

### 23. Simple Harmonic Motion. (91), (93)

A particle is said to execute Simple Harmonic Motion if it moves in a straight line such that its acceleration is always directed towards a fixed point in the line and is proportional to the distance of the particle from the fixed point.

Let  $O$  be the fixed point on a line  $BOA$  and  $P$  be the position of the particle at time  $t$  where  $OP = x$ , so that the acceleration of the particle in the sense  $OP$  is  $\ddot{x}$ .

Now the given acceleration is towards  $O$  and is proportional to  $x$ . Let it be  $\mu x$ , where  $\mu$  is constant.



Since  $\ddot{x}$  is in the direction of  $OP$  produced and  $\mu x$  is towards  $O$ , the equation of motion is

$$\ddot{x} = -\mu x.$$

$$\frac{d^2x}{dt^2} \propto x \Rightarrow \frac{d}{dx}\left(\frac{dx}{dt}\right) = -\mu x$$

Taking  $v \frac{dv}{dx}$  instead of  $\ddot{x}$ ; we can write the above equation as  $\boxed{v \frac{dv}{dx} = -\mu x}$

$$v \frac{dv}{dx} = -\mu x.$$

$$v dx = -\mu x dt \dots (1)$$

Integrating with respect to  $x$ , we get

$$\frac{v^2}{2} = -\mu \frac{x^2}{2} + C_1 \dots (2)$$

$$\frac{v^2}{2} = -\mu \frac{x^2}{2} + \frac{C}{2} \text{ where } C \text{ is a constant}$$

or

$$v^2 = -\mu x^2 + C.$$

If  $A$  be the extreme position of the particle i.e., it is at rest at  $A$  when  $x = a$ ,  $v = 0$  where  $OA = a$ , we get

$$0 = -\mu a^2 + C \quad \therefore C = \mu a^2$$

$$C = \frac{1}{2} k_1 a^2$$

$$\frac{v^2}{2} = -\frac{\mu x^2}{2} + \frac{1}{2} \mu a^2 \Rightarrow v^2 = -\mu x^2 + \mu a^2$$

$$v^2 = \mu(a^2 - x^2)$$

Hence

$$v^2 = \mu(a^2 - x^2),$$

i.e.,

If the particle moves from  $A$  towards  $O$ ,  $v$  is negative

Hence

$$\dot{x} = v = -\sqrt{\mu} \sqrt{a^2 - x^2}$$

or

$$\frac{dx}{dt} = -\sqrt{\mu} \sqrt{a^2 - x^2}$$

or

Integrating we get  $\sqrt{\mu} \cdot t = \cos^{-1} \frac{x}{a} + C_1$  where  $C_1$  is a constant

Initially at  $A$ ,  $t=0$ ,  $x=a$  i.e., the particle started from  $A$ ,  
then

$$0 = \cos^{-1} 1 + C_1 \quad \therefore C_1 = 0.$$

Hence

$$\sqrt{\mu} \cdot t = \cos^{-1} \frac{x}{a}$$

or

$$x = a \cos \sqrt{\mu} \cdot t.$$

If the particle moves from  $O$  towards  $A$ ,  $v$  is positive

so that

$$\dot{x} = \sqrt{\mu} \sqrt{a^2 - x^2}$$

or

$$\sqrt{\mu} \cdot dt = \frac{dx}{\sqrt{a^2 - x^2}}.$$

Integrating, we get,  $\sqrt{\mu} \cdot t = \sin^{-1} \frac{x}{a} + C_2$  where  $C_2$  is a constant

If the particle starts from  $O$ ,  $t=0$ ,  $x=0$ ,  $0 = \sin^{-1} 0 + C_2$

$\therefore$

$$C_2 = 0,$$

so that

$$x = a \sin \sqrt{\mu} \cdot t.$$

Thus the solution of (1) is  $x = a \cos \sqrt{\mu} \cdot t$  or  $x = a \sin \sqrt{\mu} \cdot t$  according as the starting point is  $A$  or  $O$ .

From (2),

$$v = 0 \text{ when } x = \pm a.$$

Thus if  $B$  is a point on the other side of  $O$  such that  $OB = a$ , the particle comes to rest also at  $B$ . When  $x=0$ ,  $v = \pm \sqrt{\mu} \cdot a$  i.e., at  $O$ , the velocity is  $\sqrt{\mu} \cdot a$ .

Consider the solution  $x = a \cos \sqrt{\mu} \cdot t$ .

The motion starts from  $A$  under an attraction towards  $O$ .

When the particle reaches  $O$ ,  $x=0$ .  $\therefore \cos \sqrt{\mu} \cdot t = 0 \quad \therefore \sqrt{\mu} \cdot t = \frac{\pi}{2}$  i.e.,  $t = \frac{\pi}{2\sqrt{\mu}}$  is the time required in moving from  $A$  to  $O$ .



As the particle reaches  $O$ , the attraction ceases but the particle has a velocity  $\sqrt{\mu} \cdot a$  towards the negative side of  $O$  hence the particle passes  $O$  and moves towards the negative side. As soon as the particle comes to the left side of  $O$ , attraction changes direction and becomes towards  $O$ ; hence the velocity will go on decreasing as the particle moves towards the left, till at  $B$ , the velocity becomes zero so that the particle stops. But the particle is being attracted towards  $O$  hence starts moving towards  $O$  and reaches  $O$  with a velocity  $\sqrt{\mu} \cdot a$ , due to which it passes  $O$  and moves towards  $A$  and again stops at  $A$  where its velocity becomes zero. The motion is then repeated. Thus the motion is from  $A$  to  $B$  and back to  $A$  and so on. The motion is *oscillatory*. Time from  $O$  to  $B$  is equal to that from  $A$  to  $O$  hence the **period** i.e., the time from  $A$  to  $B$  and back to  $A$  is

$$4 \cdot \frac{\pi}{2\sqrt{\mu}} = \frac{2\pi}{\sqrt{\mu}}.$$

The distance  $a (=OA)$  i.e., the distance of the centre from one of the positions of rest is called the **amplitude**.

Thus the period which is equal to  $\frac{2\pi}{\sqrt{\mu}}$  is independent of the amplitude i.e., whatever be the amplitude the period is the same. Thus the simple harmonic motion is oscillatory and periodic, the period being independent of amplitude.

The **frequency** is the number of complete oscillations in one second, so that if  $n$  be the frequency and  $T$  the periodic time,

$$n = \frac{1}{T} = \frac{\sqrt{\mu}}{2\pi}.$$

The equation (1), namely  $\ddot{x} = -\mu x$ , can be solved as a differential equation. The most general solution of this equation is

$$x = A \cos \sqrt{\mu} \cdot t + B \sin \sqrt{\mu} \cdot t \quad \dots(5)$$

$A, B$  are constants to be determined from initial conditions. In the first case when the motion starts from  $A$ , the initial conditions are  $t=0, x=a, \dot{x}=0$ .

Now  $t=0, x=a$  gives  $a=A$ .

Differentiating (5),  $\dot{x} = -A\sqrt{\mu} \sin \sqrt{\mu} \cdot t + B\sqrt{\mu} \cos \sqrt{\mu} \cdot t \quad \dots(6)$

The condition  $t=0, \dot{x}=0$  gives  $0=0+B\sqrt{\mu} \therefore B=0$ .

Hence the solution is  $x = a \cos \sqrt{\mu} t$ .

In the second case when the motion starts from  $O$ , the first condition is  $t=0, x=0$ .  $\therefore 0=A, A=0$ .

Hence  $x = B \sin \sqrt{\mu} t$ .

To determine  $B$ , we must know the velocity of projection from  $O$ .



Let us take the case of a particle, projected from  $A$  with velocity  $V$  along  $OA$  produced, so that the initial conditions are  $t=0$ ,  $x=a$ ,  $\dot{x}=V$ .

Hence from (5) and (6), we get

$$a=A$$

$$V=B\sqrt{\mu}. \quad \therefore B=\frac{V}{\sqrt{\mu}}$$

Hence the solution is  $x=a \cos \sqrt{\mu}t + \frac{V}{\sqrt{\mu}} \sin \sqrt{\mu}t$ .

Also the general solution of (1) can be written as

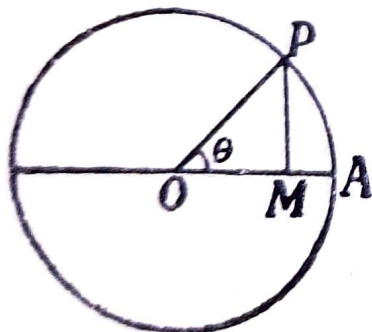
$$x=a \cos (\sqrt{\mu}t + \xi)$$

This is periodic with period  $\frac{2\pi}{\sqrt{\mu}}$ .

The quantity  $\xi$  is called the **epoch**, the angle  $\sqrt{\mu}t + \xi$  is called the **argument**. The particle is at its maximum distance at time  $t_0$  where  $\sqrt{\mu}t_0 + \xi = 0$  i.e.,  $t_0 = -\frac{\xi}{\sqrt{\mu}}$ . Hence the time that has elapsed since the particle was at its maximum distance is equal to  $t - t_0 = t + \frac{\xi}{\sqrt{\mu}} = \frac{\sqrt{\mu}t + \xi}{\sqrt{\mu}}$ . This is the **phase** at time  $t$ .

✓ A geometrical representation of the S.H.M.

Let a particle  $P$  move on a circle with constant angular velocity  $\omega$  and let  $M$  be the foot of the perpendicular from  $P$  on any diameter  $OA$ . If  $a$  be the radius of the circle, the only acceleration of  $P$  is  $\omega^2 a$  towards  $O$ .



If  $\angle AOP = \theta$  and  $OM = x$ , the component of this acceleration along  $OA$

$$= \omega^2 a \cdot \cos \theta = \omega^2 a \cdot \frac{x}{a} = \omega^2 x \text{ towards } O.$$

Hence the equation of motion of the point  $M$  is

$$\ddot{x} = -\omega^2 x.$$

This is S.H.M.

Thus if a particle describes a circle with constant angular velocity the foot of the perpendicular from it on any diameter executes simple harmonic motion.