

Numerical solution for the propagation of light along a spatial direction

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Abstract

In this short note, I introduce the equations pertaining to the numerical solution of the equations governing an electromagnetic wave propagating along one spatial direction.

1 The equations and their discretisation

We consider the system of partial differential equations:

$$\begin{aligned}\frac{\partial E}{\partial t}(t, z) &= -\frac{1}{\varepsilon_0} \frac{\partial H}{\partial z}(t, z), \\ \frac{\partial H}{\partial t}(t, z) &= -\frac{1}{\mu_0} \frac{\partial E}{\partial z}(t, z).\end{aligned}$$

Our task is to solve for the fields E and H in these equations over a closed square in (t, z) -space defined by $[0, T] \times [0, z_{\max}]$. Since the solution we seek is numerical, and for stability reasons, we will subdivide $[0, T]$ in $2N$ equal parts, respectively $[0, z_{\max}]$ in $2K$ equal parts, whose length we denote by Δt , respectively Δz . More specifically

$$\Delta t = \frac{T}{2N}, \quad \Delta z = \frac{z_{\max}}{2K}$$

We then want to define the fields over an “intertwined grid”, i.e.,

$$\begin{aligned}E &: \{\Delta t, 3\Delta t, 5\Delta t, \dots, (2N-1)\Delta t\} \times \{0, 2\Delta z, 4\Delta z, \dots, 2K\Delta z\} \longrightarrow \mathbb{R}, \\ H &: \{0, 2\Delta t, 4\Delta t, \dots, 2N\Delta t\} \times \{\Delta z, 3\Delta z, 5\Delta z, \dots, (2K-1)\Delta z\} \longrightarrow \mathbb{R}.\end{aligned}$$

In more concrete notation, we will have values for these fields at the following points:

$$\begin{aligned}E((2n+1)\Delta t, 2k\Delta z), \quad n &= 0, 1, \dots, N-1, \quad k = 0, 1, \dots, K, \\ H(2n\Delta t, (2k+1)\Delta z), \quad n &= 0, 1, \dots, N, \quad k = 0, 1, \dots, K-1.\end{aligned}$$

We’re going to be interested in the finite differences:

$$\begin{aligned}\frac{\partial E}{\partial t}((2n+1)\Delta t, 2k\Delta z) &\approx \frac{E((2n+1)\Delta t, 2k\Delta z) - E((2n-1)\Delta t, 2k\Delta z)}{\Delta t}, \quad n = 1, \dots, N-1, \quad k = 0, 1, \dots, K, \\ \frac{\partial E}{\partial z}((2n+1)\Delta t, 2k\Delta z) &\approx \frac{E((2n+1)\Delta t, 2k\Delta z) - E((2n+1)\Delta t, 2(k-1)\Delta z)}{\Delta z}, \quad n = 0, 1, \dots, N-1, \quad k = 1, \dots, K, \\ \frac{\partial H}{\partial t}(2n\Delta t, (2k+1)\Delta z) &\approx \frac{H(2n\Delta t, (2k+1)\Delta z) - H(2(n-1)\Delta t, (2k+1)\Delta z)}{\Delta t}, \quad n = 1, \dots, N, \quad k = 0, 1, \dots, K-1, \\ \frac{\partial H}{\partial z}(2n\Delta t, (2k+1)\Delta z) &\approx \frac{H(2n\Delta t, (2k+1)\Delta z) - H(2n\Delta t, (2k-1)\Delta z)}{\Delta z}, \quad n = 0, 1, \dots, N, \quad k = 1, \dots, K-1.\end{aligned}$$

Inserting these approximations into the differential equation and simplifying, taking into account the domains of n and k , we get:

$$\begin{aligned}E((2n+1)\Delta t, 2k\Delta z) &= E((2n-1)\Delta t, 2k\Delta z) - \frac{\Delta t}{\varepsilon_0 \Delta z} (H(2n\Delta t, (2k+1)\Delta z) - H(2n\Delta t, (2k-1)\Delta z)), \\ H(2n\Delta t, (2k+1)\Delta z) &= H(2(n-1)\Delta t, (2k+1)\Delta z) - \frac{\Delta t}{\mu_0 \Delta z} (E((2n+1)\Delta t, 2k\Delta z) - E((2n+1)\Delta t, 2(k-1)\Delta z)),\end{aligned}$$

both of which hold for $n = 1, \dots, N-1$ and $k = 1, \dots, K-1$. Renormalising $\tilde{E} := \sqrt{\frac{\varepsilon_0}{\mu_0}} E$, we get

$$\begin{aligned}\tilde{E}((2n+1)\Delta t, 2k\Delta z) &= \tilde{E}((2n-1)\Delta t, 2k\Delta z) - c \frac{\Delta t}{\Delta z} (H(2n\Delta t, (2k+1)\Delta z) - H(2n\Delta t, (2k-1)\Delta z)), \\ H(2n\Delta t, (2k+1)\Delta z) &= H(2(n-1)\Delta t, (2k+1)\Delta z) - c \frac{\Delta t}{\Delta z} (\tilde{E}((2n+1)\Delta t, 2k\Delta z) - \tilde{E}((2n+1)\Delta t, 2(k-1)\Delta z)),\end{aligned}$$

We take into account the stability conditions

$$\Delta z \leq \frac{\lambda}{10}, \quad \text{e.g.,} \quad \Delta z = \frac{\lambda}{10} - \varepsilon,$$

and then the *Courant condition*

$$c \frac{\Delta t}{\Delta z} \leq \frac{1}{2} \implies \Delta t \leq \frac{1}{2c} \Delta z, \quad \text{e.g.,} \quad \Delta t = \frac{1}{2c} \Delta z - \varepsilon'$$