第十二章 电磁感应 电磁场

§ 12.1 法拉第电磁感应定律

一、电磁感应现象

闭合导体回路所包围的面积内的磁(感应)通量发生变化时,回路中就会产生感应电流

二、法拉第电磁感应定律

$$\Phi_m$$
变化 — 感应电动势 $\mathbf{\epsilon}_{\vec{B}}$ — $I_{\vec{B}}$ $\mathbf{\epsilon}_i = -\frac{d\Phi}{dt}$

说明:

1)
$$\varepsilon_i$$
 的数值为 $\left| \frac{d\Phi}{dt} \right|$

N匝线圈串联
$$\Psi = N\Phi$$
 磁通链数 $\longrightarrow \varepsilon_i = -N \frac{d\Phi}{dt} = -\frac{d\Psi}{dt}$ 每匝通量 Φ

2) 负号反映 ϵ_i 的方向 负号表示感应电动势总是反抗磁通的变化

楞次定律: ——判断感应电流方向的法则 (P150)

- *闭合回路中感应电流产生的通量总是阻碍引起感应电流的通量的变化。
- *感应电流的效果,总是反抗引起感应电流的原因。

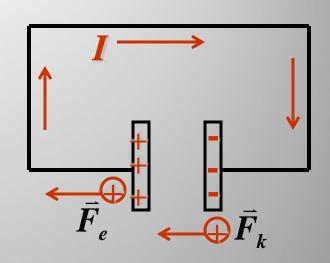
§ 12.2 动生电动势和感生电动势

一、动生电动势

1. 电动势概念:

静电力
$$\bar{F}_e$$
: 使+ q 从高电位 \rightarrow 低电位 $\bar{F}_e = \bar{E}q$

非静电力 \vec{F}_k : 使+q从低电位 \rightarrow 高电位 $\vec{F}_k = \vec{E}_k q$



电源: 提供非静电力的装置

$$+q$$
绕闭合回路一周: $A = \int_L (\vec{F}_k + \vec{F}_e) \cdot d\vec{l} = \int_L q(\vec{E}_k + \vec{E}) \cdot d\vec{l} = \int_L q\vec{E}_k \cdot d\vec{l}$

电源电动势:单位正电荷绕闭合回路一周,

非静电力所做的功

$$\varepsilon = \frac{A}{q} = \oint_L \vec{E}_k \cdot d\vec{l} = \int_{-}^{+} \vec{E}_k \cdot d\vec{l}$$

2. 动生电动势的产生及意义:

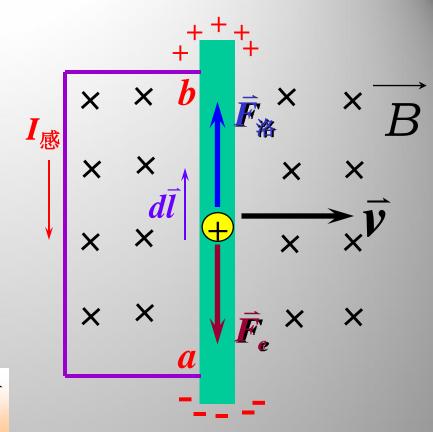
当 $\vec{F}_e = \vec{F}_{\text{A}} \rightarrow$ 形成稳定的电动势 ε_i

产生动生电动势的非静电力:

$$\vec{F}_k = \vec{F}_{\mbox{\tiny AB}} = q \left(\vec{v} \times \vec{B} \right)$$

非静电场强: $\bar{E}_k = \frac{\bar{F}_k}{q} = \bar{v} \times \bar{B}$

定义:
$$\varepsilon_{\bar{z}j} = \int_a^b \vec{E}_k \cdot d\vec{l} = \int_a^b (\vec{v} \times \vec{B}) \cdot d\vec{l}$$



例1: 细导体棒 OA 绕 O 点以 ω 转动,

求: OA 垂直长直电流 I 时刻的感应电动势。

解:
$$\varepsilon_{\bar{z}_{\bar{z}_{\bar{z}_{\bar{z}}}}} = \int (\bar{v} \times \bar{B}) \cdot d\bar{l}$$

 $d\vec{l}$ 指向: $O \rightarrow A$, $d\vec{l} = d\bar{x}$

$$\varepsilon_{OA} = \int_{0}^{A} (\vec{v} \times \vec{B}) \cdot d\vec{x}$$
$$= \int_{0}^{A} vB \cos \pi dx = -\int_{a}^{2a} vB dx$$

$$=-\int_{a}^{2a}\omega(x-a)\frac{\mu_{0}I}{2\pi x}dx$$

$$=-\frac{\mu_o I \omega a}{2\pi} \left(1-\ln 2\right) < 0$$

$$\varepsilon_{OA} < 0$$
,表明 ε_{ob} 从 $A \to O$

通常取dī指向与v×B相同

3. 动生电动势的计算:

(1). 由定义:
$$\varepsilon_{\bar{z}\bar{d}} = \int_a^b (\bar{v} \times \bar{B}) \cdot d\bar{l}$$

(2). 法拉弟定律: $\begin{cases} \text{补回路} \rightarrow \varepsilon = \left| \frac{d\Phi}{dt} \right| \\ \text{楞次定律判定方向} \end{cases}$

(感应电流方向)

例2: 长为 L 的导线在均匀磁场 B 中绕o点以 ω 匀速转动

求: 导线中的动生电动势

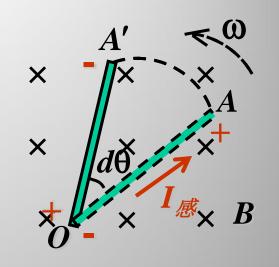
t时刻: OA

 $t \mapsto A$ t + dt 时刻: OA' 组成回路: OAA'O

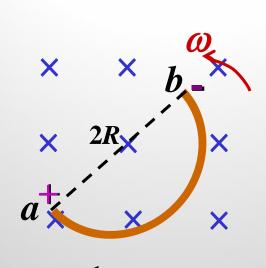
$$d\Phi = BdS = B \cdot \frac{1}{2} L \cdot Ld\theta = \frac{1}{2} BL^2 d\theta$$

$$\varepsilon = \left| \frac{d\Phi}{dt} \right| = \frac{1}{2} B L^2 \frac{d\theta}{dt} = \frac{1}{2} B L^2 \omega$$

OA'导线内: ε_{d} 从 $A' \to O$

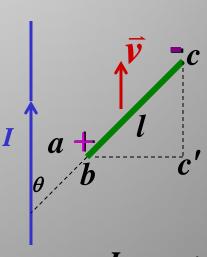


讨论:
$$\begin{array}{cccc}
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 $\varepsilon_{ba} = \frac{1}{2}B\omega(2R)^2$

$$I \uparrow a + \frac{\vec{v}}{l}$$



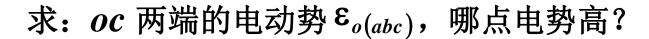
$$= \frac{\mu_0 I v}{2\pi} \ln \frac{a+l}{a} \qquad \varepsilon = \frac{\mu_0 I v}{2\pi} \ln \frac{a+2I}{a}$$

$$\varepsilon_{cb} = \varepsilon_{c'b} = \frac{\mu_0 I v}{2\pi} \ln \frac{a + l \sin \theta}{a}$$

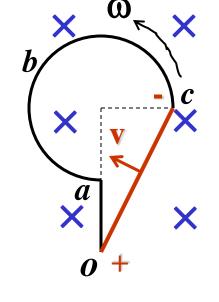
例3: (自测P22)

导线
$$oa = R$$
, (abc) 为 $\frac{3}{4}$ 圆弧,

在匀强磁场B内以ω绕o转动时,



$$\varepsilon_{o(abc)} + \varepsilon_{\overline{co}} = \varepsilon_i = \frac{d\Phi}{dt} = 0$$



$$\mathbf{\varepsilon}_{o(abc)} = -\mathbf{\varepsilon}_{\overline{co}} = \mathbf{\varepsilon}_{\overline{oc}} = \frac{1}{2} B \omega (\overline{oc})^2 = \frac{1}{2} B \omega (\sqrt{5}R)^2 = \frac{5}{2} B \omega R^2$$

O点电势最高

例:长直载流导线 I 旁,一矩形线圈以v 匀速向右移动

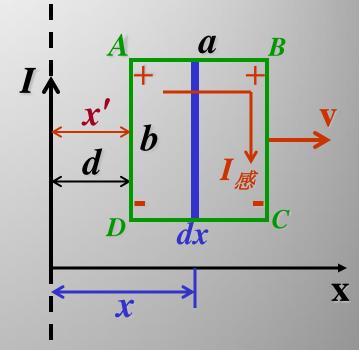
求:到图示位置时刻,线圈中的 ϵ_{d}

解: 1)
$$\epsilon_{\text{动}}$$
 方向: $\bar{v} \times \bar{B}$

$$\varepsilon_{DA} = \int_{D}^{A} \left(\vec{v} \times \vec{B} \right) \cdot d\vec{l} = \frac{\mu_{0}I}{2\pi d} v \int_{D}^{A} dl = \frac{\mu_{0}I}{2\pi d} vb$$

$$\varepsilon_{CB} = \frac{\mu_0 I}{2\pi (d+a)} vb$$

$$\therefore \varepsilon = \varepsilon_{DA} - \varepsilon_{CB} = \frac{\mu_0 I v b}{2\pi} \left(\frac{1}{d} - \frac{1}{d+a} \right)$$



2)
$$\varepsilon = \left| \frac{d\Phi}{dt} \right| \quad \Phi = \int \vec{B} \cdot d\vec{S} = \int_{x'}^{x'+a} \frac{\mu_0 I}{2\pi x} b dx = \frac{\mu_0 I b}{2\pi} \left[\ln(x'+a) - \ln x' \right]$$

$$\varepsilon = \left| \frac{d\Phi}{dt} \right| = \left| \frac{d\Phi}{dx'} \cdot \frac{dx'}{dt} \right|_{\frac{dx'}{dt} = v, x' = d} \qquad \therefore \varepsilon = \frac{\mu_0 I v b}{2\pi} \left(\frac{1}{d} - \frac{1}{d+a} \right)$$

二、感生电场(涡旋电场) 感生电动势

- 1. 涡旋电场 (感生电场) 麦克斯韦假设:
- 1) 变化的磁场会在空间激发出一个源于变化磁场的电场 →→ 点生电场
- 2) 感生电场的电力线是闭合的。

$$\begin{cases} \oint ar{E}_{ec{m{B}}} \cdot dar{l}
eq 0 \longrightarrow 涡旋电场/非保守力场 \ \oint ar{E}_{ec{m{B}}} \cdot dar{S} = 0 \longrightarrow 无源场 \end{cases}$$

3) 产生感生电动势的非静电力就是感生电场力

$$\int_{L} \bar{E}_{\bar{s}} \cdot d\bar{l} = -\int_{S} \frac{\partial \bar{B}}{\partial t} \cdot d\bar{S} \quad S: 是静止导体回路所包围的面积$$

- •空间某处 B 随时间变化,周围就会激发出感生电场,其电力线是围绕 \bar{B} 线的同心圆,方向由楞次定律判断
- \bullet 空间电场 $\left\{\begin{array}{l}$ 由静止电荷按库仑定律激发: $ar{E}_{\sharp}$ 由变化磁场激发: $ar{E}_{\aleph}$
- •空间总场强: $\bar{E} = \bar{E}_{\text{#}} + \bar{E}_{\text{K}}$

场方程:
$$\oint_{L} \vec{E} \cdot d\vec{l} = \oint_{L} (\vec{E}_{\hat{\mathbb{B}}} + \vec{E}_{\hat{\mathbb{H}}}) \cdot d\vec{l} = \oint_{L} \vec{E}_{\hat{\mathbb{H}}} \cdot d\vec{l} = -\int \frac{\partial \vec{B}}{\partial t} \cdot d\vec{S}$$

$$\oint_{S} \vec{E} \cdot d\vec{S} = \oint_{S} (\vec{E}_{\hat{\mathbb{B}}} + \vec{E}_{\hat{\mathbb{H}}}) \cdot d\vec{S} = \oint_{S} \vec{E}_{\hat{\mathbb{B}}} \cdot d\vec{S} = \frac{\sum q_{i}}{\epsilon_{0}}$$

2. 感生电场和感生电动势的计算:

例5: 无限长圆柱形空间,均匀磁场 B 随 t 变化, $\frac{dB}{dt} = \mathbb{R} > 0$

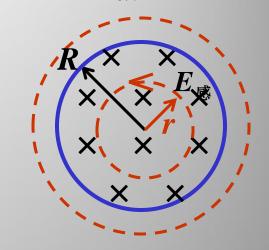
求: E_感

解: r<R:

$$\oint_{L} \vec{E} \cdot d\vec{l} = E \cdot 2\pi r$$

$$- \int_{S_{1}} \frac{\partial \vec{B}}{\partial t} \cdot d\vec{S} = -\frac{dB}{dt} \int_{S_{1}} dS = -\frac{dB}{dt} \cdot \pi r^{2}$$

$$E_{\vec{B}} = -\frac{r}{2} \cdot \frac{dB}{dt}$$



r>R:

$$\begin{cases}
\int_{L} \vec{E} \cdot d\vec{l} = E \cdot 2\pi r \\
-\int_{S_{1}} \frac{\partial \vec{B}}{\partial t} \cdot d\vec{S} = -\frac{dB}{dt} \int_{S_{1}} dS = -\frac{dB}{dt} \cdot \pi R^{2}
\end{cases} E_{\vec{B}} = -\frac{R^{2}}{2r} \cdot \frac{dB}{dt}$$

例6: 圆柱形均匀磁场 $\frac{dB}{dt} = C > 0$, 求 ε_{ab} 、 ε_{cd}

$$E_{\vec{\mathbb{R}}} = \begin{cases} -\frac{r}{2} \frac{dB}{dt} \leftarrow r < R \\ -\frac{R^2}{2r} \frac{dB}{dt} \leftarrow r > R \end{cases}$$

$$cd: \ \vec{E}_{\mathbb{R}} \perp d\vec{l} \rightarrow \varepsilon_{cd} = \int_{c}^{d} \vec{E}_{\mathbb{R}} \cdot d\vec{l} = 0$$

$$ab: \quad \varepsilon_{ab} = \int_a^b \vec{E}_{\mathbb{R}} \cdot d\vec{l} = \int_a^b E \cos\theta dl = \int_a^b \frac{r}{2} \frac{dB}{dt} \cos\theta dl$$

$$= \int_a^b \frac{h}{2} \frac{dB}{dt} \cdot dl = \frac{hab}{2} \cdot \frac{dB}{dt} > 0 \quad \text{ } \hat{\mathcal{T}} = 0 \quad \hat{\mathcal{T}} = 0 \quad \hat{\mathcal{T}} = 0$$

計 回 路
$$oabo$$
: $\varepsilon_i = \varepsilon_{oa} + \varepsilon_{ab} + \varepsilon_{bo} = \varepsilon_{ab}$

$$= -\frac{d\phi}{dt} = -\frac{dB}{dt} \cdot S_{\Delta oab} = -\frac{dB}{dt} \cdot \frac{h\overline{ab}}{2} < 0 \quad \mathcal{E}_{ab}$$

*感生电动势的计算:

1) 由定义:
$$\varepsilon_i = \int_a^b \bar{E} \cdot d\bar{l}$$
 ——已知 $\bar{E}_{\bar{B}}$ 的分布 ε_i : $a \to b$

$$\varepsilon_i$$
: $a \to b$

空间分布均匀的
$$E_{\bar{\mathbb{S}}} = \begin{cases} -\frac{r}{2} \frac{dB}{dt} \leftarrow r < R \\ -\frac{R^2}{2r} \frac{dB}{dt} \leftarrow r > R \end{cases}$$

感生电场的电力线是围绕 B 线的闭合曲线 感生电场产生的磁场一定反抗原来磁场的变化

2) 由法拉第电磁感应定律: $\varepsilon_i = -\frac{d\Phi}{dt}$

闭合回路:
$$\left| \frac{d\Phi}{dt} \right| = \left| \int \frac{d\vec{B}}{dt} \cdot d\vec{S} \right| = S \left| \frac{dB}{dt} \right|$$
 均匀磁场

$$\left\{ \begin{array}{ll} |\partial \Phi| = \left| \int \frac{d\vec{B}}{dt} \cdot d\vec{S} \right| = S \left| \frac{dB}{dt} \right| & \text{均匀磁场} \\ |\partial \Phi| = |\partial \Phi| + |\partial \Phi| & \text{0} \\ |\partial \Phi| = |\partial \Phi| & \text{0} \\ |\partial \Phi| = |\partial \Phi| & \text{0} \\ |\partial \Phi| & \text{0}$$

讨论: (1) 比较
$$\varepsilon_{ac}$$
与 ε_{cb}

$$egin{aligned} arepsilon_{ac} &= arepsilon_{acoa} = -rac{d\phi}{dt} = -S_{acoa} \cdot rac{dB}{dt}, \ \ \dot{\mathcal{T}}$$
 问: $a
ightarrow c \ &arepsilon_{cb} = arepsilon_{cboc} = -rac{d\phi}{dt} = -S_{cboc} \cdot rac{dB}{dt}, \ \ \dot{\mathcal{T}}$ 问: $c
ightarrow b \ &S_{acoa} = S_{cboc}
ightarrow arepsilon_{ac} = arepsilon_{cb} \end{aligned}$

(2) 比较
$$\varepsilon_{ef}$$
与 ε_{gh}

$$egin{align*} arepsilon_{eff} & \mathcal{F} \mathcal{E}_{efo} = -rac{d\phi}{dt} = -ec{S}_{efoe} \cdot rac{dec{B}}{dt} = S_o rac{dB}{dt} < 0$$
,方向: $f o e$ $\varepsilon_{gh} = \varepsilon_{ghog} = S_o rac{dB}{dt} < 0$,方向: $h o g$

*比较。柱内导线比较其两端与圆心连线所围面积的大小柱外导线比较其两端与圆心连线的张角的大小

例:导线 ab 长 l 在轨道上以速度 v 向右平移。轨道处于均匀磁场 B 中,B=kt (k>0) 并与轨道面法线成 60° ,t=0 时,ab 在 cd 处。求任意时刻回路中感应电动势

解:
$$\varepsilon_{\bar{\bowtie}} = \int (\bar{v} \times \bar{B}) \cdot d\bar{l}$$

$$= \int_{b}^{a} vB \sin 30^{\circ} dl = \frac{vBl}{2} = \frac{vlkt}{2}$$

$$\varepsilon_{\bar{\otimes}} = \frac{d\phi}{dt} = \int \frac{d\bar{B}}{dt} \cdot d\bar{S} = \int \frac{dB}{dt} \cos 60^{\circ} dS = \frac{1}{2} k \cdot vt \cdot l$$

$$\varepsilon = \varepsilon_{\bar{\bowtie}} + \varepsilon_{\bar{\otimes}} = klvt$$

$$\dot{\beta} = \frac{1}{2} k \cdot vt \cdot l$$

$$\dot{\beta} = \frac{1}{2} k \cdot vt \cdot l$$

$$\varepsilon = \left| \frac{d\Phi}{dt} \right| = \frac{Bl}{2} \cdot \frac{dx}{dt} + \frac{lx}{2} \cdot \frac{dB}{dt} = \frac{Blv}{2} + \frac{lxk}{2} = ktlv$$

*ab距cd为x处: $\Phi = \int \bar{B} \cdot d\bar{S} = \int_0^x B \cos 60^\circ l dx = \frac{B l x}{2}$

§ 12.3 自感和互感

- 一、自感 ((P167~172)
 - 1. 自感现象:

2. 自感系数、自感电动势

$$B \propto I \rightarrow \Phi_m \propto I \rightarrow \Psi = N\Phi_m \propto I \rightarrow \Psi = LI$$

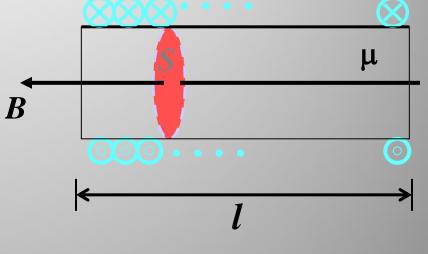
自感电动势:
$$\varepsilon_L = -\frac{d\Psi}{dt} = -L\frac{dI}{dt}$$
 ε_L 总是反抗电路中 I 的变化, 其能力与 L 有关

例 1: 长直螺旋管 l,总匝数 N,充介质 μ 。 求它的自感系数。

解: 假设螺旋管中通电流 I

$$B = \mu nI = \mu \frac{N}{l}I \rightarrow \Phi = BS = \mu \frac{N}{l}IS$$

$$L = \frac{N\Phi}{l} = \frac{\mu N^2}{l}S = \mu \frac{N^2}{l^2}Sl = \mu n^2V$$



L 的计算: 假设 $I \to B \to \Phi$, $\Psi = N\Phi \to L = \frac{\Psi}{I}$

例2: 求一环形螺线管N、h、 R_1 、 R_2 的自感系数。

解:
$$\int_{L} \vec{H} \cdot d\vec{l} = \sum I$$

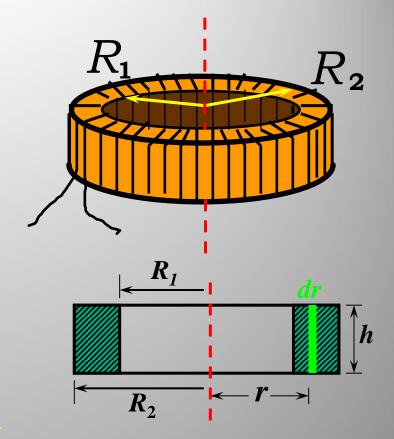
$$H \cdot 2\pi r = NI$$

$$H = \frac{NI}{2\pi r} \to B = \frac{\mu NI}{2\pi r}$$

$$\mathbf{\Phi} = \int_{S} \vec{B} \cdot d\vec{S}$$

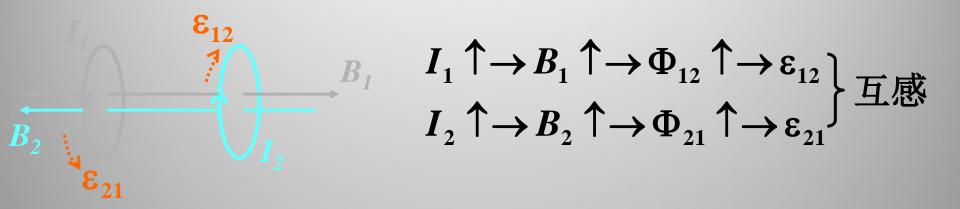
$$= \int_{R_1}^{R_2} \frac{\mu NI}{2\pi r} \cdot h dr = \frac{\mu NIh}{2\pi r} \ln \frac{R_2}{R_1}$$

$$L = \frac{\Psi}{I} = \frac{N\Phi}{I} = \frac{\mu N^2 h}{2\pi r} \ln \frac{R_2}{R_1}$$



二、互感 (P172~176)

1. 互感现象: 两个载流回路相互激发感应电动势的现象



2. 互感系数、互感电动势:

$$\begin{cases} \Phi_{12} \propto I_1 \to \Psi_{12} = N_2 \Phi_{12} \to \Psi_{12} = M_{12} I_1 \\ \Phi_{21} \propto I_2 \to \Psi_{21} = N_1 \Phi_{21} \to \Psi_{21} = M_{21} I_2 \end{cases}$$

互感系数:
$$M_{12} = \frac{\Psi_{12}}{I_1}$$
 $M_{12} = M_{21} = M$ 反映两个线圈相互感应的能力

取决于两个线圈自身的形状、大小、 相对位置及周围介质情况

互感电动势:
$$\varepsilon_{12} = -\frac{d\Psi_{12}}{dt} = -M_{12}\frac{dI_1}{dt} = -M\frac{dI_1}{dt}$$

$$\varepsilon_{21} = -\frac{d\Psi_{21}}{dt} = -M_{21}\frac{dI_2}{dt} = -M\frac{dI_2}{dt}$$

3. 互感系数、互感电动势的计算:

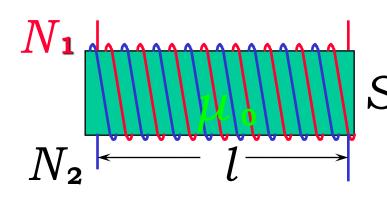
假定回路1通电流
$$I_1 \to B_1 \to \Psi_{12} \to M = \frac{\Psi_{12}}{I_1} \to \varepsilon_{12} = -M \frac{dI_1}{dt}$$
 假定回路2通电流 $I_2 \to B_2 \to \Psi_{21} \to M = \frac{\Psi_{21}}{I_2} \to \varepsilon_{21} = -M \frac{dI_2}{dt}$

例3:有两个直长螺线管,它们绕在同一个圆柱面上

$$(\mu_o, N_1, N_2, S, l)$$

求: 互感系数

$$I_1 \to B_1 = \mu n I_1 = \mu \frac{N_1}{I} I_1$$



$$\Psi_1 = N_1 \Phi_1 = N_1 B_1 S \rightarrow L_1 = \frac{\Psi_1}{I_1} = \frac{\mu N_1^2}{l} S$$

同理
$$\rightarrow L_2 = \frac{\Psi_2}{I_2} = \frac{\mu N_2^2}{l} S$$

$$\Psi_{12} = N_2 \Phi_{12} = N_2 B_1 S \rightarrow M_{12} = \frac{\Psi_{12}}{I_1} = \frac{\mu N_1 N_2}{l} S = M_{21} = \frac{\Psi_{21}}{I_2}$$

此时线圈1的磁通全部通过线圈2——无磁漏

在一般情况下: $M = K \sqrt{L_1 L_2}$

K称为耦合系数 $0 \leq K \leq 1$

耦合系数的大小反映了两个回路磁场耦合松紧的程度。一般情况下都有漏磁,耦合系数小于1。

例:将 L_I 、 L_2 、M两线圈串联。顺接:两线圈的磁通互相加强;反接:磁通相互削弱。计算两种接法下线圈的等效总自感。

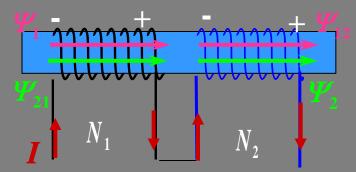
1) 顺接:

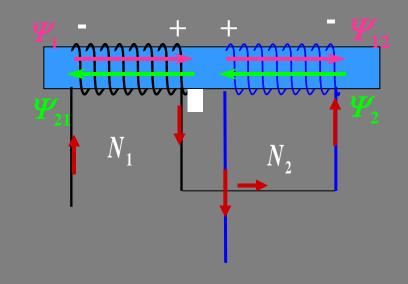
$$\Psi = \Psi_1 + \Psi_{21} + \Psi_2 + \Psi_{12}
= L_1 I + M I + L_2 I + M I
L = \frac{\Psi}{I} = L_1 + L_2 + 2M$$

2) 反接:

$$\mathcal{Y} = \mathcal{Y}_1 - \mathcal{Y}_{21} + \mathcal{Y}_2 - \mathcal{Y}_{12}$$
$$= L_1 I - MI + L_2 I - MI$$

$$L = \frac{\Psi}{I} = L_1 + L_2 - 2M$$







§ 12.4 磁场的能量

自感线圈:建立磁场过程 储能过程 磁场能量?

一、自感线圈的储能(自感磁能)(P176):

$$\begin{aligned}
\varepsilon + \varepsilon_L &= IR \to \varepsilon = -\varepsilon_L + IR \\
dA &= Pdt = \varepsilon Idt = -\varepsilon_L Idt + I^2 Rdt
\end{aligned}$$

$$\int_0^{t_0} \varepsilon Idt = \int_0^{t_0} -\varepsilon_L Idt + \int_0^{t_0} I^2 Rdt$$

电源做功

克服 ε_L做功 电阻释放 焦耳热

$$W_{m} = \int_{0}^{t_{0}} -\varepsilon_{L} I dt = \int_{0}^{t_{0}} L \frac{dI}{dt} I dt = \int_{0}^{I_{0}} L I dI = \frac{1}{2} L I_{0}^{2}$$

自感线圈的能量:线圈储藏的磁场能量

或电流稳定前,电源克服 ε_L 所做的功

二、磁场能量:

长直螺旋管:
$$L = \mu n^2 V$$

 $B = \mu nI \rightarrow I = \frac{B}{\mu n}$ $W_m = \frac{1}{2}LI^2 = \frac{1}{2}\frac{B^2}{\mu}V$

磁能密度:
$$W_m = \frac{W_m}{V} = \frac{1}{2} \frac{B^2}{\mu} = \frac{1}{2} \mu H^2 = \frac{1}{2} BH$$

磁场能量:
$$W_m = \int_V w_m dV = \int_V \frac{1}{2} BH dV = \frac{1}{2} LI^2$$

例1: 长a,宽b,间隔为d的平行金属板 (d << a, b)

求: (1) 自感系数 L, (2) 通电流 I后,板间的 W_m

解:
$$B = \mu_0 i = \mu_0 \frac{I}{h}$$

$$W_{m} = \frac{1}{2} \frac{B^{2}}{\mu_{0}} V = \frac{1}{2} \frac{B^{2}}{\mu_{0}} abd = \frac{\mu_{0} I^{2} ad}{2b}$$

$$L = \frac{2W_m}{I^2} = \frac{\mu_0 ad}{h} \quad \leftarrow L = \frac{\Phi}{I} \leftarrow \Phi = BS$$

例2: 同轴电缆,内充介质µ,通

电流I。求单位长度L及 W_m 。

解: 1) $I \rightarrow B \rightarrow$ 单位长度 Φ_m

$$L = \frac{\Phi_m}{I} \to W_m = \frac{1}{2}LI^2$$

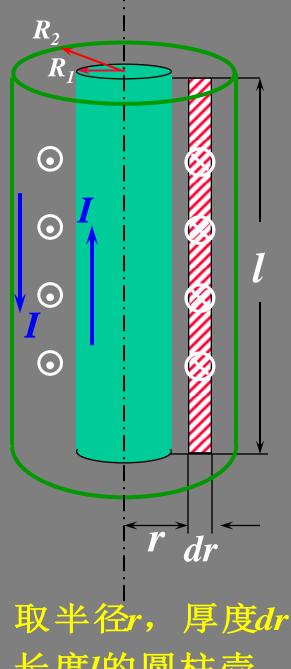
2) 先求 W_m

$$\oint \vec{H} \cdot d\vec{l} = I \to H = \frac{I}{2\pi r}$$

$$B = \mu H = \frac{\mu I}{2\pi r}$$
 $(R_1 < r < R_2)$

$$W_{m} = \int \frac{1}{2}BHdV = \int_{R_{1}}^{R_{2}} \frac{1}{2} \frac{B^{2}(r)}{l} \cdot 2\pi r \cdot l \cdot dr \leftarrow \Re l = 1$$

$$W_m = \frac{1}{2}LI^2 \to L$$



长度I的圆柱壳