

QuTiP lecture: Quantum Monte-Carlo Trajectories

Author: J.R. Johansson, robert@riken.jp

<http://dml.riken.jp/~rob/>

Latest version of this ipython notebook lecture is available at: <http://github.com/jrjohansson/qutip-lectures>

The example in this lecture is based on an example by P.D. Nation.

```
In [1]: # setup the matplotlib graphics library and configure it to show figures inline in
        the notebook
        %pylab inline
```

```
Welcome to pylab, a matplotlib-based Python environment [backend:
module://IPython.zmq.pylab.backend_inline].
For more information, type 'help(pylab)'.
```

```
In [2]: # make qutip available in the rest of the notebook
        from qutip import *

        #qutip.settings.qutip_graphics=False
        #qutip.settings.qutip_gui=False
```

Introduction to the Quantum Monte-Carlo trajectory method

The Quantum Monte-Carlo trajectory method is an equation of motion for a single realization of the state vector $|\psi(t)\rangle$ for a quantum system that interacts with its environment. The dynamics of the wave function is given by the Schrodinger equation,

$$\frac{d}{dt} |\psi(t)\rangle = -\frac{i}{\hbar} H_{\text{eff}} |\psi(t)\rangle$$

where the Hamiltonian is an effective Hamiltonian that, in addition to the system Hamiltonian $H(t)$, also contains a non-Hermitian contribution due to the interaction with the environment:

$$H_{\text{eff}}(t) = H(t) - \frac{i\hbar}{2} \sum_n c_n^\dagger c_n$$

... incomplete ...

$$\delta p = \delta t \sum_n \langle \psi(t) | c_n^\dagger c_n | \psi(t) \rangle$$

$$|\psi(t + \delta t)\rangle = c_n |\psi(t)\rangle / \langle \psi(t) | c_n^\dagger c_n | \psi(t) \rangle^{1/2}$$

Decay of a single-photon Fock state in a cavity

This is a Monte-Carlo simulation showing the decay of a cavity Fock state $|1\rangle$ in a thermal environment with an average occupation number of $n = 0.063$.

Here, the coupling strength is given by the inverse of the cavity ring-down time $T_c = 0.129$.

The parameters chosen here correspond to those from S. Gleyzes, et al., Nature 446, 297 (2007), and we will carry out a simulation that corresponds to these experimental results from that paper:



Problem parameters

```
In [3]: N = 4                # number of basis states to consider
        kappa = 1.0/0.129    # coupling to heat bath
        nth = 0.063          # temperature with <n>=0.063

        tlist = linspace(0,0.6,100)
```

Create operators, Hamiltonian and initial state

Here we create QuTiP Qobj representations of the operators and state that are involved in this problem.

```
In [4]: a = destroy(N)      # cavity destruction operator
        H = a.dag() * a     # harmonic oscillator Hamiltonian
        psi0 = basis(N,1)   # initial Fock state with one photon: |1>
```

Create a list of collapse operators that describe the dissipation

```
In [5]: # collapse operator list
        c_op_list = []

        # decay operator
        c_op_list.append(sqrt(kappa * (1 + nth)) * a)
```

```
# excitation operator
c_op_list.append(sqrt(kappa * nth) * a.dag())
```

Run Monte-Carlo simulation

Here we start the Monte-Carlo simulation, and we request expectation values of photon number operators with 1, 5, 15, and 904 trajectories (compare with experimental results above).

```
In [31]: ntraj = [1, 5, 15, 904] # list of number of trajectories to avg. over

mc = mcsolve(H, psi0, tlist, c_op_list, [a.dag()*a], ntraj)

# get expectation values from mc data (need extra index since ntraj is list)
ex1  = mc.expect[0][0] # for ntraj=1
ex5  = mc.expect[1][0] # for ntraj=5
ex15 = mc.expect[2][0] # for ntraj=15
ex904 = mc.expect[3][0] # for ntraj=904
```

Lindblad master-equation simulation and steady state

For comparison with the averages of single quantum trajectories provided by the Monte-Carlo solver we here also calculate the dynamics of the Lindblad master equation, which should agree with the Monte-Carlo simulations for infinite number of trajectories.

```
In [32]: # run master equation to get ensemble average expectation values
me = mesolve(H, psi0, tlist, c_op_list, [a.dag()*a])

# calculate final state using steadystate solver
final_state = steadystate(H, c_op_list) # find steady-state
fexpt = expect(a.dag()*a, final_state) # find expectation value for particle
number
```

Plot the results

```
In [36]: import matplotlib.font_manager
leg_prop = matplotlib.font_manager.FontProperties(size=10)

fig, axes = subplots(4, 1, sharex=True, figsize=(8,12))

fig.subplots_adjust(hspace=0.1) # reduce space between plots

for idx, n in enumerate(ntraj):

    axes[idx].plot(tlist, mc.expect[idx][0], 'b', lw=2)
    axes[idx].plot(tlist, me.expect[0], 'r--', lw=1.5)
    axes[idx].axhline(y=fexpt, color='k', lw=1.5)

    axes[idx].set_yticks(linspace(0, 2, 5))
    axes[idx].set_ylim([0, 1.5])
    axes[idx].set_ylabel(r'$\left<N\right>$', fontsize=14)

    if idx == 0:
        axes[idx].set_title("Ensemble Averaging of Monte Carlo Trajectories")
        axes[idx].legend(('Single trajectory', 'master equation', 'steady state'),
prop=leg_prop)
```

```

else:
    axes[idx].legend('%d trajectories' % n, 'master equation', 'steady
state'), prop=leg_prop)

axes[3].axis.set_major_locator(MaxNLocator(4))
axes[3].set_xlabel('Time (sec)', fontsize=14)

```

Out [36]: <matplotlib.text.Text at 0xb46df50>

