

RECAP S1:

TOWARD A VARIATIONAL ESTIMATION OF THE STEADY
STATE IN OPEN QUANTUM SYSTEMS USING GENERAL-
IZED GIBBS ENSEMBLE

FEBRUARY 26, 2024

OPEN QUANTUM SYSTEMS

■ Lindblad equation:

$$\dot{\rho} = \mathcal{L} \rho \equiv -i [H, \rho] + \sum_{k=1}^{N^2-1} \gamma_k \left(L_k \rho L_k^\dagger - \frac{1}{2} \{L_k^\dagger L_k, \rho\} \right) \quad (1)$$

■ Steady state:

$$\dot{\rho} \equiv 0$$

■ thermalisation performed by Gibbs ensemble (GE):

$$\rho_{\text{GE}} = \frac{e^{-\beta \hat{H}}}{\text{Tr} \left(e^{-\beta \hat{H}} \right)} \quad (2)$$

GENERALISED GIBBS ENSEMBLES (GGE)

■ Von Neumann entropy

$$S(\rho) = -\text{Tr}(\rho \ln \rho) \quad (3)$$

■ definition GGE (Max Entropy Principle):

$$\rho_{\text{GGE}} = \text{argmax}_{\rho} \{S(\rho) : \forall A \in \mathcal{C} : \text{Tr}(A \rho) = \text{Tr}(A \rho(0))\} \quad (4)$$

■ General form of GGE (= most general form of any steady state):

$$\rho_{\text{GGE}} = \frac{e^{-\sum_{A \in \mathcal{C}} \lambda_A A}}{\text{Tr} \left(e^{-\sum_{A \in \mathcal{C}} \lambda_A A} \right)} \quad (5)$$

VARIATIONNAL ESTIMATION OF THE STEADY STATE: PRINCIPLE

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1. Start from GGE ansatz with predefined conserved quantities:

$$\rho_{\text{GGE}} = \frac{e^{-\sum_{A \in \mathcal{C}} \lambda_A A}}{\text{Tr} \left(e^{-\sum_{A \in \mathcal{C}} \lambda_A A} \right)}$$

2. Vary the Lagrange multipliers $\{\lambda_A\}_{A \in \mathcal{C}}$ to minimize $\|\dot{\rho}\|$

APPLICATION: THE DAMPED HARMONIC OSCILLATOR

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- Lindblad equation ($\bar{n} = (\exp \beta \omega - 1)^{-1}$):

$$\begin{aligned}\mathcal{L}_H \rho = & -i [H, \rho] \\ & + \gamma(\bar{n} + 1) \left(a \rho a^\dagger - \frac{1}{2} \{a^\dagger a, \rho\} \right) \\ & + \gamma \bar{n} \left(a^\dagger \rho a - \frac{1}{2} \{a a^\dagger, \rho\} \right)\end{aligned}\tag{6}$$

- Different Hamiltonians (ϵ small) and ansätze:

$$H_0 = a^\dagger a$$

$$H_1 = a^\dagger a + \epsilon (a + a^\dagger)$$

$$H_2 = a^\dagger a + \epsilon a^\dagger a^\dagger a a$$

$$\rho_{\text{GGE } 1} = \frac{e^{-\lambda_0 a^\dagger a}}{\text{Tr} (e^{-\lambda_0 a^\dagger a})}$$

$$\rho_{\text{GGE } 2} = \frac{e^{-\lambda_0 a^\dagger a - \lambda_1 a^\dagger a^\dagger a a}}{\text{Tr} (e^{-\lambda_0 a^\dagger a - \lambda_1 a^\dagger a^\dagger a a})}$$

RESULTS

RESULTS - HO 1

■ Hamiltonian:

$$H_0 = a^\dagger a$$

■ Ansatz:

$$\rho_{\text{GGE}1} = \frac{e^{-\lambda_0 a^\dagger a}}{\text{Tr}(e^{-\lambda_0 a^\dagger a})}$$

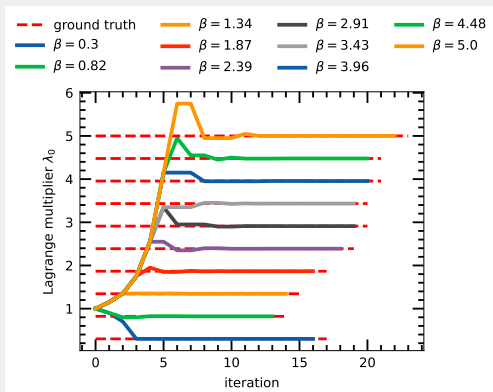


Figure: Optimization of the GGE: Lagrange multiplier as a function of the number of iterations.

RESULTS - HO 2

$$H_0 = a^\dagger a$$

$$\rho_{\text{GGE}2} = \frac{e^{-\lambda_0 a^\dagger a - \lambda_1 a^\dagger a^\dagger a a}}{\text{Tr} (e^{-\lambda_0 a^\dagger a - \lambda_1 a^\dagger a^\dagger a a})}$$

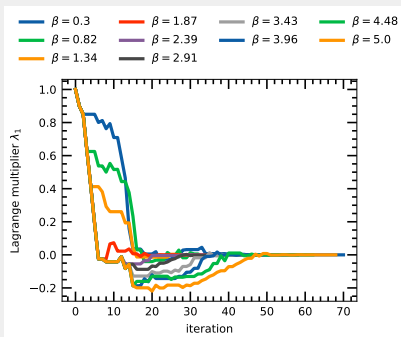
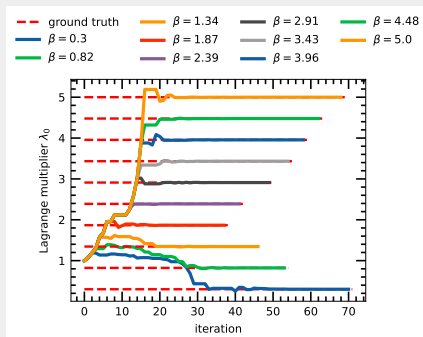


Figure: λ_0 as a function of the number of iterations.

Figure: λ_1 as a function of the number of iterations.

RESULTS - DRIVE

$$H_1 = a^\dagger a + \epsilon (a + a^\dagger)$$

$$\rho_{\text{GGE}2} = \frac{e^{-\lambda_0 a^\dagger a - \lambda_1 a^\dagger a^\dagger a a}}{\text{Tr} (e^{-\lambda_0 a^\dagger a - \lambda_1 a^\dagger a^\dagger a a})}$$

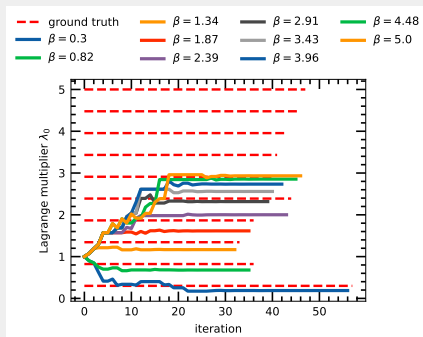


Figure: λ_0 as a function of the number of iterations.

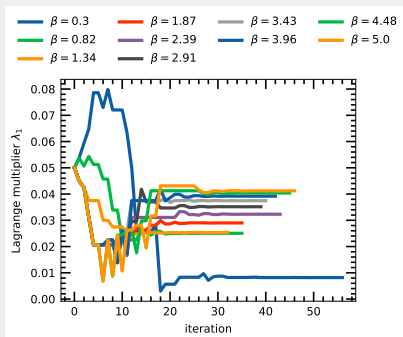


Figure: λ_1 as a function of the number of iterations.

RESULTS - KERR

$$H_2 = a^\dagger a + \epsilon a^\dagger a^\dagger a a$$

$$\rho_{\text{GGE}2} = \frac{e^{-\lambda_0 a^\dagger a - \lambda_1 a^\dagger a^\dagger a a}}{\text{Tr} (e^{-\lambda_0 a^\dagger a - \lambda_1 a^\dagger a^\dagger a a})}$$

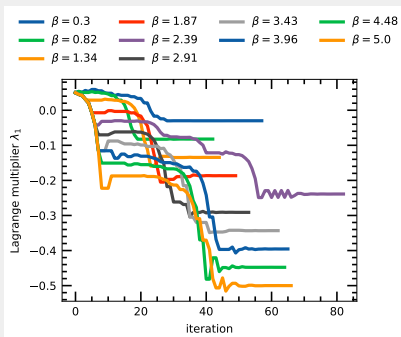
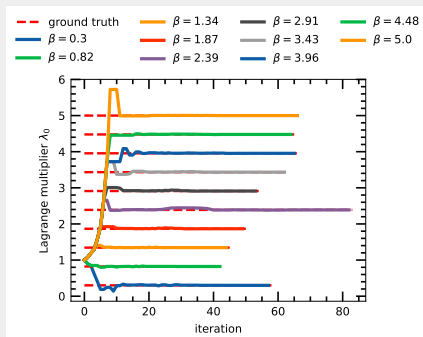


Figure: λ_0 as a function of the number of iterations.

Figure: λ_1 as a function of the number of iterations.

RESULTS

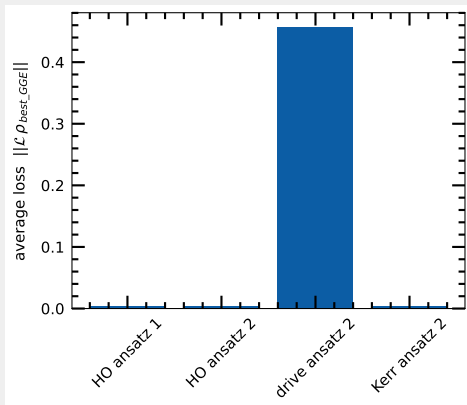






Figure: Comparison of the losses of the optimized GGEs in the various scenarios. The loss is the average value of $\|\mathcal{L} \rho_{GGE}\|$ calculated across all true temperatures.

OUTLOOK


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- Use a variational state less informed by the analytical solution (= more general way to express conserved quantities ?)
- Generalize the method in the multi-body case

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