

2. Quantum Computational Models

(1)

(warwick.ac.uk/qinfo)

ANIMESH DATTA

Church-Turing thesis: is a hypothesis about the nature of computable functions.

The heuristic agreement is that all algorithmically computable functions are Turing-computable.



Crudely speaking, this is a classical computer with unlimited memory.

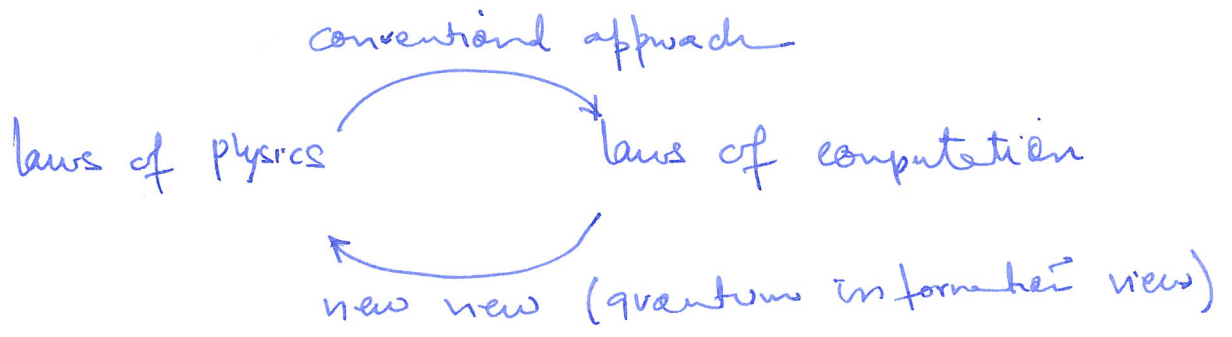
But it addresses and suggests connection to the space and time requirements for computations. It asks fundamental questions such as

1. What is physical computation?
2. What is the potential and limitation of physical computation?
3. What is meant by the universality of computation?

The long-standing open question is its derivation from the laws of physics. A modern, sharper conjecture:

~~Church~~ Church-Turing-Deutsch Principle:

Every physical process can be simulated efficiently by a universal computing device.



→ Nature is a massive computing device

→ The laws of physics (including the undiscovered theories of the future) accommodate the notion of universal computation

Current Status of observation: Classical Computers do not seem to be able to simulate quantum systems efficiently

Possible resolutions:

- ① A quantum Turing machine should be a universal computing model
- ② An efficient classical simulation of quantum systems is indeed possible.

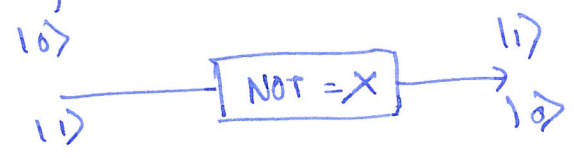
The following models are known to be universal computing devices/systems:

- Quantum Turing machine
- Quantum circuit model
 - quantum cellular automata
 - Topological computation
 - Holonomic computation
- Quantum adiabatic computation
- Quantum random walk
- Measurement-based quantum computation
 - teleportation-based
 - Entanglement-based.

Quantum circuits:

bit strings \longleftrightarrow state vectors
 gates \longleftrightarrow unitary operations
 measurements \longleftrightarrow Born rule for probs.

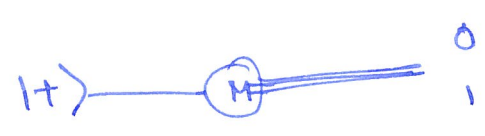
Since vectors ^{and operators} need a basis, we will conventionally use the eigenstates of a Pauli-Z operator, $|0\rangle$ and $|1\rangle$.



$$X|a\rangle = |\bar{a}\rangle = |1 \oplus a\rangle$$

$$|a\rangle \xrightarrow{Z} (-1)^a |a\rangle \leftarrow (\text{Eigen states})$$

Measurement:



$$p_0 = |\langle 0|+\rangle|^2 = 1/2$$

$$p_1 = |\langle 1|+\rangle|^2 = 1/2$$

In general:

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle \xrightarrow{U} \alpha U|0\rangle + \beta U|1\rangle$$

For the conservation of probabilities:

$$\begin{pmatrix} \alpha^* & \beta^* \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = 1$$

$$\begin{pmatrix} \alpha^* & \beta^* \end{pmatrix} U^\dagger U \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = 1$$

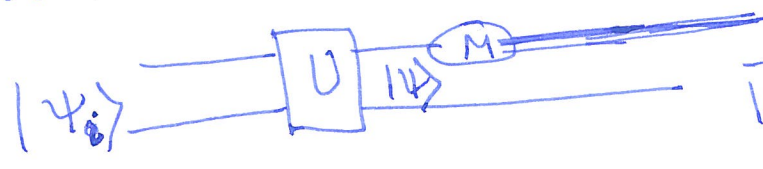
$$\Rightarrow U^\dagger U = I \quad \therefore (U \text{ is unitary})$$

$$= \alpha (U_{11}|0\rangle + U_{21}|1\rangle) + \beta (U_{12}|0\rangle + U_{22}|1\rangle)$$

$$= \begin{pmatrix} U_{11} & U_{12} \\ U_{21} & U_{22} \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix}$$

A general measurement:

$$|\psi\rangle = C_{00}|00\rangle + C_{01}|01\rangle + C_{10}|10\rangle + C_{11}|11\rangle$$



$$0: p_0 = |C_{00}|^2 + |C_{01}|^2$$

$$1: p_1 = |C_{10}|^2 + |C_{11}|^2$$

$$|\psi_0\rangle = \frac{C_{00}|0\rangle + C_{01}|1\rangle}{\sqrt{p_0}}$$

$$|\psi_1\rangle = \frac{C_{10}|0\rangle + C_{11}|1\rangle}{\sqrt{p_1}}$$

$$p_0 = |\langle 00 | U | \psi_i \rangle|^2 + |\langle 01 | U | \psi_i \rangle|^2$$

$$= \langle \psi_i | U^\dagger | 00 \rangle \langle 00 | U | \psi_i \rangle + \langle \psi_i | U^\dagger | 01 \rangle \langle 01 | U | \psi_i \rangle$$

$$= \langle \psi_i | U^\dagger (P_0 \otimes I) U | \psi_i \rangle$$

where $P_0 = |0\rangle\langle 0| \otimes I$

$$= |00\rangle\langle 00| + |01\rangle\langle 01|$$

and the state after the

measurement is $\frac{1}{\sqrt{p_0}} (P_0 \otimes I) U | \psi_{in} \rangle$

So, what's going on?

1. An arbitrary classical gate M becomes a linear operator

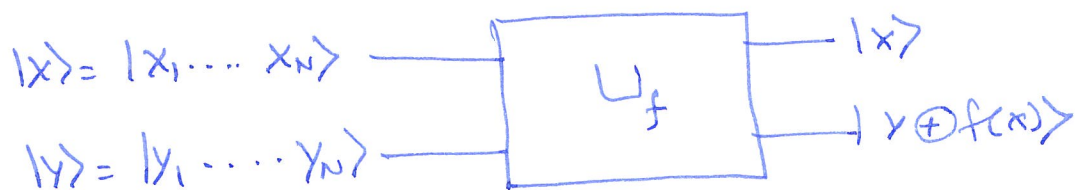
$$A|x\rangle = |Mx\rangle, \quad \langle y | A|x\rangle = \delta_{y, Mx}$$

2. A reversible classical gate becomes a permutation unitary

3. There are many more unitaries than reversible classical gates.

4. Born rule for measurements.

Function evaluation in quantum circuits:



If f is a Boolean function, U_f is a controlled bit flip op.

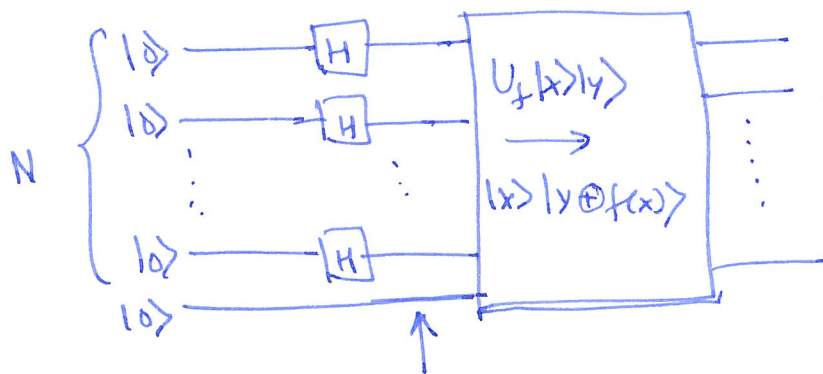
The control qubits are $|x\rangle = |x_1, \dots, x_n\rangle$, the value of the control is determined by $f(x)$.

If $f(x)=0$, the target qubit remains unchanged,

If $f(x)=1$, " " " " " flips.

$$U_f |x\rangle |y\rangle = |x\rangle |y \oplus f(x)\rangle = |x\rangle \otimes X^{f(x)} |y\rangle.$$

Quantum parallelism: f on an N -bit Boolean function.



$$|\psi\rangle = \frac{1}{\sqrt{2^N}} \sum_x |x\rangle |f(x)\rangle$$

All values of $f(x)$ calc. in parallel

$$|\psi\rangle = \frac{1}{\sqrt{2^N}} \sum_x |x\rangle |0\rangle$$

$$\begin{aligned} H|0\rangle &= \frac{|0\rangle + |1\rangle}{\sqrt{2}} \\ H|1\rangle &= \frac{|0\rangle - |1\rangle}{\sqrt{2}} \end{aligned}$$

$H^{\otimes N}$ - Walsh-Hadamard transform

$$H^{\otimes N} |0\rangle^{\otimes N} = \left[\frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \right]^{\otimes N} = \frac{1}{\sqrt{2^N}} \sum_x |x\rangle \rightarrow \text{all possible bit strings}$$

For any bit b :

$$H|b\rangle = \frac{1}{\sqrt{2}} (|0\rangle + (-1)^b |1\rangle) = \frac{1}{\sqrt{2}} \sum_x (-1)^{bx} |x\rangle$$

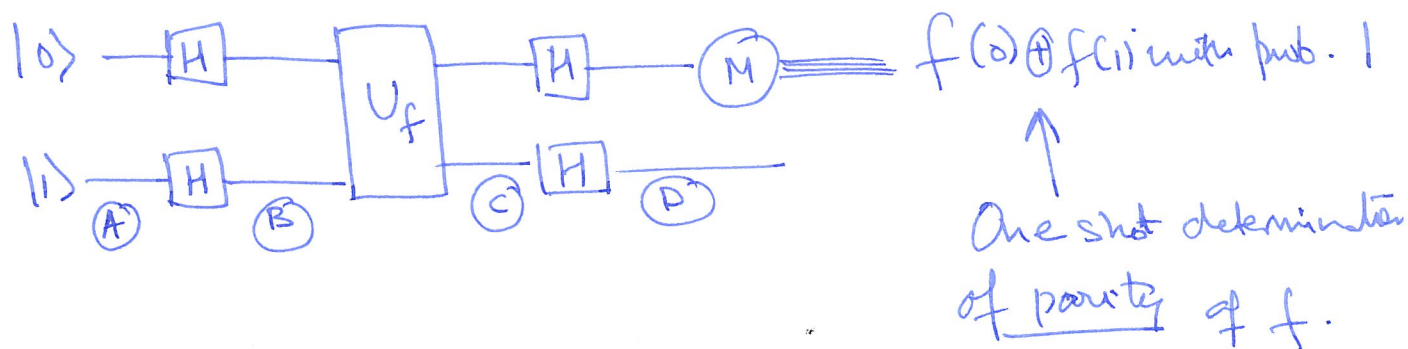
For any bit string s :

$$H^{\otimes N} |s\rangle = \frac{1}{\sqrt{2^N}} \sum_{x_1, \dots, x_n} (-1)^{x_1 s_1 + \dots + x_n s_n} |x_1 \dots x_n\rangle = \frac{1}{\sqrt{2^N}} \sum_x (-1)^{x \cdot s} |x\rangle$$

Deutsch's algorithm:

$$f: \{0,1\} \rightarrow \{0,1\}$$

1-bit Boolean function.



(A): $|\psi_A\rangle = |0\rangle \otimes |1\rangle$

(B): $|\psi_B\rangle = \frac{1}{2} (|0\rangle + |1\rangle) (|0\rangle - |1\rangle)$

(C): $|\psi_C\rangle = \frac{1}{2} (|0\rangle |f_0\rangle - |0\rangle |\bar{f}_0\rangle + |1\rangle |f_1\rangle - |1\rangle |\bar{f}_1\rangle)$

1 qubit
medium 2.

$$= \begin{cases} \frac{1}{2} (|0\rangle + |1\rangle) (|0\rangle - |1\rangle) & : f(0) = f(1) = 0 \\ \frac{1}{2} (|0\rangle - |1\rangle) (|0\rangle - |1\rangle) & : f(0) = 0, f(1) = 1 \\ -\frac{1}{2} (|0\rangle - |1\rangle) (|0\rangle - |1\rangle) & : f(0) = 1, f(1) = 1 \\ -\frac{1}{2} (|0\rangle + |1\rangle) (|0\rangle - |1\rangle) & : f(0) = f(1) = 1 \end{cases}$$

$$= \frac{1}{2} (-1)^{f(0)} (|0\rangle + (-1)^{f(0) \oplus f(1)} |1\rangle) (|0\rangle - |1\rangle).$$

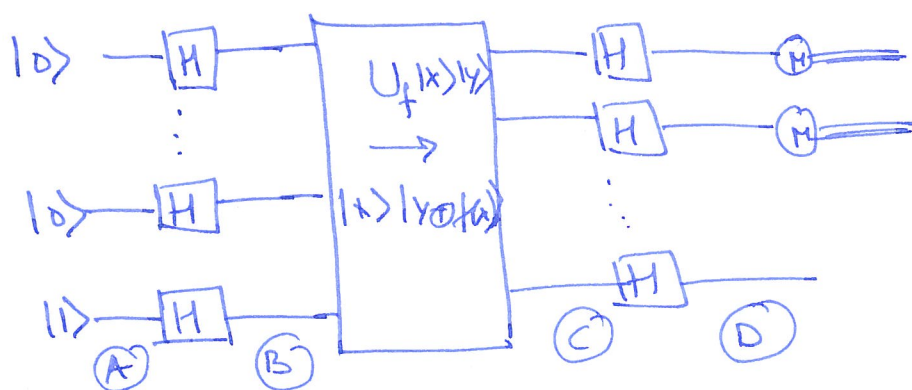
(D): $|\psi_D\rangle = (-1)^{f(0)} |f(0) \oplus f(1)\rangle |1\rangle$

Any classical algorithm to determine the parity of f , or whether f is constant or balanced, needs two runs. Deutsch does it in one.

Deutsch-Jozsa algorithm: To determine whether (7)

$f : \{0,1\}^N \rightarrow \{0,1\}$ is constant or balanced.

The number of classical calls to be sure is $\frac{2^N}{2} = 2^{N-1}$.



Quantum strategy needs just one shot.

Exponential adv. but needs an oracle U_f .

(A): $|\psi_A\rangle = |0\rangle^{\otimes N} \otimes |1\rangle$

(B): $|\psi_B\rangle = H^{\otimes N} |0\rangle^{\otimes N} \otimes H|1\rangle$

$$= \frac{1}{\sqrt{2^N}} \sum_x |x\rangle \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

(C) $|\psi_C\rangle = \frac{1}{\sqrt{2^N}} \sum_x |x\rangle \frac{1}{\sqrt{2}} (|f(x)\rangle - |\overline{f(x)}\rangle)$

$$= \left(\frac{1}{\sqrt{2^N}} \sum_x (-1)^{f(x)} |x\rangle \right) \otimes \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

$|\phi\rangle \equiv$ function value in phase

ancilla qubit uncorrelated

$$\langle \phi^{\text{const}} | \phi^{\text{balanced}} \rangle = \pm \sum_x (-1)^{f(x)} = 0$$

(D) $|\psi_D\rangle = \left(\frac{1}{\sqrt{2^N}} \sum_{x,y} (-1)^{f(x) \oplus x \cdot y} |y\rangle \right) (|1\rangle)$

constant f : $|x\rangle = \pm |0\rangle^{\otimes N}$

balanced f :

$$N \otimes \langle 0|x \rangle = \frac{1}{2^N} \sum_x (-1)^{f(x)} = 0$$