



# A relook at the GZK effect

**Greisen, PRL (1966), Zatsepin and Kuzmin, JETP Letters (1966)**

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Sheeba Shafaq (2022)



Samiran Roy (2020)

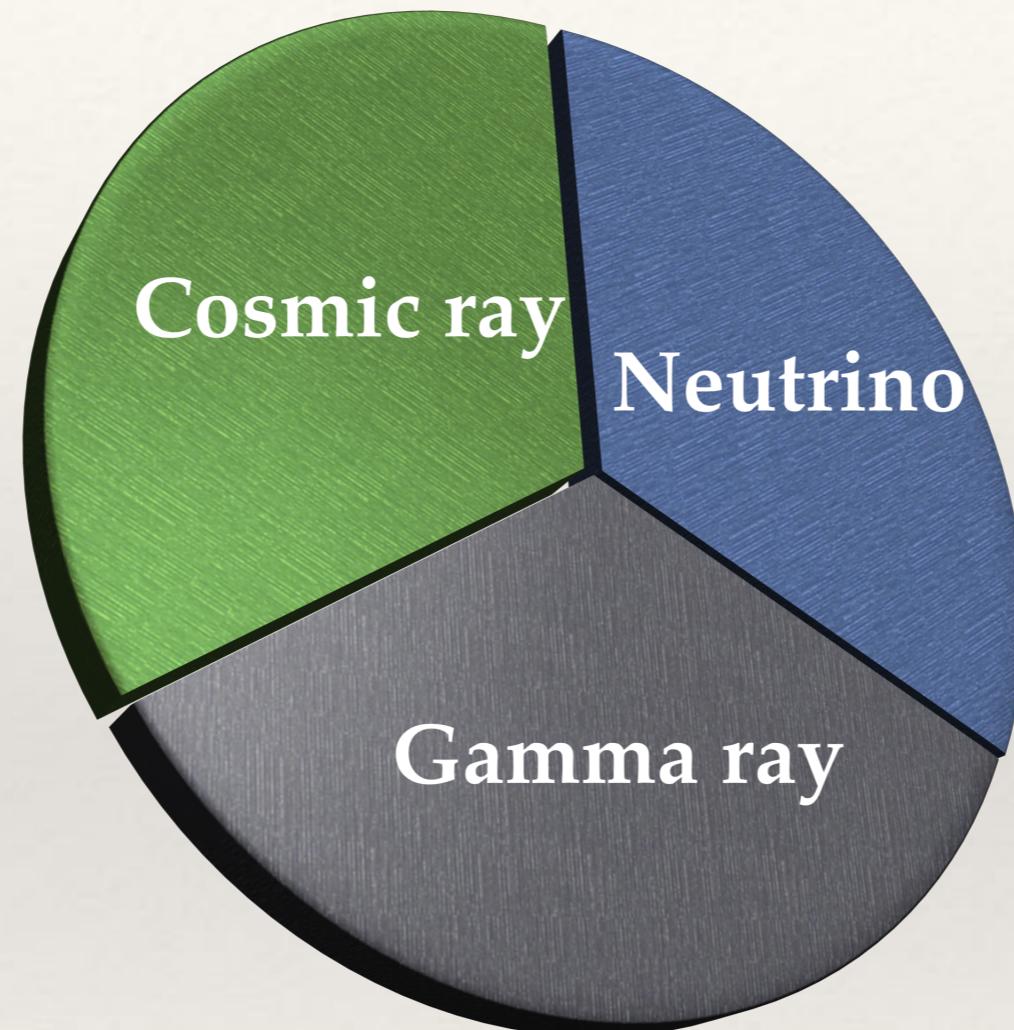


JNU-HBNI MoU

Jogesh Rout (2020)



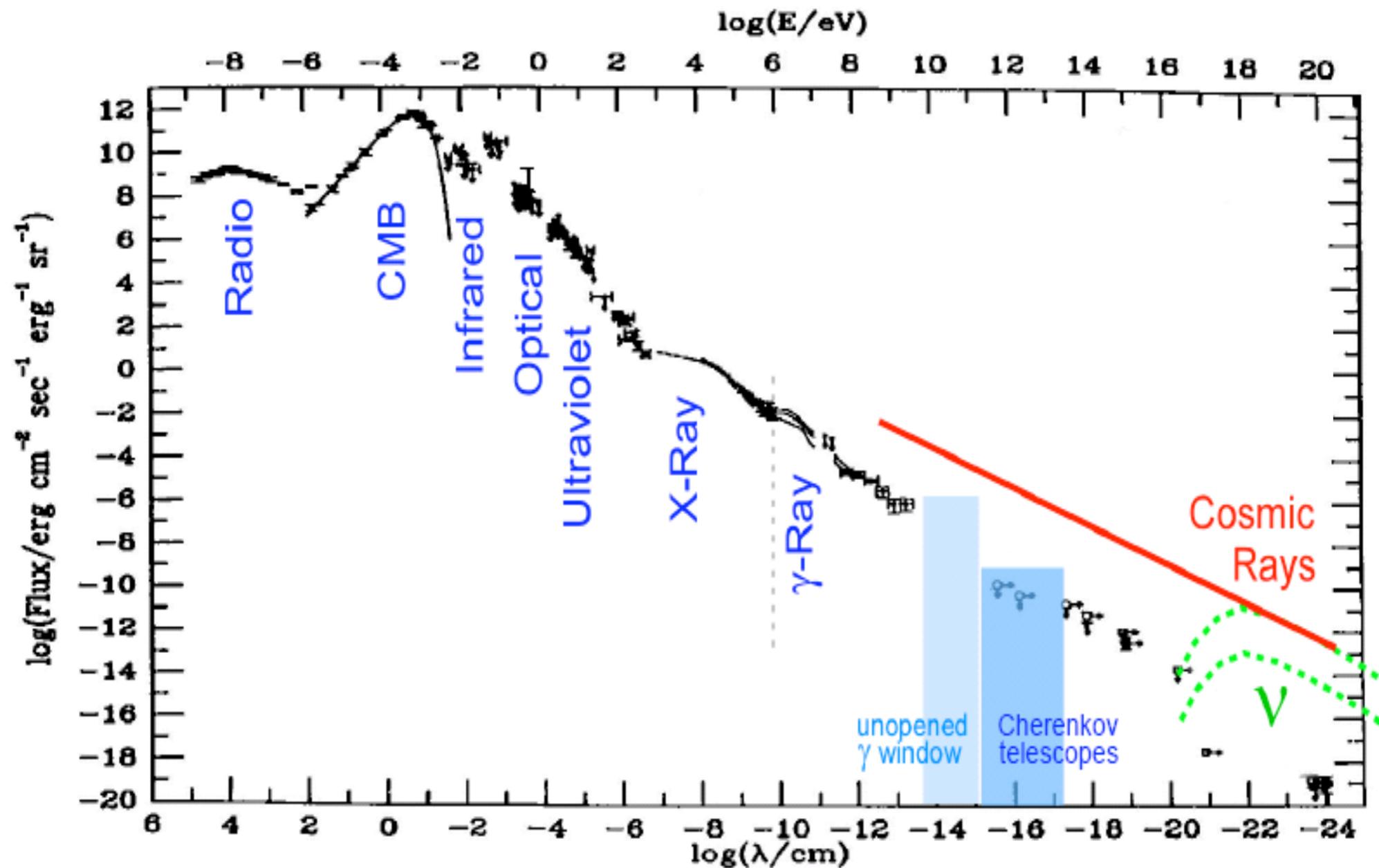
# The three messengers at extreme energies



work with Prantik Sarmah and Sovan Chakraborty (IIT Guwahati),

JCAP 01, 058 (2024)

We cannot see the deep universe at energies  $>$  few TeV, since photons are attenuated through  $\gamma\gamma \rightarrow e^+e^-$  on cosmic radiation backgrounds

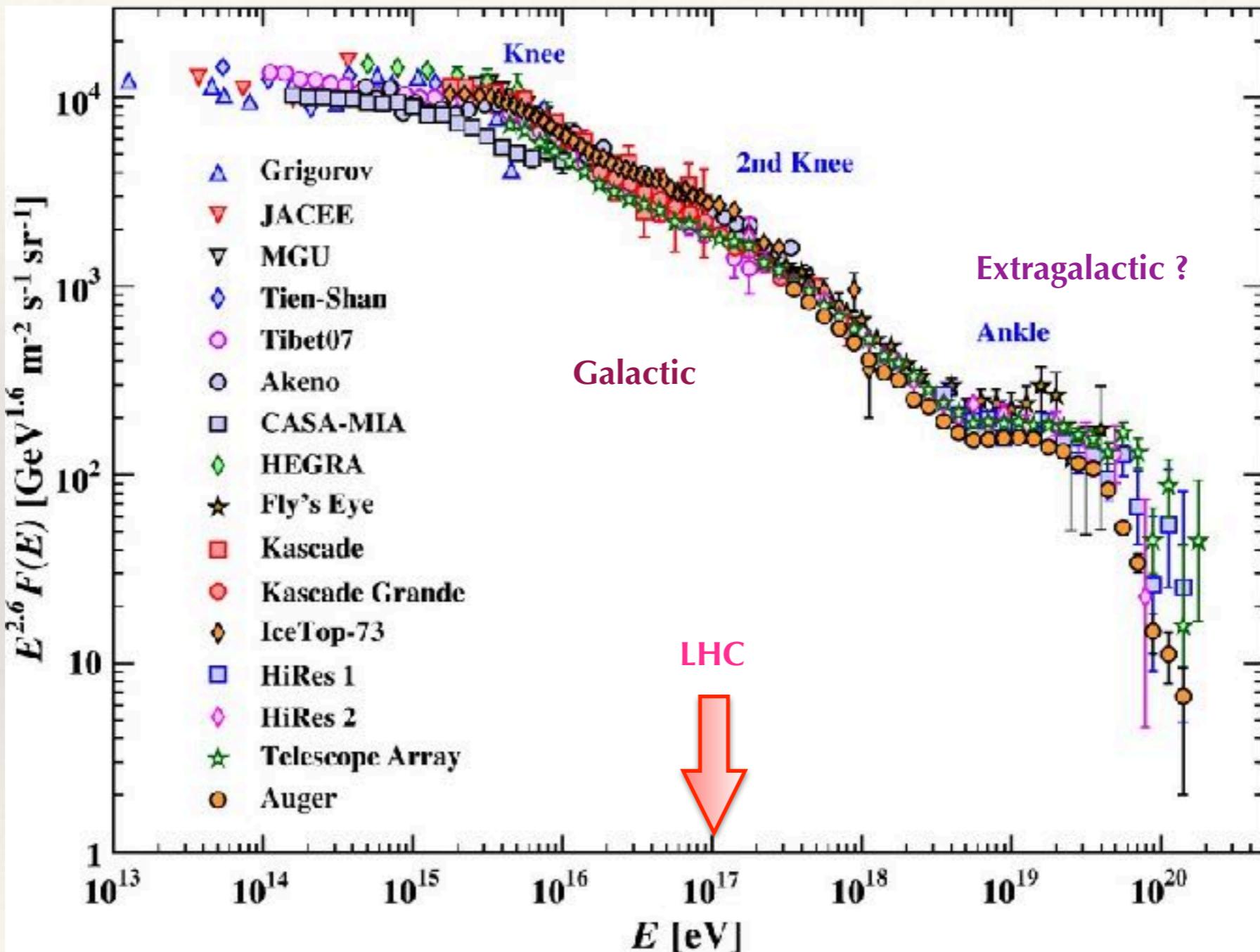


But using cosmic rays we can ‘see’ up to  $\sim 6 \times 10^{10}$  GeV before they get attenuated through photopion interactions on the CMB

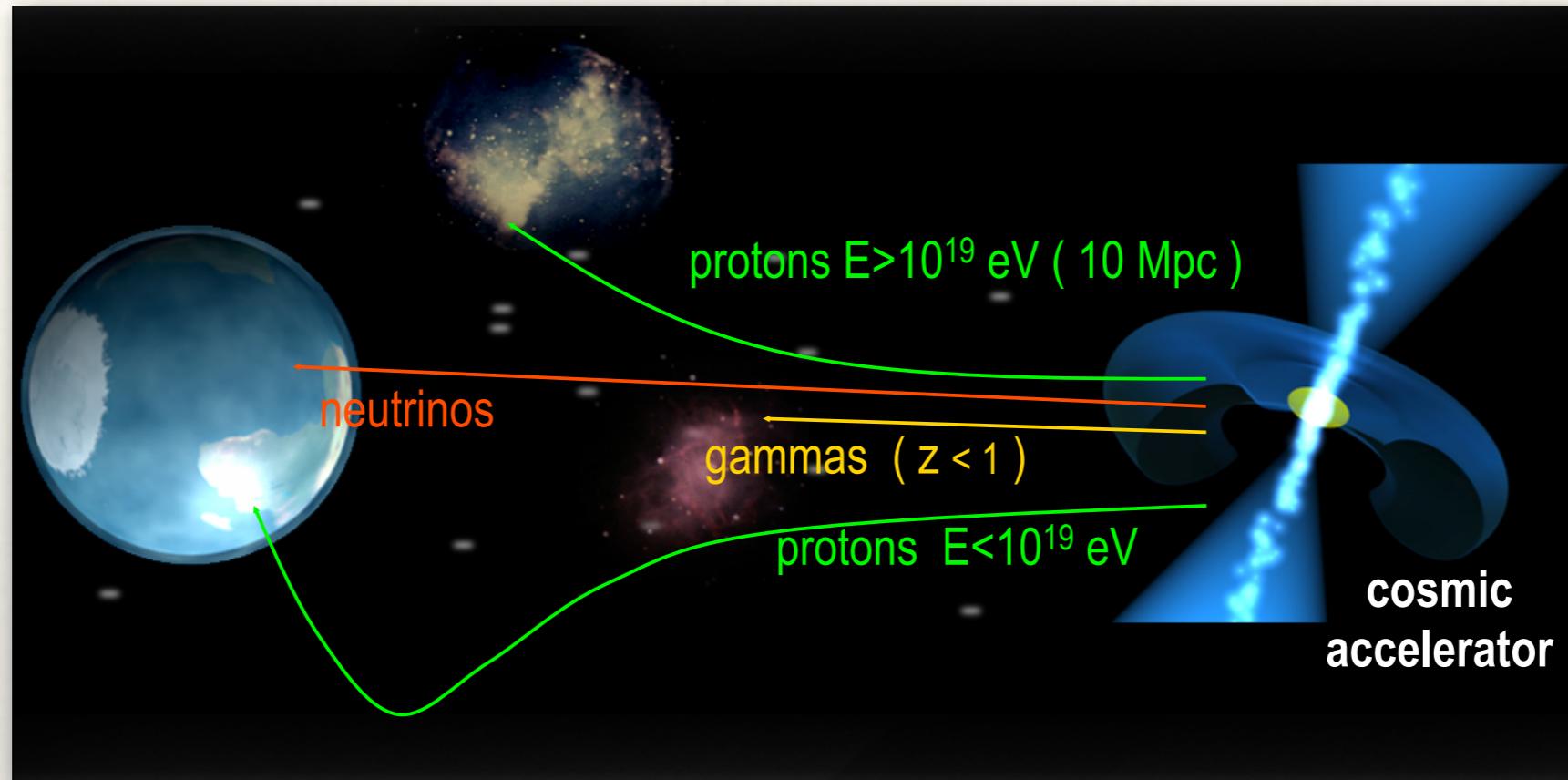
... and the universe is ~transparent to neutrinos at effectively all energies

# Cosmic ray spectrum

Courtesy : Subir Sarkar



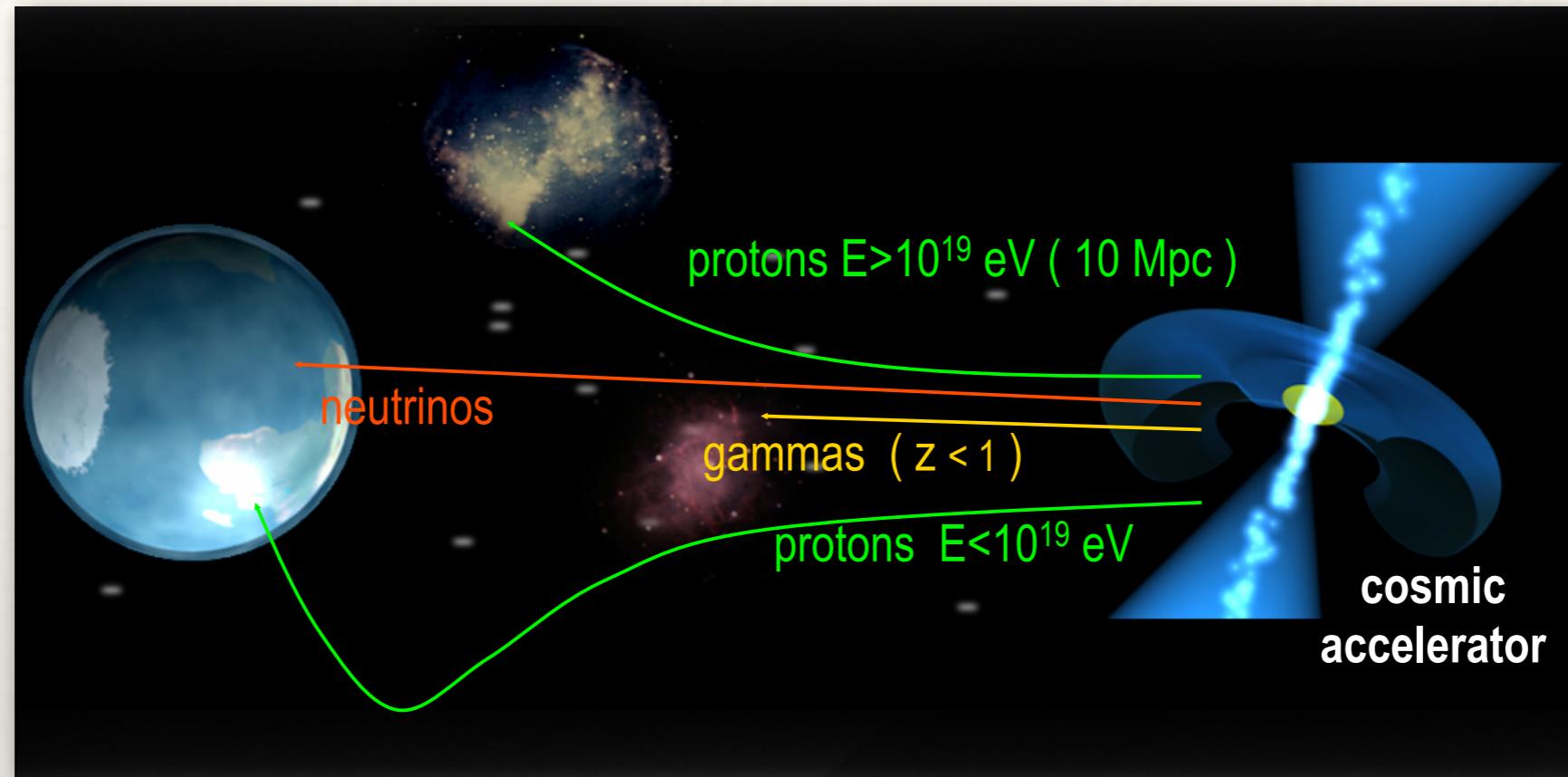
# Cosmic accelerators



- protons/nuclei: deflected by magnetic fields, absorbed on radiation (GZK)
- photons: absorbed on radiation/dust; reprocessed at source
- neutrinos: neither absorbed nor bent, straight path from source

# Cosmic accelerators

Courtesy : Montaruli, SSI 2010



*Neutrinos : can reliably lead to the discovery of such point sources*

$$\gamma + \gamma_{2.7K}$$

$$l_\gamma = \frac{1}{\sigma_{p-\gamma_{2.7K}} \times n_\gamma} \sim \frac{1}{5 \times 10^{-28} \text{ cm}^2 \times 400 \text{ cm}^{-3}} = 10 \text{ Mpc}$$

$$\nu + \nu_{1.95K} \rightarrow Z + X$$

$$l_\nu = \frac{1}{\sigma_{res} \times n} = \frac{1}{5 \times 10^{-31} \text{ cm}^2 \times 112 \text{ cm}^{-3}} = 6 \text{ Gpc}$$

# A typical cosmic accelerator

## Pion photoproduction

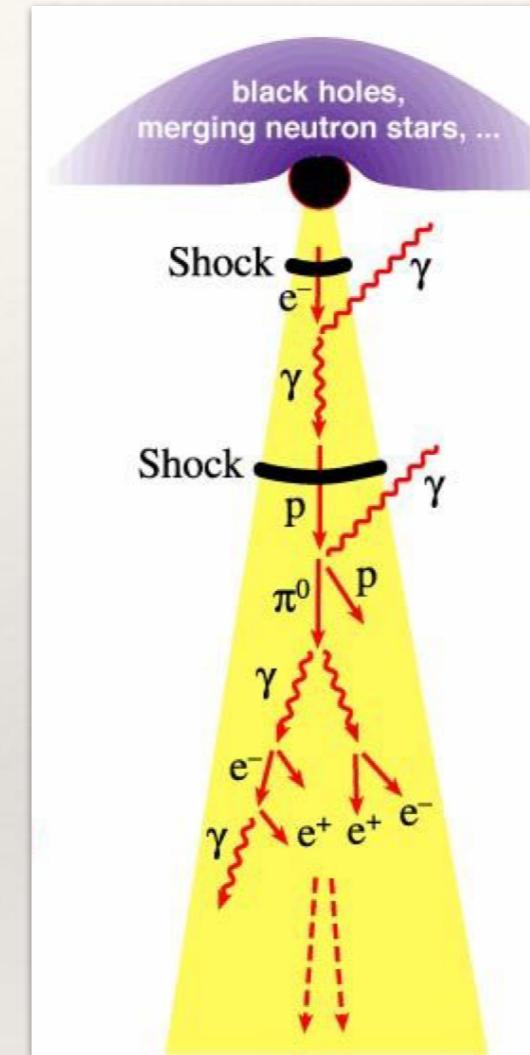


## Weak decays

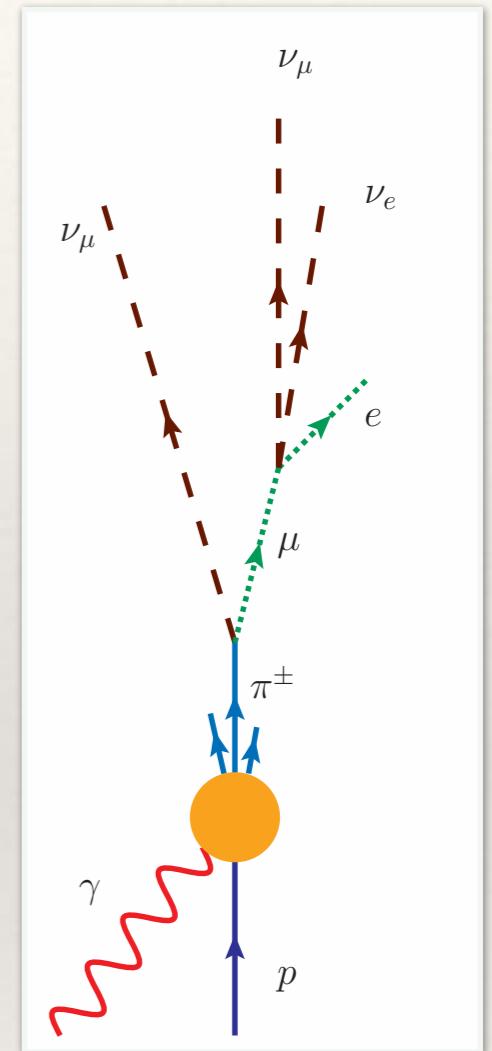


1:2:0

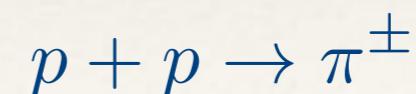
If  $n$  exists the source



Ref: Halzen, ICRC'07



## Hadron-hadron interactions



# A typical cosmic accelerator

## Pion photoproduction

$$p + \gamma \rightarrow \Delta_{1232} \rightarrow p + \pi^0 \quad \text{BR}=2/3$$

$$p + \gamma \rightarrow \Delta_{1232} \rightarrow n + \pi^+ \quad \text{BR}=1/3$$

## Weak decays

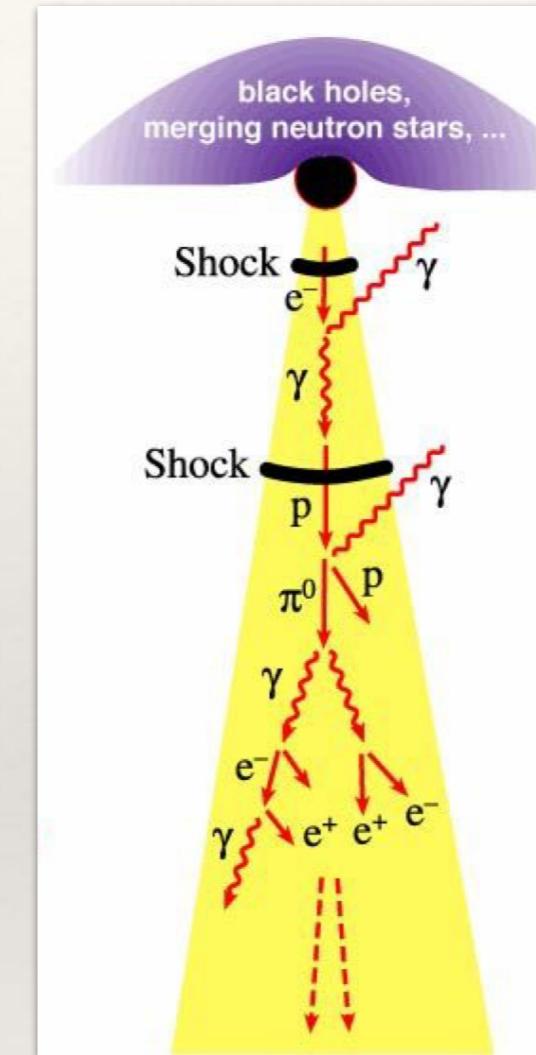
$$\pi^+ \rightarrow \mu^+ + \nu_\mu$$

$$\mu^+ \rightarrow e^+ + \nu_e + \bar{\nu}_\mu$$

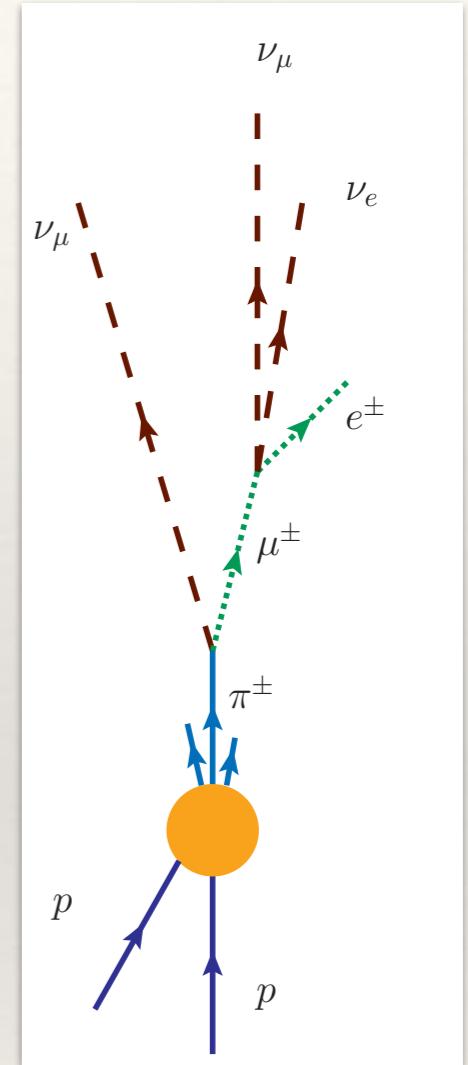
1:2:0

If  $n$  exists the source

$$n \rightarrow p + e^- + \bar{\nu}_e$$



Ref: Halzen, ICRC'07



## Hadron-hadron interactions

$$p + p \rightarrow \pi^\pm$$

# Flavor composition

## Classification of sources

Pakvasa, MPLA23, 1313 (2008) , talk@Nusky2011 A relativistic jet [credits : Zhang and Woosley]

Source Candidates : AGN, GRB, SNR, MQ

### 1. Conventional pion beam source :

Cosmic rays (assumed to be  $p$ ) interact with  $\gamma$  or  $p$  and produce charged mesons (pions, kaons) which decay to give neutrinos via the  $\pi \rightarrow \mu \rightarrow e$  decay chain :

$$p\gamma : p + \gamma \rightarrow \pi^\pm + \text{all} \quad ; \quad pp : p + p \rightarrow \pi^\pm + \text{all}$$

### 2. Damped muon source : Muons lose energy prior to decay (depends on $E$ )

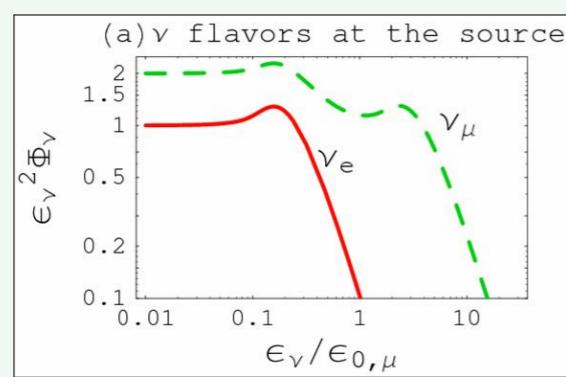
### 3. Muon decay : Muons from damped muon source decay at lower $E$

### 4. Prompt : Decay of short-lived heavy flavors (pions interact and do not decay)

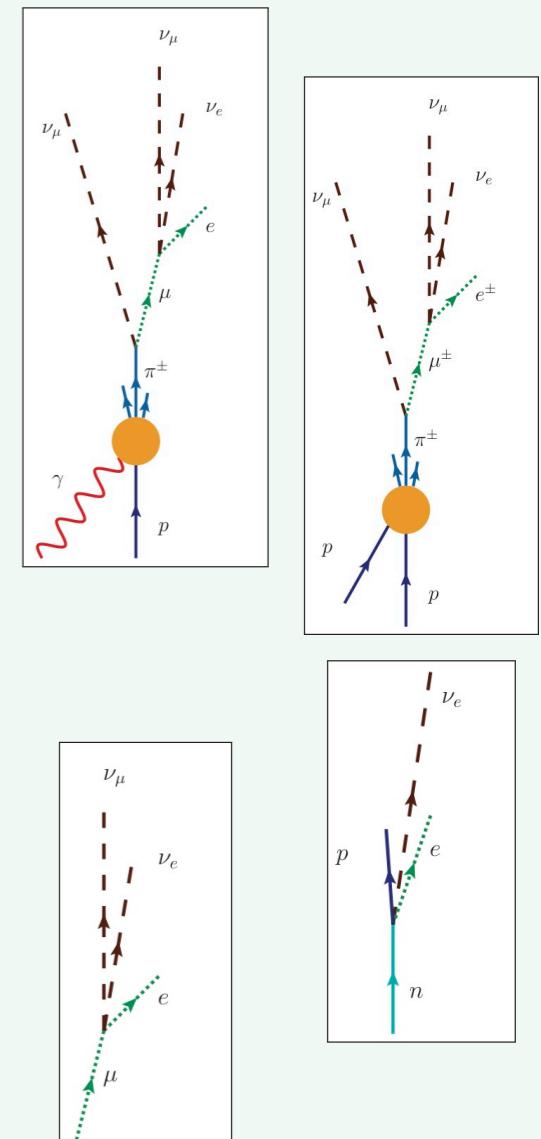
### 5. Neutron decay : Photodisintegration of heavy nuclei

- ▶ Source type can be characterized by  $\hat{X} = \Phi_e^0 / \Phi_\mu^0$  (since  $\Phi_\tau^0$  is negligible)
- ▶ However same source can mimic different source types as a function of  $E$  i.e.  $\hat{X}(E)$

Kashti and Waxman, PRL95, 181101 (2005)



Source	$\hat{X}$	$\Phi_e^0 : \Phi_\mu^0 : \Phi_\tau^0$
Pion beam	0.5	1:2:0
Neutron decay	$\gg 1$	1:0:0
Muon decay/Prompt	1	1:1:0
Damped muon	0	0:1:0



# Astrophysical parameter space

- Hillas criterion for acceleration and confinement :

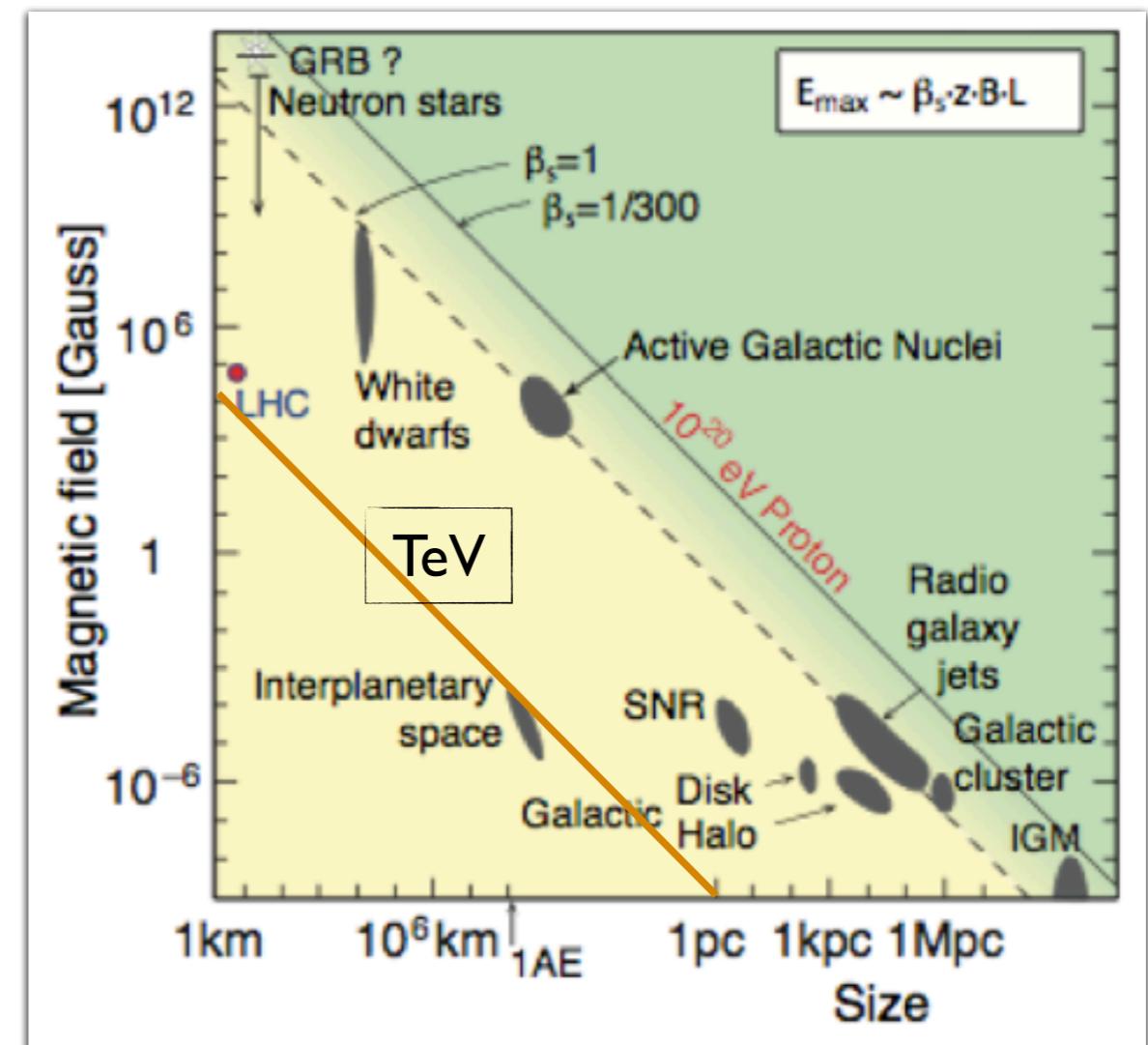
$$E \leq E_{max} = qBR$$

$$r_L < R$$

Larmor radius

size of accelerator

Ref: Hillas (1984), Boratav (2001), Hummer et al. (2010)



- constraint on B and R

- Call sources as “test points” in order to discuss E-dependent effects at source

# The GZK mechanism

Pion photoproduction

$$p + \gamma_{CMB} \rightarrow \Delta_{1232} \rightarrow n + \pi^+ (E_{\text{th}} \sim 5 \times 10^{19} \text{ eV})$$

$$\pi^+ \rightarrow \mu^+ + \nu_\mu$$

$$\mu^+ \rightarrow e^+ + \nu_e + \bar{\nu}_\mu$$

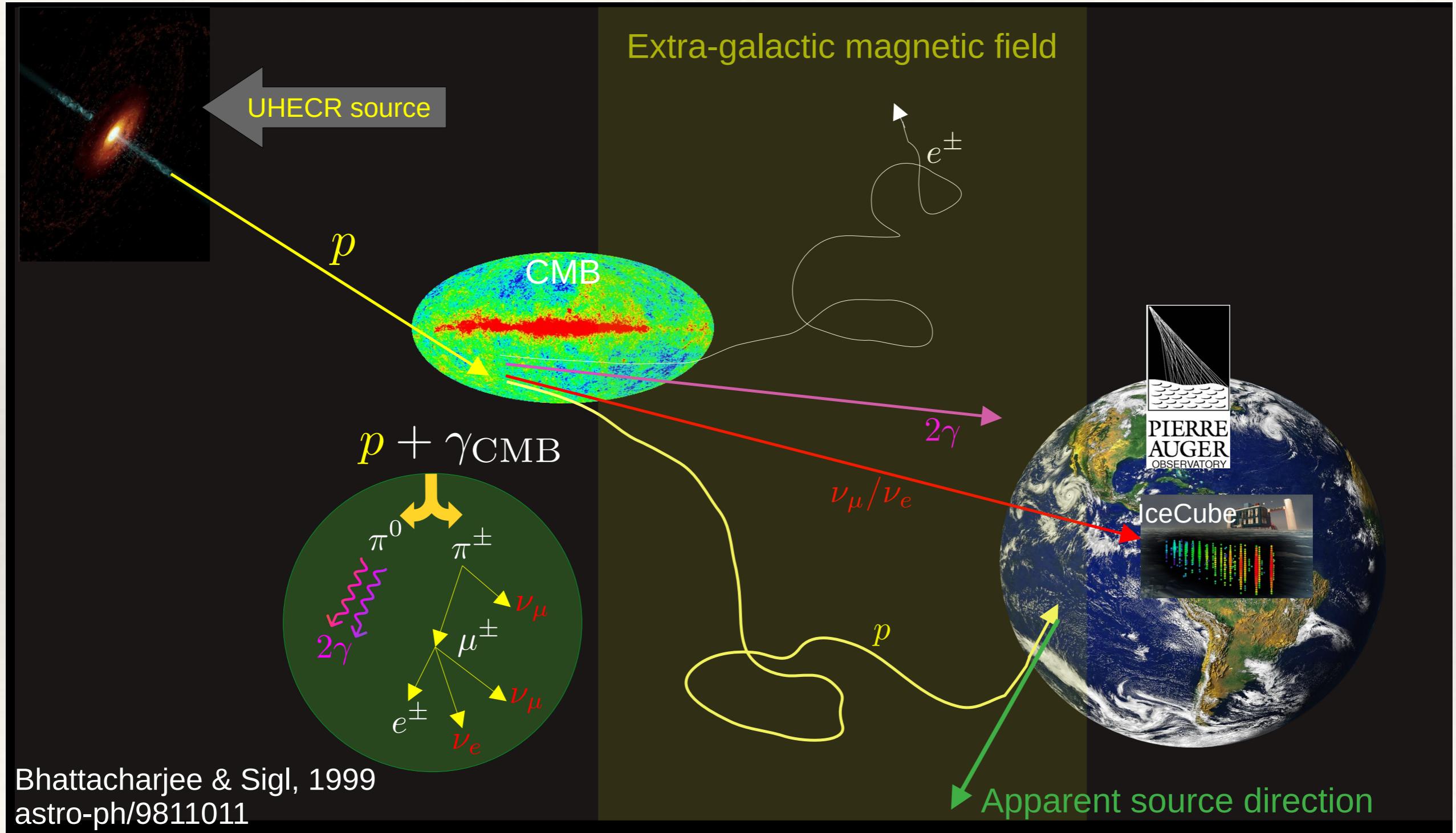
$$n \rightarrow p + e^- + \bar{\nu}_e$$

Predicts that the spectrum should be cut-off...

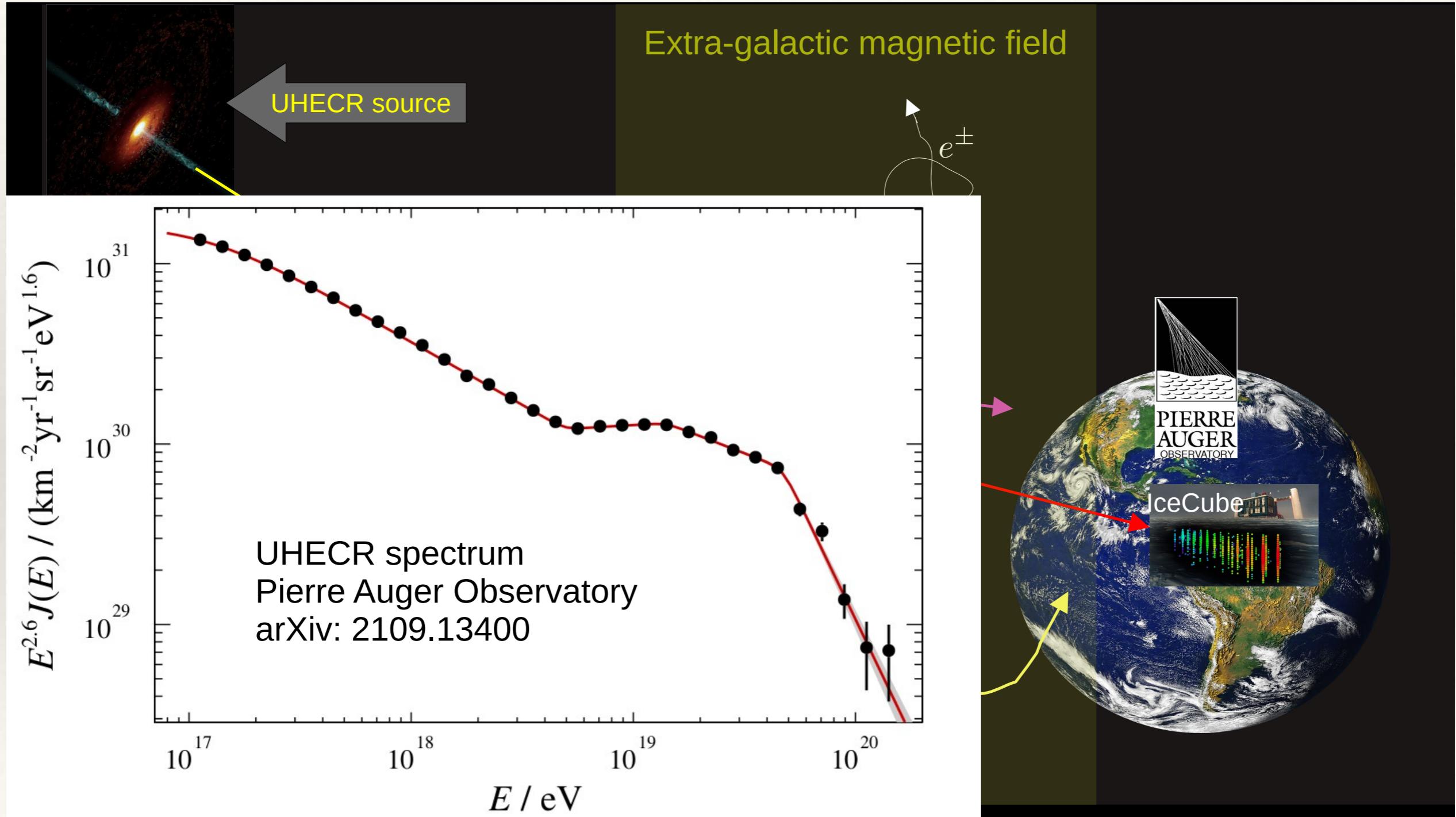
Telescope Array (SD) have seen a recent event at 244 EeV, so naturally questions arise about its origin...

R. U. Abbasi et al. (Telescope Array), Science 382, abo5095 (2023), arXiv:2311.14231 [astro-ph.HE].  
 (See also P. Sarmah, N. Das, D. Borah, S Chakraborty and P. Mehta arXiv:2406.03174 [hep-ph])

# UHECR propagation & GZK process

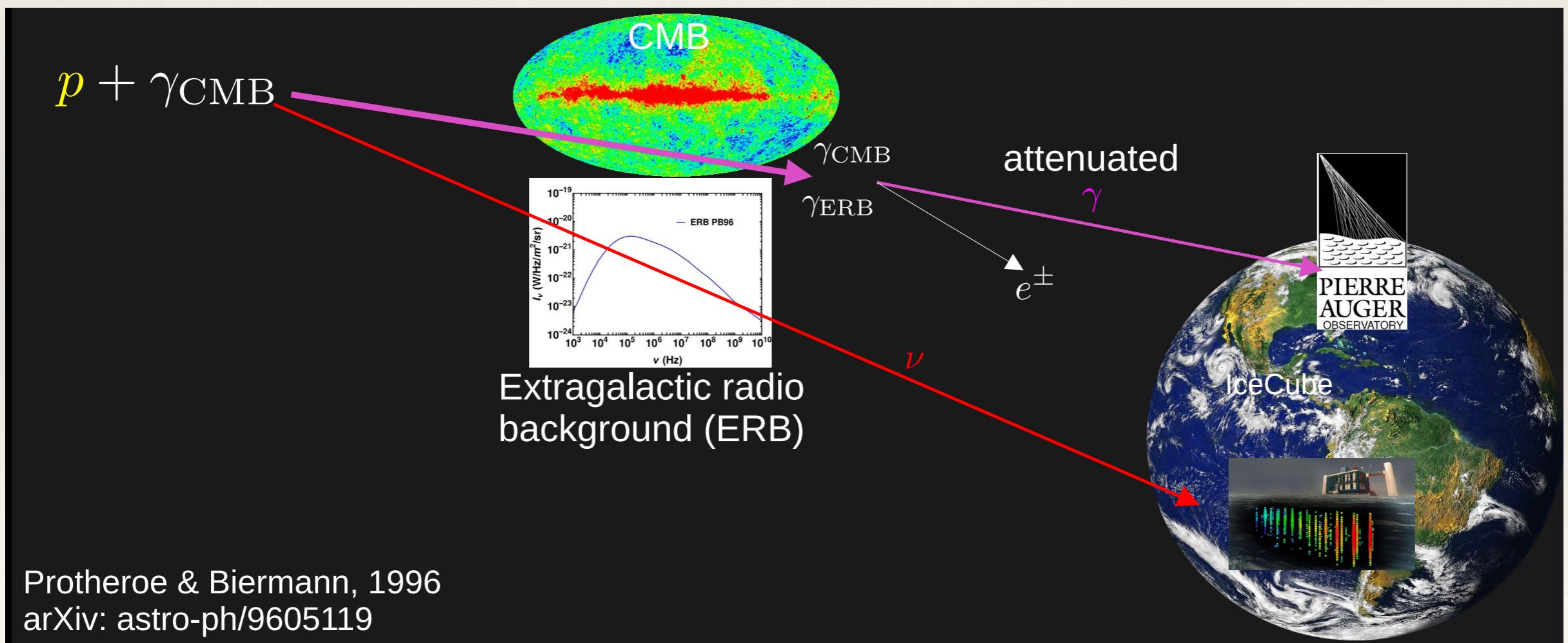


# The GZK process



# The GZK photon flux

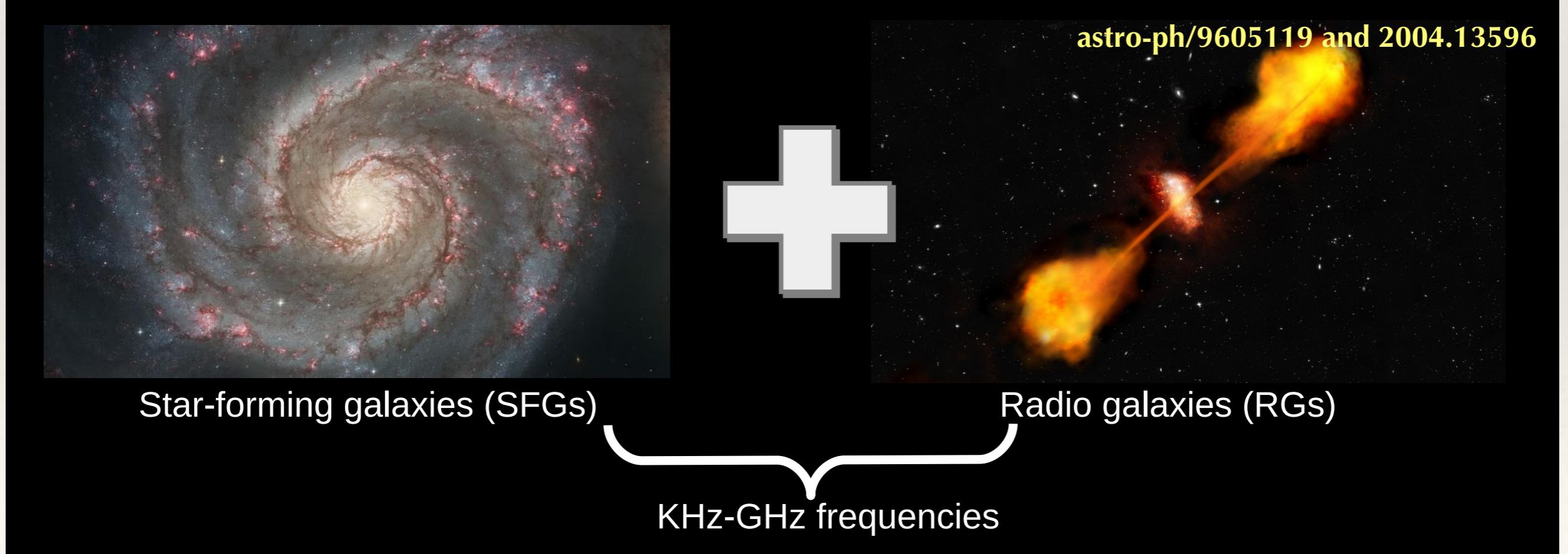
- Also affected by inverse Compton, double pair production, em cascades
- Need precise estimates of CMB and ERB
- While CMB spectrum is well measured, the ERB has large uncertainties



Protheroe & Biermann, 1996  
arXiv: astro-ph/9605119

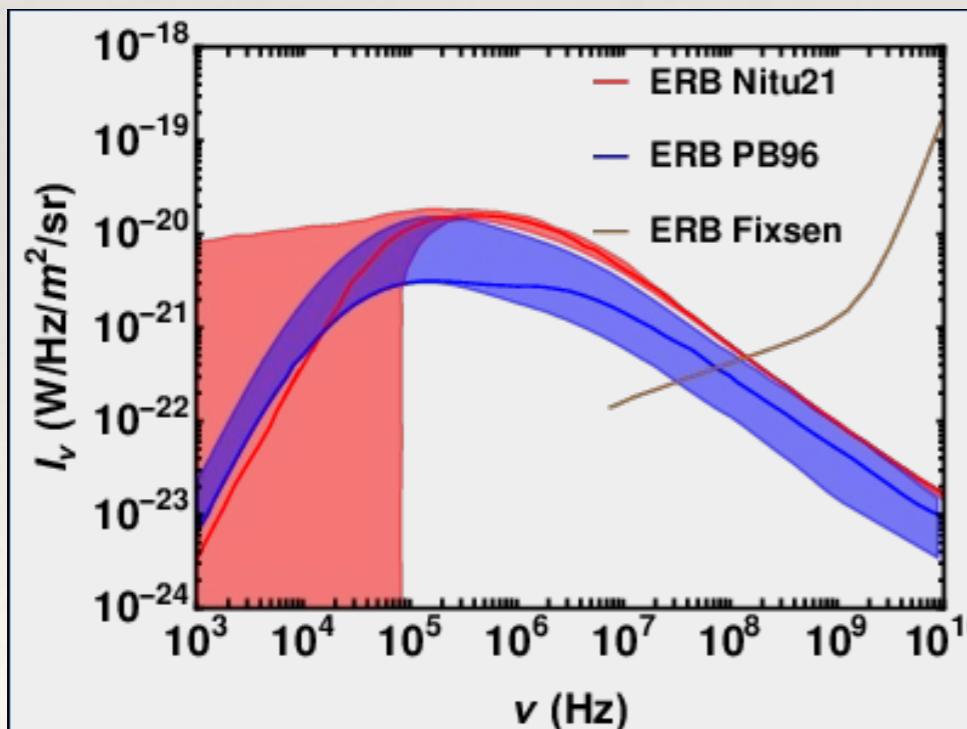
# Extragalactic Radio Background

- Free-free emission in SFGs & synchrotron radiation in RGs
- SFG contribution dominates at low frequencies (below  $\sim 100$  MHz)
- RG dominates at higher frequencies
- 2 Models : Protheroe and Biermann 1996 (PB96) & Nitu et al 2021 (Nitu21)



# Extragalactic Radio Background

- The ARCADE2 (Absolute Radiometer for Cosmology, Astrophysics, and Diffuse Emission) measured radiometric temperature at 3-90 GHz.
- Analysis (Fixsen et. al. 2011) showed an excess of temperature of about  $24.1 (\nu/310 \text{ MHz})^{-2.59}$  K over the CMB temperature  $\sim 2.73$  K
- May be due to unresolved radio sources or simplistic geometry adopted for galactic radio emission.



[astro-ph/9605119](#), [2004.13596](#), [0901.0555](#)

These uncertainties can  
impact the GZK photon flux  
prediction

# GZK flux & UHECR Source properties

- UHE Primary spectrum

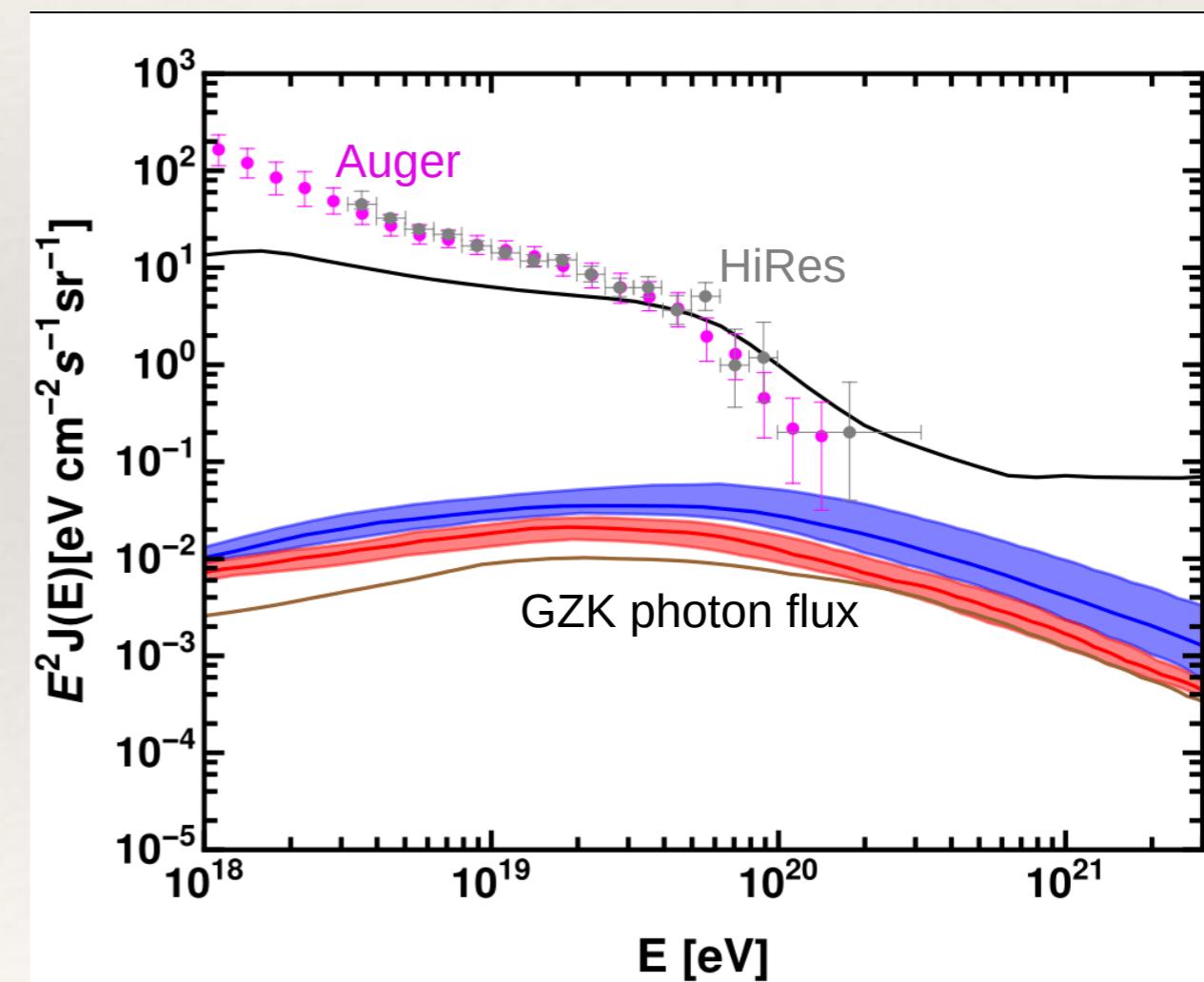
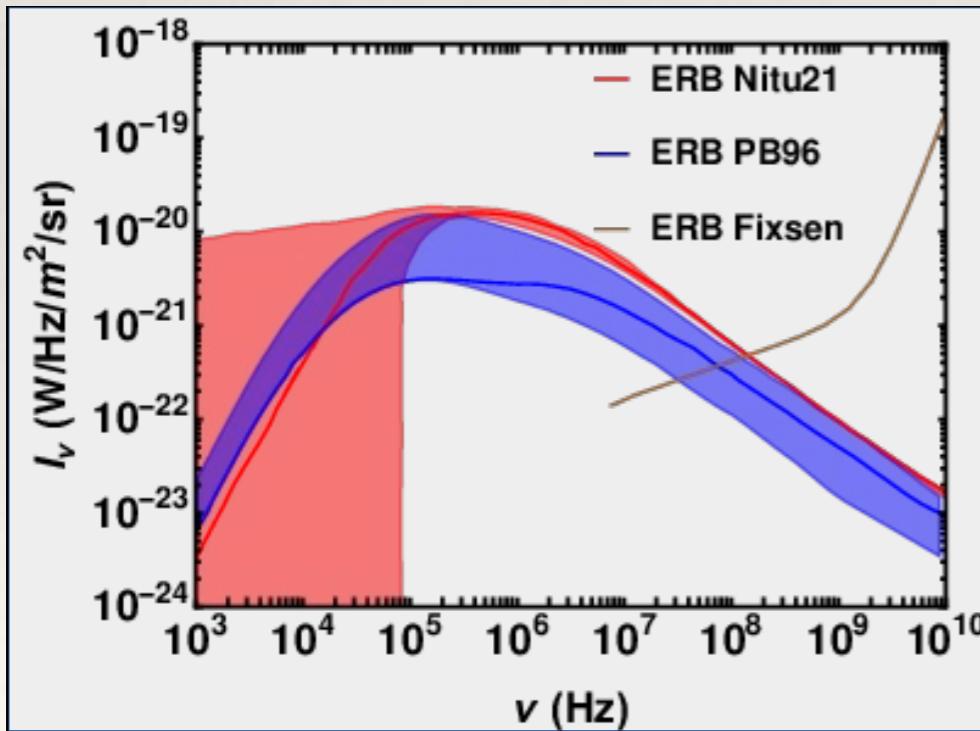
$$E^{-\alpha} \exp\left(\frac{-E}{E_{\text{cut}}}\right) \quad \begin{aligned} \alpha &: 2.2 - 2.7 \\ E_{\text{cut}} &: 5 \times 10^{20} - 10^{22} \text{ eV} \end{aligned}$$

- Composition : protons and heavy elements suggested by observations
- We take 2 extreme cases : 100% proton and 100% iron primary
- How far are the sources : 0.1 kpc to 100 Mpc

# Impact of ERB uncertainties

- ERB uncertainties give rise to about an order of magnitude of uncertainties in the GZK photon flux.
- ARCADE 2 radio excess becomes significant at lower energies for GZK photons.

astro-ph/9605119, 2004.13596, 0901.0555



# Present and Future experiments

Pierre Auger Observatory (Auger)

Giant Radio Array for  
Neutrino Detection (GRAND)

IceCube, IceCube-Gen2

Antarctic Impulsive Transient  
Antenna (ANITA)

Askaryan Radio Array (ARA)

Square Kilometre Array (SKA)

**GZK Photons**

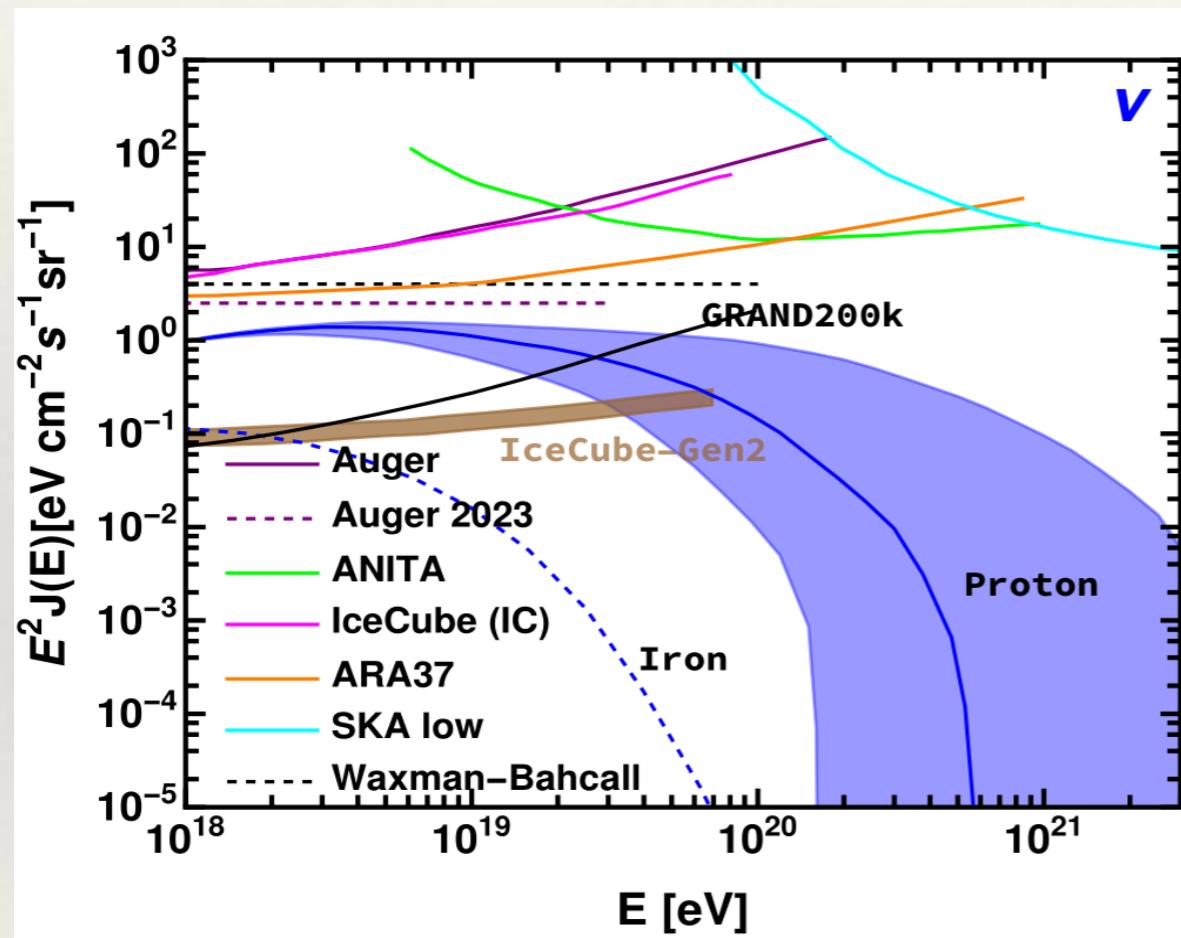
Auger

GRAND

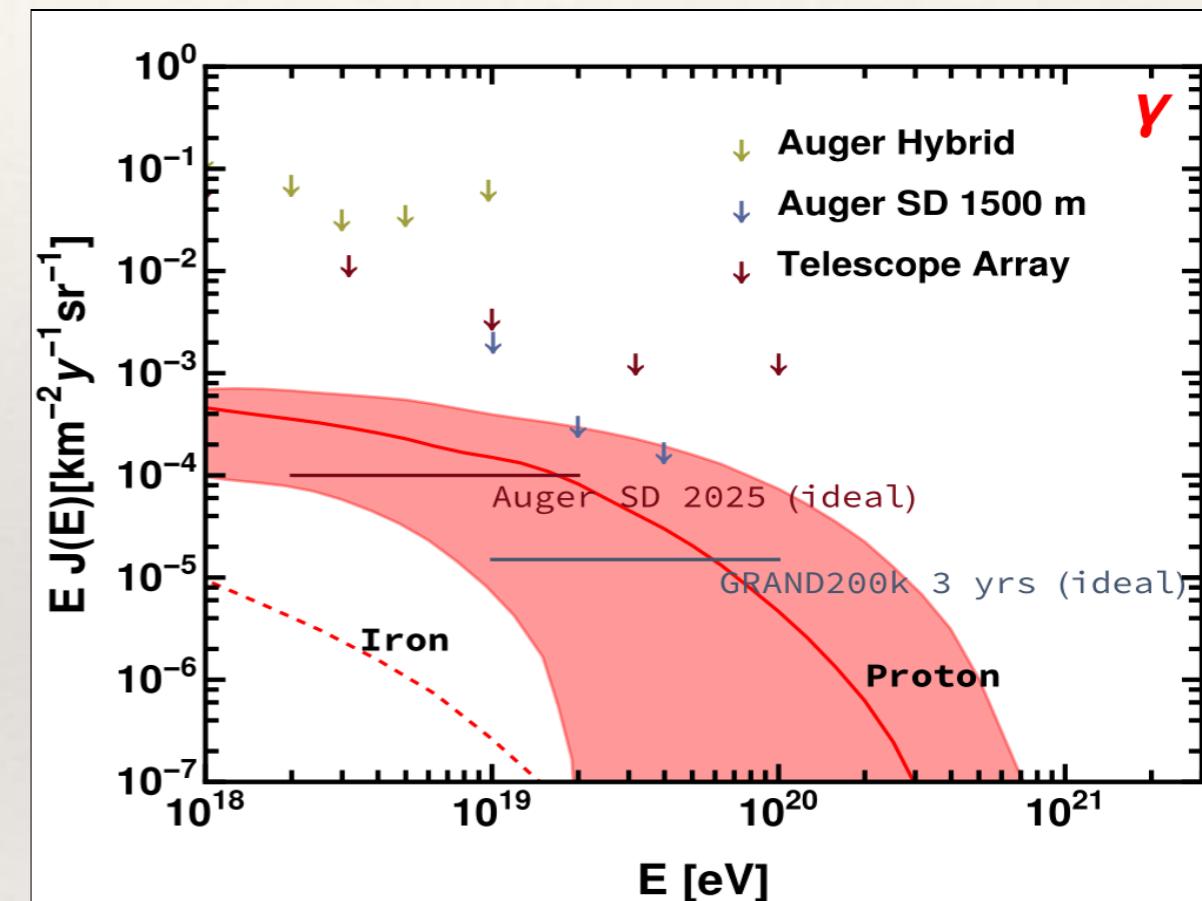
Telescope Array (TA)

# Detection Prospects

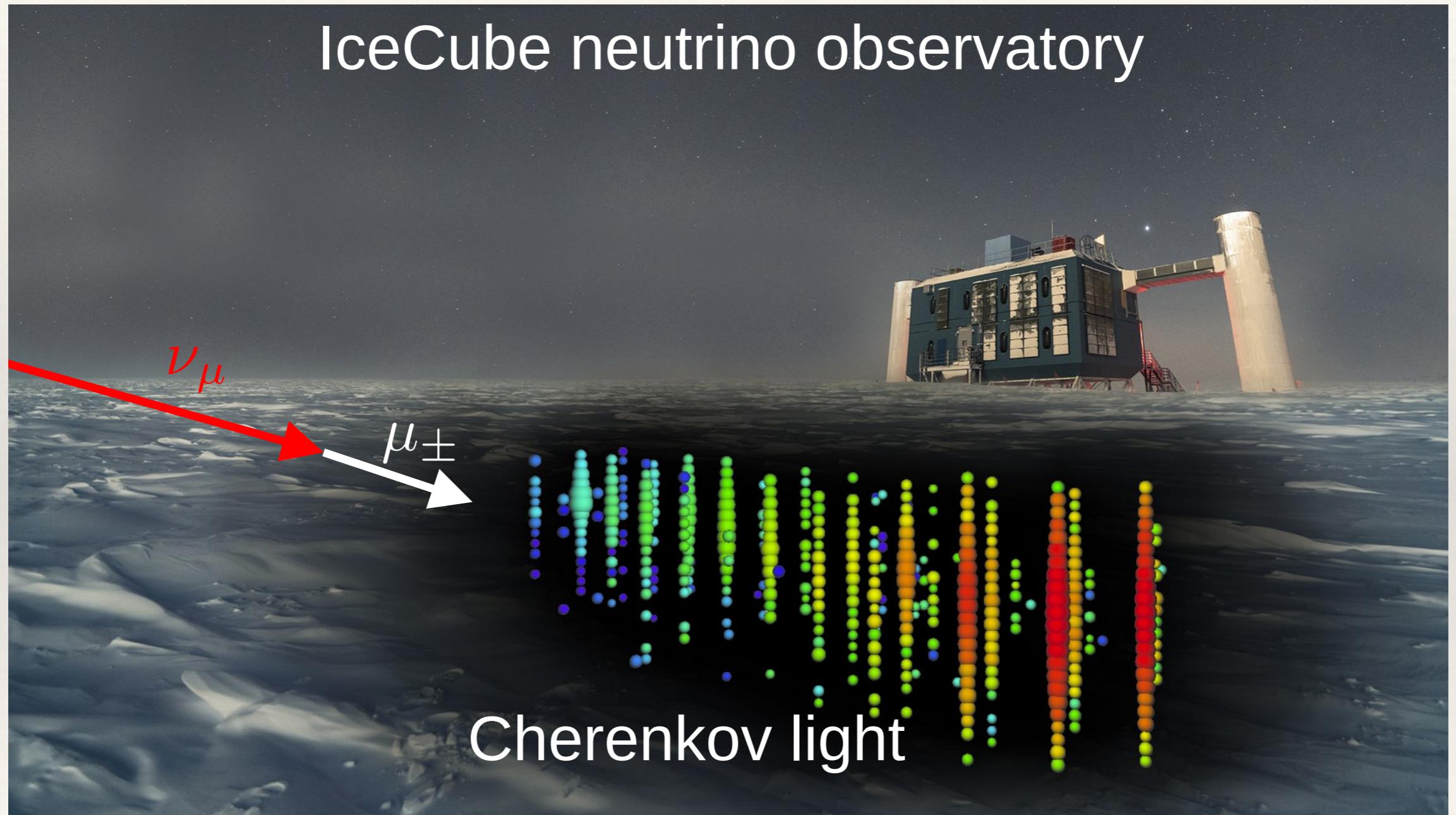
## GZK Neutrinos



## GZK Photons

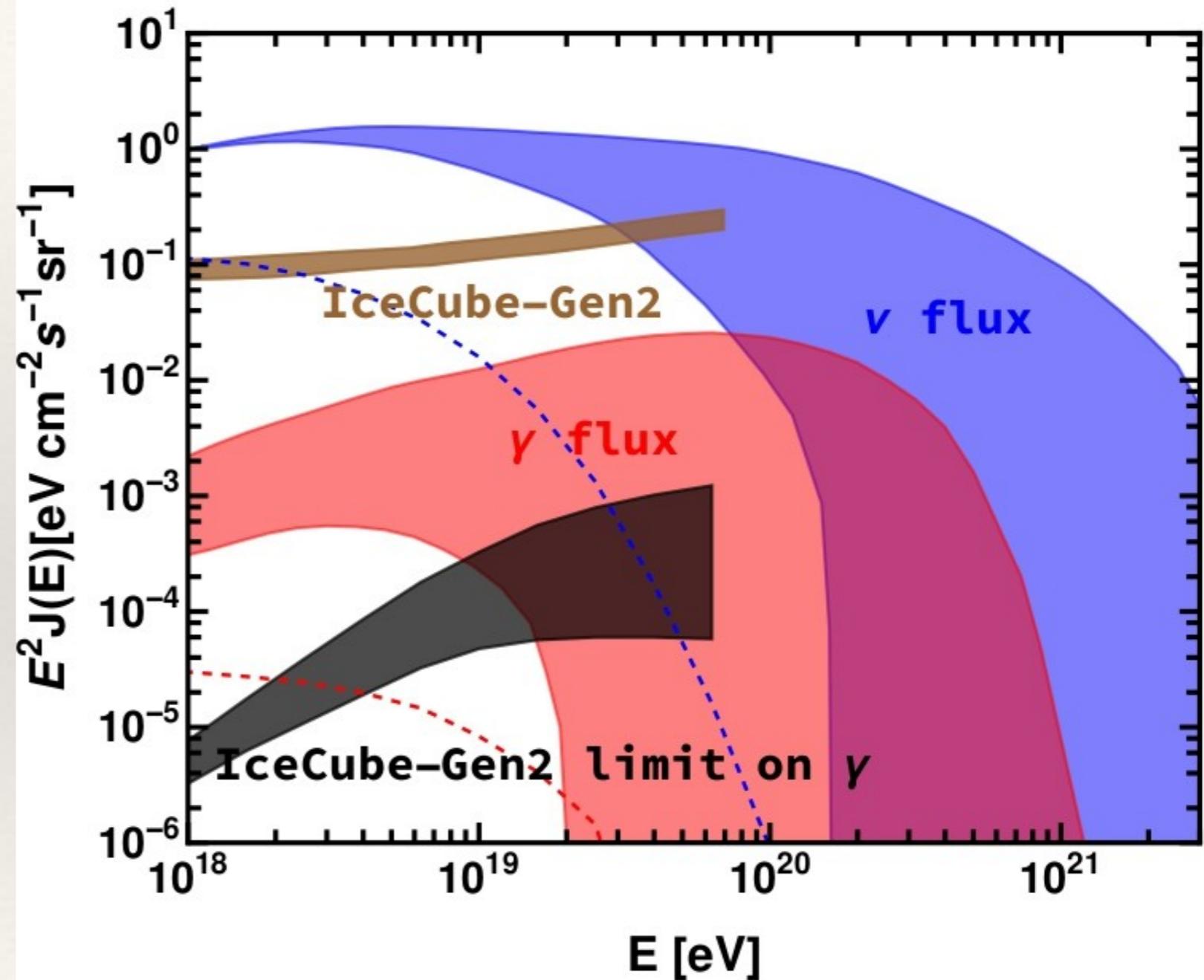


# Multi-messenger constraints from IceCube-Gen2



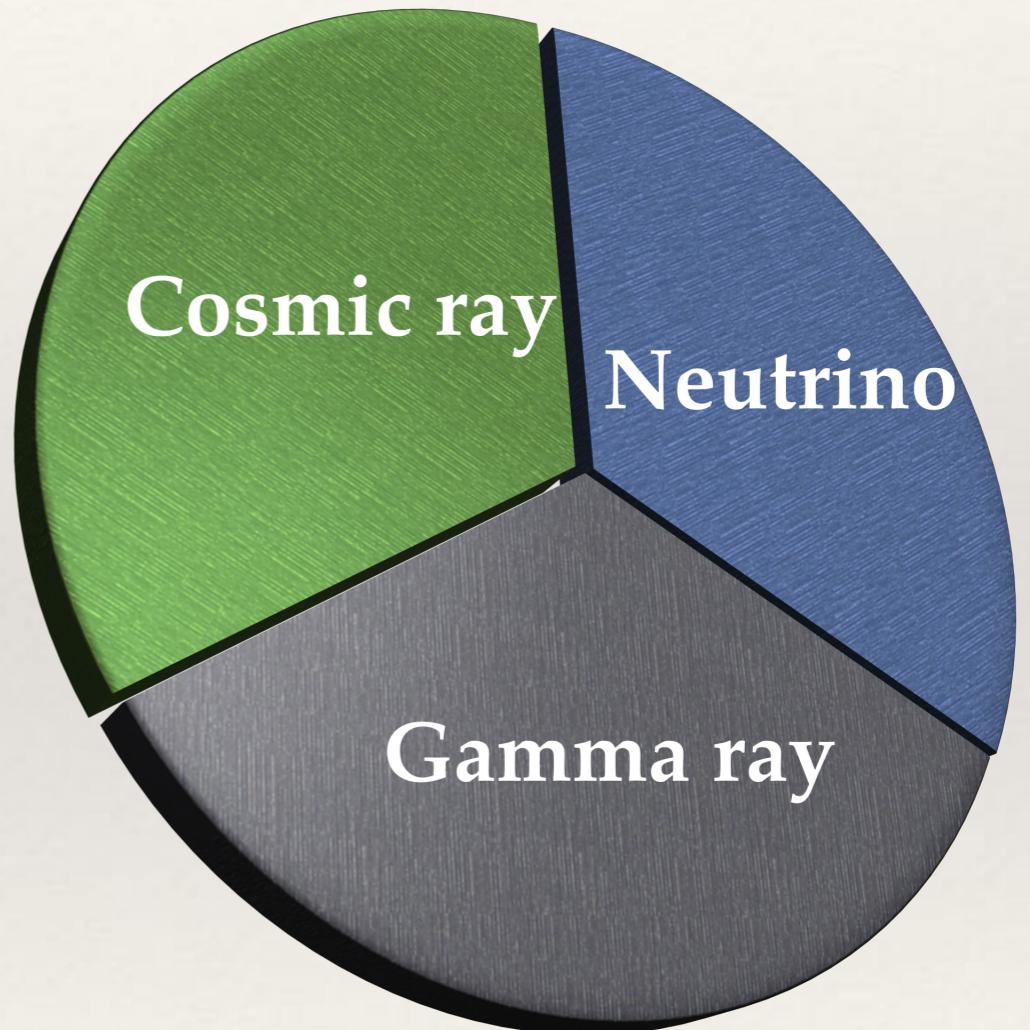
# Multi-messenger constraints from IceCube-Gen2

- GZK Photon flux limit (black band) corresponding to the IceCube-Gen2 sensitivity
- If GZK neutrinos are detected, we can expect GZK photon flux above the black band.
- Non-detection - tension b/w GZK process and UHECR data.



# The future...

- GZK ?
- Better model of ERB, Role of SKA ?
- Sources of UHECR ?
- Power of multi-messenger astronomy!





*Thanks for your attention!*