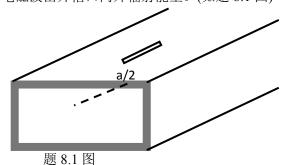
第八章习题及解答

8.1 为什么一般矩形波导测量线的纵槽开在波导的中线上?

解:因为矩形波导中的主模为 TE_{10} 模,而由 TE_{10} 的管壁电流分布可知,在波导宽边中线处只有纵向电流。因此沿波导宽边的中线开槽不会因切断管壁电流而影响波导内的场分布,也不会引起波导内电磁波由开槽口向外辐射能量。(如题 8.1 图)



8.2 下列二矩形波导具有相同的工作波长,试比较它们工作在 TM_{11} 模式的截止频率。

- (1) $a \times b = 23 \times 10 mm^2$;
- (2) $a \times b = 16.5 \times 16.5 mm^2$

解: 截止频率
$$f_c = \frac{1}{2\pi\sqrt{\mu\varepsilon}}\sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2}$$
 当介质为空气 $\sqrt{\mu\varepsilon} = \sqrt{\mu_0\varepsilon_0} = \frac{1}{c}$

(1) 当 $a \times b = 23mm \times 10mm$, 工作模式为 TM_{11} (m=1,n=1), 其截止频率为

$$f_c = \frac{3 \times 10^{11}}{2} \sqrt{\left(\frac{1}{23}\right)^2 + \left(\frac{1}{10}\right)^2} = 16.36 \text{ (GHz)}$$

(2)当 $a \times b = 16.5mm \times 16.5mm$,工作模式仍为 TM_{11} (m=1,n=1),其截止频率为的

$$f_c = \frac{3 \times 10^{11}}{2} \sqrt{\left(\frac{1}{16.5}\right)^2 + \left(\frac{1}{16.5}\right)^2} = 12.86 \text{ (GHz)}$$

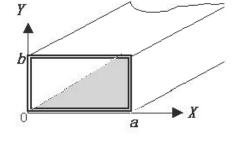
由以上的计算可知:截止频率与波导的尺寸、传输模式及波导填充的介质有关,与工作频率无关。

8.3 推导矩形波导中TE_{mn}模的场分布式。

解: 对于 TE 波有
$$E_z = 0, H_z \neq 0$$

H_z应满足下面的波动方程和边界条件:

$$\begin{cases} \nabla^2 H_z + k^2 H_z = 0 \\ E_y \big|_{x=0} = 0 \\ E_y \big|_{x=a} = 0 \\ E_x \big|_{y=0} = 0 \\ E_x \big|_{y=b} = 0 \end{cases}$$
 (1)



由均匀导波系统的假设,

$$H_z(x,y,z) = H_z(x,y)e^{-\Gamma z}$$

将其代入式(1),得

$$\frac{\partial H_z^2}{\partial x^2} + \frac{\partial H_z^2}{\partial y^2} + \frac{\partial H_z^2}{\partial z} + k^2 H_z = 0$$

$$\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \left(\Gamma^2 + k^2 \right) \right] H_z(x, y) = 0$$

$$\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + h^2 \right] H_z(x, y) = 0$$
(2)

其中 $h^2 = \Gamma^2 + k^2$

该方程可利用分离变量法求解。设其解为:

$$H_z(x,y) = f(x)g(y) \tag{3}$$

将式(.3) 代入式 (2), 然后等式两边同除以f(x)g(y), 得

$$-\frac{1}{f(x)}\frac{d^{2}f(x)}{dx^{2}} = \frac{1}{g(y)}\frac{d^{2}g(y)}{dy^{2}} + h^{2}$$

上式中等式左边仅为x的涵数,等式右边仅为y的函数,要使其相等,必须各等于常数。于是,该式可分离出两个常微分方程

$$\frac{d^2f(x)}{dx^2} + k_x^2 f(x) = 0 \tag{4a}$$

$$\frac{d^{2}g(y)}{dy^{2}} + k_{y}^{2}g(y) = 0$$
 (4b)

$$k_x^2 + k_y^2 = h^2 (5)$$

式(4a)的通解为

$$f(x) = A\sin k_x x + B\cos k_x x \tag{6}$$

由于在 x=0 和 x=a 的边界上,满足

$$E_{y}\big|_{x=0}=0 \qquad \qquad E_{y}\big|_{x=a}=0$$

由纵向场与横向场的关系,得 $E_y = \frac{j\omega\mu}{k^2} \frac{\partial H_z}{\partial x}$

则在 x=0 和 x=a 的边界上, $H_z(x,y)$ 满足

$$\frac{\partial H_z}{\partial x}\Big|_{x=0} = 0 \qquad \frac{\partial H_z}{\partial x}\Big|_{x=a} = 0$$

于是将其代入式 (6)得

$$A = 0$$

 $k_x = \frac{m\pi}{a} (m = 0, 1, 2, 3.....)$

所以

$$f(x) = B\cos\frac{m\pi}{a}x$$

同理得式(4)的通解

$$g(y) = C\sin k_y y + D\cos k_y y$$

满足的边界条件为

$$\frac{\partial H_z}{\partial y}\Big|_{y=0} = 0 \qquad \qquad \frac{\partial H_z}{\partial y}\Big|_{y=b} = 0$$

于是得

$$C = 0$$

$$k_{y} = \frac{n\pi}{b} (n = 0, 1, 2, 3, \dots)$$
$$g(y) = D\cos\frac{n\pi}{b} y$$

所以,得到矩形波导中 TE 波的纵向场分量

$$H_z(x,y) = H_0 \cos\left(\frac{m\pi}{a}x\right) \cos\left(\frac{n\pi}{b}y\right)$$

式中 H⊶CD 由激励源强度决定

本征值由式
$$h^2 = k_x^2 + k_y^2 = \left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2$$

利用纵向场与横向场的关系式可求得 TE 的其他横向场分量

$$E_{x}(x,y) = \frac{j\omega\mu}{h^{2}} \left(\frac{n\pi}{b}\right) H_{0} \cos\left(\frac{m\pi}{a}x\right) \sin\left(\frac{n\pi}{b}y\right)$$

$$E_{y}(x,y) = -\frac{j\omega\mu}{h^{2}} \left(\frac{m\pi}{a}\right) H_{0} \sin\left(\frac{m\pi}{a}x\right) \cos\left(\frac{n\pi}{b}y\right)$$

$$H_{x}(x,y) = \frac{jk_{z}}{h^{2}} \left(\frac{m\pi}{a}\right) H_{0} \sin\left(\frac{m\pi}{a}x\right) \cos\left(\frac{n\pi}{b}y\right)$$

$$H_{y}(x,y) = \frac{jk_{z}}{h^{2}} \left(\frac{n\pi}{b}\right) H_{0} \cos\left(\frac{m\pi}{a}x\right) \sin\left(\frac{n\pi}{b}y\right)$$

8.4 设矩形波导中传输 TE_{10} 模,求填充介质(介电常数为 ε)时的截止频率及波导波长。

解:截止频率
$$f_c = \frac{1}{2\pi\sqrt{\mu\varepsilon}}\sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2}$$
 对于 TE_{10} $(m=1,n=0)$, 得
$$f_c = \frac{1}{2\pi\sqrt{\mu\varepsilon}}\sqrt{\left(\frac{\pi}{a}\right)^2} = \frac{1}{2a\sqrt{\mu\varepsilon}}$$
 波导波长
$$\lambda_g = \frac{2\pi}{\beta} = \frac{2\pi}{\omega\sqrt{\mu\varepsilon}}\sqrt{1 - \frac{f_c^2}{f^2}} = \frac{\lambda}{\sqrt{1 - \frac{f_c^2}{f^2}}}$$
 式中 $\lambda = \frac{2\pi}{\omega\sqrt{\varepsilon\mu}}$ 为无界空间介质中的

8.5 已知矩形波导的横截面尺寸为 $a \times b = 23 \times 10 mm^2$,试求当工作波长 $\lambda = 10 mm$ 时,波导中能传输哪些波型? $\lambda = 30 mm$ 时呢?

解:波导中能传输的模式应满足条件

$$\lambda < (\lambda_c)_{mn}$$
 (工作波长小于截止波长)
或 $f > (f_c)_{mn}$ (工作頻率大于截止頻率)
在矩形波导中截止波长为
$$\lambda_c = \frac{2\pi}{\sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2}}$$

曲传输条件
$$\lambda < \frac{2}{\sqrt{\left(\frac{m}{23}\right)^2 + \left(\frac{n}{10}\right)^2}}$$

当
$$\lambda$$
=10 mm 时上式可写为 $n<10$ $\left[\left(\frac{2}{10}\right)^2-\left(\frac{m}{23}\right)^2\right]^{\frac{1}{2}}$

能满足传输条件的 m 和 n 为

- (1) m=0,n<2 有以下波型 TE₀₁
- (2) m=1,n<1.95 有以下波型 TE₁₀,TE₁₁,TM₁₁
- (3) m=2,n<1.8 有以下波型 TE₂₀,TE₂₁,TM₂₁
- (4) m=3,n<1.5 有以下波型 TE₃₀,TE₃₁,TM₃₁
- (5) m=4,n<0.95 有以下波型 TE₄₀

当
$$\lambda$$
=30 mm 时,应满足 $n<10$ $\left[\left(\frac{2}{30}\right)^2-\left(\frac{m}{23}\right)^2\right]^{\frac{1}{2}}$

- (1) m=0,n<0.66 (无波型存在)
- (2) m=1,n<0.5有以下波型 TE₁₀
- (3) m=2, 不满足条件。

故此时只能传输TE₁₀模

8.6 一矩形波导的横截面尺寸为 $a \times b = 23 \times 10 mm^2$ 由紫铜制作,传输电磁波的频率为 f = 10 GHz。试计算:

- (1) 当波导内为空气填充,且传输TE₁₀波时,每米衰件多少分贝?
- (2) 当波导内填充以 $\varepsilon_{\rm r}=2.54$ 的介质,仍传输 ${\rm TE}_{10}$ 波时,每米衰件多少分贝?

解: 当波导内为空气填充时,其工作波长为
$$\lambda = \frac{v}{f} = \frac{3 \times 10^8}{10 \times 10^9} = 3 \times 10^{-2} = 3 (cm)$$

当波导内填充以 $\varepsilon_{\rm r}=2.54$ 的介质时。其工作波长为

$$\lambda = \frac{v}{f} = \frac{3 \times 10^8}{\sqrt{2.54} \times 10 \times 10^9} = 1.88 \times 10^{-2} = 1.88 (cm)$$

波导壁的表面电阻 $R_s = \sqrt{\frac{\pi f \mu}{\gamma}}$

查表得紫铜的电导率 $\gamma = 5.8 \times 10^7 (S/m)$,于是

$$R_s = \sqrt{\frac{3.14 \times 10 \times 10^9 \times 4 \times 3.14 \times 10^{-7}}{5.8 \times 10^7}} = 0.0261(\Omega)$$

矩形波导中传输 TE₁₀ 波时,由导体引起的衰减为

$$\alpha_c = \frac{R_s}{b\eta \sqrt{1 - \left(\frac{\lambda}{2a}\right)^2}} \left[1 + 2\frac{b}{a} \left(\frac{\lambda}{2a}\right)^2 \right]$$

(1) 当波导内为空气填充,
$$\eta = \sqrt{\frac{\mu_0}{\varepsilon_0}} \approx 377(\Omega)$$
,得

$$\alpha_{c} = \frac{R_{s}}{b\eta\sqrt{1 - \left(\frac{\lambda}{2a}\right)^{2}}} \left[1 + 2\frac{b}{a}\left(\frac{\lambda}{2a}\right)^{2}\right]$$

$$= \frac{0.0261}{10 \times 10 \times 377\sqrt{1 - \left(\frac{30}{2 \times 23}\right)^{2}}} \left[1 + \frac{20}{23}\left(\frac{30}{2 \times 23}\right)^{2}\right]$$

$$= 0.011(Np/m)$$

用分贝表示

$$\alpha_c = 0.011 \times 8.686 = 0.094 (dB/m)$$

(2) 当波导内填充以 $\varepsilon_{\rm r}=2.54$ 的介质时

$$\alpha_{c} = \frac{R_{s}\sqrt{\varepsilon_{r}}}{b\eta_{0}\sqrt{1-\left(\frac{\lambda}{2a}\right)^{2}}} \left[1+2\frac{b}{a}\left(\frac{\lambda}{2a}\right)^{2}\right]$$

$$= \frac{0.0261\times\sqrt{2.54}}{10\times10\times377\sqrt{1-\left(\frac{1.88}{2\times23}\right)^{2}}} \left[1+\frac{20}{23}\left(\frac{1.88}{2\times23}\right)^{2}\right]$$

$$= 0.013(Np/m)$$

用分贝表示

$$\alpha_c = 0.013 \times 8.686 = 0.113 (dB/m)$$

8.7 试设计 $\lambda = 10cm$ 的矩形波导。材料用紫铜,内充空气,并且要求 TE_{10} 模的工作 频率至少有 30% 的安全因子,即 $0.7f_{c2} \ge f \ge 1.3f_{c1}$,此处 f_{c1} 和 f_{c2} 分别表示 TE_{10} 波和相 邻高阶模式的截止频率。

解; 由题给: $0.7f_{c2} \ge f \ge 1.3f_{c1}$

即
$$0.7(f_c)_{\mathrm{TE}_{20}} \ge f \ge 1.3(f_c)_{\mathrm{TE}_{10}}$$

若用波长表示,上式变为 $\frac{0.7}{\left(\lambda_{c}\right)_{\text{TE}_{10}}} \geq \frac{1}{\lambda} \geq \frac{1.3}{\left(\lambda_{c}\right)_{\text{TE}_{10}}}$

即

$$\frac{0.7}{a} \ge \frac{1}{10}$$

$$\frac{1.3}{2a} \le \frac{1}{10}$$

由此可得 $6.5 \le a \le 7$

选择: a = 6.8(cm)

为防止高次模 TE_{01} 的出现,窄边 b 的尺寸应满足 $\lambda > (\lambda_c)_{TE_{ac}} = 2b$

即
$$0 < b < 5(cm)$$

考虑到传输功率容量和损耗情况,一般选取 $b = (0.4 \sim 0.5)a$

故设计的矩形波导尺寸为 $a \times b = 6.8 \times 3.4 (cm^2)$

8.8 矩形波导的前半段填充空气,后半段填充介质(介电常数为 ε),问当 TE_{10} 波从空气段入射介质段时,反射波场量和透射波场量各为多大?

解: 由反射系数
$$\rho = \frac{|\mathbf{E}_r|}{|\mathbf{E}_t|} = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1} = \frac{Z_2 - Z_1}{Z_2 + Z_1}$$

得 $|\mathbf{E}_r| = \rho |\mathbf{E}_i|$

即反射波场量的大小为入射波场量的户倍

由透射系数
$$\tau = \frac{|\mathbf{E}_t|}{|\mathbf{E}_t|} = \frac{2\eta_2}{\eta_2 + \eta_1} = \frac{2Z_2}{Z_2 + Z_1}$$

得 $|\mathbf{E}_t| = \tau |\mathbf{E}_i|$

即透射波场量的大小为入射波场量的 τ 倍。因此只须求出 ρ 和 τ 即可得到解答。

矩形波导中
$$TE_{10}$$
模的波阻抗为
$$Z_{TE_{10}} = \frac{\eta}{\sqrt{1 - \left(\frac{\lambda}{\lambda_c}\right)^2}}$$

其中
$$\eta = \sqrt{\frac{\mu}{\varepsilon}} = \frac{\eta_0}{\sqrt{\varepsilon_r}}$$

$$\lambda = \frac{\lambda_0}{\sqrt{\varepsilon_r}}$$

当介质为空气时,得 $Z_1 = \frac{\eta_0}{\sqrt{1 - \left(\frac{\lambda_0}{\lambda_2}\right)^2}}$

当介质的介电常数为 $\varepsilon = \varepsilon_0 \varepsilon_r$ 时,得

$$\mathcal{Z}_2 = \frac{\eta}{\sqrt{1 - \left(\frac{\lambda}{\lambda_c}\right)^2}} = \frac{\eta_0}{\sqrt{\varepsilon_r} \sqrt{1 - \frac{1}{\varepsilon_r} \left(\frac{\lambda_0}{\lambda_c}\right)^2}} = \frac{\eta_0}{\sqrt{\varepsilon_r - \left(\frac{\lambda_0}{\lambda_c}\right)^2}}$$

$$\frac{\eta_0}{\sqrt{\varepsilon_r - \left(\frac{\lambda_0}{\lambda_c}\right)^2}} - \frac{\eta_0}{\sqrt{1 - \left(\frac{\lambda_0}{\lambda_c}\right)^2}} = \frac{\sqrt{1 - \left(\frac{\lambda_0}{2a}\right)^2} - \sqrt{\varepsilon_r - \left(\frac{\lambda_0}{2a}\right)^2}}{\sqrt{1 - \left(\frac{\lambda_0}{\lambda_c}\right)^2}}$$

$$\frac{\eta_0}{\sqrt{\varepsilon_r - \left(\frac{\lambda_0}{\lambda_c}\right)^2}} + \frac{\eta_0}{\sqrt{1 - \left(\frac{\lambda_0}{\lambda_c}\right)^2}} = \frac{\sqrt{1 - \left(\frac{\lambda_0}{2a}\right)^2} - \sqrt{\varepsilon_r - \left(\frac{\lambda_0}{2a}\right)^2}}{\sqrt{1 - \left(\frac{\lambda_0}{\lambda_c}\right)^2}}$$

$$\tau = \frac{2Z_2}{Z_2 + Z_1} = \frac{2\eta_0}{\sqrt{\varepsilon_r - \left(\frac{\lambda_0}{\lambda_c}\right)^2}} + \frac{2\eta_0}{\sqrt{1 - \left(\frac{\lambda_0}{\lambda_c}\right)^2}} = \frac{2\sqrt{1 - \left(\frac{\lambda_0}{2a}\right)^2}}{\sqrt{1 - \left(\frac{\lambda_0}{2a}\right)^2}} + \sqrt{\varepsilon_r - \left(\frac{\lambda_0}{2a}\right)^2}$$

8.9 试推导在矩形波导中传输 TE_{mn} 波时的传输功率。

解:波导中传输的功率可由波导横截面上坡印廷矢量的积分求得

$$P = \operatorname{Re} \frac{1}{2} \int_{s} \mathbf{E} \times \mathbf{H}^{*} \cdot d\mathbf{S} = \frac{1}{2Z_{\operatorname{TE}_{mn}}} \int_{s} \left| \mathbf{E} \right|^{2} dS = \frac{Z_{\operatorname{TE}_{mn}}}{2} \int_{s} \left| \mathbf{H} \right|^{2} dS$$
$$= \frac{1}{2Z_{\operatorname{TE}}} \int_{0.0}^{b} \int_{0.0}^{a} \left(\left| E_{x} \right|^{2} + \left| E_{y} \right|^{2} \right) dx dy$$

式中 \mathbf{E} 和 \mathbf{H} 分别为波导横截面内的电场强度和磁场强度, Z_{TE} 为波阻抗。

式中 E 和 H 分别为波导横截面内的电场强度和磁场强度,
$$Z_{TE_{man}}$$
 为波阴矩形波导中 $E_x(x,y) = \frac{j\omega\mu}{h^2} \left(\frac{n\pi}{b}\right) H_0 \cos\left(\frac{m\pi}{a}x\right) \sin\left(\frac{n\pi}{b}y\right)$
得 $|E_x(x,y)| = E_m \left(\frac{n\pi}{b}\right) \cos\left(\frac{m\pi}{a}x\right) \sin\left(\frac{n\pi}{b}y\right)$
 $E_y(x,y) = -\frac{j\omega\mu}{h^2} \left(\frac{m\pi}{a}\right) H_0 \sin\left(\frac{m\pi}{a}x\right) \cos\left(\frac{n\pi}{b}y\right)$
对中 $E_m \left(\frac{m\pi}{a}\right) \sin\left(\frac{m\pi}{a}x\right) \cos\left(\frac{n\pi}{b}y\right)$
式中 $E_m = \frac{\omega\mu}{h^2} H_0$
于是
$$P = \frac{1}{2Z_{TE_{man}}} \int_0^b \left(|E_x|^2 + |E_y|^2 \right) dx dy$$

$$= \frac{E_m^2 \left(\frac{n\pi}{b}\right)^2}{2Z_{TE_{man}}} \int_0^b \sin^2\left(\frac{n\pi}{b}y\right) dy \int_0^a \cos^2\left(\frac{m\pi}{a}x\right) dx$$

$$+ \frac{E_m^2 \left(\frac{m\pi}{a}\right)^2}{2Z_{TE_{man}}} \int_0^b \cos^2\left(\frac{n\pi}{b}y\right) dy \int_0^a \sin^2\left(\frac{m\pi}{a}x\right) dx$$

$$= \frac{E_m^2 \left(\frac{n\pi}{a}\right)^2}{2Z_{TE_{man}}} \int_0^b \cos^2\left(\frac{n\pi}{b}y\right) dy \int_0^a \sin^2\left(\frac{m\pi}{a}x\right) dx$$

$$= \frac{E_m^2 \left(\frac{n\pi}{a}\right)^2}{2Z_{TE_{man}}} \int_0^b \cos^2\left(\frac{n\pi}{b}y\right) dy \int_0^a \sin^2\left(\frac{m\pi}{a}x\right) dx$$

$$= \frac{E_m^2 \left(\frac{n\pi}{a}\right)^2}{2Z_{TE_{man}}} \int_0^b \cos^2\left(\frac{n\pi}{b}y\right) dy \int_0^a \sin^2\left(\frac{m\pi}{a}x\right) dx$$

$$= \frac{E_m^2 \left(\frac{n\pi}{b}\right)^2}{2Z_{TE_{man}}} N_m N_n E_m^2 \left[\left(\frac{n\pi}{b}\right)^2 + \left(\frac{m\pi}{a}\right)^2\right]$$

$$= \frac{ab}{8Z_{TE_{man}}} E_m^2 h^2 N_m N_n$$

$$\exists \Phi N_m = \begin{cases} 1 & (m \neq 0) \\ 2 & (m = 0) \end{cases}, N_n = \begin{cases} 1 & (n \neq 0) \\ 2 & (m = 0) \end{cases}$$

8.10 试设计一工作波长 $\lambda = 5cm$ 的圆柱形波导, 材料用紫铜, 内充空气, 并要求 TE_{11} 波的工作频率应有一定的安全因子。

解: TE_{11} 模是圆柱形波导中的主模,为保证单模传输,应使工作频率大于 TE_{11} 模的截止频率而小于第一次高模 TM_{01} 的截止频率,即

$$\lambda < (\lambda_{c})_{TE_{11}} \left(= \frac{2\pi a}{1.841} \right)$$

$$\lambda > (\lambda_{c})_{TM_{11}} \left(= \frac{2\pi a}{2.405} \right)$$

于是得
$$\frac{2\pi a}{2.405} < \lambda < \frac{2\pi a}{1.841}$$

圆柱形波导的半径 a 应满足 $\frac{\lambda}{2.61} > a > \frac{\lambda}{3.41}$

选择 $a = \frac{\lambda}{3} = \frac{5}{3}$ (cm)

8.11 求圆柱形波导中 TE_{0n} 波的传输功率。

解 : 传 输 功 率

$$P = \text{Re} \frac{1}{2} \int_{S} \mathbf{E} \times \mathbf{H}^{*} \cdot d\mathbf{S} = \frac{1}{2Z_{\text{TE}_{0n}}} \int_{S} |\mathbf{E}|^{2} dS = \frac{1}{2Z_{\text{TE}_{0n}}} \int_{0}^{2\pi} \int_{0}^{a} (|E_{r}|^{2} + |E_{\varphi}|^{2}) r dr d\varphi$$

柱形波导中的 TE_{on} 模的场分量

$$\begin{aligned} \left| E_r \right| &= 0 \\ \left| E_{\varphi} \right| &= \frac{\omega \mu}{k_c} H_0 J_{\mathrm{m}}(hr) = E_0 J_{\mathrm{m}}(hr) \end{aligned}$$

由贝塞尔函数的递推公式 $J_{m}(hr) = \frac{m}{k_{c}r} J_{m}(hr) - J_{m+1}(hr)$

因为 m=0 则
$$J_{\rm m}^{'}(hr) = J_{\rm 0}^{'}(hr) = -J_{\rm 1}(hr)$$

所以
$$\left|E_{\varphi}\right| = E_0 J_1(hr)$$

$$P = \frac{2\pi}{2Z_{\text{TE}_{0n}}} E_0^2 \int_0^a J_1^2 (hr) r dr$$

而

$$\int_{0}^{a} J_{1}^{2}(hr)rdr = \frac{1}{k_{c}^{2}} \int_{0}^{a} (hr) J_{1}^{2}(hr) d(hr)$$
$$= \frac{a^{2}}{2} \left[J_{1}^{2}(ha) - J_{0}(ha) J_{2}(ha) \right]$$

由电场切向分量连续的边界条件可知 $E_{\varphi}(r=a)=0$

$$J_0(ha) = -J_1(ha) = 0$$

$$J_2(ha) = \frac{2}{ha}J_1(ha) - J_0(ha) = -J_0(ha)$$

故
$$\int_{0}^{a} J_{1}^{2}(hr)rdr = \frac{a^{2}}{2}J_{0}^{2}(ha)$$

则圆柱形波导中 TE_{0n} 的传输功率为 $P = \frac{\pi a^2}{2Z_{TE_0}} E_0^2 J_0^2 (ha)$

8.12 试求圆波导中 TE_{0n} 模由于管壁不是完纯导体而引起的衰减 α_{c} 。

解:波导中由于管壁不是完纯导体而引起的衰减 $lpha_{
m c}=rac{P_l}{2P}$

式中: P_1 表示波导中单位长度的损耗功率; P表示传输功率。

$$P_l = \frac{1}{2} \int_s R_s \left| \mathbf{J}_s \right|^2 dS$$

R。为导体的表面电阻。而

$$\begin{aligned} \left| \mathbf{J}_{s} \right|_{r=a} &= \left| -\mathbf{a}_{\mathbf{r}} \times \left(\mathbf{a}_{r} H_{r} + \mathbf{a}_{\varphi} H_{\varphi} + \mathbf{a}_{z} H_{z} \right) \right|_{r=a} \\ &= \left| \left(-\mathbf{a}_{z} H_{\varphi} + \mathbf{a}_{\varphi} H_{z} \right) \right|_{r=a} \end{aligned}$$

由圆波导中 $\mathrm{TE}_{\mathrm{mn}}$ 的场分量表示式可知,当 $\mathrm{m=0}$ 时 $H_{\mathrm{o}}=0$,得

$$\left|\mathbf{J}_{s}\right|_{r=a} = \left|\mathbf{a}_{\varphi}H_{z}\right|_{r=a} = \left|H_{z}\right|_{r=a} = H_{0}J_{0}(ha)$$

则

$$P_{l} = \frac{1}{2} \int_{s} R_{s} |\mathbf{J}_{s}|^{2} dS = \frac{1}{2} \int_{0}^{2\pi} R_{s} |J_{\varphi}|^{2} r d\varphi$$
$$= \frac{1}{2} R_{s} 2\pi a H_{0}^{2} J_{0}^{2} (ha)$$

由上题得TEon模的传输功率

$$P = \frac{\pi a^2}{2Z_{TE}} E_0^2 J_0^2 (ha) = \frac{\pi a^2}{2Z_{TE}} \left(\frac{\omega \mu}{h}\right)^2 H_0^2 J_0^2 (ha)$$

故

$$\alpha_{c} = \frac{P_{l}}{2P} = \frac{\frac{1}{2}R_{s}2\pi aH_{0}^{2}J_{0}^{2}(ha)}{2\frac{\pi a^{2}}{2Z_{TE_{0n}}}\left(\frac{\omega\mu}{h}\right)^{2}H_{0}^{2}J_{0}^{2}(ha)}$$
$$= \frac{R_{s}Z_{TE_{0n}}}{a}\left(\frac{h}{\omega\mu}\right)^{2}$$

因为

$$Z_{\text{TE}_{0n}} = \frac{\eta}{\sqrt{1 - \left(\frac{f_c}{f}\right)^2}}; \qquad h = \omega_c \sqrt{\mu \varepsilon}$$

所以TE_{0n}模的衰减常数为

$$\alpha_{c} = \frac{R_{s}}{a} \frac{\eta}{\sqrt{1 - \left(\frac{f_{c}}{f}\right)^{2}}} \left(\frac{\omega_{c}\sqrt{\mu\varepsilon}}{\omega\mu}\right)$$

$$= \frac{R_{s}}{a} \frac{\eta}{\sqrt{1 - \left(\frac{f_{c}}{f}\right)^{2}}} \left(\frac{f_{c}}{f}\right)^{2} \left(\frac{1}{\eta}\right)^{2}$$

$$= \frac{R_{s}\left(\frac{f_{c}}{f}\right)^{2}}{a\eta\sqrt{1 - \left(\frac{f_{c}}{f}\right)^{2}}}$$

8.13 已知在圆柱形波导中, TM_{mn} 波由于壁面不完纯而引起的衰减常数为

$$\alpha_c = \frac{R_s / a\eta}{\sqrt{1 - \left(\frac{f_c}{f}\right)^2}}$$

求证:衰减的最小值出现在 $f = \sqrt{3} f_c$ 处。

证: 因为

$$\alpha_c = \frac{R_s / a\eta}{\sqrt{1 - \left(\frac{f_c}{f}\right)^2}}$$

而导体的表面电阻 $R_s = \sqrt{\frac{\pi \mu f}{\gamma}} = N\sqrt{f}$

因此 α_c 的最小值可由 $\frac{d\alpha}{df} = 0$ 求得

$$\begin{split} \frac{d\alpha}{df} &= \frac{d}{df} \left[\frac{N}{a\eta} \left(\frac{f^3}{f^2 - f_c^2} \right)^{\frac{1}{2}} \right] \\ &= \frac{N}{a\eta} \frac{1}{2} \left(\frac{f^3}{f^2 - f_c^2} \right)^{-\frac{1}{2}} \cdot \left[\frac{3f^2 \left(f^2 - f_c^2 \right) - 2f \cdot f^3}{\left(f^2 - f_c^2 \right)^2} \right] \\ &= \frac{N}{2a\eta} \cdot \left[\frac{\left(3f^4 - 3f^2 f_c^2 - 2f^4 \right) \left(f^2 - f_c^2 \right)^{\frac{1}{2}}}{f^{3/2} \left(f^2 - f_c^2 \right)^2} \right] = \frac{N}{2a\eta} \cdot \frac{\left(f^4 - 3f^2 f_c^2 \right)}{f^{3/2} \left(f^2 - f_c^2 \right)^{\frac{3}{2}}} = 0 \end{split}$$

得
$$Nf^2(f^2-3f_c^2)=0$$

$$\mathbb{P} \quad f^2 - 3f_c^2 = 0$$

所以
$$f = \sqrt{3} f_c$$

8.14 设计一矩形谐振腔, 使在 1 及 1.5GHz 分别谐振于两个不同模式上。

解: 矩形谐振腔的谐振频率为
$$f_{mnl} = v \left[\left(\frac{m}{2a} \right)^2 + \left(\frac{n}{2b} \right)^2 + \left(\frac{l}{2d} \right)^2 \right]^{\frac{1}{2}}$$

若使在 1 及 1.5GHz 分别谐振于矩形谐振腔的 TE_{101} 及 TE_{102} 两个不同模式上,则它们的谐振频率分别为

$$f_{101} = 3 \times 10^8 \left[\left(\frac{1}{2a} \right)^2 + \left(\frac{1}{2d} \right)^2 \right]^{\frac{1}{2}} = 1 \times 10^9$$

$$f_{102} = 3 \times 10^8 \left[\left(\frac{1}{2a} \right)^2 + \left(\frac{1}{d} \right)^2 \right]^{\frac{1}{2}} = 1.5 \times 10^9$$

$$\left(\frac{1}{2a} \right)^2 + \left(\frac{1}{2d} \right)^2 = \left(\frac{10}{3} \right)^2 \qquad (1)$$

$$\left(\frac{1}{2a} \right)^2 + \left(\frac{1}{d} \right)^2 = \left(\frac{15}{3} \right)^2 \qquad (2)$$
将以上二式相减得 $\frac{1}{d^2} \left(1 - \frac{1}{4} \right) = \left(\frac{15}{3} \right)^2 - \left(\frac{10}{3} \right)^2 = 13.9$
可得 $d = \sqrt{\frac{3}{4 \times 13.9}} = 0.23(m)$
将其代入式 (2) 得 $\left(\frac{1}{2a} \right)^2 = 25 - \left(\frac{100}{23} \right)^2 = 6.1$
所以 $a = \sqrt{\frac{1}{4 \times 6.1}} = 0.20(m)$
尺寸 b 可取为 $b = \frac{a}{2} = 0.10(m)$

于是该矩形谐振腔的尺寸为 $a \times b \times d = 0.20 \times 0.10 \times 0.23 (m^3)$

8.15 由空气填充的矩形谐振腔,其尺寸为 a=25mm,b=12.5mm,d=60mm,谐振于 TE102 模式,若在腔内填充介质,则在同一工作频率将谐振一 TE103 模式,求介质的相对介电常数 ε_r 应为多少?

解: 矩形谐振腔的谐振频率为
$$f_{mnl} = v \left[\left(\frac{m}{2a} \right)^2 + \left(\frac{n}{2b} \right)^2 + \left(\frac{l}{2d} \right)^2 \right]^{\frac{1}{2}}$$
 当填充介质为空气时 $v = c = 3 \times 10^8 \left(m/s \right)$ TE102 模的谐振频率为

$$f_{102} = 3 \times 10^8 \left[\left(\frac{10^3}{2 \times 25} \right)^2 + \left(\frac{2 \times 10^3}{2 \times 60} \right)^2 \right]^{\frac{1}{2}}$$
$$= 7.8 \times 10^9 \left(H_z \right)$$

当填充介质的介电常数为 ε_r 时, $v = \frac{c}{\sqrt{\varepsilon_r}}$,TE103 模的谐振频率为

$$f_{103} = \frac{3 \times 10^8}{\sqrt{\varepsilon_r}} \left[\left(\frac{10^3}{2 \times 25} \right)^2 + \left(\frac{3 \times 10^3}{2 \times 60} \right)^2 \right]^{\frac{1}{2}}$$

 $f_{103} = f_{102} = 7.8 \times 10^9$

得
$$\varepsilon_r = \left(\frac{3 \times 10^8}{7.8 \times 10^9}\right)^2 \left[\left(\frac{10^3}{50}\right)^2 + \left(\frac{3 \times 10^3}{120}\right)^2 \right] = 1.52$$

平行双线传输线的线间距 D=8cm, 导线的直径 d=1cm, 周围是空气, 试计算: (1) 分布电感和分布电容; (2) f=600MHz 时的相位系数和特性阻抗 $(R_1 = 0, G_1 = 0)$ 。

解: (1) 双线传输线分布电容
$$C_1 = \frac{\varepsilon \pi}{\ln \frac{2D}{d}} = \frac{\pi \varepsilon_0}{\ln 16} = 10 \left(\text{pF/m} \right)$$

分布电感
$$L_1 = \frac{\mu_0}{\pi} \ln \frac{2D}{d} = \frac{4\pi \times 10^{-7}}{\pi} \times 2.7726 = 1.11 \left(\mu H/m \right)$$

$$f = 6 \times 10^8 \text{Hz}$$

(2)
$$f = 6 \times 10^8 \text{ Hz}$$
$$\beta = \omega \sqrt{L_1 C_1} = 2\pi \times 10^8 \times \sqrt{10^{-11} \times 1.11 \times 10^{-6}} = 12.86 \text{ (rad/m)}$$

$$Z_0 = \sqrt{\frac{L_1}{C_1}} = \sqrt{\frac{1.11 \times 10^{-6}}{10^{-11}}} = 333(\Omega)$$

8.17 同轴线的外导体半径 b = 23mm,内导体半径 a = 10mm,填充介质分别为空气 和 $\varepsilon_r = 2.25$ 的无耗介质, 试计算其特性组抗。

解:(1)填充空气时

$$C_{1} = \frac{2\pi\varepsilon_{0}}{\ln\frac{b}{a}} = \frac{2\pi \times 8.85 \times 10^{-12}}{\ln\frac{23}{10}} = 6.68 \times 10^{-11} (\text{F/m})$$

$$L_{1} = \frac{\mu_{0}}{2\pi} \ln\frac{b}{a} = \frac{4\pi \times 10^{-7}}{2\pi} \ln\frac{23}{10} = 1.67 \times 10^{-7} (\text{H/m})$$

$$Z_{0} = \sqrt{\frac{L_{1}}{C_{1}}} = \frac{\eta_{0}}{2\pi} \ln\frac{b}{a} = \frac{120\pi}{2\pi} \ln 2.3 = 50 (\Omega)$$

特性阻抗

(2)
$$\varepsilon_r = 2.25 \,\text{ft}$$
, $\eta = \frac{\eta_0}{\sqrt{\varepsilon_r}} = \frac{120\pi}{\sqrt{2.25}}$

$$Z_0' = \sqrt{\frac{L_1'}{C_1'}} = \frac{\eta}{2\pi} \ln \frac{b}{a} = \frac{Z_0}{\sqrt{2.25}} = 33.32(\Omega)$$

- 8. 18 在构造均匀传输线时,用聚乙烯($\varepsilon_r = 2.25$)作为电介质。假设不计损耗。
- (1) 对于 300Ω 的平行双线, 若导线的半径为 0.6mm, 则线间距应选多少?
- (2) 对于 75Ω 的同轴线,若内导体的半径为 0.6mm,则外导体的半径应选多少? \mathbf{M} : (1) 双线传输线,设 \mathbf{a} 为导体半径, \mathbf{D} 为线间距,则

$$C_1 = \frac{\varepsilon \pi}{\ln \frac{D}{a}} \qquad L_1 = \frac{\mu_0}{\pi} \ln \frac{D}{a}$$

$$Z_0 = \sqrt{\frac{L_1}{C_1}} = \frac{\eta_0}{\sqrt{\varepsilon_r}\pi} \ln \frac{D}{a} 300(\Omega)$$
$$\ln \frac{D}{a} = \frac{300}{120} \times \sqrt{2.25} = 3.75$$

则线间距

$$D = 42.5 \times 0.6 = 25.5 (mm)$$

(2) 同轴线传输线,设 a 为内导体半径, b 为外导体内半径,则

$$C_1 = \frac{2\varepsilon\pi}{\ln\frac{b}{a}} \qquad L_1 = \frac{\mu_0}{2\pi} \ln\frac{b}{a}$$

$$Z_0 = \sqrt{\frac{L_1}{C_1}} = \frac{\eta_0}{2\sqrt{\varepsilon_r}\pi} \ln\frac{b}{a} = 75(\Omega)$$

$$\ln\frac{b}{a} = \frac{75 \times 2 \times \sqrt{2.25}}{120} = 1.875$$

则外导体的内半径 $b = 6.516 \times 0.6 = 3.91 (mm)$

- **8.19** 试以传输线输入端电压 U_1 和电流 I_1 以及传输线的传播系数 Γ 和特性阻抗 Z_0 表示线上任意一点的电压分布U(z)和电流分布I(z)。
 - (1) 用指数形式表示;
 - (2) 用双曲函数表示。

解: 传输线上电压和电流的通解形式为

$$U(z) = A_1 e^{\Gamma z} + A_2 e^{-\Gamma z}$$
$$I(z) = \frac{1}{Z_0} \left(A_1 e^{\Gamma z} - A_2 e^{-\Gamma z} \right)$$

式中传播系数 Γ 和特性阻抗 Z_0 分别为

$$\Gamma = \sqrt{\left(R_1 + j\omega L_1\right)\left(G_1 + j\omega C_1\right)}$$

$$Z_0 = \sqrt{\frac{R_1 + j\omega L_1}{G_1 + j\omega C_1}}$$

对于输入端: z=l

$$U_{1} = A_{1}e^{\Gamma l} + A_{2}e^{-\Gamma l}$$

$$I_{1} = \frac{1}{Z_{0}} \left(A_{1}e^{\Gamma l} - A_{2}e^{-\Gamma l} \right)$$

联立求解得

$$A_{1} = \frac{1}{2} (U_{1} + I_{1} Z_{0}) e^{-\Gamma I}$$

$$A_{2} = \frac{1}{2} (U_{1} - I_{1} Z_{0}) e^{\Gamma I}$$

可得

$$U(z) = \frac{1}{2} (U_1 + I_1 Z_0) e^{\Gamma(z-l)} + \frac{1}{2} (U_1 - I_1 Z_0) e^{-\Gamma(z-l)}$$
$$I(z) = \frac{1}{2Z_0} (U_1 + I_1 Z_0) e^{\Gamma(z-l)} - \frac{1}{2Z_0} (U_1 - I_1 Z_0) e^{-\Gamma(z-l)}$$

用双曲函数表示

$$U(z) = \frac{1}{2} U_1 \left(e^{\Gamma(z-l)} + e^{-\Gamma(z-l)} \right) + \frac{I_1 Z_0}{2} \left(e^{\Gamma(z-l)} - e^{-\Gamma(z-l)} \right)$$

$$= U_1 \text{ch}\Gamma(z-l) + I_1 Z_0 \text{sh}\Gamma(z-l)$$

$$I(z) = \frac{1}{2Z_0} U_1 \left(e^{\Gamma(z-l)} + e^{-\Gamma(z-l)} \right) + \frac{I_1}{2} \left(e^{\Gamma(z-l)} - e^{-\Gamma(z-l)} \right)$$

$$= I_1 \text{ch}\Gamma(z-l) + \frac{U_1}{Z_0} \text{sh}\Gamma(z-l)$$

8.20 一根特性阻抗为 50Ω 、长度为 2m 的无损耗传输线工作于频率 $200 \mathrm{MHz}$,终端接有阻抗 $Z_{\mathrm{L}} = 40 + j30\Omega$,试求其输入阻抗。

解:无损耗线得输入阻抗 $Z_{\rm in}=Z_0+rac{Z_{\rm L}+jZ_0 aneta z}{Z_0+jZ_{\rm L} aneta z}$

而

$$\beta z = \frac{2\pi}{\lambda} \times 2$$

$$\lambda = \frac{c}{f} = \times \frac{3 \times 10^8}{210^8} = 1.5(m)$$

所以

$$\beta z = \frac{4}{1.5}\pi = 480^{\circ}$$

$$\tan 480^{\circ} = -1.732$$

则

$$Z_{in} = 50 + \frac{(40 + j30) + j50 \times (-1.732)}{50 + j(40 + j30) \times (-1.732)}$$
$$= 26.32 - j9.87(\Omega)$$

- **8.21** 一根 75 Ω 的无损耗线,终端接有负载阻抗 $Z_L = R_L + jX_L$ 。
- (1) 欲使线上的电压驻波比等于 3,则 RL 和 XL 有什么关系?
- (2) 若 $R_L = 150\Omega$, 求XL等于多少?
- (3) 求在(2)情况下, 距负载最近的电压最小点位置。

解: (1) 由驻波比 S 与反射系数 ρ 的关系 $|\rho| = \frac{S-1}{S+1} = \frac{1}{2}$

$$\overrightarrow{m} \quad \left| \rho \right| = \left| \frac{Z_{L} - Z_{0}}{Z_{L} + Z_{0}} \right|$$

即

$$\left|\rho\right| = \left[\frac{\left(R_{L} - Z_{0}\right)^{2} + X_{L}^{2}}{\left(R_{L} + Z_{0}\right)^{2} + X_{L}^{2}}\right]^{\frac{1}{2}} = \frac{1}{2}$$

$$4\left(R_{L} - Z_{0}\right)^{2} + 4X_{L}^{2} = \left(R_{L} + Z_{0}\right)^{2} + X_{L}^{2}$$

解得

$$X_{L} = \pm Z_{0} \sqrt{-\left(\frac{R_{L}}{Z_{0}}\right)^{2} + \frac{10}{3} \left(\frac{R_{L}}{Z_{0}}\right) - 1}$$
$$= \pm 75 \sqrt{-\left(\frac{R_{L}}{75}\right)^{2} + \frac{10}{3} \left(\frac{R_{L}}{75}\right) - 1}$$

(2) 将 $R_L = 150\Omega$ 代如上式,得

$$X_{L} = \pm 75 \sqrt{-\left(\frac{150}{75}\right)^{2} + \frac{10}{3}\left(\frac{150}{75}\right) - 1}$$
$$= \pm 96.82(\Omega)$$

(3) 终端反射系数

$$\rho_2 = \frac{\left(R_L - Z_0\right) + X_L}{\left(R_L + Z_0\right) + X_L}$$

$$= \frac{\left(150 - 75\right) + j96.82}{\left(150 + 75\right) + j96.82}$$

$$= 0.4375 + j0.242$$

$$= 0.5e^{j\theta_2}$$

式中

$$\theta_2 = \arctan \frac{0.242}{0.4375}$$
$$= \arctan 0.0553$$
$$\approx 29^0$$

传输线的电压分布

$$U(z) = Ae^{j\beta z} + \rho_2 Ae^{-j\beta z}$$

$$= Ae^{j\beta z} \left(1 + \rho_2 e^{-2j\beta z} \right)$$

$$= Ae^{j\beta z} \left(1 + \left| \rho_2 \right| e^{j\theta_2} e^{-2j\beta z} \right)$$

$$= Ae^{j\beta z} \left[1 + \left| \rho_2 \right| e^{j(2\beta z - \theta_2)} \right]$$

电压的幅值

$$|U(z)| = |Ae^{j\beta z}| |1 + |\rho_2|e^{j(2\beta z - \theta_2)}|$$

$$= |A| + \sqrt{1 + |\rho_2|^2 + 2|\rho_2|\cos(2\beta z - \theta_2)}$$

波节点出现在 $\cos(2\beta z - \theta_2) = -1$

第一波节点出现在 $2\beta z_1 - \theta_2 = 180^{\circ}$

即
$$\frac{4 \times 180^{0}}{\lambda} z_{1} - 29^{0} = 180^{0}$$

解得
$$z_1 = \frac{180^0 + 29^0}{4 \times 180^0} \lambda = 0.29\lambda$$

8.22 考虑一根无损耗传输线,

- (1) 当负载阻抗 $Z_L = (40 + j30)\Omega$ 时,欲使线上驻波比最小,则线的特性阻抗应为多少?
- (2) 求出该最小的驻波比及相应的电压反射系数。
- (3) 确定距负载最近的电压最小点位置。

解: (1) 因为
$$|\rho| = \frac{S-1}{S+1}$$
 得 $S = \frac{1+|\rho|}{1-|\rho|}$

驻波比 S 要最小,就要求反射系数 $|\rho|$ 最小,而 $|\rho| = \left[\frac{(R_{\rm L} - Z_0)^2 + X_{\rm L}^2}{(R_{\rm L} + Z_0)^2 + X_{\rm L}^2}\right]^{\frac{1}{2}}$

其最小值可由 $\frac{d|\rho|}{dZ_0} = 0$ 求得 $Z_0^2 = R_L^2 + X_L^2 = 40^2 + 30^2$

故 $Z_0 = 50(\Omega)$

(2) 将 $Z = 50(\Omega)$ 代入反射系数公式,得

$$\left| \rho \right|_{\min} = \left[\frac{\left(R_{L} - Z_{0} \right)^{2} + X_{L}^{2}}{\left(R_{L} + Z_{0} \right)^{2} + X_{L}^{2}} \right]^{\frac{1}{2}}$$

$$= \left[\frac{\left(40 - 50 \right)^{2} + 30^{2}}{\left(40 + 50 \right)^{2} + 30^{2}} \right]^{\frac{1}{2}}$$

$$= \frac{1}{3}$$

最小驻波比为 $S_{\min} = \frac{1+\left|\rho\right|_{\min}}{1-\left|\rho\right|_{\min}} = \frac{1+\frac{1}{3}}{1-\frac{1}{3}} = 2$

(3) 终端反射系数

$$\rho_2 = \frac{(R_L - Z_0) + jX_L}{(R_L + Z_0) + jX_L}$$
$$= \frac{(40 - 50) - j30}{(40 + 50) - j30}$$
$$= 0.333e^{-j90^0}$$

由上题的结论,电压的第一个波节点 z_1 应满足 $2\frac{2\pi}{\lambda}z_1-\theta_2=180^0$

即
$$\frac{4 \times 180^{\circ}}{\lambda} z_1 + 90^{\circ} = 180^{\circ}$$

解得 $z_1 = \frac{180^{\circ} - 90^{\circ}}{4 \times 180^{\circ}} \lambda = 0.125\lambda$

8.23 有一段特性阻抗为 $Z_0 = 500\Omega$ 的无损耗线, 当终端短路时, 测的始端的阻抗为 250Ω 的感抗, 求该传输线的最小长度, 如果该线的终端为开路, 长度又为多少?

解: (1) 终端短路线的输入阻抗为 $Z_{in} = jZ_0 \tan \beta z$

$$j500 \tan \beta z = j250$$

 $\beta z = \arctan 0.5 = 26.57^{\circ}$

将
$$\beta z = \frac{2\pi}{\lambda}$$
代入上式得传输线的长度为 $z = \frac{26.57^{\circ}}{2 \times 180^{\circ}} \lambda = 0.074 \lambda$

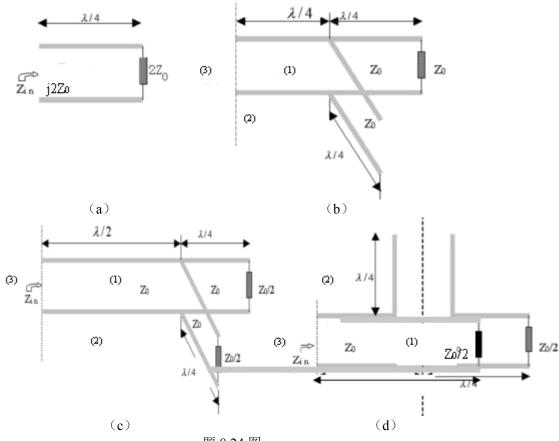
(2) 终端开路线的输入阻抗为
$$Z_{in} = \frac{Z_0}{j \tan \beta z}$$

即

$$500 = -250 \tan \beta z$$
$$\beta z = 116.57^{\circ}$$

将
$$\beta z = \frac{2\pi}{\lambda}$$
代入上式得传输线的长度为 $z = \frac{116.57^{\circ}}{2 \times 180^{\circ}} \lambda = 0.324 \lambda$

8.24 求如题 8.24 图示的分布参数电路的输入阻。



题 8.24 图

解: 设传输线无损耗,则输入阻抗为
$$Z_{\rm in}=Z_0+rac{Z_{\rm L}+jZ_0 aneta z}{Z_0+jZ_{\rm L} aneta z}$$

当传输线长度
$$z = \frac{\lambda}{4}$$
时 $Z_{in} \left(\frac{\lambda}{4}\right) = \frac{Z_0^2}{Z_L}$ ($\frac{\lambda}{4}$ 阻抗变换性)

当传输线长度
$$z = \frac{n\lambda}{2}$$
 时 $Z_{in} \left(\frac{n\lambda}{2} \right) = Z_{L}$ ($\frac{\lambda}{2}$ 阻抗还原性)

(a)
$$Z_{in} \left(\frac{\lambda}{4} \right) = \frac{Z_0^2}{Z_1} = -j0.5Z_0$$

(b) 支节①
$$Z_{\text{inl}} \left(\frac{\lambda}{4} \right) = \frac{Z_0^2}{Z_{11}} = \frac{Z_0^2}{Z_0} = Z_0$$

支节②
$$Z_{\text{in2}}\left(\frac{\lambda}{4}\right) = \frac{Z_0^2}{Z_{\text{L2}}} = \frac{Z_0^2}{\infty} = 0$$

支节③ $Z_{\text{L3}} = Z_{\text{in1}} //Z_{\text{in2}} = 0$
 $Z_{\text{in}} = Z_{\text{in3}}\left(\frac{\lambda}{4}\right) = \frac{Z_0^2}{Z_{\text{L3}}} = \frac{Z_0^2}{0} = \infty$

(c) 支节① $Z_{\text{in1}}\left(\frac{\lambda}{4}\right) = \frac{Z_0^2}{Z_{\text{L1}}} = \frac{Z_0^2}{\frac{1}{2}Z_0} = 2Z_0$

支节② $Z_{\text{in2}}\left(\frac{\lambda}{4}\right) = \frac{Z_0^2}{Z_{\text{L2}}} = \frac{Z_0^2}{\frac{1}{2}Z_0} = 2Z_0$

支节③ $Z_{\text{L3}} = Z_{\text{in1}} //Z_{\text{in2}} = Z_0$
 $Z_{\text{in}} = Z_{\text{in3}}\left(\frac{\lambda}{4}\right) = \frac{Z_0^2}{Z_{\text{L3}}} = \frac{Z_0^2}{Z_0} = Z_0$

(d) 支节① $Z_{\text{in1}}\left(\frac{\lambda}{2}\right) = \frac{1}{2}Z_0$

支节② $Z_{\text{in2}}\left(\frac{\lambda}{4}\right) = \frac{Z_0^2}{Z_{\text{L2}}} = \frac{Z_0^2}{\infty} = 0$

支节③
$$Z_{L3} = Z_{in1} + Z_{in2} = \frac{Z_0}{2}$$

$$Z_{L3} = Z_{in1} + Z_{in2} = \frac{Z_0}{2}$$

$$Z_{\text{in}} = Z_{\text{in3}} \left(\frac{\lambda}{4}\right) = \frac{Z_0^2}{Z_{L3}} = \frac{Z_0^2}{Z_0/2} = 2Z_0$$

8.24 求题 8.24 图中各段的反射系数及驻波系数。

解:: 终端反射系数
$$\rho_2 = \frac{Z_L - Z_0}{Z_L + Z_0}$$
 反射系数 $\rho = \rho_2 e^{-2j\beta z}$

驻波系数 $S = \frac{1 + |\rho_2|}{1 - |\rho_2|}$

ます②
$$S = \frac{1+|\rho_2|}{1-|\rho_2|} = \frac{1+0}{1-0} = 1$$

支节② $\rho_2 = \frac{Z_L - Z_0}{Z_L + Z_0} = \frac{\infty - Z_0}{\infty + Z_0} = 1$
 $\rho = \rho_2 e^{-2j\beta z} = e^{-2j\beta z}$
 $S = \frac{1+|\rho_2|}{1-|\rho_2|} = \frac{1+1}{1-1} = \infty$

支节③ $\rho_2 = \frac{Z_L - Z_0}{Z_L + Z_0} = \frac{0 - Z_0}{0 + Z_0} = -1 = e^{j\pi}$
 $\rho = \rho_2 e^{-2j\beta z} = e^{j(\pi - 2\beta z)}$
 $S = \frac{1+|\rho_2|}{1-|\rho_2|} = \frac{1+1}{1-1} = \infty$

(c) 艾节①、② $\rho_2 = \frac{Z_L - Z_0}{Z_L + Z_0} = \frac{\frac{1}{2}Z_0 - Z_0}{\frac{1}{2}Z_0 + Z_0} = \frac{-1}{3} = \frac{1}{3}e^{j\pi}$
 $\rho = \rho_2 e^{-2j\beta z} = \frac{1}{3}e^{j(\pi - 2\beta z)}$
 $S = \frac{1+|\rho_2|}{1-|\rho_2|} = \frac{1+\frac{1}{3}}{1-\frac{1}{3}} = 2$

支节③ $\rho_2 = \frac{Z_L - Z_0}{Z_L + Z_0} = \frac{\frac{1}{2}Z_0 - Z_0}{\frac{1}{2}Z_0 + Z_0} = 0$
 $\rho = \rho_2 e^{-2j\beta z} = 0$
 $S = \frac{1+|\rho_2|}{1-|\rho_2|} = \frac{1+0}{1-0} = 1$

(d) 支节① $\rho_2 = \frac{Z_L - Z_0}{Z_L + Z_0} = \frac{\frac{1}{2}Z_0 - Z_0}{\frac{1}{2}Z_0 + Z_0} = -\frac{1}{3} = \frac{1}{3}e^{j\pi}$
 $\rho = \rho_2 e^{-2j\beta z} = \frac{1}{3}e^{j(\pi - 2\beta z)}$
 $S = \frac{1+|\rho_2|}{1-|\rho_2|} = \frac{1+\frac{1}{3}}{1-\frac{1}{3}} = 2$

支节② $\rho_2 = \frac{Z_L - Z_0}{Z_L + Z_0} = \frac{\infty - Z_0}{\infty + Z_0} = 1$
 $\rho = \rho_2 e^{-2j\beta z} = e^{-2j\beta z}$
 $S = \frac{1+|\rho_2|}{1-|\rho_2|} = \frac{1+1}{1-1} = \infty$

支帯③
$$\rho_2 = \frac{Z_L - Z_0}{Z_L + Z_0} = \frac{\frac{1}{2}Z_0 - Z_0}{\frac{1}{2}Z_0 + Z_0} = -\frac{1}{3} = \frac{1}{3}e^{j\pi}$$

$$\rho = \rho_2 e^{-2j\beta z} = \frac{1}{3}e^{j(\pi - 2\beta z)}$$

$$S = \frac{1 + |\rho_2|}{1 - |\rho_2|} = \frac{1 + \frac{1}{3}}{1 - \frac{1}{3}} = 2$$