

Fourier Series

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- Fourier series and Fourier Transform provide us a set of important tools for signals and systems analysis.
- The main idea is to represent the signals as linear combination of basic signals.
- The basic signal has two properties:-
 - 1- can be used to represent broad range of signals.
 - 2- It's response/output is simple to compute.

Historical Perspectives:-

- The concept of using "trigonometric sums"- that is, sums of harmonically related sines and cosines or periodic complex exponentials - to describe periodic phenomena goes back at least as far as the "Babylonians" \rightarrow prediction of astronomical events
- The modern history begins with "Euler" in 1748 with examination of the vibrating string; specifically he noted that if the vertical deflection of the string at any time instant can be represented as linear combination of Normal Modes of this string then we can do at any subsequent time
- in 1753, "Bernoulli" argued on physical grounds all physical motions of string could be represented by linear combinations of normal modes, But he did not mathematically.
- in 1759, "Lagrange" strongly criticized the use of trigonometric sum in the examination of vibrating string, more specifically signals with corner, he said cannot be represented.
- Around 1800, Fourier presented his ideas and with the claim that harmonically-related sinusoids can be used to describe more and more phenomena like heat propagation, temperature distribution, and he claimed that "any" periodic signal could be represented by such series.

- In 1829 Dirichlet provided precise conditions under which periodic signal could be represented by a Fourier series. Then Fourier extend his claim to aperiodic signals.
- in 1965, The fast Fourier transform FFT was introduced.

- For LTI system, if the input $x(t) = e^{st}$, then the output $y(t) = \int_{-\infty}^{+\infty} h(\tau) x(t-\tau) d\tau = \int_{-\infty}^{+\infty} h(\tau) e^{s(t-\tau)} d\tau$

$$\therefore y(t) = \int_{-\infty}^{+\infty} h(\tau) e^{st} \cdot e^{-s\tau} d\tau = e^{st} \int_{-\infty}^{+\infty} h(\tau) e^{-s\tau} d\tau$$

$$= e^{st} \cdot H(s) \quad \left\{ \begin{array}{l} \text{i.e. } e^{st} \text{ in input appears } e^{st} \text{ in output} \\ \text{too but multiplied by a variable gain} \end{array} \right\}$$

then e^{st} is called eigenfunction of the LTI systems

with eigenvalue is $H(s) = \int_{-\infty}^{+\infty} h(t) e^{-st} dt$

then if $x(t) = a_1 e^{s_1 t} + a_2 e^{s_2 t} + a_3 e^{s_3 t}$

$$\therefore y(t) = a_1 H_1(s_1) e^{s_1 t} + a_2 H_2(s_2) e^{s_2 t} + a_3 H_3(s_3) e^{s_3 t}$$

Fourier analysis involves restricting $s = j\omega$

Fourier claim that any periodic signal $x(t)$ could be represented as $x(t) = \sum_{k=-\infty}^{+\infty} a_k e^{jkw_0 t} \rightarrow ①$

as the set $\phi_k(t) = e^{jkw_0 t}$ are harmonically related signals and they are periodic with common period T , which is the fundamental period of $x(t) \Rightarrow$ i.e. $x(t) = x(t+T)$

multiply both sides of eqn ① with $e^{-jn\omega_0 t}$

$$\therefore x(t) e^{-jn\omega_0 t} = \sum_{k=-\infty}^{+\infty} a_k \cdot e^{jkw_0 t} \cdot e^{-jn\omega_0 t}, \omega_0 = \frac{2\pi}{T}$$

integrating both sides from 0 to $T = \frac{2\pi}{\omega_0}$

$$\therefore \int_0^T x(t) e^{-jn\omega_0 t} dt = \int_0^T \left\{ \sum_{k=-\infty}^{+\infty} a_k \cdot e^{jkw_0 t} \cdot e^{-jn\omega_0 t} \right\} dt$$

$$= \sum_{k=-\infty}^{+\infty} a_k \int_0^T e^{j(k-n)\omega_0 t} dt$$

$$\int_0^T e^{j(k-n)\omega_0 t} dt = \int_0^T \cos((k-n)\omega_0 t) dt + j \int_0^T \sin((k-n)\omega_0 t) dt$$

[3]

For $k \neq n$, $\cos(k-n)\omega_0 t$ and $\sin(k-n)\omega_0 t$ are periodic signals with $\omega = (k-n)\omega_0$. When comparing with general form $\cos \omega t$ or $\sin \omega t$

with fundamental period $T_0 = \frac{2\pi}{\omega} = \frac{2\pi}{|(k-n)\omega_0|} = \frac{2\pi}{|\omega_0|} \cdot \frac{1}{|k-n|}$
 $= T \cdot \frac{1}{|k-n|}$; with T is the original fundamental period of $x(t)$,
as k is integer and n is integer

$$\Rightarrow (k-n) \text{ is integer} \Rightarrow T_0 = \frac{T}{\text{integer}}$$

$$\therefore T = \text{integer} \cdot T_0$$

$\therefore \int^T$ is integrating $\cos(k-n)\omega_0 t$ and $\sin(k-n)\omega_0 t$ are integrations over integer number of periods of both \cos and $\sin = \text{zeros}$.

$$\therefore \int_0^T e^{j(k-n)\omega_0 t} dt = \int_0^T \cos(k-n)\omega_0 t dt + j \int_0^T \sin(k-n)\omega_0 t dt$$

$$= \begin{cases} 0 & k \neq n \\ T & k = n \end{cases}$$

as if $k = n$ $\int_0^T e^{j(k-n)\omega_0 t} dt = \int_0^T \cos(\theta) dt + j \int_0^T \sin(\theta) dt$
 $= \int_0^T 1 \cdot dt + j \cdot \theta$, as $\cos \theta = 1$ and $\sin \theta = \theta$
 $= t |_0^T = T - 0 = T$

$$\therefore \int_0^T x(t) e^{-j\omega_0 n t} dt = \sum_{k=-\infty}^{+\infty} q_k \cdot \begin{cases} T & k = n \\ 0 & k \neq n \end{cases}$$

$$\therefore a_n = \frac{1}{T} \int_0^T x(t) e^{-j\omega_0 n t} dt = q_n \cdot T$$

it is valid for any integration over any complete period.

$$\therefore q_k = \frac{1}{T} \int_0^T x(t) e^{-j\omega_0 k t} dt \quad \begin{array}{l} \xrightarrow{\text{called analysis equation}} \\ \text{OR Fourier series coefficients} \\ \text{OR spectral coefficients} \end{array}$$

$$\text{if } x(t) = \sum_{k=-\infty}^{+\infty} q_k e^{j\omega_0 k t} \quad \begin{array}{l} \xrightarrow{\text{called synthesis equation}} \\ \text{OR Fourier series equation} \end{array}$$

$$\text{as } q_K = \frac{1}{T} \int_T x(t) e^{-jk\omega_0 t} dt$$

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$$\Rightarrow \bar{x}_0 = \frac{1}{T} \int_T x(t) dt \quad \text{if } K=0$$

- = average value of $x(t)$ over one period.
- = dc component or constant component.

a_1 and a_{-1} are called coefficients of first harmonic,
 a_2 and a_{-2} " " " " " second " .
 \vdots
 a_k and a_{-k} " " " " " k -th "

Ex 1:

$$x(t) = \sin \omega_0 t$$

$$\therefore x(t) = \frac{1}{2j} \left\{ e^{j\omega_0 t} - e^{-j\omega_0 t} \right\}$$

$$= \frac{1}{2j} e^{j(\omega_0)t} - \frac{1}{2j} e^{j(-\omega_0)t}$$

$$\text{as } x(t) = \sum_{k=-\infty}^{+\infty} q_k e^{jk\omega_0 t} \Rightarrow \therefore q_1 = \frac{1}{z_j} \text{ and } q_{-1} = \frac{-1}{z_j}$$

and $a_K = 0$; $K \neq +$ or -1

Ex 2:

$$x(t) = 1 + \sin \omega_0 t + 2 \cos \omega_0 t + \cos(2\omega_0 t + \frac{\pi}{4})$$

$$\therefore x(t) = 1 + \frac{1}{2j} \left\{ e^{jw_0 t} - e^{-jw_0 t} \right\} + 2 \cdot \frac{1}{2} \left\{ e^{jw_0 t} + e^{-jw_0 t} \right\} + \frac{1}{2} \left\{ e^{j(2w_0 + \frac{\pi}{4})} + e^{-j(2w_0 + \frac{\pi}{4})} \right\}$$

$$\therefore x(t) = 1 + \left(1 + \frac{1}{2j}\right) e^{j\omega_0 t} + \left(1 - \frac{1}{2j}\right) e^{-j\omega_0 t} + \left(\frac{1}{2} e^{j\frac{\pi}{4}}\right) e^{j2\omega_0 t} + \left(\frac{1}{2} e^{-j\frac{\pi}{4}}\right) e^{-j2\omega_0 t}$$

$$\therefore q_0 = 1 = \text{dc component}$$

$$a_1 = 1 + \frac{1}{2}j = 1 - \frac{1}{2}j$$

$$q_{-1} = 1 - \frac{1}{z} j = 1 + \frac{1}{z} j$$

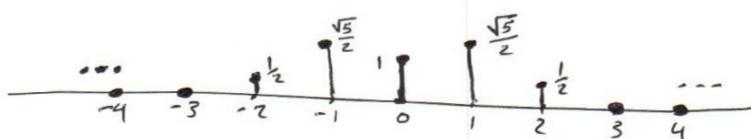
$$z_2 = \frac{1}{2} e^{j\frac{\pi}{4}} = \frac{1}{2} \left\{ \cos \frac{\pi}{4} + j \sin \frac{\pi}{4} \right\} = \frac{1}{2} \left\{ \frac{1}{\sqrt{2}} + j \frac{1}{\sqrt{2}} \right\} = \frac{\sqrt{2}}{4} (1+j)$$

$$a_{-2} = \frac{1}{2} e^{j\frac{\pi}{4}} = \frac{1}{2} \left\{ \cos \frac{\pi}{4} - j \sin \frac{\pi}{4} \right\} = \frac{1}{2} \left\{ \frac{1}{\sqrt{2}} - j \frac{1}{\sqrt{2}} \right\} = \frac{\sqrt{2}}{4} (1-j)$$

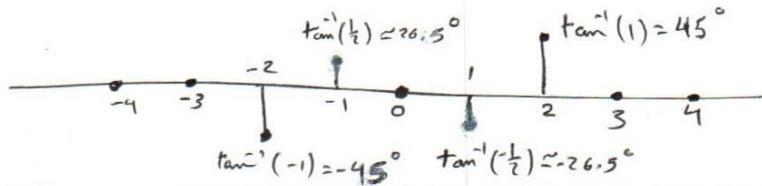
$$a_K = 0 \quad , \quad |K| \geq 2$$

magnitudes

19_K 1



Phase $\propto q_k$



Ex 3

$$x(t) = \begin{cases} 1 & ; |t| < T_1 \\ 0 & ; T_1 < |t| < \frac{T}{2} \end{cases}$$

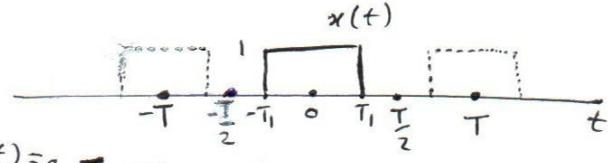
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as T is the fundamental period of $x(t)$

$$q_0 = \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} x(t) dt$$

$$= \frac{1}{T} \int_{-T_1}^{T_1} dt = \frac{1}{T} \{ t \Big|_{-T_1}^{T_1} \} \text{ as } x(t) = 0 \quad T_1 < |t| < \frac{T}{2}$$

$$= \frac{1}{T} \{ T_1 - (-T_1) \} = \frac{2T_1}{T}$$



$$q_K = \frac{1}{T} \int_T x(t) e^{-jK\omega_0 t} dt$$

$$= \frac{1}{T} \int_{-T_1}^{T_1} e^{-jK\omega_0 t} dt = \frac{1}{T} \left\{ \frac{e^{-jK\omega_0 T_1}}{-jK\omega_0} - \frac{e^{jK\omega_0 T_1}}{-jK\omega_0} \right\}$$

$$= \frac{1}{jK\omega_0 T} \left\{ e^{jK\omega_0 T_1} - e^{-jK\omega_0 T_1} \right\} = \frac{2j \sin(K\omega_0 T_1)}{jK\omega_0 T}$$

$$\therefore q_K = \frac{2j \sin(K\omega_0 T_1)}{jK\omega_0 \cdot \frac{2\pi}{\omega_0}} = \frac{\sin(K\omega_0 T_1)}{K\pi}; K \neq 0$$

$$\Rightarrow T q_K = \frac{2 \sin(\omega T_1)}{\omega} \Big|_{\omega = K\omega_0} \equiv \text{Samples of envelope } T q_K$$

$$T = 4T_1 \Rightarrow q_0 = \frac{2T_1}{4T_1} = \frac{1}{2}$$

$$T_1 = \frac{T}{4} = \frac{2\pi}{\omega_0} \cdot \frac{1}{4} = \frac{\pi}{2\omega_0}$$

$$K\omega_0 T_1 = K\omega_0 \cdot \frac{\pi}{2\omega_0} \cdot \frac{1}{4} = K \cdot \frac{\pi}{8} \stackrel{!}{=} \text{every 4-points} \equiv \text{one cycle} \stackrel{!}{=} 2\pi$$

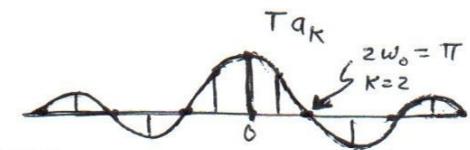
$$\frac{\sin(K\frac{\pi}{2})}{K\pi}$$

$$T = 8T_1 \Rightarrow q_0 = \frac{2T_1}{8T_1} = \frac{1}{4}$$

$$T_1 = \frac{T}{8} = \frac{2\pi}{\omega_0} \cdot \frac{1}{8} = \frac{\pi}{4\omega_0}$$

$$K\omega_0 T_1 = K\omega_0 \cdot \frac{\pi}{4\omega_0} \cdot \frac{1}{8} = K \cdot \frac{\pi}{32} \stackrel{!}{=} \text{every 8-points} \equiv \text{one cycle} \stackrel{!}{=} 2\pi$$

$$\therefore wT_1 = w \cdot \frac{\pi}{2\omega_0}$$



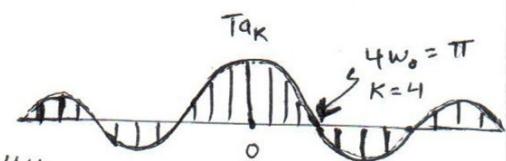
$$\therefore wT_1 = \pi \text{ if } w = 2\omega_0$$

$$T = 16T_1 \Rightarrow q_0 = \frac{2T_1}{16T_1} = \frac{1}{8}$$

$$T_1 = \frac{T}{16} = \frac{2\pi}{\omega_0} \cdot \frac{1}{16} = \frac{\pi}{8\omega_0}$$

$$K\omega_0 T_1 = K\omega_0 \cdot \frac{\pi}{8\omega_0} \cdot \frac{1}{8} = K \cdot \frac{\pi}{64} \stackrel{!}{=} \text{every 16-points} \equiv \text{one cycle} \stackrel{!}{=} 2\pi$$

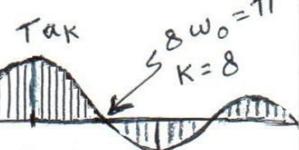
$$\therefore wT_1 = w \cdot \frac{\pi}{8\omega_0}$$



$$\therefore wT_1 = \pi \text{ if } w = 4\omega_0$$

$$\frac{1}{8} q_K$$

$$\therefore wT_1 = w \cdot \frac{\pi}{8\omega_0}$$

as T increases \Rightarrow Samples becomes more closer.

- Convergence of Fourier Series

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$$a_k = \frac{1}{T} \int_T x(t) e^{-j k \omega_0 t} dt$$

Recall: any complex number $\beta = a+ib = |z|e^{j\theta}$
does NOT go to ∞ unless its magnitude $|z| = \infty$

\therefore as the value $\{x(t) e^{-j k \omega_0 t}\}$ is a complex number

\therefore it is not go to ∞ unless $|x(t)| = \infty$

\therefore if $|x(t)| < \infty, \forall t$

$$\Rightarrow a_k = \frac{1}{T} \int_T x(t) e^{-j k \omega_0 t} dt < \infty$$

\therefore ONE sufficient conditions for Fourier series to converge

is $\boxed{\int_T |x(t)|^2 dt < \infty} \equiv \{|x(t)| < \infty \forall t\}$

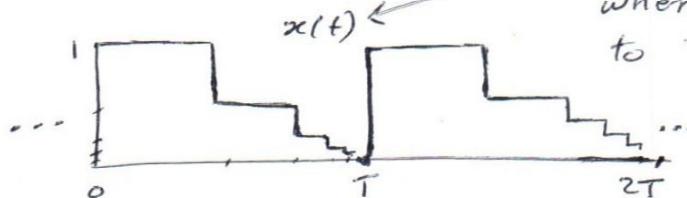
i.e. $x(t)$ has a finite energy over a period.

Dirichlet conditions are as follow :-

① $\int_T |x(t)| dt < \infty \equiv$ absolutely integrable over any period.

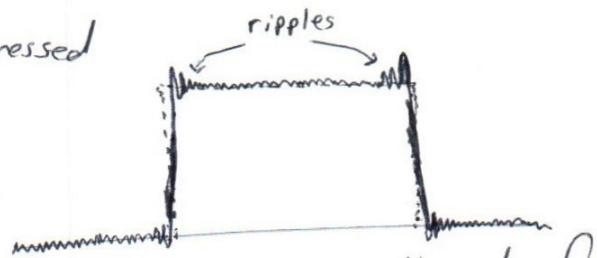
② $x(t)$ has a finite number of maxima and minima during any single period \equiv counter example $x(t) = \frac{1}{t}; 0 < t \leq 1$
with period = 1

③ $x(t)$ has a finite number of discontinuities over any finite interval of time \equiv counter example the value of $x(t)$ decreases by a factor of 2 whenever the distance from t to T decreases by a factor of 2



- Gibbs phenomenon :-

- An American physicist, Albert Michelson, tested the convergence of Fourier series through an experiment
- He uses the approximation formula $x_N(t) = \sum_{k=-N}^{+N} a_k e^{j k \omega_0 t}$ in a trial to test the position of discontinuity of the square wave.
- He found that around the position of discontinuity there are ripples that do not vanish as he increases N . instead the ripples are compressed but its amplitude does not decreases.
- it is called Gibbs phenomenon as the mathematical physicist "Josiah Gibbs" is who explained it mathematically. He found the height of ripples = 0.09% of the height of discontinuity.



Properties of continuous-time Fourier series

① Linearity :-

$$\begin{array}{ccc} x(t) & \xleftrightarrow{\text{FS}} & a_k \\ y(t) & \xleftrightarrow{\text{FS}} & b_k \end{array}$$

$$\text{let } w(t) = m x(t) + n y(t) \xleftrightarrow{\text{FS}} c_k$$

$$c_k = \frac{1}{T} \int_T w(t) e^{-jk\omega_0 t} dt$$

$$= \frac{1}{T} \int_T \{m x(t) + n y(t)\} e^{-jk\omega_0 t} dt$$

$$= \frac{1}{T} \int_T m x(t) e^{-jk\omega_0 t} dt + \frac{1}{T} \int_T n y(t) e^{-jk\omega_0 t} dt$$

$$= m \left\{ \frac{1}{T} \int_T x(t) e^{-jk\omega_0 t} dt \right\} + n \left\{ \frac{1}{T} \int_T y(t) e^{-jk\omega_0 t} dt \right\}$$

$$\therefore c_k = m a_k + n b_k$$

② Time-shifting :- if $x(t) \xleftrightarrow{FS} a_k$ [8]

$$\text{and } y(t) = x(t-t_0) \xleftrightarrow{FS} b_k$$

$$\therefore b_k = \frac{1}{T} \int_T y(t) e^{-jk\omega_0 t} dt = \frac{1}{T} \int_T x(t-t_0) e^{-jk\omega_0 t} dt$$

$$\text{let } \tau = t - t_0 \Rightarrow d\tau = dt \Rightarrow t = \tau + t_0$$

$$\therefore b_k = \frac{1}{T} \int_T x(\tau) e^{-jk\omega_0 (\tau+t_0)} d\tau$$

$$= \frac{1}{T} \int_T x(\tau) e^{-jk\omega_0 \tau} e^{-jk\omega_0 t_0} d\tau$$

$$= e^{-jk\omega_0 t_0} \cdot \left\{ \frac{1}{T} \int_T x(\tau) e^{-jk\omega_0 \tau} d\tau \right\} \quad \begin{matrix} \text{as } t \text{ is constant} \\ \text{w.r.t } \tau \end{matrix}$$

$$\therefore b_k = e^{-jk\omega_0 t_0} \cdot a_k$$

$$\therefore \boxed{x(t-t_0) \xleftrightarrow{FS} e^{-jk\omega_0 t_0} a_k}$$

~~Note magnitude does not change just the phase that is changed~~

③ Time-Reversal :- if $x(t) \xleftrightarrow{FS} a_k$

$$\text{and } y(t) = x(-t) \xleftrightarrow{FS} b_k$$

$$\therefore b_k = \frac{1}{T} \int_T y(t) e^{-jk\omega_0 t} dt = \frac{1}{T} \int_T x(-t) e^{-jk\omega_0 t} dt$$

$$\text{let } \tau = -t \Rightarrow d\tau = -dt \text{ and } t = -\tau$$

$$\therefore b_k = \frac{1}{T} \int_T x(\tau) e^{+jk\omega_0 \tau} (-d\tau) \quad \begin{matrix} \{ \text{with boundaries of} \\ \text{integration is still w.r.t } t \} \end{matrix}$$

$$= \frac{1}{T} \int_T x(\tau) e^{jk\omega_0 \tau} d\tau \quad \begin{matrix} \{ \text{with boundaries of integration} \\ \text{is replaced with boundaries} \\ \text{of } \tau \text{ and } \tau = -t \} \end{matrix}$$

as if t is from $-\frac{T}{2}$ to $\frac{T}{2}$

τ will be from $\frac{T}{2}$ to $-\frac{T}{2}$

and by making τ to range from $-\frac{T}{2}$ to $\frac{T}{2}$

You will multiply by (-1) which will be cancelled with the $-ve$ associated with $(-d\tau)$

$$\therefore \boxed{b_k = a_{-k} \xleftrightarrow{FS} x(-t)} \quad \begin{matrix} \{ \text{i.e. time-reversal} \\ \text{make time-reversal} \\ \text{in the coefficients} \} \end{matrix}$$

Another proof

$$\text{if } x(t) = \sum_{k=-\infty}^{+\infty} a_k e^{jk\omega_0 t} \Rightarrow x(t) = \sum_{k=-\infty}^{+\infty} a_k e^{-jk\omega_0 t}$$

$$\text{let } m = -k \Rightarrow x(-t) = \sum_{m=+\infty}^{-\infty} a_m e^{jm\omega_0 t} = \sum_{m=-\infty}^{+\infty} a_m e^{jm\omega_0 t}$$

$$\therefore \boxed{x(t) \xleftrightarrow{FS} b_k = a_{-k}}$$

- As a consequence of time-reversal property :-

(A) if $x(t)$ is even $\Rightarrow x(t) = x(-t)$

(B) if $x(t)$ is odd $\Rightarrow x(t) = -x(-t)$

(4) Time-scaling :-

if $x(t) = \sum_{k=-\infty}^{+\infty} q_k e^{jk(\omega_0)t}$

then $x(\alpha t) = \sum_{k=-\infty}^{+\infty} q_k e^{jk(\alpha \omega_0)t}$

which means as the fundamental period becomes $(\frac{T}{\alpha})$
then the fundamental frequency becomes $(\alpha \omega_0)$
but the coefficients of Fourier series not changed

$x(\alpha t) \xleftrightarrow{FS} q_k$

(5) Multiplication :-

if $x(t) \xleftrightarrow{FS} q_k$
 $y(t) \xleftrightarrow{FS} b_k$

then $x(t) \cdot y(t) \xleftrightarrow{FS} c_k$

$$x(t) \cdot y(t) = \left\{ \sum_{k=-\infty}^{+\infty} q_k e^{jk\omega_0 t} \right\} \left\{ \sum_{m=-\infty}^{+\infty} b_m e^{jm\omega_0 t} \right\}$$

$\therefore c_k = \frac{1}{T} \int_T x(t) \cdot y(t) e^{-jk\omega_0 t} dt$

$$= \frac{1}{T} \int_T \left\{ \sum_{m=-\infty}^{+\infty} q_m e^{jm\omega_0 t} \right\} \left\{ \sum_{l=-\infty}^{+\infty} b_l e^{jl\omega_0 t} \right\} e^{-jk\omega_0 t} dt$$

$$= \frac{1}{T} \sum_{m=-\infty}^{+\infty} q_m \sum_{l=-\infty}^{+\infty} b_l \int_T e^{j(m+l-k)\omega_0 t} dt$$

remember $\int_T e^{j(m+l-k)\omega_0 t} dt = \begin{cases} T & ; k = m + l \\ 0 & ; \text{otherwise} \end{cases}$
as $e^{jG} = \cos G + j \sin G$

and integrating the cos or sin over one period = zero

$\therefore c_k = \sum_{m=-\infty}^{+\infty} q_m \sum_{l=-\infty}^{+\infty} b_l \cdot \frac{1}{T} \cdot \begin{cases} T & \text{if } k = m + l \\ 0 & \text{or } l = k - m \end{cases}$

$\therefore c_k = \sum_{m=-\infty}^{+\infty} q_m b_{k-m} \equiv c(n) = \sum_{k=-\infty}^{+\infty} q(k) b(n-k)$

i.e. $c_k \equiv$ discrete-time convolution between q_k and b_k

⑥ Conjugation ::

$$\text{if } x(t) \xleftrightarrow{\text{FS'}} q_k$$

$$x^*(t) \xleftrightarrow{\text{FS'}} b_k$$

$$\therefore b_k = \frac{1}{T} \int_T x^*(t) e^{-jkw_0 t} dt$$

$$x(t) = \sum_{k=-\infty}^{+\infty} q_k e^{jkw_0 t}$$

$$\therefore x^*(t) = \left\{ \sum_{k=-\infty}^{+\infty} q_k e^{jkw_0 t} \right\}^*$$

$$\text{Recall } \{ \bar{z}_1 + \bar{z}_2 \}^* = \bar{z}_1^* + \bar{z}_2^*$$

$$\therefore x^*(t) = \sum_{k=-\infty}^{+\infty} \{ q_k e^{jkw_0 t} \}^*$$

$$\text{Recall } \{ \bar{z}_1 \bar{z}_2 \}^* = \bar{z}_1^* \bar{z}_2^*$$

$$\therefore x^*(t) = \sum_{k=-\infty}^{+\infty} q_k^* e^{-jkw_0 t}$$

$$\text{let } m = -k \Rightarrow x^*(t) = \sum_{m=+\infty}^{-\infty} q_{-m}^* e^{jmw_0 t}$$

by comparing to equation # ① $\Rightarrow b_k = a_{-k}^*$

if $x(t)$ is real $\Rightarrow x^*(t) = x(t) \Rightarrow q_k = q_{-k}^* \text{ or } q_{-k} = a_k^*$
i.e. conjugate symmetric and $|q_k| = |q_{-k}|$ and q is real

if $x(t)$ is real and even i.e. $q_k = q_{-k} \Rightarrow q_k = q_k^*$

if $x(t)$ is real and odd i.e. Fourier series coefficients are real and even

$q_k = -q_{-k} \Rightarrow q_k^* = -q_k$ and $q_{-k}^* = -q_{-k} \Rightarrow q_k$ pure imaginary and odd

⑦ Parseval's Relation :: The average in one period of periodic signal

$$x(t) \equiv \frac{1}{T} \int_T |x(t)|^2 dt$$

$$= \frac{1}{T} \int_T |x(t)|^2 dt = \frac{1}{T} \int_T \left| \sum_{k=-\infty}^{+\infty} q_k e^{jkw_0 t} \right|^2 dt$$

$$= \sum_{k=-\infty}^{+\infty} |q_k|^2 \left[\frac{1}{T} \int_T |e^{jkw_0 t}|^2 dt \right]$$

$$\text{as } |A \cdot B| = |A| \cdot |B|$$

$$\therefore \frac{1}{T} \int_T |x(t)|^2 dt = \sum_{k=-\infty}^{+\infty} |q_k|^2 \left[\frac{1}{T} \int_T dt \right] \text{ as } |e^{j\theta}|^2 = 1$$

$$= \sum_{k=-\infty}^{+\infty} |q_k|^2 \left[\frac{1}{T} \cdot T \right]$$

$$\therefore \frac{1}{T} \int_T |x(t)|^2 dt = \sum_{k=-\infty}^{+\infty} |q_k|^2$$

as $|q_k|$ is the average power in k th harmonic

\Rightarrow The total average power in a periodic signal equals the sum of the average powers in all of its harmonic components ~~of~~

⑧ Periodic convolution: if $x(\tau) \xleftrightarrow{FS} a_k$
and $y(\tau) \xleftrightarrow{FS} b_k$

$$\therefore x(\tau) = \sum_m a_m e^{jm\omega_0\tau}$$

$$\text{and } y(\tau) = \sum_n b_n e^{jn\omega_0\tau}$$

$$\therefore y(t-\tau) = \sum_n b_n e^{jn\omega_0(t-\tau)} = \sum_n b_n e^{jn\omega_0t} e^{-jn\omega_0\tau}$$

$$\therefore x(\tau) \cdot y(t-\tau) = \left\{ \sum_m a_m e^{jm\omega_0\tau} \right\} \left\{ \sum_n b_n e^{jn\omega_0t} e^{-jn\omega_0\tau} \right\}$$

$$= \sum_m a_m \sum_n b_n e^{jn\omega_0t} \cdot e^{j(m-n)\omega_0\tau}$$

$$\therefore \int_T x(\tau) y(t-\tau) d\tau = \int_T \sum_m a_m \sum_n b_n e^{jn\omega_0t} e^{j(m-n)\omega_0\tau} d\tau$$

$$= \sum_m a_m \sum_n b_n e^{jn\omega_0t} \int_T e^{j(m-n)\omega_0\tau} d\tau$$

$$= \sum_m a_m \sum_n b_n e^{jn\omega_0t} \cdot \begin{cases} T & \text{if } n=m \\ 0 & \text{if } n \neq m \end{cases}$$

$$= T \sum_m a_m b_m e^{jm\omega_0t}$$

$$\therefore \int_T x(\tau) y(t-\tau) d\tau = T \sum_m a_m b_m e^{jm\omega_0t} \xleftrightarrow{FS} c_k$$

$$\therefore c_k = \frac{1}{T} \int_T \left\{ T \sum_m a_m b_m e^{jm\omega_0t} \right\} e^{-jk\omega_0t} dt$$

$$= \frac{1}{T} \cancel{\int_T} \sum_m a_m b_m \cdot \int_T e^{j(m-k)\omega_0t} dt$$

$$= \sum_m a_m b_m \cdot \begin{cases} T & \text{if } m=k \\ 0 & \text{if } m \neq k \end{cases}$$

$\therefore \boxed{c_k = T a_k b_k}$



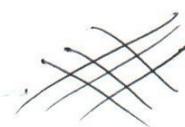
⑨ Differentiation property: if $x(t) \xleftrightarrow{FS} a_k$

$$\therefore x(t) = \sum_K a_K e^{jk\omega_0 t}$$

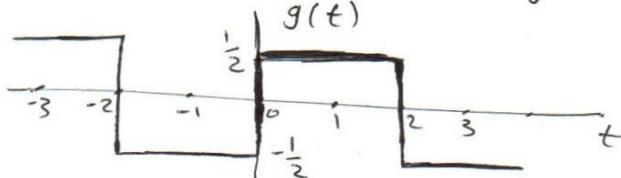
$$\therefore \frac{d}{dt} x(t) = \frac{d}{dt} \left\{ \sum_K a_K e^{jk\omega_0 t} \right\} = \sum_K a_K \cdot jk\omega_0 \cdot e^{jk\omega_0 t}$$

$$\Rightarrow \frac{d}{dt} x(t) = \sum_K jk\omega_0 a_K \cdot e^{jk\omega_0 t}$$

$$\therefore \boxed{\frac{d}{dt} x(t) \xleftrightarrow{FS} jk\omega_0 \cdot a_K}$$

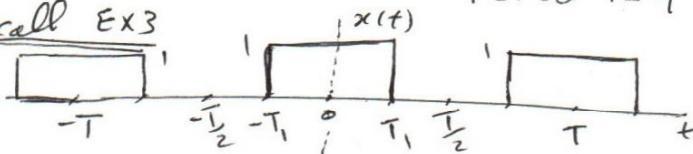


Ex 4:- If $g(t)$ is as shown in Figure



$g(t)$ is periodic with period $T=4$

Recall Ex 3,



if $T=4$ and $T_1=1$

$$\therefore g(t) = x(t-1) - \frac{1}{2}$$

Recall $x(t) \xleftrightarrow{\text{F.S.'}} q_K = \begin{cases} \frac{\sin(k\omega_0 T_1)}{k\pi} & ; k \neq 0 \\ \frac{2T_1}{T} & ; k=0 \end{cases}$

$$\therefore a_K = \begin{cases} \frac{2}{4} = \frac{1}{2} & ; k=0 \\ \frac{\sin(k\omega_0)}{k\pi} & ; k \neq 0 \end{cases}$$

using F.S' properties

$$x(t) \xleftrightarrow{\text{F.S.'}} q_K$$

$$\therefore x(t-1) \xleftrightarrow{\text{F.S.'}} e^{-jk\omega_0} q_K = c_K$$

$$\omega_0 = \frac{2\pi}{T} = \frac{2\pi}{4} = \frac{\pi}{2}$$

$$\therefore x(t-1) \xleftrightarrow{\text{F.S.'}} \left[e^{-jk\frac{\pi}{2}} \cdot q_K \right] = c_K$$

If $g(t) = -\frac{1}{2} \xleftrightarrow{\text{F.S.'}} d_K$

$$\therefore d_K = \frac{1}{4} \int_{-T}^T (-\frac{1}{2}) e^{-j k \omega_0 t} dt$$

$$= -\frac{1}{8} \cdot \begin{cases} 4 & ; k=0 \\ 0 & ; k \neq 0 \end{cases}$$

$$\therefore d_K = \begin{cases} -\frac{1}{2} & ; k=0 \\ 0 & ; k \neq 0 \end{cases}$$

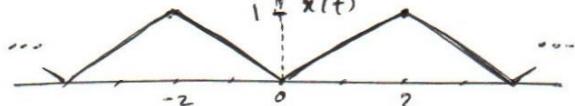
$$\therefore g(t) \xleftrightarrow{\text{F.S.'}} b_K = c_K + d_K$$

$$\therefore b_K = a_K \cdot e^{-jk\frac{\pi}{2}} + d_K$$

$$= \begin{cases} 0 + \frac{\sin(k\omega_0)}{k\pi} \cdot e^{-jk\frac{\pi}{2}} & ; k \neq 0 \\ -\frac{1}{2} + \frac{1}{2} & ; k=0 \end{cases}$$

$$\therefore b_K = \begin{cases} \frac{\sin(k\cdot\frac{\pi}{2})}{k\pi} \cdot e^{-jk\frac{\pi}{2}} & ; k \neq 0 \\ 0 & ; k=0 \end{cases}$$

Ex 5:- If $x(t)$ is as shown in Figure [12]



is periodic with period $T=4$

Note: $\leftarrow g(t) = \frac{d}{dt} x(t)$
of Ex 4 if $g(t) \xleftrightarrow{\text{F.S.'}} b_K$

and $x(t) \xleftrightarrow{\text{F.S.'}} e_K$
using F.S' properties

$$\therefore b_K = j k \omega_0 \cdot e_K = \frac{j k \pi}{2} \cdot e_K$$

$$\therefore e_K = \frac{2 b_K}{j k \pi} = \frac{2 \sin(k\frac{\pi}{2}) \cdot e}{j k \pi \cdot k \pi} ; k \neq 0$$

For $k=0 \Rightarrow e_K = \text{average of one period of } x(t)$

$$\therefore e_K = \begin{cases} \frac{2 \sin(k\frac{\pi}{2})}{j (k\pi)^2} \cdot e^{-jk\frac{\pi}{2}} & ; k \neq 0 \\ \frac{1}{2} & ; k=0 \end{cases}$$

Ex 6:- $x(t) = \sum_{k=-\infty}^{+\infty} \delta(t - kT)$

$$\begin{aligned} q_K &= \frac{1}{T} \int_T x(t) e^{-jk\omega_0 t} dt \\ &= \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} \delta(t) e^{-jk(\frac{2\pi}{T})t} dt \end{aligned}$$

$$\therefore q_K = \frac{1}{T}$$

If $g(t) \xrightarrow{\text{F.S.'}}$

and if $g(t) = \frac{d}{dt} g(t)$
 $\therefore q(t) \xrightarrow{\text{F.S.'}}$

$$\therefore q(t) = x(t+T_1) - x(t-T_1)$$

$$\text{If } q(t) \xleftrightarrow{\text{F.S.'}} b_K$$

$$\therefore b_K = e^{jk\omega_0 T_1} q_K - e^{-jk\omega_0 T_1} q_K$$

$$\therefore b_K = q_K [e^{jk\omega_0 T_1} - e^{-jk\omega_0 T_1}]$$

$$\therefore b_K = \frac{1}{T} \cdot 2j \sin(k\omega_0 T_1)$$

$$\text{as } q(t) = \frac{d}{dt} g(t)$$

$$\therefore b_K = j k \omega_0 \cdot c_K \quad \text{as } g(t) \xleftrightarrow{\text{F.S.'}} c_K$$

$$\therefore c_K = \frac{b_K}{j k \omega_0} = \frac{2j \sin(k\omega_0 T_1)}{j k \omega_0 T} ; k \neq 0$$

$$c_0 = \text{average value of one period} = \frac{2T_1}{T}$$

$$\therefore c_K = \begin{cases} \frac{\sin(k\omega_0 T_1)}{k\pi} & ; k \neq 0 \\ \frac{2T_1}{T} & ; k=0 \end{cases}$$

Ex 7 :- Given the following facts about a signal $x(t)$

1- $x(t)$ is a real signal

2- $x(t)$ is periodic with period $T=4$ and it has Fourier series coefficients a_k

3- $a_k = 0$ for $|k| > 1$

4- The signal with Fourier coefficients $b_k = e^{-jk\frac{\pi}{2}} \cdot a_k$ is odd.

$$5- \frac{1}{4} \int_4 |x(t)|^2 dt = \frac{1}{2}$$

Determine $x(t)$ within a sign factor?

Solution

Fact #3

The non-zero coefficients are :- a_0 , a_1 , and a_{-1}

$$\therefore x(t) = a_0 + a_1 e^{j t \cdot w_0} + a_{-1} e^{j(-1) \cdot w_0}$$

Fact #2

$$T=4 \Rightarrow w_0 = \frac{2\pi}{4} = \frac{\pi}{2}$$

$$\therefore x(t) = a_0 + a_1 e^{jt\frac{\pi}{2}} + a_{-1} e^{-jt\frac{\pi}{2}}$$

Fact #1

as $x(t)$ is real $\Rightarrow a_{-k} = a_k^*$

$\therefore a_{-1} = a_1^*$ and a_0 is real

$$\therefore x(t) = a_0 + a_1 e^{jt\frac{\pi}{2}} + a_1^* e^{-jt\frac{\pi}{2}}$$

Fact #4

- From time-reversal property

$$\text{if } x(t) \xleftrightarrow{\text{FS}} a_k$$

$$\text{then } x(-t) \xleftrightarrow{\text{FS}} a_{-k}$$

- From time-shifting property

$$x(t-t_0) \xleftrightarrow{\text{FS}} e^{-jkw_0 t_0} a_k$$

$$\therefore x(-(t-1)) \xleftrightarrow{\text{FS}} b_k = e^{-jk\frac{\pi}{2}} \cdot a_{-k}$$

as $y_1(t) = x(-t)$ time-reversal

then $y_1(t) = y_1(t-1)$ time-shift

$$= x(-(t-1)) = x(-t+1)$$

$$\therefore x(-t+1) \xleftrightarrow{\text{FS}} e^{-jk\frac{\pi}{2}} \cdot a_{-k}$$

- (From Fact #4) $y(t) = x(-t+1)$ is odd

- From fact #1 as $x(t)$ is real 13

$\Rightarrow x(-t+1)$ is also real.

$\therefore x(-t+1)$ is real and odd

$\therefore b_k$ is pure-imaginary and odd

$$\Rightarrow b_0 = 0$$

$$\text{and } b_{-1} = -b_1$$

Fact #5

average power over one

$$\text{period of } x(t) = \frac{1}{2}$$

- As the time-reversal and time-shift cannot change the average power

\therefore the average power of $x(t)$ = the average power of $x(-t+1)$

then Fact #5 holds for $x(-t+1)$ too

$$\therefore \frac{1}{4} \int_4 |x(-t+1)|^2 dt = \frac{1}{2}$$

(From Parseval's relation)

$$\therefore \sum_k |b_k|^2 = \frac{1}{2}$$

$$\text{as } b_0 = 0$$

$$\therefore |b_1|^2 + |b_{-1}|^2 = \frac{1}{2}$$

as b_k is odd $\Rightarrow b_1 = -b_{-1}$

$$\therefore |b_1|^2 + |-b_1|^2 = 2|b_1|^2 = \frac{1}{2}$$

$$\therefore |b_1|^2 = \frac{1}{4} \Rightarrow |b_1| = \frac{1}{2}$$

as b_k is pure imaginary conjugation property

$\therefore b_1$ is either $\frac{1}{2}j$ or $-\frac{1}{2}j$

$$\therefore b_k = e^{-jk\frac{\pi}{2}} \cdot a_k \Rightarrow b_0 = a_0 = 0 \quad \text{Fact #4}$$

$$\text{if } k=-1 \Rightarrow b_{-1} = e^{-j\frac{\pi}{2}} a_1 \Rightarrow a_1 = -b_1 \cdot e^{-j\frac{\pi}{2}}$$

$$\Rightarrow a_1 = (-b_1) \cdot (-j) = j b_1$$

$$\text{Case 1: } \therefore b_1 = \frac{1}{2}j \Rightarrow a_1 = -\frac{1}{2}$$

$$\Rightarrow x(t) = \frac{1}{2} \left\{ e^{j\frac{\pi}{2}} + e^{-j\frac{\pi}{2}} \right\} = -\cos(\frac{\pi}{2})$$

$$\text{Case 2: } \therefore b_1 = -\frac{1}{2}j \Rightarrow a_1 = \frac{1}{2}$$

$$\Rightarrow x(t) = \frac{1}{2} \left\{ e^{j\frac{\pi}{2}} + e^{-j\frac{\pi}{2}} \right\} = \cos(\frac{\pi}{2})$$