

SUPERPOSITION OF E. FIELDS

Since force follows superposition,

$$\sum \vec{F} = \vec{F}_1 + \vec{F}_2 + \vec{F}_3 + \dots = q_1 \vec{E}_1 + q_2 \vec{E}_2 + \dots$$

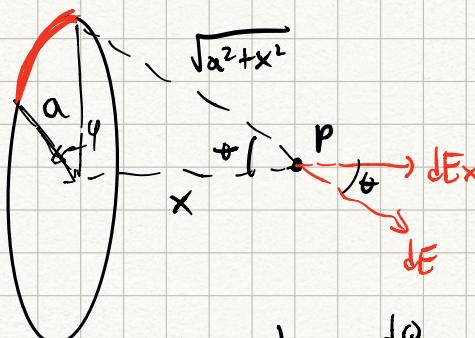
$$\sum \vec{E} = \frac{\vec{E}_1 + \vec{E}_2 + \vec{E}_3 + \dots}{q_1} = \frac{\vec{E}_1}{q_1} + \frac{\vec{E}_2}{q_2} + \dots = \vec{E}_1 + \vec{E}_2 + \vec{E}_3 + \dots$$

λ = linear charge density ($\frac{C}{m}$)

σ = surface charge density ($\frac{C}{m^2}$)

ρ = volume charge density ($\frac{C}{m^3}$)

ex. field across ring of charge



$$(\omega) = \frac{x}{\sqrt{a^2 + x^2}}$$

$$dE = \frac{1}{4\pi\epsilon_0} \frac{dQ}{x^2 + a^2}$$

$$dQ = \lambda ds = \lambda (a d\phi) \Rightarrow Q = \int_0^{2\pi} \lambda a d\phi = 2\pi\lambda a$$

$$dE_x = \frac{1}{4\pi\epsilon_0} \left(\frac{\lambda a d\phi}{x^2 + a^2} \right) (\omega) \theta$$

$$= \frac{1}{4\pi\epsilon_0} \left(\frac{\lambda a x}{(x^2 + a^2)^{3/2}} \right) d\phi$$

$$E = \int dE_x = \int_0^{2\pi} \frac{\lambda a x}{4\pi\epsilon_0 (x^2 + a^2)^{3/2}} d\phi = \frac{\lambda a x}{4\pi\epsilon_0 (x^2 + a^2)^{3/2}} (2\pi) = \frac{Qx}{4\pi\epsilon_0 (x^2 + a^2)^{3/2}}$$

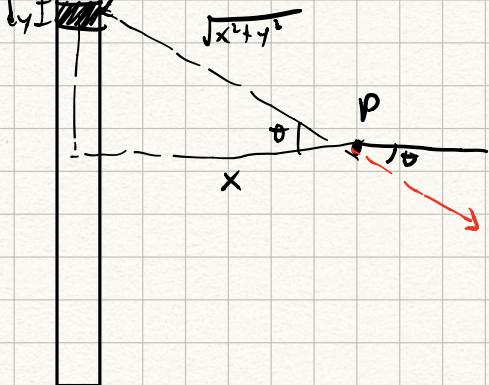
alternatively

$$dE_x = \frac{1}{4\pi\epsilon_0} \left(\frac{\lambda x}{(x^2 + a^2)^{3/2}} \right) ds, \quad E = \int_0^{2\pi a} \frac{\lambda x}{4\pi\epsilon_0 (x^2 + a^2)^{3/2}} ds$$

ex. charged like segment

$$\boxed{dQ}$$

$$\lambda = \frac{Q}{2a}, \quad dQ = \lambda dy = \frac{Q}{2a} dy$$



$$dE = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} = \frac{1}{4\pi\epsilon_0} \frac{Q}{2a} \frac{1}{(x^2+y^2)} dy$$

$$dE_x = \frac{1}{4\pi\epsilon_0} \left(\frac{Q}{2a}\right) \left(\frac{1}{x^2+y^2}\right) \left(\frac{x}{\sqrt{x^2+y^2}}\right) dy$$

$$E_x = \frac{1}{4\pi\epsilon_0} \left(\frac{Q}{2a}\right) \int_{-a}^{a} \frac{x}{(x^2+y^2)^{3/2}} dy$$

$$= \frac{1}{4\pi\epsilon_0} \left(\frac{Q}{2a}\right) \left(\frac{y}{x\sqrt{x^2+y^2}} \Big|_{-a}^a\right)$$

$$= \frac{1}{4\pi\epsilon_0} \left(\frac{Q}{2a}\right) \left(\frac{2a}{x\sqrt{x^2+a^2}}\right)$$

$$= \frac{Q}{4\pi\epsilon_0 x \sqrt{x^2+a^2}}$$

$$x = a \tan u \Rightarrow dx = a \sec^2 u du$$

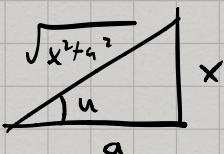
$$= \int \frac{1}{(a^2 + a^2 \tan^2 u)^{3/2}} dx$$

$$= \frac{1}{a^2} \int \frac{1}{\sec^3 u} (a \sec^2 u) du$$

$$= \frac{1}{a^2} \int \cos u du$$

$$\therefore \frac{\sin u}{a^2} = \frac{1}{a^2} \left(\frac{x}{\sqrt{x^2+a^2}} \right) + C$$

$$\frac{x}{a} = \tan u$$



$$\sin u = \frac{x}{\sqrt{x^2+a^2}}$$

if segment is very short $x \gg a$,

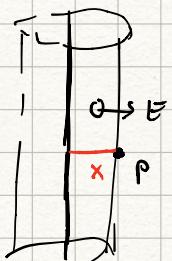
$$\vec{E} = \frac{Q}{4\pi\epsilon_0 x^2}$$

if segment is very long, $a \gg x$,

$$\vec{E} = \frac{Q}{4\pi\epsilon_0 \times \sqrt{x^2+a^2}} = \frac{1}{2\pi\epsilon_0} \left(\frac{Q}{2a}\right) \left(\frac{1}{x\sqrt{(\frac{x}{a})+1}}\right) = \frac{1}{2\pi\epsilon_0} \frac{\lambda}{x} = \frac{\lambda}{2\pi\epsilon_0 x}$$

$$\text{so } \vec{E} = \frac{\lambda}{2\pi\epsilon_0 x} \quad (\text{infinite line of charge})$$

- using Gauss's Law

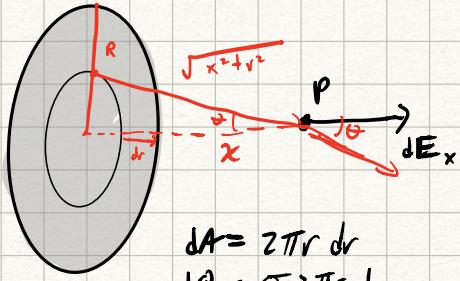


$$\phi = \frac{Q_{\text{enc}}}{\epsilon_0} = \oint \vec{E} \cdot d\vec{A} = \vec{E} \cdot A$$

$$\frac{Q}{\epsilon_0} = \vec{E} (2\pi \times h)$$

$$\lambda = \frac{Q}{h} \Rightarrow \vec{E} = \frac{\lambda}{2\pi\epsilon_0 x}$$

Ex: field of uniformly charged disk



$$dA = 2\pi r dr$$

$$dQ = \sigma 2\pi r dr$$

$$dE_{\text{ring}} = \frac{1}{4\pi\epsilon_0} \left(\frac{\lambda r d\phi}{x^2 + r^2} \right) \left(\frac{x}{\sqrt{x^2 + r^2}} \right)$$

$$= \frac{(2\pi r \lambda) x}{4\pi\epsilon_0 (x^2 + r^2)^{3/2}}$$

$$= \frac{\sigma x}{4\pi\epsilon_0 (x^2 + r^2)^{3/2}}$$

$$dE_{\text{disk}} = \frac{x dQ}{4\pi\epsilon_0 (x^2 + r^2)^{3/2}} = \frac{x}{4\pi\epsilon_0} \left(\frac{\sigma 2\pi r}{(x^2 + r^2)^{3/2}} \right) dr$$

$$E_{\text{disk}} = \frac{\sigma x}{4\pi\epsilon_0} \int_0^R \frac{r}{(x^2 + r^2)^{3/2}} dr \quad u = r^2 + x^2 \quad du = 2r dr$$

$$= \frac{\sigma x}{4\epsilon_0} \int_0^R u^{-3/2} du$$

$$= \frac{\sigma x}{4\epsilon_0} \left(-2u^{-1/2} \Big|_0^R \right)$$

$$= -\frac{\sigma x}{2\epsilon_0} \left(\frac{1}{\sqrt{R^2 + x^2}} \right) \Big|_0^R$$

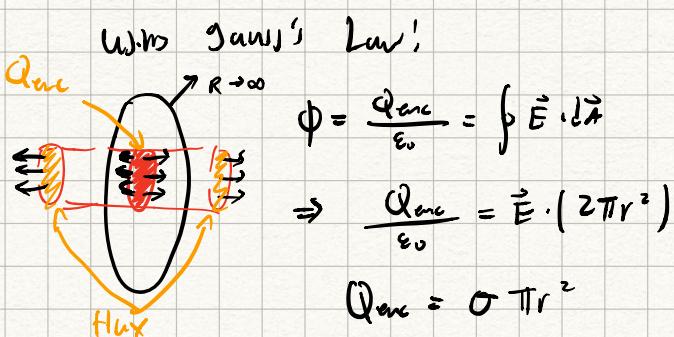
$$= \frac{\sigma x}{2\epsilon_0} \cdot \left(\frac{1}{\sqrt{R^2 + x^2}} - \frac{1}{x} \right)$$

$$= \frac{\sigma x}{2\epsilon_0} \left(\frac{1}{x} - \frac{1}{\sqrt{R^2 + x^2}} \right)$$

$$= \frac{\sigma}{2\epsilon_0} \left(1 - \frac{1}{\sqrt{\frac{R^2}{x^2} + 1}} \right)$$

If disk is really large,

$$E_{\text{disk}} = \frac{\sigma}{2\epsilon_0}$$



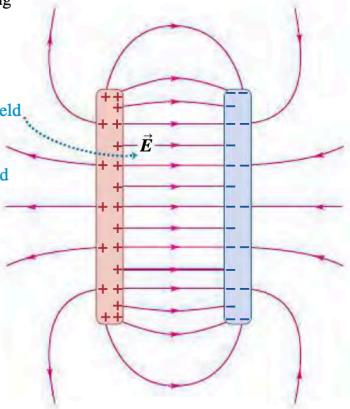
$$\phi = \frac{Q_{\text{pillbox}}}{\epsilon_0} = \int \vec{E} \cdot d\vec{A}$$

$$\Rightarrow \frac{Q_{\text{pillbox}}}{\epsilon_0} = \vec{E} \cdot (2\pi r^2)$$

$$Q_{\text{pillbox}} = \sigma \pi R^2$$

(a) Realistic drawing

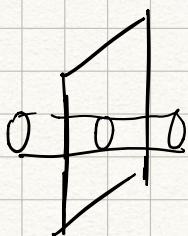
Between the two plates the electric field is nearly uniform, pointing from the positive plate toward the negative one.



$$\Rightarrow \frac{\sigma \pi r^2}{\epsilon_0} = E (2\pi r)$$

$$E = \frac{\sigma}{2\epsilon_0}$$

infinite place



$$\phi = \frac{Q_{enc}}{\epsilon_0} = \vec{E} \cdot \vec{A}$$

$$\Rightarrow \frac{Q_{enc}}{\epsilon_0} = \vec{E} \cdot (2\pi r^2)$$

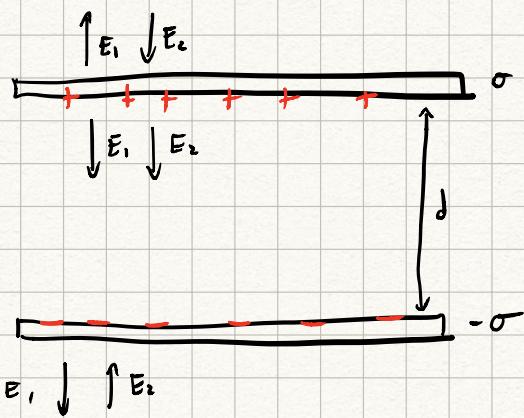
$$Q_{enc} = \sigma \pi r^2$$

$$\Rightarrow \frac{\sigma \pi r^2}{\epsilon_0} = E (2\pi r^2)$$

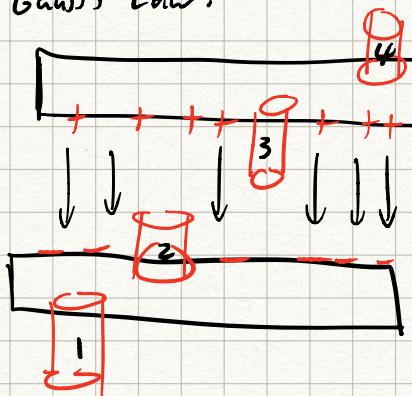
$$\Rightarrow E = \frac{\sigma}{2\epsilon_0}$$

* electric field for all 2D planes we $E = \frac{\sigma}{2\epsilon_0}$

Ex field of two oppositely charged conducting sheets, ∞ in length



using Gauss's Law:



The electric field from a single plane is

$$\vec{E} = \frac{\sigma}{2\epsilon_0}, |\vec{E}| = E_1$$

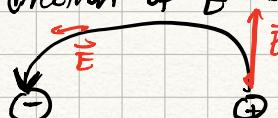
$$\sum \vec{E} = \vec{E}_1 + \vec{E}_2 = \begin{cases} 0 & \text{when above } \sigma \\ \frac{\sigma}{\epsilon_0} & \text{in between} \\ 0 & \text{when below } -\sigma \end{cases}$$

* Electric field lines are not trajectories

- \vec{E} is tangent to the field line at that point
- $\vec{F} = q\vec{E}$ is exerted on the particle
- however, since moving along curved path, acceleration cannot be tangent to that path.

* electric field := electric field lines

- magnitude of \vec{E} is usually different along the same line
- direction of \vec{E} is tangent to the



$$1: \phi = \frac{Q_{enc}}{\epsilon_0} = \vec{E} \cdot \vec{A}$$

$$\frac{\sigma}{\epsilon_0} = E (2\pi r^2)$$

$$2: \phi = -\frac{\sigma (\pi r^2)}{\epsilon_0} = -E (\pi r^2) \quad E \text{ only top}$$

$$E = \frac{\sigma}{\epsilon_0}$$

$$3: \phi = \frac{\sigma \pi r^2}{\epsilon_0} = E (\pi r^2)$$

$$E = \frac{\sigma}{\epsilon_0}$$

$$4: E = 0$$