习题课: \vec{E} 、U的计算

- 一、 \vec{E} 的计算
- 1. 由定义求
- 2. 由点电荷(或典型电荷分布) \vec{E} 公式 和叠加原理求
- 3. 由高斯定理求
- 4. 由 \vec{E} 与 U的关系求

典型静电场:

点电荷:
$$\vec{E} = \frac{qr}{4\pi\varepsilon_0 r^3}$$

均匀带电圆环轴线上:
$$\vec{E} = \frac{1}{4\pi\varepsilon_0} \frac{qxi}{(R^2 + x^2)^{\frac{3}{2}}}$$

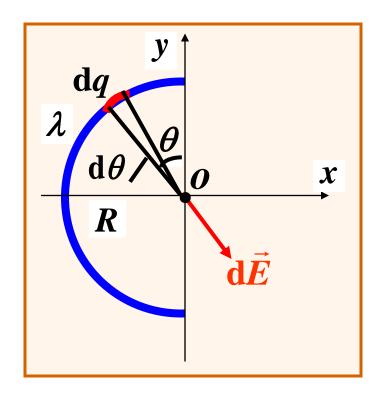
无限长均匀带电直线:
$$E=rac{\lambda}{2\piarepsilon_0 r}$$
 ($oxedsymbol{\perp}$ 带电直线)

均匀带电球面:
$$ec{E}_{
m p}=0$$
 , $ec{E}_{
m h}=rac{qr}{4\piarepsilon_0 r^3}$

无限大均匀带电平面:
$$E=rac{\sigma}{2arepsilon_0}$$
 (丄带电平面)

习题1: 求半径 R 的带电半圆环环心处的电场强度

- 1. 均匀带电,线密度为入
- 2. 上半部带正电,下半部带负电,线密度为礼
- 3. 非均匀带电,线密度为 $\lambda = \lambda_0 \sin \theta$



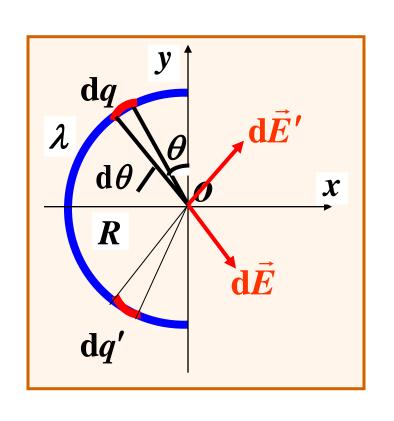
思路: 叠加法

$$dq \rightarrow d\vec{E} \rightarrow \vec{E}$$

解:建立如图所示坐标系

1.
$$dq = \lambda R d\theta$$

$$dE = \frac{dq}{4\pi\varepsilon_0 R^2}; 沿径向$$



用分量积分:

$$\mathrm{d}E_x = \mathrm{d}E \cdot \sin\theta$$

 由对称性: $E_y = \int \mathrm{d}E_y = 0$
 $E = E_x = \int \mathrm{d}E_x = \int \mathrm{d}E \cdot \sin\theta$

$$=\int_{0}^{\pi} \frac{\lambda \sin \theta d\theta}{4\pi \varepsilon_{0} R} = \frac{\lambda}{2\pi \varepsilon_{0} R}$$

$$\therefore \vec{E}_o = E_x \vec{i} + E_y \vec{j} = \frac{\lambda \vec{i}}{2\pi \varepsilon_0 R}$$

$$\begin{array}{c|c}
dq & y \\
\lambda & \theta \\
0 & x \\
\hline
R & d\vec{E}' \\
-\lambda & d\vec{q}'
\end{array}$$

2.
$$dq = \lambda R d\theta$$
 $dE = \frac{\lambda R d\theta}{4\pi\varepsilon_0 R^2}$;沿径向
 $dq' = -\lambda R d\theta$
 $dE' = \frac{-\lambda R d\theta}{4\pi\varepsilon_0 R^2}$;沿径向
由对称性: $E_x = \int dE_x = 0$

$$dE_y = -dE \cdot \cos \theta$$

$$E_o = E_y = \int dE_y = 2 \int_0^{\pi/2} dE_y = 2 \int_0^{\pi/2} -\frac{\lambda \cos \theta d\theta}{4\pi \varepsilon_0 R} = -\frac{\lambda}{2\pi \varepsilon_0 R}$$

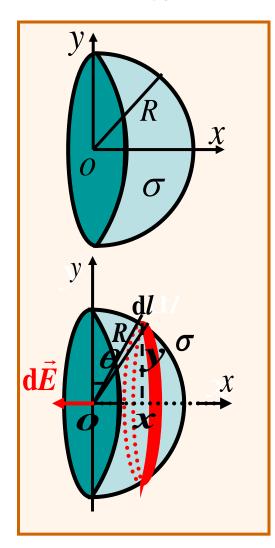
$$\therefore \vec{E}_o = E_x \vec{i} + E_y \vec{j} = -\frac{\lambda \vec{j}}{2\pi \varepsilon_0 R}$$

$$\begin{array}{c|c}
dq & y \\
\lambda & \theta \\
d\vec{E}' \\
R & d\vec{E} \\
dq'
\end{array}$$

3.
$$dq = \lambda R d\theta = \lambda_0 \sin \theta R d\theta$$
 $dE = \frac{dq}{4\pi\varepsilon_0 R^2}$; 沿径向
 $dq' = \lambda_0 \sin(\pi - \theta) R d\theta$
 $= \lambda_0 \sin \theta R d\theta = dq$

$$\begin{split} \mathrm{d}E_x &= \mathrm{d}E \cdot \sin\theta \quad y \, \bar{\beta} \, \mathrm{向} \, \underline{\beta} \, \bar{\eta} \, \bar{\eta} \, \underline{\kappa} \, \underline{k} \, ! \quad E_y = \int \mathrm{d}E_y = 0 \\ E &= E_x = \int \mathrm{d}E_x^{-1} \int_0^\pi \frac{\lambda_0 \sin^2\theta \mathrm{d}\theta}{4\pi\varepsilon_0 R} = \frac{\lambda}{8\varepsilon_0 R} \\ & \therefore \vec{E}_o = E_x \vec{i} + E_y \vec{j} = \frac{\lambda \vec{i}}{8\varepsilon_0 R} \end{split}$$

习题2(P_{235} 9.6): 求均匀带电半球面(已知R, σ) 球心处电场。



解:将半球面视为由许多圆环拼成,任取一个圆环:

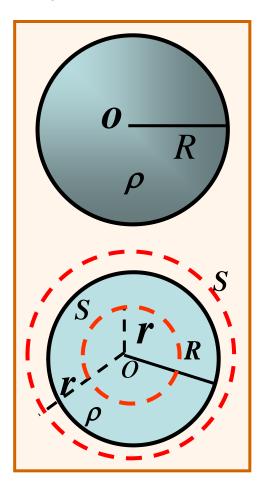
$$dq = \sigma \cdot 2\pi y dl = \sigma \cdot 2\pi R \cos\theta \cdot R d\theta$$

$$dE = \frac{x dq}{4\pi \varepsilon_0 (y^2 + x^2)^{\frac{3}{2}}}$$

$$= \frac{R\sin\theta dq}{4\pi\varepsilon_0 R^3} = \frac{\sigma\cos\theta\sin\theta}{2\varepsilon_0} d\theta^{ - 3\pi - 3\pi } \dot{\beta}$$

因为各圆环在o点处 $d\vec{E}$ 同向,可直接积分。

习题3(P_{236} 9.14): 求半径为R,电荷体密度 $\rho = k/r$ (k为常数, $r \le R$)的带电球体内外的场强。



思考: 选用哪种方法求解更方便?

 $\rho = k/r$ 未破坏电场分布的球对称性。用高斯定理求解方便。

解:高斯面为:半径为 r的同心球面S

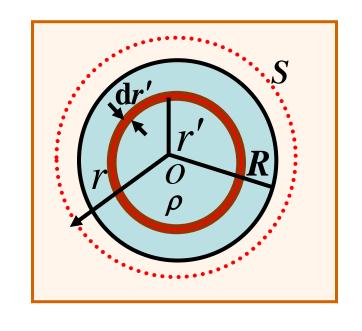
$$\oint_{S} \vec{E} \cdot d\vec{S} = E \cdot 4\pi r^{2}$$

$$\sum q_{\beta} = \rho \cdot V = \frac{\kappa}{r} \cdot \frac{4}{3} \pi r^3$$

对否?

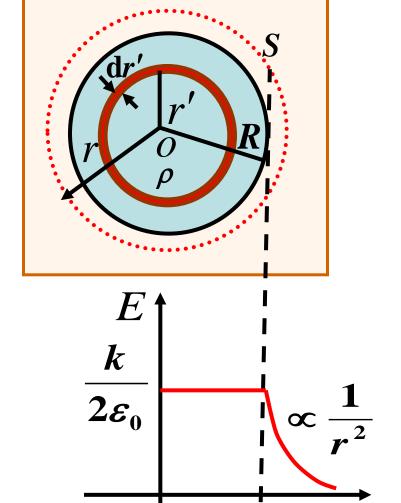


$$dq = \rho dV = \frac{k}{r'} \cdot 4\pi r'^2 dr'$$



$$r > R$$
: $\sum q_{|\gamma|} = \int dq = \int_0^R \frac{k}{r'} \cdot 4\pi r'^2 dr' = 2\pi k R^2$

$$r < R: \qquad \sum q_{\mid h \mid} = \int \mathrm{d}q = \int \frac{k}{r'} \cdot 4\pi r'^2 \mathrm{d}r' = 2\pi k r^2$$



由高斯定理得:

$$E \cdot 4\pi r^2 = \frac{1}{\varepsilon_0} \sum q_{\mid \gamma \mid}$$

$$\therefore E_{\bowtie} = \frac{2\pi kr^2}{4\pi\varepsilon_0 r^2} = \frac{k}{2\varepsilon_0}$$

$$E_{\text{sh}} = \frac{2\pi kR^2}{4\pi\varepsilon_0 r^2} = \frac{kR^2}{2\varepsilon_0 r^2}$$
 沿径向

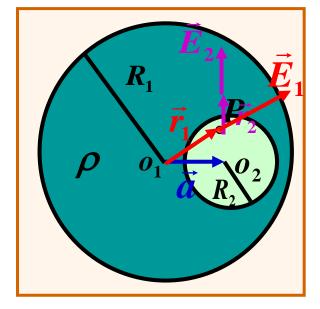
一、产的计算

习题4(P_{236} 9.13):在半径 R_1 ,体电荷密度 ρ 的均匀带电球体内挖去一个半径 R_2 的球形空腔。空腔中心 o_2 与带电

球体中心 o_1 相距为 $a[(R_2+a) < R_1]$,求空腔内任一点电场。

思考: 选用何种方法求解?

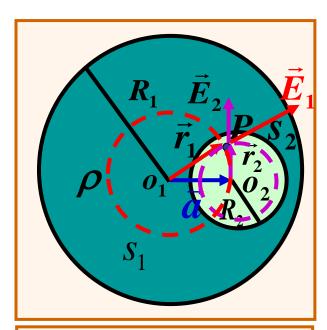
挖去空腔 —— 失去球对称性,用补偿法恢复对称性!



设:

半径为 R_1 均匀带电实心球体在P点的场强: \vec{E}_1 \vec{E}_1 、 \vec{E}_2 半径为 R_2 均匀带电实心球体在P点的场强: \vec{E}_2 均可由高斯定理求出

所求场强 $\vec{E}_P = \vec{E}_1 - \vec{E}_2$



$$E_1 \cdot 4\pi r_1^2 = \frac{1}{\varepsilon_0} \rho \cdot \frac{4}{3}\pi r_1^3 \qquad \vec{E}_1 = \frac{\rho \vec{r}_1}{3\varepsilon_0}$$

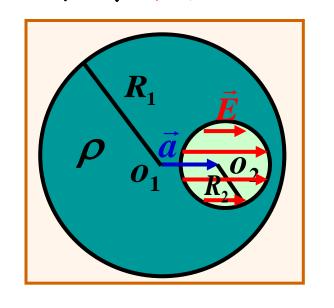
$$E_2 \cdot 4\pi r_2^2 = \frac{1}{\varepsilon_0} \rho \cdot \frac{4}{3}\pi r_2^3 \quad \vec{E}_2 = \frac{\rho \vec{r}_2}{3\varepsilon_0}$$

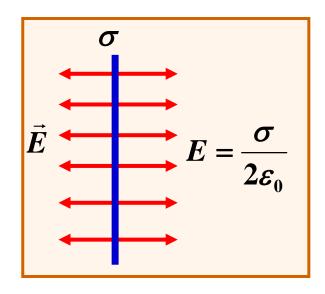
$$\begin{array}{c|c}
R_1 & \vec{E} \\
\hline
\rho & \vec{a} & \vec{R_2} \\
\hline
\end{array}$$

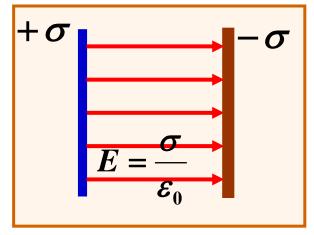
$$\vec{E}_P = \vec{E}_1 - \vec{E}_2 = \frac{\rho}{3\varepsilon_0} (\vec{r}_1 - \vec{r}_2) = \frac{\rho a}{3\varepsilon_0}$$

腔内为平行于
$$\overrightarrow{o_1o_2} = \overrightarrow{a}$$
 的均匀电场!

请总结获得均匀电场的方法:







••••

= U =

注意

1. 场强积分法
$$U_a = \int\limits_a^{\mathop{\approx} h \atop a} \vec{E} \cdot \mathbf{d} \vec{l}$$

适合:场强已知或 容易求到的情况.

- (1) 积分与路径无关,可依题意选最简便的积分路径。
- (2) \vec{E} 为路径上各点总场,若各区域 \vec{E} 表达式不同, 应分段积分.
- (3) 积分值与零势点选取有关 。选取原则: 电荷有限分布选 $U_{\scriptscriptstyle \infty}$ =0; 电荷无限分布选 $U_{\scriptscriptstyle ARM}$ =0

2. 叠加法 思路: $dq \rightarrow dU \rightarrow U = \int dU$

适合:场强不知或不容易求到的情况.

注意:应用典型带电体的电势公式时需选取相同的零势点 典型带电体的电势:

点电荷:

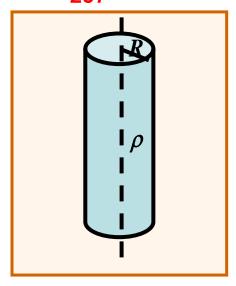
$$U = \frac{q}{4\pi\varepsilon_0 r}$$

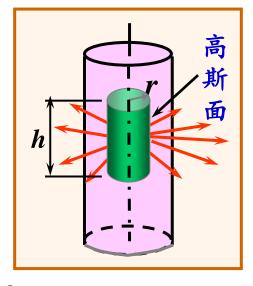
均匀带电圆环轴线上: $U = \frac{q}{4\pi\varepsilon_0(R^2 + x^2)^{1/2}}$

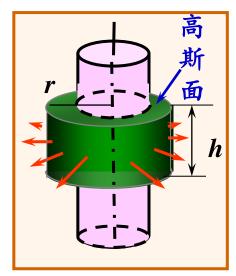
均匀带电球面:

$$U_{\text{Pl}} = \frac{q}{4\pi\varepsilon_0 R} \qquad U_{\text{Pl}} = \frac{q}{4\pi\varepsilon_0 r}$$

习题 $5(P_{237}, 9.19)$:求无限长均匀带电圆柱体 (R, ρ) 电势分布





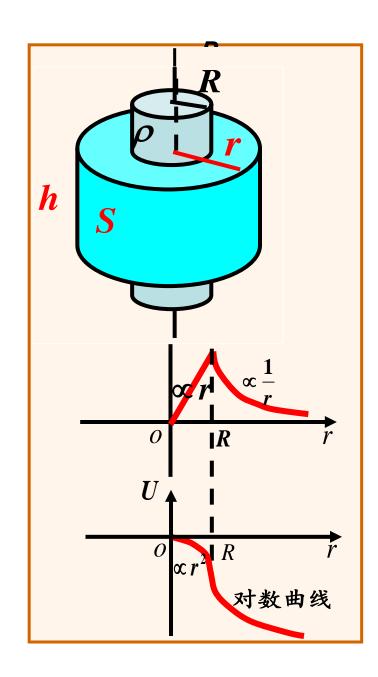


解:场强积分法:
$$\oint_s \vec{E} \cdot d\vec{S} = E \cdot 2\pi rh = \frac{1}{\varepsilon_0} \sum q_{\text{Pl}}$$

$$r \leq R: \qquad \sum q_{\mid j \mid} = \rho \cdot \pi r^2 h \qquad \vec{E}_{\mid j \mid} = \frac{\rho \vec{r}}{2\varepsilon_0}$$

$$r \ge R$$
:
$$\sum q_{\bowtie} = \rho \cdot \pi R^2 h \qquad \vec{E}_{\bowtie} = \frac{\rho R^2 \vec{r}}{2\varepsilon_0 r^2}$$

二、U的计算



令
$$r = 0$$
处 $U = 0$,沿径向积分。
$$U_{\text{內}} = \int_{r}^{0} \vec{E}_{\text{內}} \cdot d\vec{r} = \int_{r}^{0} \frac{\rho \vec{r} \cdot d\vec{r}}{2\varepsilon_{0}}$$

$$= \frac{\rho}{2\varepsilon_{0}} \int_{r}^{0} r dr = -\frac{\rho r^{2}}{4\varepsilon_{0}}$$

$$U_{\text{內}} = \int_{r}^{R} \vec{E}_{\text{內}} \cdot d\vec{r} + \int_{R}^{0} \vec{E}_{\text{內}} \cdot d\vec{r}$$

$$= \int_{r}^{R} \frac{\rho R^{2} dr}{2\varepsilon_{0} r} + \int_{R}^{0} \frac{\rho r dr}{2\varepsilon_{0}}$$

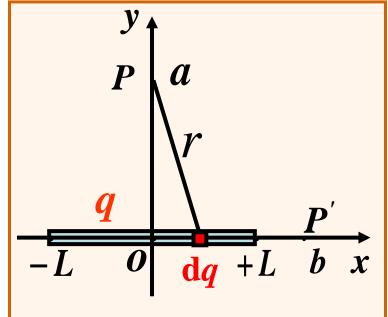
$$= \frac{\rho R^{2}}{2\varepsilon_{0}} \ln \frac{R}{r} - \frac{\rho R^{2}}{4\varepsilon_{0}}$$

$$= U_{\text{內}} + \frac{\rho R^{2}}{2\varepsilon_{0}} \ln \frac{R}{r} - \frac{\rho R^{2}}{4\varepsilon_{0}}$$

习题 $6(P_{237}, 9.23)$:电量q均匀分布在长为2L的细棒上。

求:(1)细棒中垂面上距细棒中心a处P点的电势。

(2) 细棒延长线上距细棒中心b处P 点的电势。



解:将带电细棒视为点电荷集合用叠加法。建立如图所示坐标系:

$$dq = \frac{q}{2L} dx \quad \diamondsuit U_{\infty} = 0$$

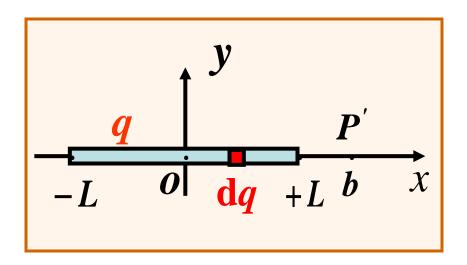
$$\frac{P'}{dq + L b x} \qquad (1) dU = \frac{dq}{4\pi\varepsilon_0 r} = \frac{qdx}{8\pi\varepsilon_0 L(x^2 + a^2)^{\frac{1}{2}}}$$

$$U_{P} = \int dU = \int_{-L}^{L} \frac{q dx}{8\pi \varepsilon_{0} L(x^{2} + a^{2})^{\frac{1}{2}}} = \frac{q}{4\pi \varepsilon_{0} L} \ln \frac{L + \sqrt{a^{2} + L^{2}}}{a}$$

二、U的计算

(2)求细棒延长线上距细棒中心b处P/点的电势

$$dU = \frac{dq}{4\pi\varepsilon_0(b-x)}$$
$$= \frac{qdx}{8\pi\varepsilon_0 L(b-x)}$$



$$U_{P'} = \int dU = \int_{-L}^{L} \frac{q dx}{8\pi \varepsilon_0 L(b-x)}$$

$$= \frac{q}{8\pi\varepsilon_0 L} \ln \frac{b+L}{b-L}$$

作业

- 1. No.7(希望在作业题纸中选择、填空、判断各题的相应位置处写出其关键步骤);
- 2. 自学本章各例题并完成书上的习题(对照书后的参考答案自己订正)。

第十二周星期三交作业

