$$[F \cdot d\vec{l} = \vec{F} \cdot \vec{dl}] \text{ and work done}$$

- since, to move point charge from A to B, work is done externally.
- $\therefore W = -F \cdot \vec{dl}$
- if ' dW ' is the workdone to move charge at distance dl , then,

$$dW = -F \cdot dl$$

$$W = - \int_A^B F \cdot dl$$

- Electric force on test charge ' Q_t ' is given by

$$\Rightarrow F = E Q_t$$

$$\Rightarrow W = - \int_A^B E Q_t dl$$

$$W = - Q_t \int_A^B \vec{E} \cdot d\vec{l}$$

$$\frac{W}{Q_t} = - \int_A^B \vec{E} \cdot d\vec{l}$$

$$V = - \int_A^B \vec{E} \cdot d\vec{l}$$

we have electric field intensity due to point charge.

$$\text{is, } \vec{E} = \frac{1}{4\pi\epsilon_0} \cdot \frac{Q}{r^2} \cdot \vec{a}_r$$

$$\therefore V = - \int_A^B \frac{Q}{4\pi\epsilon_0 r^2} \cdot \vec{a}_r \cdot d\vec{l}$$

$d\vec{l}$ in spherical coordinates is:

$$d\vec{l} = dr \vec{a}_r + r d\theta \vec{a}_\theta + r \sin\theta d\phi \vec{a}_\phi$$

$$V_{AB} = \int_A^B \frac{Q}{4\pi\epsilon_0 r^2} dr$$

$$= - \int_{r_A}^{r_B} \frac{Q}{4\pi\epsilon_0 r^2} dr$$

$$V_{AB} = \frac{-Q}{4\pi\epsilon_0} \left[-\frac{1}{r} \right]_{r_A}^{r_B}$$

$$V_{AB} = \frac{Q}{4\pi\epsilon_0} \left[\frac{1}{r_B} - \frac{1}{r_A} \right]$$

$$V_{AB} = \frac{Q}{4\pi\epsilon_0 r_B} - \frac{Q}{4\pi\epsilon_0 r_A}$$

$$\therefore V_{AB} = V_B - V_A$$

If point 'A' is at infinity then $r_A \rightarrow \infty$

⇒ If potential at any particular point is measured with reference to a point at infinity, then the measured potential is called "absolute potential".

$$V_{AB} = \frac{Q}{4\pi\epsilon_0 r_B} - 0$$

$$V_{AB} = \frac{Q}{4\pi\epsilon_0 r_B}$$

If source charge 'Q' point is not at origin but somewhere with position vector 'r'

$$\text{Then } V = \frac{Q}{4\pi\epsilon_0 |(r-r')|}$$

If there are 'N' no. of source charges then

$$V = V_1 + V_2 + V_3 + \dots + V_N$$

$$V = \frac{Q}{4\pi\epsilon_0 |(r-r_1)|} + \frac{Q}{4\pi\epsilon_0 |(r-r_2)|} + \dots + \frac{Q}{4\pi\epsilon_0 |(r-r_N)|}$$

$$V = \frac{1}{4\pi\epsilon_0} \sum_{i=1}^N \frac{Q}{|r-r_i|}$$

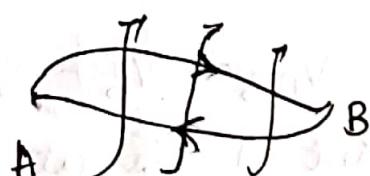
NOTE:

Electric potential is a scalar quantity

units: Volts or J/coulomb.

Relation b/w E & V - Maxwell's Equation:

The potential difference b/w points A & B is independent of path taken.



$$\therefore V_{AB} = -V_{BA}$$

$$\Rightarrow V_{AB} + V_{BA} = 0$$

$$-\int_A^B \vec{E} \cdot d\vec{l} - \int_B^A \vec{E} \cdot d\vec{l} = 0$$

$$-\oint \vec{E} \cdot d\vec{l} = 0$$

It shows that the integral of \vec{E} along a closed path is '0'

On applying "Stokes' theorem:

$$\oint \vec{E} \cdot d\vec{l} = \iint \nabla \times \vec{E} \cdot d\vec{s} = 0$$

$$\Rightarrow \nabla \times \vec{E} = 0$$

Any vector field that satisfy the above eqn is said to be conservative or irrotational.

∴ static electric field is a conservative field or irrotational.

→ The above eqn is also referred to as "Maxwell's Eqn" for static electric field.

$$dV = - \int_A^B \vec{E} \cdot d\vec{l}$$

$$dV = - [E_x dx + E_y dy + E_z dz]$$

$$L.H.S = dV$$

$$= \frac{\partial V}{\partial x} dx + \frac{\partial V}{\partial y} dy + \frac{\partial V}{\partial z} dz$$

$$= \left[\frac{\partial V}{\partial x} \vec{a}_x + \frac{\partial V}{\partial y} \vec{a}_y + \frac{\partial V}{\partial z} \vec{a}_z \right] \cdot (dx \vec{a}_x + dy \vec{a}_y + dz \vec{a}_z)$$

Equating LHS & RHS:

$$\nabla V \cdot d\vec{l} = - \vec{E} \cdot d\vec{l}$$

$$\boxed{\vec{E} = - \nabla V}$$

Energy in Static electric field :

To determine the energy present in the assembly of charges, 1st determine the amount of work done to assemble them if Q_1 , Q_2 & Q_3 are charges placed in empty space at points P_1 , P_2 & P_3 respectively. If W_1 , W_2 & W_3 are the works done to transfer these charges from infinity then

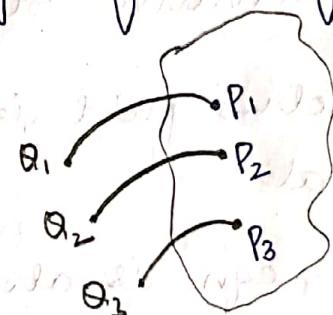
$$W_1 = 0$$

$$W_2 = Q_2 V_{\infty}$$

$$W_3 = Q_3 (V_{31} + V_{32})$$

$$W = W_1 + W_2 + W_3 \rightarrow ①$$

→ If the charges are brought to the region in Reverse Order



$$W_3 = 0$$

$$W_2 = Q_2 V_{\infty}$$

$$W_1 = Q_1 (V_{13} + V_{12})$$

$$W = W_1 + W_2 + W_3 \rightarrow ②$$

$$① + ②$$

$$\Rightarrow 2W = 0 + Q_2 V_{\infty} + Q_3 [V_{31} + V_{32}] + 0 + Q_2 V_{\infty} + Q_1 [V_{13} + V_{12}]$$

$$W = \frac{1}{2} [Q_1 (V_{13} + V_{12}) + Q_2 (V_{21} + V_{23}) + Q_3 (V_{31} + V_{32})]$$

$$W = \frac{1}{2} [Q_1 V_1 + Q_2 V_2 + Q_3 V_3]$$

if 'N' no. of charges were moved from ∞ to region

$$W = \frac{1}{2} [Q_1 V_1 + Q_2 V_2 + \dots + Q_N V_N]$$

$$W = \frac{1}{2} \sum_{k=1}^N Q_k V_k$$

from the definition, energy present = Work done

$$\therefore W_E = \frac{1}{2} \sum_{k=1}^N Q_k V_k$$

Electric Dipole:

Electric dipole is an arrangement of two charges equal in magnitude but of opposite polarity separated by a small distance.

Consider '+Q' charge at pt. B

'-Q' charge at pt. A

separated by a distance 'd'

Potential due to charge at pt. B

at a distant point 'P' is given

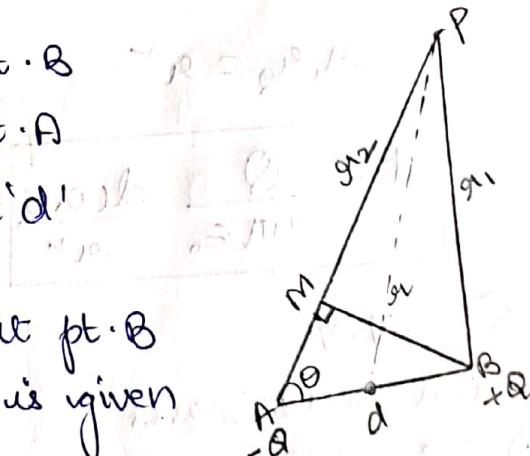
$$V_B = \frac{Q}{4\pi \epsilon_0 r_1}$$

Potential due to charge at pt. A

at a distant pt 'P' is given by

$$V_A = \frac{-Q}{4\pi \epsilon_0 r_2}$$

Total potential due to charge at pt 'P' is $V = V_A + V_B$



$$V = \frac{Q}{4\pi\epsilon_0 r_1} - \frac{Q}{4\pi\epsilon_0 r_2}$$

$$= \frac{Q}{4\pi\epsilon_0} \left[\frac{1}{r_1} - \frac{1}{r_2} \right]$$

$$= \frac{Q}{4\pi\epsilon_0} \left[\frac{r_2 - r_1}{r_1 r_2} \right]$$

from Δ^k A.M.D

$$\cos\theta = \frac{AM}{AB}$$

$$AM = AB \cos\theta$$

$$AM = r_2 - r_1 = d \cos\theta$$

if 'd' is small & 'p' is large distance, then

$$r_1 \approx r_2 \approx r$$

$$r_1, r_2 = r^2$$

$$\therefore V = \frac{Q}{4\pi\epsilon_0} \frac{dr \cos\theta}{r^2}$$

$$\text{W.R.T } \vec{E} = -\nabla \cdot V$$

In spherical coordinates

$$\nabla = \frac{\partial}{\partial r} \vec{a}_r + \frac{1}{r} \frac{\partial}{\partial \theta} \vec{a}_\theta + \frac{1}{r \sin\theta} \frac{\partial}{\partial \phi} \vec{a}_\phi$$

$$\Rightarrow \vec{E} = -\frac{\partial V}{\partial r} \vec{a}_r - \frac{1}{r} \frac{\partial V}{\partial \theta} \vec{a}_\theta - \frac{1}{r \sin\theta} \frac{\partial V}{\partial \phi} \vec{a}_\phi$$

$$\Rightarrow \vec{E} = \frac{Qd}{4\pi\epsilon_0 r^3} \left[\frac{2 \cos\theta}{r^3} \vec{a}_r + \frac{\sin\theta}{r^3} \vec{a}_\theta \right]$$

$$\therefore \vec{E} = \frac{Qd}{4\pi\epsilon_0 r^3} \left[2 \cos\theta \vec{a}_r + \sin\theta \vec{a}_\theta \right]$$

Current & current density:

The current through a given area is the electric charge passing through the area per unit time

$$\text{i.e } I = \frac{dQ}{dt}$$

(OR)

Current is defined as the rate of flow of charge.

If the "dI" is the current flowing through a planar differential surface "ds", then current density is defined as "current passing through unit normal surface area" i.e

$$\vec{J}_n = \frac{\vec{I}}{S}$$

$$J_n = \frac{dI}{ds} \Rightarrow I = \int_S dI = \int_S \vec{J}_n ds$$

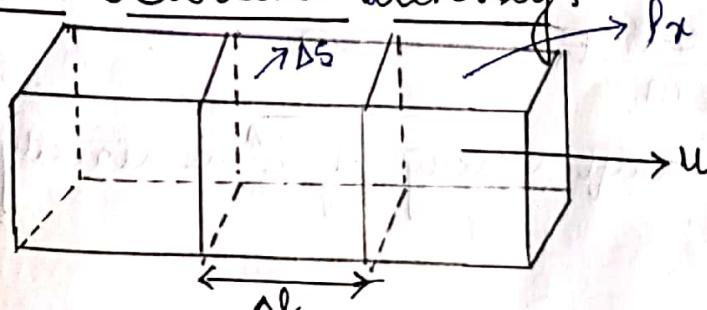
Depending upon the medium through which current passes, current produced. There are 3 types of current densities.

Convection current density

Conduction

Displacement

Convection Current density:



Total current flowing through surface 'S'

$$I = \int J \cdot ds$$

$$\Delta I = \frac{\Delta Q}{\Delta t}$$

$$\begin{aligned} dQ &= j_v dv \\ &= j_v ds dl \\ &= dl \Rightarrow dy \\ \rightarrow dQ &= j_v ds dy \end{aligned}$$

$$\rightarrow \Delta I = \frac{\Delta Q}{\Delta t} = j_v \cdot \Delta s \frac{\Delta y}{\Delta t}$$

$$\Delta I = j_v \Delta s \cdot u_y$$

The current density at a given pt is the current through a unit normal area at that pt.

Conduction current density:

when an electric field 'E' is applied, the force on an electron with charge -e is

$$F = -eE \quad \text{--- ①} \quad (\because F = QE \text{ or } E = \frac{F}{Q})$$

If an \bar{e}^n which mass 'm' is moving in an electric field 'E' with an average drift velocity (u), acc to Newton's law; the average change in momentum of free e^- must match the applied force.

$$\frac{mu}{2} = -eE$$

$$\text{①} \quad u = -\frac{e \gamma}{m} E \quad \text{--- ②}$$

from above eqn, the drift velocity of \bar{e}^n is directly proportional to applied field.

If there are 'n' e^{-n} per unit volume, the electric charge density given by

$$g_v = -ne \quad \text{--- (3)}$$

$$J = g_v \cdot u = \frac{\tau ne^2}{m} E = \sigma E \quad \text{or} \quad J = \sigma E \quad \text{--- (4)}$$

where $\sigma = \frac{\tau ne^2}{m}$ is the conductivity of the conductor

$$\therefore [J = \sigma E]$$

Continuity Equation:

The principle of charge conservation; the time rate of decrease of charge within a given volume must be equal to the net outward current flow through the surface of the volume.

current I_{out} coming out of closed surface is :

$$I_{\text{out}} = \oint \mathbf{J} \cdot d\mathbf{s} = -\frac{dQ_{\text{in}}}{dt} \quad \text{--- (1)}$$

where Q_{in} is the total charge enclosed by the surface.

& J - conduction current density.

→ By divergence theorem:

$$\oint \mathbf{J} \cdot d\mathbf{s} = \int \nabla \cdot \mathbf{J} dv \quad \text{--- (2)}$$

$$\text{but } -\frac{dQ_{\text{in}}}{dt} = -\frac{d}{dt} \int g_v dv = -\int \frac{\partial g_v}{\partial t} dv \quad \text{--- (3)}$$

subⁿ ② & ③ in ①

$$\Rightarrow \int \nabla \cdot J \, dV = - \int \frac{\partial \rho_v}{\partial t} \, dV$$

④)

$$\boxed{\nabla \cdot J = - \frac{\partial \rho_v}{\partial t}}$$

This is called "Continuity Equation"

Relaxation Time:

The Maxwell's first eqn is given by:

$$\boxed{f_v = \nabla \cdot D}$$

$$W.K.T \quad J = \sigma E$$

$$f_v = \nabla \cdot \epsilon E \quad [\because D = \epsilon E]$$

$$\frac{f_v}{\epsilon} = \nabla \cdot E$$

$$\frac{\sigma}{\epsilon} \frac{f_v}{\epsilon} = \nabla \cdot \sigma E \Rightarrow \frac{\sigma f_v}{\epsilon} = \nabla \cdot J$$

From continuity eqn

$$\nabla \cdot J = - \frac{\partial f_v}{\partial t}$$

$$\frac{\partial f_v}{\partial t} = - \frac{\sigma f_v}{\epsilon} \Rightarrow \frac{\partial f_v}{f_v} = - \frac{\sigma}{\epsilon} \frac{1}{t}$$

$$\therefore \boxed{f_v = f_{v_0} e^{-\frac{t}{T_r}}}$$

where $\boxed{T_r = \frac{\epsilon}{\sigma}}$ is known as "Relaxation Time" or Rearrangement Time

- It is the time taken by the charge placed in interior of a material to drop to $\frac{1}{e}$ of its initial value.

POISSON'S EQUATIONS

Maxwell's 1st eqn is given by : $\rho_V = \nabla \cdot D$

$$\rho_V = \nabla \cdot \epsilon E \quad (\because D = \epsilon E)$$

$$\rho_V = \epsilon \nabla \cdot E$$

$$\rho_V = \epsilon \nabla \cdot (-\nabla V)$$

$$\frac{\rho_V}{\epsilon} = -\nabla^2 V$$

(or)

$$\nabla^2 V = -\frac{\rho_V}{\epsilon}$$

which is known as "POISSON'S LAW".

A special case of above eqn occurs when $\rho_V = 0$
(i.e. for a free charge region)

$$\therefore \nabla^2 V = 0 \rightarrow \text{which is known as "LAPLACE Eqn".}$$

LAPLACE eqn in cartesian coordinate system is :

$$\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0$$

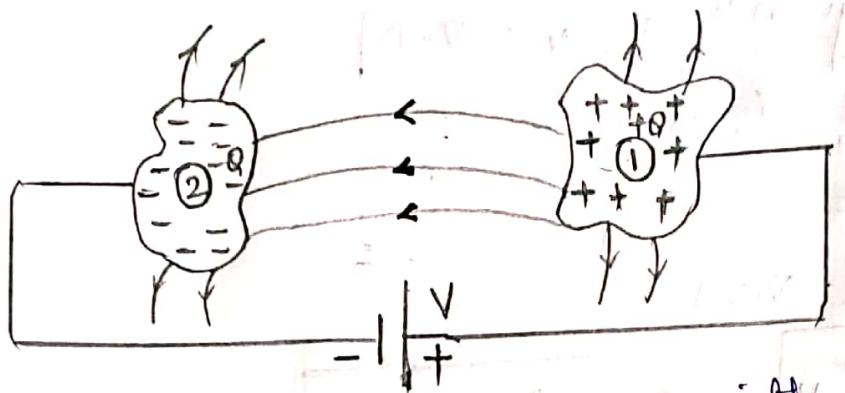
CAPACITOR & CAPACITANCE :-

Capacitor is an arrangement of two or more conductors carrying equal but opposite charge, hence flux lines start from conductor with positive charge & terminate at conductor with negative charge.

It stores energy in the form of Electric field denoted by 'C'

measured in "Farads" or "Coloumbs/volts".

Procedure to Calculate Capacitance:



- consider two conductors with +ve charge on conductor -1 & -ve charge on conductor 2
- The potential diff b/w two conductors i.e :

$$V = V_1 - V_2$$

$$V = - \int_{\text{2}}^{\text{1}} E \cdot dl$$

The capacitance 'C' of the capacitor is defined as the ratio of magnitude of charge on one of the plates to the potential diff b/w them

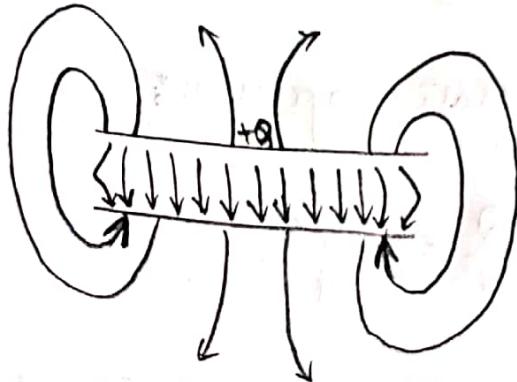
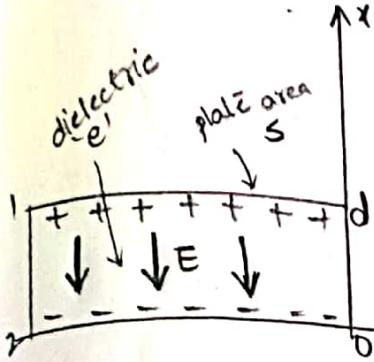
$$\text{i.e } C = \frac{Q}{V}$$

$$C = \frac{1}{\epsilon_0} \int E \cdot ds$$

Capacitance is calculated with two approaches:

- ① Assume 'Q' & calculate 'E' using Coulomb's or Gauss law & then calculate 'V' using $V = - \int \vec{E} \cdot d\vec{l}$
- ② Assume 'V' & calculate 'Q' in terms of 'V' using Laplace eqn: $\nabla^2 V = 0$

parallel plate capacitor:



assuming $+Q$ & $-Q$ charges on plates 1 & 2 respectively such that $\rho_s = \frac{Q}{S} \rightarrow ①$

The electric field intensity \vec{E} due to 1st plate in -ve x-direction is

$$E_1 = -\frac{\rho_s}{2\epsilon_0} \hat{a}_x$$

||| E_2 in +ve x-direction $E_2 = -\frac{\rho_s}{2\epsilon_0} \cdot \hat{a}_x$

∴ Total electric field Intensity \vec{E} due to 1 & 2 is

$$\vec{E} = \vec{E}_1 + \vec{E}_2$$

$$\vec{E} = -\frac{\rho_s}{2\epsilon_0} \hat{a}_x - \frac{\rho_s}{2\epsilon_0} \hat{a}_x$$

$$\vec{E} = -\frac{\rho_s}{\epsilon_0} \hat{a}_x$$

$$\therefore \vec{E} = -\frac{Q}{\epsilon_0 S} \hat{a}_x \quad (\text{From } ① \rho_s = \frac{Q}{S})$$

$$V = - \int_2^1 \vec{E} \cdot d\vec{l} = \int_0^d \left(-\frac{Q}{\epsilon_0 S} \hat{a}_x \right) dx \hat{a}_x$$

$$V = \frac{Qd}{\epsilon_0 S}$$

For a 11^{th} plate capacitor:

$$C = \frac{Q}{V} = \frac{\epsilon_0 S}{d}$$

If dielectric ϵ between the plates is not free space then capacitance is given by:

$$C = \frac{\epsilon S}{d}$$