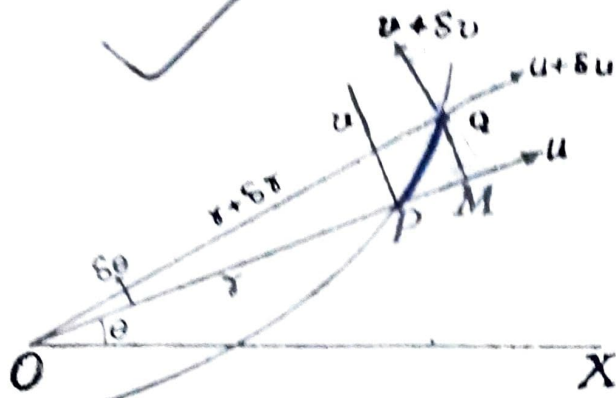


1.4. Radial and Transverse Components of Velocity and Acceleration:

A particle is moving in a plane curve, to find components of velocity and acceleration at time t along and perpendicular to the radius vector drawn from a fixed point in the plane.



Take the fixed point O as the pole and line OX as the initial line. Let P be the position of the particle at time t , its coordinates be (r, θ) and Q be the position at time $t + \delta t$, its coordinates $(r + \delta r, \theta + \delta \theta)$, so that the chord PQ is the displacement in time δt . Draw QM perpendicular from Q to OP so that PM and QM are the components of the displacement PQ along and perpendicular to OP . Let u, v be the components of velocity along and perpendicular to OP .

Then

$$u = \lim_{\delta t \rightarrow 0} \frac{\text{displacement along } OP \text{ in time } \delta t}{\delta t}$$

$$= \lim_{\delta t \rightarrow 0} \frac{PM}{\delta t} = \lim_{\delta t \rightarrow 0} \frac{OM - OP}{\delta t}$$

$$= \lim_{\delta t \rightarrow 0} \frac{OQ \cos \delta \theta - OP}{\delta t}$$

$$= \lim_{\delta t \rightarrow 0} \frac{(r + \delta r) \cos \delta \theta - r}{\delta t}$$

$$= \lim_{\delta t \rightarrow 0} \frac{(r + \delta r) 1 - r}{\delta t}, \text{ small quantities of above the first order being neglected.}$$

$$= \lim_{\delta t \rightarrow 0} \frac{\delta r}{\delta t} = \frac{dr}{dt} = \dot{r} \quad \text{radial velocity}$$

$$v = \lim_{\delta t \rightarrow 0} \frac{\text{displacement perp. to } OP \text{ in time } \delta t}{\delta t}$$

$$= \lim_{\delta t \rightarrow 0} \frac{QM}{\delta t} = \lim_{\delta t \rightarrow 0} \frac{OQ \sin \delta \theta}{\delta t}$$

$$= \lim_{\delta t \rightarrow 0} \frac{(r + \delta r) \sin \delta \theta}{\delta t} = \lim_{\delta t \rightarrow 0} \frac{(r + \delta r) \sin \delta \theta}{\delta \theta} \cdot \frac{\delta \theta}{\delta t}$$

$$= \lim_{\delta t \rightarrow 0} (r + \delta r) \frac{\delta \theta}{\delta t} \text{ because } \lim_{\delta \theta \rightarrow 0} \frac{\sin \delta \theta}{\delta \theta} = 1$$

$$= \lim_{\delta t \rightarrow 0} \frac{r \delta \theta}{\delta t}, \text{ neglecting the other term}$$

$$= r \frac{d\theta}{dt} = r\dot{\theta} \quad \text{transverse velocity}$$

Thus the components of velocity along and perpendicular to the radius vector are \dot{r} and $r\dot{\theta}$, in the senses in which r and θ increase. These are called the radial and transverse or cross-radial components of velocity.

Now let the components of velocity along and perpendicular to OQ be $u + \delta u$, $v + \delta v$; (u, v) being those along and perpendicular to OP .

Thus the change of velocity along OP in time δt

$$= (u + \delta u) \cos \delta \theta - (v + \delta v) \sin \delta \theta - u$$

$$= (u + \delta u) \cdot 1 - (v + \delta v) \cdot \delta \theta - u, \quad \text{neglecting higher powers of } \delta \theta$$

$$= \delta u - v \delta \theta, \quad \text{neglecting the other term.}$$

Similarly the change of velocity perpendicular to OP in time δt

$$= (u + \delta u) \sin \delta \theta + (v + \delta v) \cos \delta \theta - v$$

$$= (u + \delta u) \delta \theta + (v + \delta v) \cdot 1 - v, \text{ as before}$$

$$= u \delta \theta + \delta v, \quad \text{neglecting the other term.}$$

Therefore, radial acceleration

$$= \lim_{\delta t \rightarrow 0} \frac{\text{change of velocity along } OP \text{ in time } \delta t}{\delta t}$$

$$= \lim_{\delta t \rightarrow 0} \frac{\delta u - v \delta \theta}{\delta t}$$

$$= \frac{du}{dt} - v \frac{d\theta}{dt}$$

$$= \frac{d}{dt} \left(\frac{dr}{dt} \right) - r \frac{d\theta}{dt} \cdot \frac{d\theta}{dt}$$

$$= \frac{d^2 r}{dt^2} - r \left(\frac{d\theta}{dt} \right)^2 = \boxed{\ddot{r} - r\dot{\theta}^2}$$

get

Transverse acceleration = $\lim_{\delta t \rightarrow 0} \frac{\text{change of velocity perp. to } OP \text{ in time } \delta t}{\delta t}$

$$\begin{aligned}
 &= \lim_{\delta t \rightarrow 0} \frac{u\delta\theta + \delta v}{\delta t} \\
 &= \lim_{\delta t \rightarrow 0} \frac{d\theta}{u \frac{dt}{dt}} + \frac{dv}{dt} \\
 &= \frac{dr}{dt} \cdot \frac{d\theta}{dt} + \frac{d}{dt} \left(r \frac{d\theta}{dt} \right) \\
 &= \frac{dr}{dt} \cdot \frac{d\theta}{dt} + \frac{dr}{dt} \cdot \frac{d\theta}{dt} + r \frac{d^2\theta}{dt^2} \\
 &= r \frac{d^2\theta}{dt^2} + 2 \frac{dr}{dt} \cdot \frac{d\theta}{dt} \\
 &= r\ddot{\theta} + 2\dot{r}\dot{\theta} \\
 &= \frac{1}{r} \frac{d}{dt} \left(r^2 \frac{d\theta}{dt} \right) \\
 &= \boxed{\frac{1}{r} \frac{d}{dt} (r^2 \dot{\theta})}.
 \end{aligned}$$

also

Thus the components of acceleration are $\dot{r} - r\dot{\theta}^2$ along OP in the sense r increasing and $\frac{1}{r} \frac{d}{dt} (r^2\dot{\theta})$ perp. to OP in the sense θ increasing.

Cor. If the particle describes a circle of radius a , then $r = a$, $\dot{r} = 0$, so that $u = \dot{r} = 0$, $v = r\dot{\theta} = a\dot{\theta}$;

Radial acc. $= \dot{r} - r\dot{\theta}^2 = 0 - a\dot{\theta}^2 = -a\dot{\theta}^2$, i.e., $a\dot{\theta}^2$ towards the centre

and Transverse acc. $= r\ddot{\theta} + 2\dot{r}\dot{\theta} = a\ddot{\theta}$, i.e., tangentially.
Alternative method. We have

$$\begin{aligned}
 &\therefore x = r \cos \theta, y = r \sin \theta \\
 &\dot{x} = \dot{r} \cos \theta - r \sin \theta \dot{\theta} \\
 &\dot{y} = \dot{r} \sin \theta + r \cos \theta \dot{\theta} \\
 &\ddot{x} = \ddot{r} \cos \theta - 2\dot{r}\dot{\theta} \sin \theta - r\dot{\theta}^2 \cos \theta - r\ddot{\theta} \sin \theta \\
 &\ddot{y} = \ddot{r} \sin \theta + 2\dot{r}\dot{\theta} \cos \theta - r\dot{\theta}^2 \sin \theta + r\ddot{\theta} \cos \theta \quad \dots(1) \\
 &\text{Radial velocity} = \dot{x} \cos \theta + \dot{y} \sin \theta = \dot{r}, \quad \dots(2) \\
 &\text{Transverse velocity} = \dot{y} \cos \theta - \dot{x} \sin \theta = r\dot{\theta}, \quad \text{from (1)} \\
 &\text{Radial acceleration} = \ddot{x} \cos \theta + \ddot{y} \sin \theta = \ddot{r} - r\dot{\theta}^2, \quad \text{from (1)} \\
 &\text{Transverse acceleration} = \ddot{y} \cos \theta - \ddot{x} \sin \theta = 2\dot{r}\dot{\theta} + r\ddot{\theta}, \quad \text{from (2)} \\
 &= -\frac{1}{r} \cdot \frac{d}{dt} (r^2\dot{\theta}).
 \end{aligned}$$

Ex. 4. If the radial and transverse velocities of a particle are always proportional to each other, show that the path is an equiangular spiral.

Here $\frac{dr}{dt} = kr \frac{d\theta}{dt}$ where k is some constant

$$\text{or } \frac{dr}{r} = k d\theta.$$

Integrating we get, $\log r = k\theta + C$ where C is some constant

$r = ae^{k\theta}$ where a is also a constant.

This is an equiangular spiral.

Ex. 5. The velocities of a particle along and perpendicular to the radius from a fixed origin are λr and $\mu\theta$; find the path and show that the acceleration, along and perpendicular to the radius vector, are

$$\lambda^2 r - \frac{\mu^2 \theta^2}{r} \text{ and } \mu\theta \left(\lambda + \frac{\mu}{r} \right).$$

$$\text{Here } \frac{dr}{dt} = \lambda r \text{ and } r \frac{d\theta}{dt} = \mu\theta.$$

Dividing we get, $\frac{rd\theta}{dr} = \frac{\mu\theta}{\lambda r}$ or $\frac{\mu}{\lambda} \frac{dr}{r^2} = \frac{d\theta}{\theta}.$

Integrating we get

$$-\frac{\mu}{\lambda} \cdot \frac{1}{r} = \log \theta + C, \text{ where } C \text{ is a constant.}$$

This gives the path.

Radial acceleration $= \ddot{r} - r\dot{\theta}^2.$

$$= \lambda \dot{r} - r \cdot \frac{\mu^2 \theta^2}{r^2}, \quad \therefore \dot{r} = \lambda r \quad \therefore \ddot{r} = \lambda \dot{r} = \lambda^2 r$$

$$= \lambda^2 r - \frac{\mu^2 \theta^2}{r}, \quad \text{and } \dot{\theta} = \frac{\mu\theta}{r}.$$

Transverse acceleration

$$= 2\dot{r}\dot{\theta} + r\ddot{\theta} \\ = 2\lambda r \cdot \frac{\mu\theta}{r} + r\ddot{\theta}$$

By differentiating $r\dot{\theta} = \mu\theta$

we get $r\ddot{\theta} + \dot{r}\dot{\theta} = \mu\dot{\theta}$

i.e.,

$$r\ddot{\theta} = \mu \cdot \frac{\mu\theta}{r} - \lambda r, \quad \frac{\mu\theta}{r} = \frac{\mu\theta}{r} (\mu - \lambda r).$$

Hence transverse acc. $= 2\lambda\mu\theta + \frac{\mu\theta}{r} (\mu - \lambda r)$

$$= \mu\theta \left(2\lambda + \frac{\mu}{r} - \lambda \right) = \mu\theta \left(\lambda + \frac{\mu}{r} \right).$$

Ex. 6. A straight line of constant length moves with its ends on two fixed rectangular axes OX , OY and P is the foot of the perpendicular from O on the straight line. Show that the velocity of P perpendicular to OP is $OP \cdot \dot{\theta}$ and along OP is $2CP \cdot \dot{\theta}$ where C is the middle point of the line and θ is the angle COX .

Since AOB is a right angled triangle, C is the middle point of the hypotenuse, $OC = CA = CB = a$ if $AB = 2a$.



$$\angle COA = \theta = \angle CAO$$

$$\angle POX = 90 - \theta$$

$$\text{and } PO = OC \cos (90 - 2\theta) = a \sin 2\theta$$

\therefore Polar co-ordinates of P are
 $(a \sin 2\theta, 90 - \theta)$

Also $CP = a \cos 2\theta$

\therefore Velocity of P along OP

$$= \frac{d}{dt} (a \sin 2\theta) = 2a \cos 2\theta \dot{\theta}$$

$$= 2 \cdot CP \dot{\theta}.$$

$$\text{Velocity of } P \text{ perp. to } OP = a \sin 2\theta \frac{d}{dt} (90 - \theta)$$

$$= -OP \cdot \dot{\theta},$$

$$\text{i.e., velocity along } PA = OP \cdot \dot{\theta}.$$

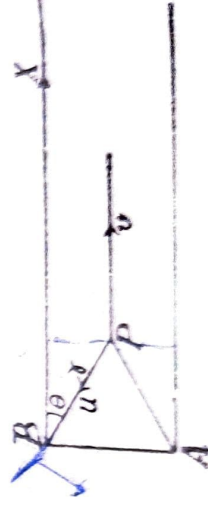
Ex. 7. A boat which is rowed with constant velocity U starts from a point A on the bank of a river which flows with a constant velocity V ; and it points always towards a point B on the other bank exactly opposite to A ; find the equation of the path of the boat. If $V = U$, show that the path is a parabola whose focus is B .

Let P be the position of the boat at time t . It has two velocities, U towards B and V downstream. Take the opposite bank as

initial line and B as pole and P the point (r, θ) , its components of velocity are \dot{r} along BP produced and $r\dot{\theta}$ perpendicular to BP .

Hence resolving along and perp. to BP , we get

$$\dot{r} = V \cos \theta - U$$



$$\text{and } r\dot{\theta} = -V \sin \theta$$

$$\text{Dividing we get } \frac{r\dot{\theta}}{dr} = -\frac{V \sin \theta}{V \cos \theta - U}$$

$$\text{or } \frac{dr}{r} = \left(-\cot \theta + \frac{U}{V} \operatorname{cosec} \theta \right) d\theta.$$

Integrating we get $\log r = \log \operatorname{cosec} \theta + \frac{U}{V} \log \tan \theta/2 + C_1$ where C_1 is constant.

This gives the path.

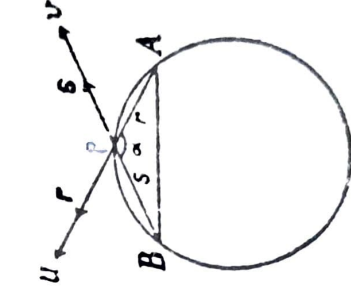
If $V = U$, it becomes $\log r = \log \operatorname{cosec} \theta + \log \tan \theta/2 + C_1$

$$\text{or } \log r = \log \left(\frac{1}{2 \cos^2 \theta/2} \right) + C_1$$

or $2r \cos^2 \theta/2 = A$ where A is also a constant

i.e., $\frac{A}{r} = 2 \cos^2 \theta/2 = 1 + \cos \theta$ which is a parabola, with B as focus.

Ex. 8. A and B are two fixed points on the circumference of a circle and the distances from A and B of any other point P on the circumference are r and s respectively. If u and v are the components of P 's velocity, as it moves round the circumference, along AP and BP respectively prove that α being the angle APB



$$u \sin^2 \alpha = \dot{r} - \dot{s} \cos \alpha,$$

$$v \sin^2 \alpha = \dot{s} - \dot{r} \cos \alpha$$

and deduce that $ur + vs = 0$

Resolving along AP ,

$$\dot{r} = u - v \cos (180 - \alpha) = u + v \cos \alpha$$

and resolving along BP

$$\dot{s} = v + u \cos (180 - \alpha) = v + u \cos \alpha$$

$$\therefore \dot{r} - \dot{s} \cos \alpha = u \sin^2 \alpha$$

$$\text{and } \dot{s} - \dot{r} \cos \alpha = v \sin^2 \alpha.$$

$$\text{Also } (ur + vs) \sin^2 \alpha = r\dot{r} - r\dot{s} \cos \alpha + s\dot{s} - s\dot{r} \cos \alpha$$

$$= \frac{1}{2} \frac{d}{dt} (r^2 + s^2) - \cos \alpha \frac{d}{dt} (rs)$$

$$= \frac{1}{2} \frac{d}{dt} (r^2 + s^2 - 2rs \cos \alpha)$$

$$= \frac{1}{2} \frac{d}{dt} AB = 0, \text{ hence}$$