

Unit - 2

ELECTROMAGNETIC THEORY

Basic laws of electrostatic & magnetostatic.

1) Gauss law of electrostatic.

→ Total Electric flux enclosed by closed surface
is $\frac{1}{\epsilon_0}$ times the total charge enclosed
by that surface. i.e.

$$\oint \vec{E} \cdot d\vec{s} = \frac{q}{\epsilon_0}$$

Where
area
in ;
incl
to
near

integrated
form

2) Gauss law of magnetostatic

→ The magnetic flux entering in a closed surface is always equal to the magnetic flux leaving the surface of same volume. i.e. net flux to a closed surface must be zero.

$$\phi_B = 0$$

or,

$$\oint \vec{B} \cdot d\vec{s} = 0$$

3) Faraday's law for electromagnetic Induct.

Whenever the magnetic flux links with a circuit is change and emf is induced in the circuit. The magnitude of induced emf is directly proportional to negative rate of variation of magnetic flux.

$$e = - \frac{d\phi_B}{dt}$$

$$\text{or, } e = - \frac{\partial \phi_B}{\partial t}$$

$$\text{or, } \oint \vec{B} \cdot d\vec{l} = - \frac{\partial \phi_B}{\partial t}$$

integral form

...4) Ampere's circuital law

Acc. to Ampere's law the line integral of magnetic field be along the closed curve is equal to μ_0 times the net current to the area bounded by the curve.

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I$$

$$\oint H \cdot dI = \Phi_B$$

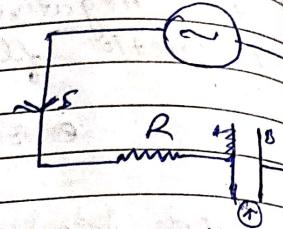
Displacement Current

Conduction Current

$$i_c = \frac{\partial q}{\partial t}$$

displacement current $D = \frac{q}{A} \Rightarrow q = A D$
Density

(vector diff.) $i_c = \frac{\partial (AD)}{\partial t}$



$$P_c = A \frac{\partial D}{\partial t}$$

$$\text{Ans} : i_d = A \frac{\partial D}{\partial t}$$

Maxwell proved that changing EF in vacuum or in dielectric also produces a magnetic field. So a changing EF is equivalent to a current which flows as long as electric field is changing and produces the same magnetic effect as ordinary conduction current. This is known as displacement current.

$$J_d = \frac{i_d}{A} = \frac{\partial D}{\partial t}$$

Note 5-

The accumulation of charges on condenser plates produces a potential difference between the two plates of the capacitor. The potential difference is responsible for the flow of current between the capacitor plates. This is known as displacement current.

Maxwell Equations

in medium

$\mathbf{D} = \epsilon \mathbf{E}$

$\mathbf{D} = \epsilon_0 \mathbf{E}$

$\mathbf{B} = \mu \mathbf{H}$

$\mathbf{B} = \mu_0 \mathbf{H}$

$\mathbf{J}_c = \sigma \mathbf{E}$

$\mathbf{J}_c = 0$

$J_r = \rho V$

Stanadard meaning of used symbols

\mathbf{D} → displacement vector or electric flux density.

\mathbf{E} → Electric field intensity.

\mathbf{B} → Magnetic flux density.

\mathbf{H} → Magnetic field intensity.

μ → Permeability of the medium.

$$\mu = \frac{\mu}{\mu_0} \Rightarrow \mu = \mu_0 \mu_s$$

ϵ_0 → Permeability of the medium in space.

σ → Conductivity of medium.

ρ → Volume charge density.

J_c → Conduction current density.

$J_c \rightarrow$ Conventional current density.
 $V \rightarrow$ Volume.

Maxwell's First eqⁿ:

Acc to the Gauss law of electrostatic.

$$\phi_E = \frac{q}{\epsilon}$$

$$\int \vec{E} \cdot d\vec{s} = \frac{q}{\epsilon} \quad \text{--- (1)}$$

$$q = \int_V \rho dv \quad \text{--- (2)}$$

So eq (1) becomes

$$\boxed{\int \vec{E} \cdot d\vec{s} = \frac{1}{\epsilon} \int_V \rho dv} \rightarrow \text{Integral form.}$$

Using Gauss divergence theorem

$$\int \vec{E} \cdot d\vec{s} = \int_V (\nabla \cdot \vec{E}) dv. \quad \left[\because \int_A d\vec{s} = \int_V \right]$$

$$\Rightarrow \int_V (\nabla \cdot \vec{E}) dv = \frac{1}{\epsilon} \int_V \rho dv$$

Comparing on both sides.

dot product divergence
cross " curl "

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$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon}$$

$$\text{div } \vec{E} = \frac{\rho}{\epsilon}$$

$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon}$$

$$\sigma \cdot \vec{E} = \rho$$

$$\nabla \cdot D = \rho \quad \therefore D = \epsilon F$$

$$\text{div } D = \rho \rightarrow \text{diff. form.}$$

$$\text{div } \vec{D}' = \rho$$

$$\nabla \cdot \vec{D}' = \rho$$

$$\text{div } \vec{E}' = \rho$$

Maxwell's Second Equation

Acc. to Coulomb law of Magnetostatic.

$$\phi_B = 0$$

$$\text{or, } \oint \vec{B} \cdot d\vec{s} = 0 \xrightarrow{\text{Integral form of } B \text{ eq}}$$

Using Gauss divergence theorem.

$$\oint \vec{B} \cdot d\vec{s} = \int (\nabla \cdot B) dv$$

$$\Rightarrow \int_V (\nabla \cdot B) dv = 0$$

$$\Rightarrow \nabla \cdot B = 0 \quad \text{Diff form of } B \text{ eq}$$

$$\text{div } B = 0$$

Maxwell Third Equation

Acc. to Faraday's Law of Electromagnetic Induction $e = -\frac{d\phi_B}{dt}$ → (1)

The induced emf $e = \oint \mathbf{E} \cdot d\mathbf{l}$

$$\oint \phi_B = \oint \mathbf{B} \cdot d\mathbf{s}$$

$$\text{so, } \oint \mathbf{E} \cdot d\mathbf{l} = - \frac{\oint \mathbf{B} \cdot d\mathbf{s}}{\partial t} \quad \text{[Integral form]}$$

Using Stoke's Theorem:

$$\oint \mathbf{E} \cdot d\mathbf{l} = \oint (\nabla \times \mathbf{E}) \cdot d\mathbf{s}$$

$$\Rightarrow \oint (\nabla \times \mathbf{E}) \cdot d\mathbf{s} = - \oint \frac{\partial \mathbf{B}}{\partial t} \cdot d\mathbf{s}$$

$$\Rightarrow \boxed{\nabla \times \mathbf{E} = - \frac{\partial \mathbf{B}}{\partial t}}$$

$$\boxed{\text{curl } \mathbf{E} = - \frac{\partial \mathbf{B}}{\partial t}} \quad \text{Puff form}$$

Maxwell Fourth Eq⁴ (modified form of Ampere's law)

Acc. to Ampere's law

$$\oint \mathbf{B} \cdot d\mathbf{l} = \mu_0 I \quad \text{---(1)}$$

But acc. to Maxwell: $\oint \mathbf{B} \cdot d\mathbf{l} = \mu_0 (\mathbf{J}_c + \mathbf{J}_d)$

where $\mathbf{J}_c = \oint \mathbf{J}_c \cdot d\mathbf{s} = \oint \mathbf{J} \cdot d\mathbf{s}$

$$\mathbf{J}_d = \oint \mathbf{J}_d \cdot d\mathbf{s}$$

Put the values of \mathbf{J}_c & \mathbf{J}_d in eq (1)

$$\oint \mathbf{B} \cdot d\mathbf{l} = \mu_0 (\mathbf{J}_c + \mathbf{J}_d)$$

$$= \mu_0 [\oint \mathbf{J} \cdot d\mathbf{s} + \oint \mathbf{J}_d \cdot d\mathbf{s}]$$

$$\boxed{\oint \mathbf{B} \cdot d\mathbf{l} = \mu_0 \oint (\mathbf{J} + \mathbf{J}_d) \cdot d\mathbf{s}}$$

or

$$\boxed{\oint \mathbf{B} \cdot d\mathbf{l} = \oint (\mathbf{J} + \mathbf{J}_d) \cdot d\mathbf{s}}$$

Integral
form.

$$\boxed{\oint \mathbf{B} \cdot d\mathbf{l} = \mu_0 \oint \left(\mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \right) \cdot d\mathbf{s}}$$

$$\boxed{\mathbf{J}_d = \frac{\partial \mathbf{D}}{\partial t}}$$

Using
Stokes law $\int \mathbf{B} \cdot d\mathbf{l} = \oint (\nabla \times \mathbf{B}) \cdot d\mathbf{s}$

$$\oint (\nabla \times \mathbf{B}) \cdot d\mathbf{s} = \mu_0 \oint \left(\mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \right) \cdot d\mathbf{s}$$

$$(\nabla \times B) = \mu_0 \left(J + \frac{\partial D}{\partial t} \right)$$

Diff. per unit

$$\nabla \times H = J + \frac{\partial D}{\partial t}$$

$$\text{curl } B = \mu_0 \left(J + \frac{\partial D}{\partial t} \right)$$

charge
i.e.

$\Rightarrow \oint$

Ci

Continuity Equation → Mathematical form
law of conservation of
charge

Acc to the law of conservation of
charges : the electric charges
can't be neither created
nor destroyed.

let us consider any region of volume
V bounded by a closed
surface S if the current
flows through the closed
surface S :

$$\frac{dQ}{dt} = i = \oint j \cdot ds \quad \text{--- (1)}$$

Acc. to the conservation principle
the current flow through
a region must be balanced
by rate of decrease of

Charge flowing through that region
i.e.

$$i = \oint J \cdot ds = -\frac{dQ}{dt} = -\int_V \frac{\partial P}{\partial t} dv$$

$$\Rightarrow \oint J \cdot ds + \int_V \frac{\partial P}{\partial t} dv = 0 \quad \text{---(2)} \quad [\because Q = \int_V P dv]$$

Using Gauss Divergence theorem.

$$\oint J \cdot ds = \int_V (\nabla \cdot J) dv$$

so, Eq (2) becomes

$$\int_V (\nabla \cdot J) dv + \int_V \left(\frac{\partial P}{\partial t} \right) dv = 0$$

$$\text{so, } \int_V \left(\nabla \cdot J + \frac{\partial P}{\partial t} \right) dv = 0$$

$$\Rightarrow \boxed{\nabla \cdot J + \frac{\partial P}{\partial t} = 0}$$

↳ continuity Eq.

for Static field

$$\frac{\partial P}{\partial t} = 0$$

$$\boxed{\nabla \cdot J = 0}$$

Another method for continuity Eq,
from maxwell fourth eq²

$$\nabla \times B = \mu_0 \left[J + \frac{\partial D}{\partial t} \right]$$

$$\nabla \times B = \mu_0 \left[J + \epsilon_0 \frac{\partial E}{\partial t} \right] \quad \left[\because D = \epsilon_0 E \right]$$

↓ → ①

Taking divergence on both sides

$$\text{div.}(\nabla \times B) = \mu_0 \text{div} J + \mu_0 \epsilon_0 \frac{\partial}{\partial t} \text{div} E$$

$$\nabla \cdot (\nabla \times B) = \mu_0 \nabla \cdot J + \mu_0 \epsilon_0 \frac{\partial}{\partial t} (\nabla \cdot E)$$

⇒ properties → $\boxed{\text{div. curl} = 0}$

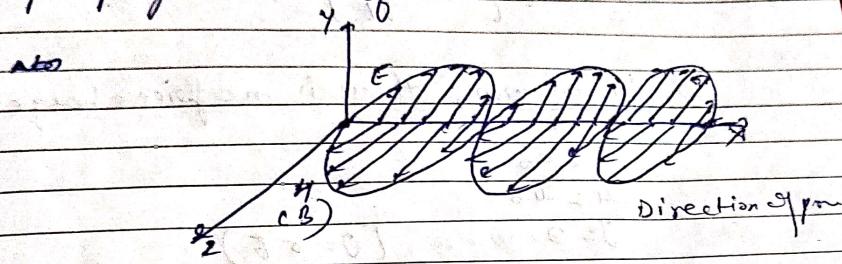
$$0 = \mu_0 \nabla \cdot J + \mu_0 \epsilon_0 \frac{\partial}{\partial t} (\nabla \cdot E) \quad \left| \begin{array}{l} \text{P}_{\text{new}} \\ \text{M}_{\text{new}} \\ J=0 \\ \nabla \cdot E=0 \end{array} \right.$$

$$\Rightarrow \mu_0 \left[\nabla \cdot J + \frac{\partial \Phi}{\partial t} \right] = 0$$

$$\boxed{\nabla \cdot J + \frac{\partial \Phi}{\partial t} = 0}$$

ELECTROMAGNETIC WAVES (E.M Wave)

Electromagnetic wave is a wave of oscillations of electric and magnetic field. In mutually perpendicular planes and these oscillations are perpendicular to the direction of propagation of wave.



NOTE :-

Acc. to Maxwell an oscillating moving charge continuously produces waves which are known as electromagnetic waves. Because oscillating charge has an accelerating motion and a magnetic field is produced around the accelerating charge. As the velocity of the charge changes the magnetic field also changes with time, thus producing an electric field.

F-TALAVE equation in Vacuum (free space)

→ due to vacuum

Maxwell's equation are

$$i) \nabla \cdot E = \frac{J}{\epsilon}$$

$$ii) \nabla \cdot B = 0$$

$$iii) \nabla \times E = -\frac{\partial B}{\partial t}$$

$$iv) \nabla \times B = \mu \left[J + \frac{\partial D}{\partial t} \right]$$

So in vacuum, there is no free charge

$$\epsilon = \epsilon_0$$

$$\mu = \mu_0$$

$$J = 0 \quad [J = \sigma E]$$

$$\rho = 0$$

So Maxwell eq reduces to in free space

$$i) \nabla \cdot E = 0 \quad (\text{non conductive})$$

$$ii) \nabla \cdot B = 0 \quad \text{HLF}$$

$$iii) \nabla \times E = -\frac{\partial B}{\partial t}$$

$$iv) \nabla \times B = \mu_0 \rho = \mu_0 \epsilon_0 \frac{\partial E}{\partial t}$$

(2)

Wave
Maxwell equation in terms of E

from Maxwell's third eq.

$$\nabla \times E = -\frac{\partial B}{\partial t}$$

Taking curl on both sides.

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$$\nabla \times (\nabla \times \vec{E}) = -\frac{\partial}{\partial t} (\vec{\nabla} \times \vec{B})$$

Using vector identity $\vec{a} \times (\vec{b} \times \vec{c}) = (\vec{a} \cdot \vec{c})\vec{b} - (\vec{a} \cdot \vec{b})\vec{c}$

$$\Rightarrow \vec{\nabla} \times (\vec{\nabla} \times \vec{E}) = (\vec{\nabla} \cdot \vec{E})\vec{\nabla} - (\vec{\nabla} \cdot \vec{\nabla})\vec{E} = -\frac{\partial}{\partial t} (\vec{\nabla} \times \vec{B})$$

$$0 = \vec{\nabla}^2 \vec{E} = -\frac{\partial}{\partial t} (\vec{\nabla} \times \vec{B}) \quad \left[\begin{array}{l} \text{---} \\ \vec{\nabla} \cdot \vec{E} = 0 \end{array} \right] \quad \text{from Maxwell}$$

first eqn J.

$$-\vec{\nabla}^2 \vec{E} = +\frac{\partial}{\partial t} (\mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t})$$

$$\Rightarrow \boxed{\vec{\nabla}^2 \vec{E} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}}$$

②

Wave eqn in terms of D

$$\text{We know } \vec{\nabla}^2 \vec{E} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2}$$

$$\& \quad D = \epsilon_0 E \Rightarrow E = \frac{D}{\epsilon_0}$$

$$\text{So, } \vec{\nabla}^2 \left[\frac{D}{\epsilon_0} \right] = \mu_0 \epsilon_0 \frac{\partial^2}{\partial t^2} \left[\frac{D}{\epsilon_0} \right]$$

$$\Rightarrow \nabla^2 D = \mu_0 \epsilon_0 \frac{\partial^2 D}{\partial t^2}$$

(3) wave equation in terms of B

from maxwell's fourth equation

$$\vec{\nabla} \times \vec{B}' = \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

Taking curl on both sides.

$$\nabla \times (\nabla \times \vec{B}) = \mu_0 \epsilon_0 \frac{\partial (\vec{\nabla} \times \vec{E})}{\partial t}$$

$$(\vec{\nabla} \cdot \vec{B}) \vec{\nabla} - (\vec{\nabla} \cdot \vec{B}) \vec{B} = \mu_0 \epsilon_0 \frac{\partial (\vec{\nabla} \times \vec{E})}{\partial t}$$

$$-(\vec{\nabla}^2) \vec{B} = \mu_0 \epsilon_0 \frac{\partial}{\partial t} \left(-\frac{\partial \vec{B}}{\partial t} \right)$$

$$+(\nabla^2) \vec{B} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{B}}{\partial t^2}$$

$$\boxed{\nabla^2 B = \mu_0 \epsilon_0 \frac{\partial^2 B}{\partial t^2}}$$

(4) wave equation in terms of H

$$\nabla \times \vec{H} = \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

$$\nabla \times (\nabla \times H) = \mu_0 \frac{\partial}{\partial t} (\nabla \times E)$$

$$(\nabla \cdot H) \nabla - (\nabla \cdot \nabla) H = \mu_0 \frac{\partial}{\partial t} \left(-\mu_0 \frac{\partial H}{\partial t} \right)$$

$$+\nabla^2 H = \mu_0 \epsilon_0 \frac{\partial^2 H}{\partial t^2}$$

$$\boxed{\nabla^2 H = \mu_0 \epsilon_0 \frac{\partial^2 H}{\partial t^2}}$$

So finally

$$\boxed{\nabla^2 \begin{bmatrix} E \\ D \\ B \\ H \end{bmatrix} = \mu_0 \epsilon_0 \frac{\partial^2}{\partial t^2} \begin{bmatrix} E \\ D \\ B \\ H \end{bmatrix} - A}$$

Biennally any wave eqt,

$$\nabla^2 A = 1 \frac{\partial^2 A}{\partial t^2} - B$$

On comparing A & B

$$\frac{1}{v^2} = \mu_0 \epsilon_0 \Rightarrow v = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$$

$$v = \frac{1}{\sqrt{4 \pi \times 10^{-7} \times 8.85 \times 10^{-12}}}$$

$$v = 3 \times 10^8 \text{ m/sec.}$$

(Hence EM waves
with speed of
light in vacuum)

v = speed in cgs

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Transverse NATURE OF E.M. wave

(Solution of Plane E.M. wave equation)

$$\nabla^2 E - \mu_0 \epsilon_0 \frac{\partial^2 E}{\partial t^2} = 0 \quad \text{--- (1)}$$

$$\nabla^2 B - \mu_0 \epsilon_0 \frac{\partial^2 B}{\partial t^2} = 0 \quad \text{--- (2)}$$

Solution of these equations are.

$$\vec{E}(r, t) = \vec{E}_0 e^{(ik \cdot r - i\omega t)} \quad \text{--- (3)}$$

$$\vec{B}(r, t) = \vec{B}_0 e^{(ik \cdot r - i\omega t)} \quad \text{--- (4)}$$

So,

$$\nabla \cdot \vec{E} = \left[\frac{\partial}{\partial x} + j \frac{\partial}{\partial y} + k \frac{\partial}{\partial z} \right] \cdot \vec{E}_0 e^{(ik \cdot r - i\omega t)}$$

$$\text{where } \vec{k} \cdot \vec{r} = (jk_x + ik_y + k_z k_z) \cdot (\hat{i}_x + \hat{j}_y + \hat{k}_z)$$

propagation
vector

$$= (k_x x + k_y y + k_z z)$$

$$\nabla \cdot \vec{E} = \left[\frac{\partial}{\partial x} + j \frac{\partial}{\partial y} + k \frac{\partial}{\partial z} \right] \cdot \left[(E_{0x} + j E_{0y} + k E_{0z}) e^{(ik \cdot r - i\omega t)} \right]$$

$$= \frac{\partial}{\partial x} \left[E_{0x} e^{(ik \cdot r + k_y y + k_z z) - i\omega t} \right] +$$

$$j \frac{\partial}{\partial y} \left[E_{0y} e^{(ik \cdot r + k_y y + k_z z) - i\omega t} \right] +$$

$$+ \frac{\partial}{\partial z} \left[E_{0z} e^{i(k_x x + k_y y + k_z z) - i\omega t} \right]$$

$$= (i k_x) E_{0x} e^{i(k_x x + k_y y + k_z z) - i\omega t} +$$

$$i k_y E_{0y} e^{i(k_x x + k_y y + k_z z) - i\omega t} + i k_z E_{0z} e^{i(k_x x + k_y y + k_z z) - i\omega t}$$

$$= i [k_x E_{0x} + k_y E_{0y} + k_z E_{0z}] e^{i(k_x x + k_y y + k_z z) - i\omega t}$$

$$= i [\vec{k} \cdot \vec{E}_0] e^{i(\vec{k} \cdot \vec{r}) - i\omega t}$$

$$\vec{\nabla} \cdot \vec{E}' = i [\vec{k} \cdot \vec{E}']$$

From Maxwell's first eqn $\vec{\nabla} \cdot \vec{E}' = 0$

$$\Rightarrow i [\vec{k} \cdot \vec{E}'] = 0$$

Similarly, $i [\vec{k} \cdot \vec{B}'] = 0$

Hence $\vec{E}' \& \vec{B}'$ are \perp to the direction of wave.

for Mutually Perpendicular of Electric &

Magnetic field.

$$\boxed{\vec{k}' \cdot \vec{E}' = \omega \vec{B}'}$$

$$\boxed{\vec{k}' \times \vec{B}' = -\omega \mu_0 \epsilon_0 \vec{E}'}$$

$$\vec{B}' = \frac{i}{\omega} (\vec{k} \times \vec{E}')$$

$$\vec{B}' = \frac{i}{\omega} (\vec{k} \times \vec{E}') \quad [\because \vec{k} \perp \vec{E}']$$

$$\omega \vec{H}' = \frac{i}{\omega} (\vec{k} \times \vec{E}')$$

$$\vec{H}' = \frac{K}{\omega_{HO}} (\hat{n} \times \vec{E}')$$

$$\left[\frac{\omega}{K} = c \Rightarrow \right]$$

$$\vec{H}' = \frac{1}{c \mu_0} (\hat{n} \times \vec{E}')$$

$$\left[c = \frac{1}{\mu_0 \epsilon_0} \right]$$

in terms of magnitude,

$$H = \frac{1}{c \mu_0} E$$

$$H = \frac{\sqrt{\mu_0 \epsilon_0}}{\epsilon_0} E \Rightarrow \frac{E}{H} = \sqrt{\frac{\mu_0}{\epsilon_0}} = \sqrt{\frac{4\pi \times 10^{-7}}{8.85 \times 10^{-12}}}$$

$$= 376.72 \text{ n.}$$

$$\boxed{\frac{E}{H} = \frac{\mu_0}{\epsilon_0} = Z = 376.72 \text{ n.}}$$

wave impedance
factor

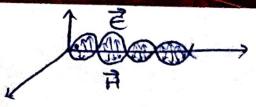
(from depth)
(sun)

(of free space)

Ans Poynting vector & Poynting theorem!

The rate of flow of wave energy through unit area of space normal to the direction of propagation of E-M wave is represented by energy - flux ~~vector~~ vector (P). This energy flux vector is also called Poynting vector.

The cross product of electric field (E) & magnetic field vector is



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Known as Poynting vector.

$$\vec{P} = \vec{E}' \times \vec{H}'$$

S.I unit of point vector is watt/m².

Poynting theorem

from Maxwell fourth eqn :

$$\nabla \times \vec{H} = \vec{J} + C \frac{\partial \vec{E}}{\partial t} \quad \text{---(1)}$$

Taking dot product of E on both sides.

$$\cdot \vec{E} \cdot (\nabla \times \vec{H}) = \vec{E} \cdot \vec{J} + C \frac{\partial \vec{E} \cdot \vec{E}}{\partial t} \quad \text{---(2)}$$

Using Vector identity :

$$\nabla \cdot (\vec{E} \times \vec{H}) = \vec{H} \cdot (\nabla \times \vec{E}) - \vec{E} \cdot (\nabla \times \vec{H})$$

from this identity.

$$\vec{E} \cdot (\nabla \times \vec{H}) = \vec{H} \cdot (\nabla \times \vec{E}) - \nabla \cdot (\vec{E} \times \vec{H})$$

Put this value in eq (2)

$$\vec{H} \cdot (\nabla \times \vec{E}) - \nabla \cdot (\vec{E} \times \vec{H}) = \vec{E} \cdot \vec{J} + C \frac{\partial \vec{E} \cdot \vec{E}}{\partial t} \quad \text{---(3)}$$

From Maxwell 3rd Eqn:

$$\nabla \times E = -\frac{\partial B}{\partial t} = -\mu \frac{\partial H}{\partial t}$$

$$\mu \left[-\mu \frac{\partial H}{\partial t} \right] - \nabla \cdot (E \times H) = E \cdot J + \epsilon \frac{\partial (E \cdot E)}{\partial t} \quad (1)$$

$$\Rightarrow -\mu \left[\mu \frac{\partial H}{\partial t} \right] - \nabla \cdot (E \times H) = E \cdot J + \mu E \cdot \frac{\partial (E \cdot E)}{\partial t}$$

$$\Rightarrow -\mu \left[\mu \frac{\partial H}{\partial t} \right] - \nabla \cdot (E \times H) = E \cdot J + \epsilon E = \frac{\partial F}{\partial t}$$

We know $\frac{\partial (H^2)}{\partial t} = 2H \frac{\partial H}{\partial t}$

$$\Rightarrow \mu \frac{\partial H}{\partial t} = \frac{1}{2} \frac{\partial (H^2)}{\partial t}$$

Similarly $E \cdot \frac{\partial E}{\partial t} = \frac{1}{2} \frac{\partial (E^2)}{\partial t}$

$$-\mu \left(\frac{1}{2} \frac{\partial H^2}{\partial t} \right) - \nabla \cdot (E \times H) = E \cdot J + \epsilon \left[\frac{1}{2} \frac{\partial (E^2)}{\partial t} \right]$$

$$-\nabla \cdot (E \times H) = E \cdot J + \frac{\partial}{\partial t} \left[\frac{\epsilon E^2}{2} + \frac{\mu H^2}{2} \right]$$

$$\nabla \cdot (E \times H) = -E \cdot J - \frac{\partial}{\partial t} \left[\frac{\epsilon E^2}{2} + \frac{\mu H^2}{2} \right] \quad (2)$$

$$\int \nabla \cdot (E \times H) dV = - \int E \cdot J dV - \frac{d}{dt} \int \left(\frac{\mu H^2}{2} + \frac{\epsilon E^2}{2} \right) dV \quad (6)$$

Using Gauss divergence theorem,

$$\int \nabla \cdot (E \times H) dV = \oint (E \times H) ds.$$

$$\boxed{\oint (E \times H) ds = - \int (E \cdot J) dV - \frac{d}{dt} \int \left(\frac{\mu H^2}{2} + \frac{\epsilon E^2}{2} \right) dV} \quad (1) \quad (2) \quad (3)$$

→ First term, from this equation
first term represent the rate
of flow of energy outward
through the surface.

→ Second term, it represent Power
dissipated (loss) in volume 'V'

$$E \cdot J = V \cdot \frac{I}{A} = \frac{V I}{A} = \text{Power loss per unit } 10^6 \text{ m}^3.$$

→ Third term, in this $\frac{d}{dt}$ → represent
the rate at which stored energy is
decreasing.

→ $\frac{\mu H^2}{2}$ → energy density due
to magnetic field.

→ $\frac{\epsilon E^2}{2}$ → energy density due to electric field.

Q1. If the Earth receives ϵ calories $\text{min}^{-1} \text{cm}^{-2}$ solar energy what are the amplitudes of electric and magnetic fields of radioactivity?

Sol:

ENERGY DENSITY OF EM WAVE IN FREE SPACE

In a small volume energy of electric field

$$U_E = \frac{1}{2} \epsilon_0 E^2 dv \quad (1)$$

In a small volume energy of M.P.

$$U_B = \frac{1}{2} \mu_0 H^2 dv \quad (2)$$

So, electric energy per unit volume
that is electric field energy density.

$$U_E = \frac{1}{2} \epsilon_0 E^2 \quad (3)$$

Magnetic energy per unit volume

$$U_B = \frac{1}{2} \mu_0 H^2 \quad (4)$$

So, total energy density

$$U = U_E + U_B \\ = \frac{1}{2} \epsilon_0 E^2 + \frac{1}{2} \mu_0 H^2$$

We know $\frac{E}{H} = \sqrt{\mu_0 / \epsilon_0}$

$$\frac{E^2}{H^2} = \frac{\mu_0}{\epsilon_0} \Rightarrow \mu_0 H^2 = \epsilon_0 E^2$$

$$U = \frac{1}{2} \epsilon_0 E^2 + \frac{1}{2} \mu_0 H^2$$

$$\boxed{U = \epsilon_0 E^2}$$

Ratio of E.f & M.f energy density.

$$\frac{U_E}{U_B} = \frac{\frac{1}{2} \epsilon_0 E^2}{\frac{1}{2} \mu_0 H^2} = \frac{\epsilon_0 E^2}{\mu_0 H^2} [\because \mu_0 H^2 = \epsilon_0 E^2]$$

$$\boxed{\frac{U_E}{U_B} = 1}$$

Hence, electric field energy density is equal to magnetic field energy density, i.e. the energy associated with the electric field is equal to the energy associated with the magnetic field.

ENERGY & MOMENTUM of EM WAVE:

Maxwell predicted that EM wave transport linear momentum in the direction of propagation.

The momentum of a particle of mass m with velocity v is given by,

$$[p = mv] \quad \textcircled{1}$$

According to Einstein mass-energy relation

$$E = mc^2 \quad \textcircled{2}$$

$$\Rightarrow m = \frac{E}{c^2}$$

So, $\vec{p} = \frac{E}{c^2} \vec{v} \quad \textcircled{3}$

~~$E = mc^2$~~

The energy density in plane electromagnetic wave is given by

$$U = E = 60 E^2 \quad \textcircled{4}$$

So, the momentum density or momentum per unit volume associated with electromagnetic wave is

$$\vec{p}' = \frac{U}{c^2} \vec{v} \quad \textcircled{5}$$

If wave is propagating in x-axis, then

$$\vec{v} = c \hat{i}$$

So, $\boxed{\vec{p}' = \frac{U}{c^2} \cdot c \hat{i} = \frac{U}{c} \hat{i}} \quad \textcircled{6}$

By Pointing vector, $\vec{S}^2 = \frac{1}{\mu_0} (\vec{E}' \times \vec{B})$

But, $\vec{E}' \times \vec{B} = \frac{\epsilon^2}{c^2} \hat{i}$.

$$\vec{S} = \frac{\epsilon^2}{c^2} \hat{i}$$

$$\Rightarrow \vec{S} = \frac{4}{\mu_0 \epsilon_0 c^2} \hat{i} \quad \left[\because \frac{4}{\epsilon_0} = \epsilon^2 \right]$$

We know $c = 1 \Rightarrow 1 = c^2$

$\therefore \vec{S} = \frac{4}{\mu_0 \epsilon_0} \hat{i}$ (in plane wave)

$$\vec{S} = (4c) \hat{i} \quad \text{and} \quad \vec{B} = \frac{\vec{S}}{c} \quad \text{in}$$

Put the values in eq. ⑥ $\vec{P} = \vec{S}$.

$$\vec{P} = 4 \hat{i}$$

$$\vec{P} = \frac{4c^2}{\mu_0 \epsilon_0} \hat{i} \quad \text{or} \quad 1 = c^2 (\vec{E}' \times \vec{B})$$

$$\Rightarrow c = \frac{1}{\sqrt{\mu_0 \epsilon_0}} \Rightarrow c^2 = \frac{1}{\mu_0 \epsilon_0}$$

$$\Rightarrow \mu_0 c^2 = \frac{1}{60}$$

$$\boxed{\vec{P} = 60 (\vec{E}' \times \vec{B})} \rightarrow \textcircled{Q}$$

$$\therefore \vec{E} \times \vec{B} = \frac{c^2}{c} \vec{i}$$

$$p = \frac{c \rho E^2}{c} \vec{j} \Rightarrow p = \rho v \vec{i} \quad \begin{cases} \rho = \text{energy density} \\ v = \text{wave velocity} \end{cases}$$

Energy Density = momentum \times wave velocity.

Radiation Pressure & Energy Density!

When EM wave strikes on a surface its momentum is changes. The rate of change of momentum is equal to the force. This force acting on unit area of the surface applied a pressure i.e. called radiation pressure.

Let a EM wave incident normally on a perfectly absorbing surface of area "A" for a time "t". If energy "U" is absorbed during this time, the momentum "p" delivered to the surface is given by $p = \frac{U}{c}$. According to Maxwell,

If "S" is the energy passing per unit area per unit time, then $U = SAT$ - ②

So, $b = \frac{C_0 A t}{c}$

$b = u A t \quad \text{--- (3)}$

So, $F = \frac{P}{t} = \frac{u A b}{t} = u A b = u A \quad \text{--- (4)}$

Radiation Pressure, $P_{\text{rad}} = \frac{F}{A}$

$P_{\text{rad}} = \frac{u A}{A}$

$P_{\text{rad}} = u$

The radiation pressure applied by a normally incident plane electromagnetic wave on a perfect absorber is equal to the energy density in that wave.

SKIN DEPTH (DEPTH OF PENETRATION)

when EM
wave travel
in conducting
medium

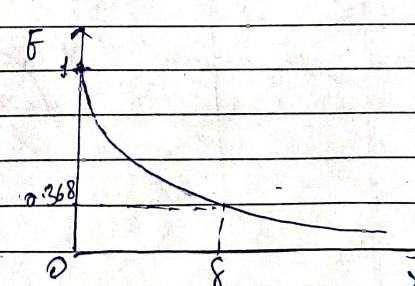
$$e^{-\alpha x} = \frac{1}{2}$$

$$e^{-\alpha x} = e^{-1}$$

$$\alpha x = 1$$

$$\text{If } x = \delta.$$

$$\Rightarrow \alpha \cdot \delta = 1 \Rightarrow \delta = \frac{1}{\alpha}$$



$\therefore \text{Skin depth} = \frac{1}{\text{Attenuation const.}}$

→ Skin depth is defined as when the EM wave travels in any medium so the value of its initial electric field is reduced to $\frac{1}{e}$ times, this depth is known as skin depth or depth of penetration.

Conductor

Skin depth

$$\delta = \sqrt{\frac{2}{\omega \mu \sigma}}$$

Dielectric / Insulator

$$\delta = \frac{2}{\omega \sqrt{\epsilon_r}}$$