

$$\text{Surface charge density} = \sigma = \frac{Q}{A}$$

$$\text{Volume charge density} = \rho = \frac{Q}{V}$$

ELECTRIC FIELD

An electric Field is the physical field that surrounds electrically charged particles and exerts force on all other charge particles in field, either attracting or repelling them.

Equation:

$$E = k \frac{Q}{r^2} = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2}$$

$$k = 9 \times 10^9 \text{ N m}^2 \text{ C}^{-2}$$

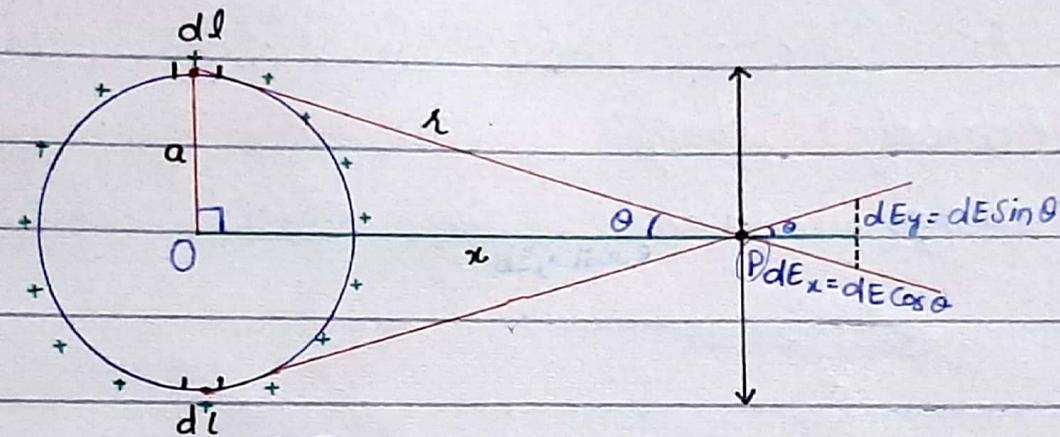
Calculation of Electric Field of Continuous Charges

Consider a Ring shaped object in which charges are distributed uniformly.

In the Ring shaped object the total distributed charge is Q . And a Point P lies at a distance

from the centre O of the Ring. The Radius of the Ring is a .

Figure:



To calculate the electric field of the Ring we have to divide the Ring into small patches.

Each patch will be equal to dl . Here r is the distance of the Point P from the patch.

Now For calculation of Electric Field on Point 'P' Resolve the point P into its components.

Total Electric Field can be given as:

$$E = \int dE$$

The Vertical Components are cancelled because they are equal in magnitude but opposite in direction.
So

$$dE = dE_x + dE_y$$

$$\therefore dE_y = 0$$

$$dE = dE_x$$

Hence

$$E = \int dE_x \quad \dots \text{(i)}$$

$$E = \int dE \cos\theta$$

$$\therefore E = \frac{1}{4\pi\epsilon_0} \frac{dQ}{r^2} = \frac{1}{4\pi\epsilon_0} \frac{\lambda d\lambda}{r^2}$$

Derivative $\frac{d}{d\lambda}$ $\int d\lambda$

$\frac{d}{d\lambda} \int d\lambda$
Area λ $\int d\lambda$
Derivative $\frac{d}{d\lambda}$ $\int d\lambda$
 λ $\int d\lambda$
Small Patched Surface λ
 λ $\int d\lambda$ Divide λ

so Now

$$E = \frac{1}{4\pi\epsilon_0} \int \frac{dQ}{r^2} \cos\theta \quad \dots \text{(ii)}$$

$$\therefore Q = \lambda l$$

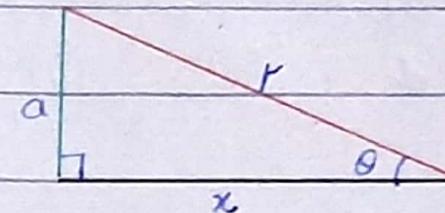
$$E = \frac{1}{4\pi\epsilon_0} \int \frac{d\lambda l}{r^2} \cos\theta \quad (\text{As } \lambda \text{ is also constant})$$

$$E = \frac{\lambda}{4\pi\epsilon_0} \int \frac{dl}{r^2} \cos\theta \quad \dots \text{(iii)}$$

From Figure we get

$$\cos \theta = \frac{R}{H} = \frac{x}{a}$$

So above Equation becomes



$$E = \frac{\lambda}{4\pi\epsilon_0} \int \frac{dl}{x^2} \cdot \frac{x}{a}$$

$$= \frac{\lambda}{4\pi\epsilon_0} \int \frac{dl \cdot x}{x^3}$$

Here $x = \text{Constant}$ Because it has not 'd' of $\frac{1}{x^3}$
∴ Derivation of $dl \sim \frac{1}{x} dl$ \therefore Derivative

$$E = \frac{\lambda x}{4\pi\epsilon_0 x^3} \int dl \dots \text{(iv)}$$

Since

$$dl = 2\pi a$$

above Eq. becomes

$$\therefore E = \frac{\lambda x (2\pi a)}{4\pi\epsilon_0 x^3}$$

$$E = \frac{x (\lambda 2\pi a)}{4\pi\epsilon_0 x^3}$$

$$\therefore \lambda 2\pi a = \text{Total charge enclosed Surface} = Q$$

$$E = \frac{Qx}{4\pi\epsilon_0 r^3} \quad \dots \quad (v)$$

From Right Triangle

$$H^2 = P^2 + B^2$$

$$r^2 = a^2 + x^2$$

$$r = \sqrt{a^2 + x^2}$$

$$r = (a^2 + x^2)^{1/2}$$

Pulling r in above Equation

$$E = \frac{Qx}{4\pi\epsilon_0 (a^2 + x^2)^{3/2}}$$

Here This is the Total Electric Field of a Ring shaped object in which charges are distributed uniformly. It mean that there is no electric Field in the centre of Surface.

$$E = \frac{1}{4\pi\epsilon_0} \frac{Qx}{(x^2 + a^2)^{3/2}}$$

$$\text{Put } x=0$$

SPECIAL CASES

Now Let's Discuss about some special cases.

1: When does the electric Field in the Ring becomes Zero.

If 'x' decreased then 'E' also decrease, when 'x' became zero; E also becomes zero.

Opp. Force \cancel{F}
Opp. Force \cancel{F}
Opp. Opposite \cancel{F}
Opp. Force \cancel{F}
Effect \cancel{F}
 \cancel{F} Cancel \cancel{F}

$$E = \frac{1}{4\pi\epsilon_0} \frac{Q(0)}{(0^2 + a^2)^{3/2}}$$

$$E = 0$$

2. When Does the Electric Field of the Ring becomes Electric Field of a Point charge.

When the Distance of Point 'P' 'x' becomes equal to the Distance of Patch 'r' from Point 'P'. Then Electric Field of Ring becomes the electric field of Point charge.

$$E = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \cdot \frac{x}{r}$$

$$\text{Put } x = r$$

$$E = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \cdot \frac{r}{r}$$

$$E = \frac{Q}{4\pi\epsilon_0 r^2}$$

3. If $x \gg a$ than $a \approx 0$
(a can be neglected)

If Point 'P' situated at the infinity Distance of 'x' Hence 'x' is too large that in front of 'x' we can neglect the Radius 'a' Because in this case 'a' does not have a significant impact on Electric Field.

$$E = \frac{1}{4\pi\epsilon_0} \frac{Qx}{(x^2 + a^2)^{3/2}}$$

$$\text{Put } a \approx 0$$

$$E = \frac{1}{4\pi\epsilon_0} \frac{Qx}{(x^2 + 0)^{3/2}}$$

$$E = \frac{1}{4\pi\epsilon_0} \frac{Qx}{(x^2)^{3/2}}$$

$$E = \frac{1}{4\pi\epsilon_0} \frac{Qx}{x^3}$$

$$E = \frac{Q}{4\pi\epsilon_0 x^2}$$

In the case of a Rod

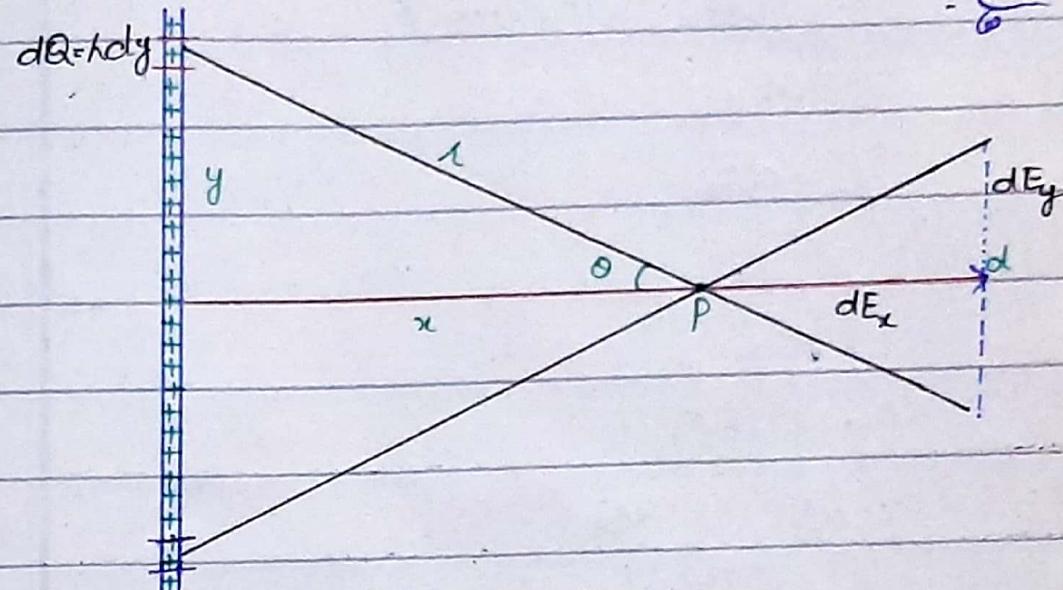
Take a Rod at which uniformly charges are distributed.

Consider a point 'P' lies at a distance of

x. To calculate the electric field on the Rod we have to divide the rod into small patches.

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Topic Digest

Figure: জ্বলন পথে পার্শ রেড



Long Rod of infinite length

To calculate the electric field we have to dissolve the

Point 'P' into its components

dE_x and dE_y .

Here dE_y is zero because they are same in magnitude but opposite direction.

So

$$dE = dE_x + dE_y \dots \underline{i}$$

$$\therefore dE_y = 0$$

$$dE = dE_x$$

Now

$$E = \int dE_x$$

$$E = \int dE \cos \theta$$

$$\therefore dE = \frac{1}{4\pi\epsilon_0} \frac{dQ}{r^2}$$

So

$$E = \int \frac{dQ}{4\pi\epsilon_0 r^2} \cos \theta \dots \underline{\underline{ii}}$$

Here $\frac{1}{4\pi\epsilon_0}$ is constant so

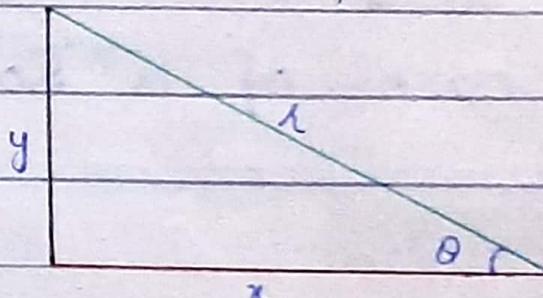
$$E = \frac{1}{4\pi\epsilon_0} \int \frac{dQ}{r^2} \cos \theta \dots \underline{\underline{ii}}$$

$$\therefore dQ = \lambda dl$$

$$E = \frac{1}{4\pi\epsilon_0} \int \frac{\lambda dl}{r^2} \cos \theta$$

As λ is constant

$$E = \frac{\lambda}{4\pi\epsilon_0} \int \frac{dl}{r^2} \cos \theta \dots \underline{\underline{iv}}$$



$$r^2 = x^2 + y^2$$

So above equation becomes

$$E = \frac{\lambda}{4\pi\epsilon_0} \int \frac{dl}{x^2 + y^2} \cos \theta \dots \underline{\underline{v}}$$

Here

From above Figure:

$$\tan\theta = \frac{y}{x} = \frac{x}{y}$$

$$\tan\theta = y$$

$$y = x \tan\theta$$

$$dy = d(x \tan\theta)$$

$$dy = x \sec^2\theta d\theta$$

From above equation. \bar{v}

$$x^2 + y^2 = x^2 + x^2 \tan^2\theta$$

$$= x^2(1 + \tan^2\theta)$$

$$x^2 + y^2 = x^2 \sec^2\theta$$

Now putting these values

$$E = \frac{\lambda}{4\pi\epsilon_0} \int \frac{x \sec^2\theta \cos\theta d\theta}{x^2 \sec^2\theta}$$

$$E = \frac{\lambda}{4\pi\epsilon_0} \int \frac{\cos\theta d\theta}{x}$$

As $\frac{1}{x}$ also constant.

So

$$E = \frac{\lambda}{4\pi\epsilon_0 x} \int \cos\theta d\theta$$

By Applying Limit

$$E = \frac{\lambda}{4\pi\epsilon_0 x} \int_{-\pi/2}^{\pi/2} \cos\theta d\theta$$

$$= \frac{\lambda}{4\pi\epsilon_0 x} \left| -\sin\theta \right|_{-\pi/2}^{\pi/2} \quad \because \int \cos\theta = \sin\theta$$

$$= \frac{\lambda}{4\pi\epsilon_0 x} \left\{ \sin\left(\frac{\pi}{2}\right) - \sin\left(-\frac{\pi}{2}\right) \right\}$$

$$= \frac{\lambda}{4\pi\epsilon_0 x} \left\{ 1 - (-1) \right\}$$

$$= \frac{\lambda}{4\pi\epsilon_0 x} (2)$$

$$E = \frac{\lambda}{2\pi\epsilon_0 x}$$

The above equation tells the electric field of an object of which angle changes while we change the length of an object.

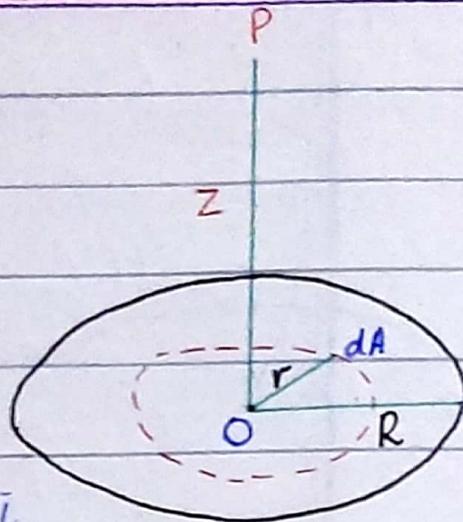
Calculation Of Electric Field Due to Surface Charge Distribution

Consider a plate object in which charges are distributed uniformly. The total radius of the object is R , and a point P lies at a distance of z from the centre of the plate.

To calculate the Electric Field of the object, first we will convert the object into small equal patches. Each patch has Area of dA and the radius of each is r

($\text{for given Area } \rightarrow \text{Electric Field } \propto \frac{1}{r^2}$ (per patch) \therefore Patches $\propto r^2 \therefore E \propto r^2$)

Figure:



As we know that

$$\text{Area} = \text{length} \times \text{width}$$

So the Area of the patch

is

$$A = 2\pi r dr$$

Surface charge Density can be given as:

$$\sigma = \frac{Q}{A}$$

As we are considering a small patch so the surface charge density of this patch

becomes

$$\sigma' = \frac{dQ}{dA}$$

As

$$dQ = \sigma dA$$

As we know that the electric Field due to ring is:

$$E = \frac{1}{4\pi\epsilon_0} \frac{Qx}{(x^2 + a^2)^{3/2}}$$

From the Figure the above equation becomes:

$$dE = \frac{1}{4\pi\epsilon_0} \frac{z dQ}{(z^2 + r^2)^{3/2}}$$

To calculate the electric field,

$$E = \int dE$$

$$= \frac{1}{4\pi\epsilon_0} \int \frac{z dQ}{(z^2 + r^2)^{3/2}}$$

$$\therefore dQ = \sigma dA$$

$$\therefore dA = 2\pi r dr$$

$$\therefore dQ = \sigma (2\pi r dr)$$

Hence

$$E = \frac{1}{4\pi\epsilon_0} \int \frac{z \sigma (2\pi r) dr}{(z^2 + r^2)^{3/2}}$$

Since $z\sigma 2\pi r$ is constant so

$$= \frac{z\sigma (2\pi r)}{4\pi\epsilon_0} \int \frac{r dr}{(z^2 + r^2)^{3/2}}$$

$$E = \frac{z\sigma}{2\epsilon_0} \int \frac{r dr}{(z^2 + r^2)^{3/2}}$$

Applying Limit $0 \rightarrow R$

So we have

$$E = \frac{z\sigma}{2\epsilon_0} \int_0^R \frac{r dr}{(z^2 + r^2)^{3/2}}$$

Now Using R

$$= \int_0^R \frac{r dr}{(z^2 + r^2)^{3/2}}$$

x-ing And Dividing by 2.

$$= \frac{1}{2} \int_0^R (z^2 + r^2)^{-3/2} 2(r dr)$$

(Power Rule)

$$= \frac{1}{2} \left| \frac{(z^2 + r^2)^{-3/2 + 1}}{-\frac{3}{2} + 1} \right|_0^R$$

$$= \frac{1}{2} \left| \frac{(z^2 + r^2)^{1/2}}{-\frac{1}{2}} \right|_0^R$$

$$E = - \left| \frac{1}{(z^2 + r^2)^{1/2}} \right|_0^R$$

\Rightarrow (by taking limit \rightarrow $r \rightarrow 0$)

$$= - \left\{ \frac{1}{(z^2 + 0^2)^{1/2}} - \frac{1}{(z^2 + R^2)^{1/2}} \right\}$$

$$= - \left\{ \frac{1}{z} - \frac{1}{(z^2 + R^2)^{1/2}} \right\}$$

$$= - \left\{ \frac{1}{z} - \frac{1}{(z^2 + R^2)^{1/2}} \right\}$$

Now Put. this Value at the place of rdr . So, we have

$$E = \frac{z\sigma}{2\epsilon_0} \left\{ \frac{1}{z} - \frac{1}{(z^2 + R^2)^{1/2}} \right\}$$

$$E = \frac{\sigma}{2\epsilon_0} \left\{ \frac{z}{z} - \frac{z}{(z^2 + R^2)^{1/2}} \right\}$$

$$E = \frac{\sigma}{2\epsilon_0} \left[1 - \frac{z}{(z^2 + R^2)^{1/2}} \right]$$

This is the total electric Field of a surface at which charges are distributed uniformly.

Special Case:

When Will be Electric Field Maximum?

If the Point 'P' lies on the centre and the distance of Point 'P' becomes zero from the centre. i.e $(z = 0)$ then Maximum electric field will be produced.

$$E = \frac{\sigma}{2\epsilon_0} \left[1 - \frac{z}{(z^2 + R^2)^{1/2}} \right]$$

As $z = 0$

so

$$E = \frac{\sigma}{2\epsilon_0} \left[1 - \frac{0}{(0^2 + R^2)^{1/2}} \right]$$

$$E = \frac{\sigma}{2\epsilon_0} \quad (1)$$

$$E = \frac{\sigma}{2\epsilon_0}$$

we can also explain the same case in this manner, that if the Radius 'R' of the object becomes too much greater than in front of 'R' 'z' becomes so small so that we can

neglect the 'z'.

$R \gg z$ then

$z \approx 0$

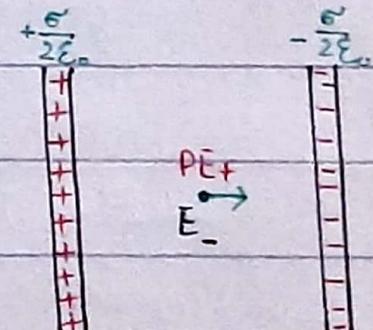
Then Maximum electric Field will be produced.

In Cases Of Point Charge

Case-I

Consider a point charge q (e.g. 6^{th} +ve Gaussian Cell)

that is placed between two charged capacitor plates. As



Here we can see that the positive capacitor plate is repelling the point charge towards negative charge plate and negative charge plate already attracting the point charge towards itself. So, At centre of capacitor plates the electric field can be given as:

$$E_{\text{TOTAL}} = E_+ + E_-$$

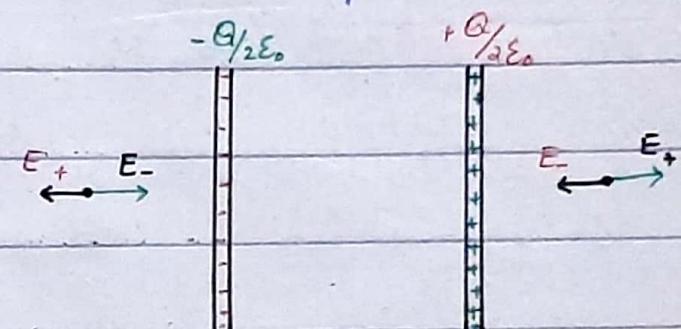
$$E_{\text{total}} = \frac{\sigma}{2\epsilon_0} + \frac{\sigma}{2\epsilon_0}$$

$$E_{\text{total}} = \frac{\sigma}{\epsilon_0}$$

This is maximum electric field that is present in the centre of plates.

Case-II

If the point charge lies at left or right side of a charged plate of capacitor. For Example:



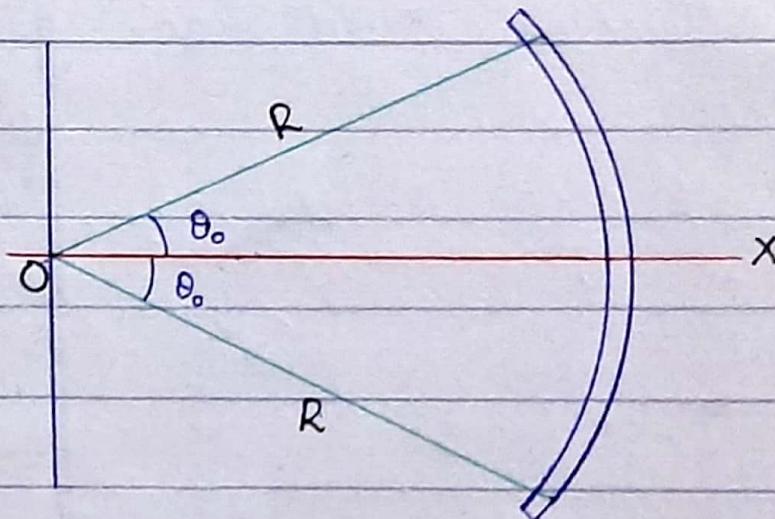
Here in such case the electric field intensity is given as:

$$E_{\text{TOTAL}} = E_+ = E_-$$

$$E = \frac{\sigma}{2\epsilon_0}$$

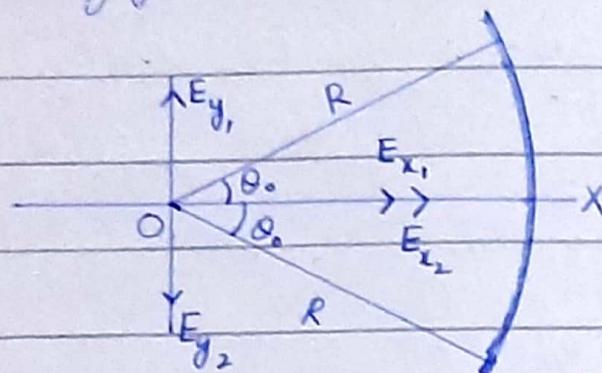
*P*ROBLEMS

1: A thin rod is bent into the shape of an arc of a circle, of radius R carries a uniform charge per unit length λ . The arc subtends a total angle $2\theta_0$. Symmetric about the x -axis as shown in figure.



Determine the electric field E at point O .

By using the figure:



By Resolving the axis

into their components.

So we have two

components of x-axis which

E_{x_1} and E_{z_2} which lies

on the same axis and

in same direction. but in case of y-axis y-axis

also have two Components E_{y_1} and E_{y_2} . But they

will be cancel each other effect because they

have same magnitude but opposite direction.

So here the total Electric Field becomes

$$E = E_x + E_y$$

$$\therefore E_y = 0$$

So

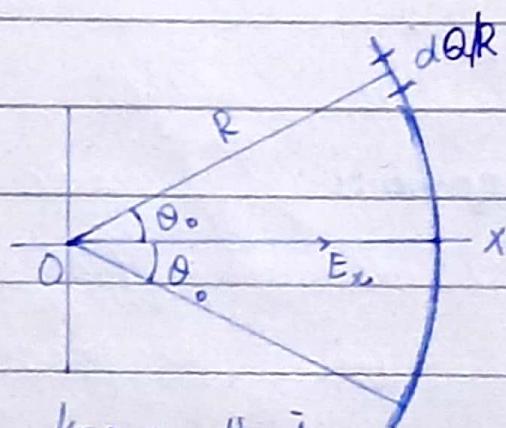
$$E = E_x$$

So

In order to calculate the electric field, we use E_x . Here we will consider a small patch of the curve. As we know that

$$E = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2}$$

From Figure



As we know that

$$\lambda = \frac{Q}{l}$$

This is the Linear charge density for a small patch it will become

$$\lambda = \frac{dQ}{l}$$

$$R \quad \therefore l = R$$

$$dQ = \lambda \cdot R$$

Now using

$$E = \int dE_x$$

$$\therefore E = \int \frac{1}{4\pi\epsilon_0} \frac{Q}{R^2}$$

In case of E_x the above eq. becomes

$$E = \int \frac{1}{4\pi\epsilon_0} \frac{\lambda R}{R^2} \cos(\theta_0) d\theta \quad \therefore dQ = \lambda R$$

$$\therefore R = r \quad \theta = 2\theta_0$$

$$= \frac{\lambda R}{4\pi\epsilon_0 R^2} (2) \int \cos(\theta_0) d\theta$$

$$= \frac{\lambda}{4\pi\epsilon_0 R} (2) \int \cos \theta_0 d\theta$$

$$= \frac{2\lambda}{4\pi\epsilon_0 R} \int \cos \theta_0 d\theta$$

$$= \frac{\lambda}{2\pi\epsilon_0 R} \int \cos \theta_0 d\theta$$

From statement it is clear that θ changes from 0 to θ_0 . So in above equation we can apply the limit as:

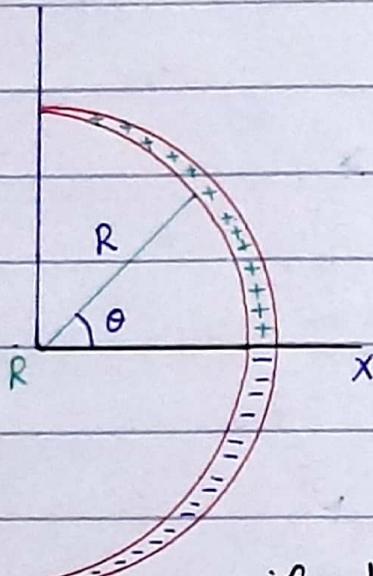
$$E = \frac{\lambda}{2\pi\epsilon_0 R} \int_0^{\theta_0} \cos\theta_0 d\theta$$

$$= \frac{\lambda}{2\pi\epsilon_0 R} \left| \sin\theta \right|_0^{\theta_0}$$

$$= \frac{\lambda}{2\pi\epsilon_0 R} \left| \sin(\theta_0) - \sin(0) \right|$$

$$E = \frac{\lambda}{2\pi\epsilon_0 R} \sin\theta_0$$

2- A thin glass rod is a semicircle of radius R as show below.

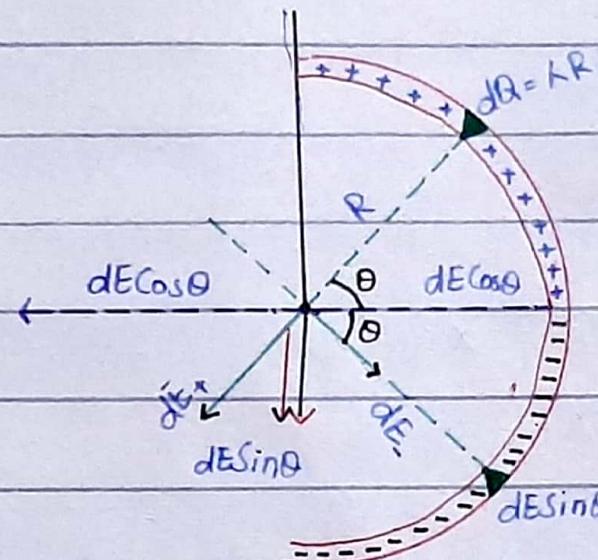


A charge is no uniformly distributed along the rod with a linear charge density given by $\lambda = \lambda_0 \sin \theta$ where λ_0 is a positive constant and point P lies at the centre of the ring.

Here In

From the figure we can understand that there is a point charge which contain positive charge. The point charge is repelling by the

positive charges and attracting by the negative charges. Here we can see the figure:



Here in the figure we can see that the point charge is moving downward towards negative charges. Means when we will resolve the axis into their components then the x-axis's components due to same magnitude but opposite direction will cancel each other effect. Hence, we will have

$$dE = dE_+ = dE_-$$

As

$$\therefore dE_{\text{netx}} = 0$$

So we have

$$dE_{\text{nety}} = 2dE \sin \theta$$

So we know

$$E = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2}$$

But In case of dE_{nety} we know

$$\lambda = \frac{dQ}{R} \Rightarrow dQ = \lambda R$$

$$r = R$$

So, by Putting the values, we have

$$E = \int dE_{\text{nety}}$$

$$= \int \frac{1}{4\pi\epsilon_0} \frac{\lambda R}{R^2} 2 \sin \theta$$

$$= \frac{2}{4\pi\epsilon_0} \frac{1}{R^2} \int \lambda \sin\theta d\theta$$

$$= \frac{1}{2\pi\epsilon_0} \frac{1}{R} \int \lambda \sin\theta d\theta$$

$$\therefore \lambda = \lambda_0 \sin\theta$$

$$= \frac{1}{2\pi\epsilon_0} \frac{\lambda_0}{R} \int (\lambda_0 \sin\theta) (\sin\theta) d\theta$$

$$= \frac{1}{2\pi\epsilon_0} \frac{\lambda_0}{R} \int \lambda_0 \sin^2\theta d\theta$$

$$= \frac{1}{2\pi\epsilon_0} \frac{(\lambda_0)}{R} \int \sin^2\theta d\theta$$

Now, apply the limit of λ_0 at $\pi/2$ in the above equation.

$$= \frac{\lambda_0}{2\pi\epsilon_0 R} \int_0^{\pi/2} \sin^2\theta d\theta$$

$$= \frac{\lambda_0}{2\pi\epsilon_0 R} \int_0^{\pi/2} \left(\frac{1}{2} - \frac{\cos 2\theta}{2} \right) d\theta$$

$$= \frac{\lambda_0}{2\pi\epsilon_0 R} \left[\frac{\theta}{2} - \frac{\sin 2\theta}{4} \right]_0^{\pi/2}$$

$$= \frac{\lambda_0}{2\pi\epsilon_0 R} \left\{ \left(\frac{\pi/2}{2} - \frac{\sin 2(\pi/2)}{4} \right) - \left(\frac{0}{2} - \frac{\sin(2(0))}{4} \right) \right\}$$

$$= \frac{\lambda_0}{2\pi\epsilon_0 R} \left\{ \frac{\pi}{4} - \frac{\sin(\pi)}{4} - (0) \right\}$$

$$= \frac{\lambda_0}{2\pi\epsilon_0 R} \left\{ \frac{\pi}{4} - 0 \right\}$$

$$= \frac{\lambda_0}{2\pi\epsilon_0 R} \left\{ \frac{\pi}{4} \right\}$$

$$E = - \frac{\lambda_0}{8\pi\epsilon_0 R}$$

Since negative sign indicates the direction of point charge to the negative charge capacitor.

CHAPTER 22

Gauss's Law

ELECTRIC FLUX

Gauss's Law involves the concept of electric flux. That can be defined as:

"Total number of electric fields passing normally through a certain area is called electric flux."

→ Electric Flux is represented by ϕ_e .

Formula:

$$\phi_e = \vec{E} \cdot \vec{A}$$
$$= EA \cos \theta$$

The electric flux depends upon

- Electric intensity
- Surface Area
- Orientation of Surface

Flux means **flow**.

- It is a **scalar quantity**. (Have no direction)
- Its SI Unit is $Nm^2 C^{-1}$

Cases

Now let's discuss some case in which we will see that whenever the electric flux will maximum and whenever the electric flux will minimum.

Case-I: Maximum Flux:

When the ^{direction of} vector **Area** is held **parallel** to the **field lines**.

Then the electric flux will be maximum.

$$\phi_e = EA \cos \theta$$

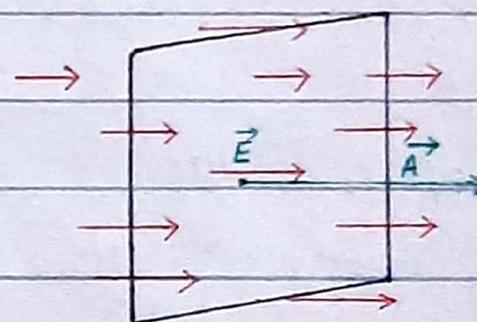
$$\because \theta = 0$$

$$\phi_e = EA \cos(0)$$

$$\therefore \cos(0) = 1$$

$$\phi_e = EA(1)$$

$$\phi_e = EA$$



Case-II: Minimum Flux:

When the ^{direction of} vector **Area** is held **perpendicular** to the **electric field lines**. Then the electric flux will be minimum. As:

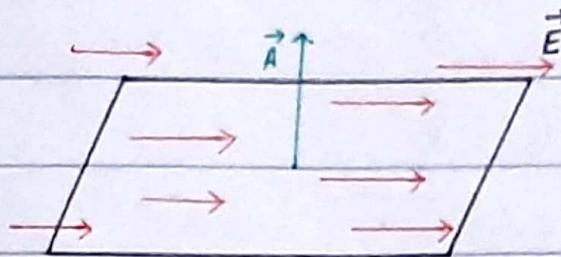
$$\phi_e = EA \cos 0^\circ$$

$$\because \theta = 90^\circ$$

$$\phi_e = EA \cos (90^\circ)$$

$$\because \cos 90^\circ = 0$$

$$\phi_e = EA(0)$$



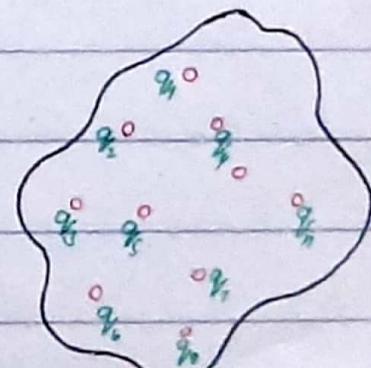
$$\phi_e = 0$$

ELECTRIC FLUX THROUGH A Surface Enclosing CHARGE

Consider an irregular surface in which point charges

Q $q_1, q_2, q_3, \dots, q_n$ are distributed arbitrarily. Each charge will act as an independent source of electric flux.

The flux through the surface due to point charge q_1 is $\phi_1 = \frac{q_1}{\epsilon_0}$



Similarly flux through other charges can be given as:

$$\phi_2 = \frac{q_2}{\epsilon_0}$$

$$\phi_3 = \frac{q_3}{\epsilon_0}$$

$$\vdots \quad \vdots$$

$$\phi_n = \frac{q_n}{\epsilon_0}$$

The total electric flux passing through the whole closed surface will be

$$\phi_1 + \phi_2 + \phi_3 + \dots + \phi_n = \frac{q_1}{\epsilon_0} + \frac{q_2}{\epsilon_0} + \dots + \frac{q_n}{\epsilon_0}$$

$$= \frac{1}{\epsilon_0} (q_1 + q_2 + \dots + q_n)$$

$$= \frac{1}{\epsilon_0} \sum_{i=1}^n q_i$$

$$\therefore \sum_{i=1}^n q_i = Q$$

$$\phi_e = \vec{E} \cdot \vec{A} = \frac{1}{\epsilon_0} Q$$

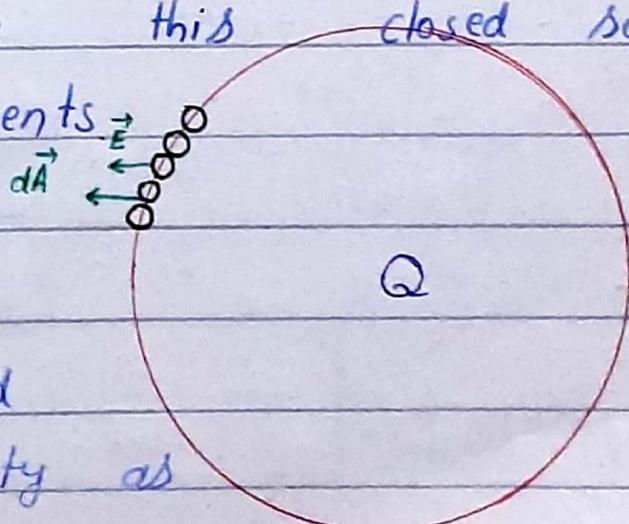
Here Q is the total charge enclosed by the surface

ELECTRIC FLUX THROUGH A CLOSED Surface

Consider a sphere of radius r due to a point charge q at its centre.

In order to calculate the electric flux through whole surface of sphere, we divide this closed surface into n number of small elements.

The electric flux through each element is given as $d\vec{A}$ and the electric field lines intensity as



E.

So the flux through an element can be given as:

$$\vec{E} \cdot d\vec{A} = \frac{1}{\epsilon_0} Q$$

$$E \cdot A = \oint \vec{E} \cdot d\vec{A}$$

$$= \oint E dA$$

$$= E \oint dA$$

$$= E(A)$$

The Area $A = 4\pi r^2$

so

$$= E(4\pi r^2)$$

Now

$$E(4\pi r^2) = \frac{1}{\epsilon_0} Q$$

$$E = \frac{1}{4\pi \epsilon_0} \frac{Q}{r^2} \Rightarrow \text{Coulomb's Law.}$$

Cases

Now let's discuss the case in which we will find that when will be the electric flux through the surface zero.

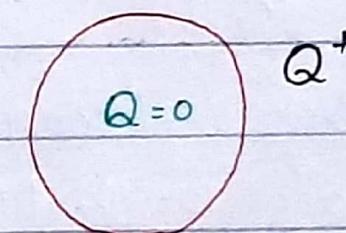
In the above figure if the charge enclosed by the surface is zero means charges lies outside the surface.

Inside $Q=0$

$$E = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2}$$

$$= \frac{0}{4\pi\epsilon_0 r^2}$$

$$E = 0$$

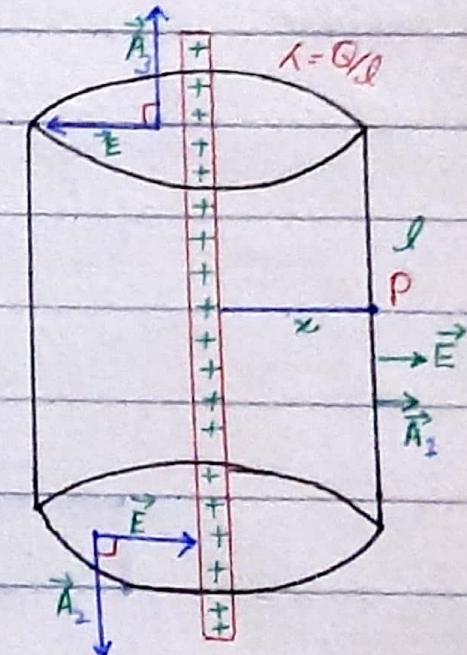


ELECTRICAL FLUX THROUGH Cylinder

Consider a cylinder in which a wire is attached to produce electric field lines. We can find out the direction of electric field lines by using Right Hand Rule.

In order to calculate the electric flux through the cylinder we take a point 'P' at a distance of 'x' from the wire.

If we want to calculate the electric flux through the cylinder then, we have to calculate the flux through three sides of the cylinder.



As we know

$$\vec{E} \cdot \vec{A} = \frac{1}{\epsilon_0} Q_{\text{enclosed}}$$

As

$$\lambda = \frac{Q}{l}$$

And

$$Q = \lambda l$$

enclosed

The electric Flux through the curved surface is

$$\phi_1 = \vec{E} \cdot \vec{A}_1 \dots (i) \quad (\text{Here } \cos 0^\circ \text{ (} \theta=0 \text{)} \text{ so Maximum output})$$

Similarly electric Flux through upper flat surface will be

$$\phi_3 = \vec{E} \cdot \vec{A}_3 \cos 90^\circ \quad (\text{Because } \vec{E} \text{ and } \vec{A} \text{ are L})$$

$$\phi_3 = 0 \dots (ii)$$

Similarly in the lower flat surface, the electric Flux will be

$$\phi_2 = \vec{E} \cdot \vec{A}_2 \cos 90^\circ$$

$$\phi_2 = 0 \dots (iii)$$

By adding the all equations we have

$$\vec{E} \cdot \vec{A}_1 + \vec{E} \cdot \vec{A}_2 + \vec{E} \cdot \vec{A}_3 = \frac{1}{\epsilon_0} Q_{\text{enclosed}}$$

$$\therefore Q_{\text{enclosed}} = \lambda l$$

$$\vec{E} \cdot \vec{A}_2 = \frac{1}{\epsilon_0} \lambda l$$

As

$$\text{Area} = \text{length} \times \text{width} = l \times 2\pi r = 2\pi r l$$

Here

Here $\therefore r = x$

$$\vec{E} (2\pi x l) = \frac{1}{\epsilon_0} \lambda l$$

$$\vec{E} = \frac{\lambda l}{2\pi \epsilon_0 x l}$$

$$\boxed{\vec{E} = \frac{\lambda}{2\pi \epsilon_0 x}}$$

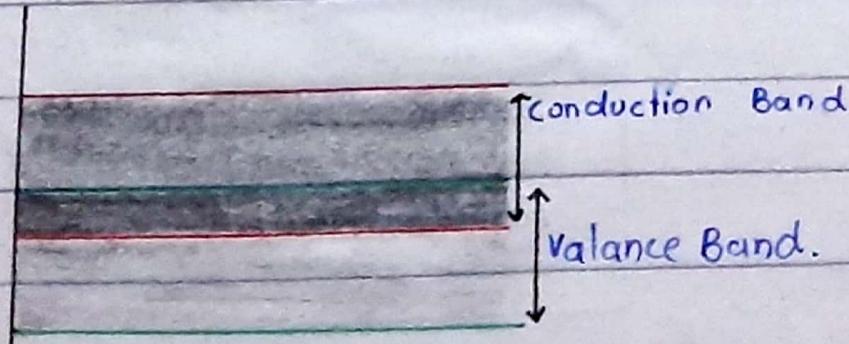
written by

Some Definitions about Practice No. 02:

Those material which conduct electricity easily, called **Conductors**: conductor. e.g copper, silver. Their conductivities are of the order of $10^7 (\Omega m)^{-1}$

Conductors are those materials which have large number of free electrons for electrical conduction. In term of Energy Band the valence and conduction bands are overlapped. i.e. there is no distinction between these two bands.

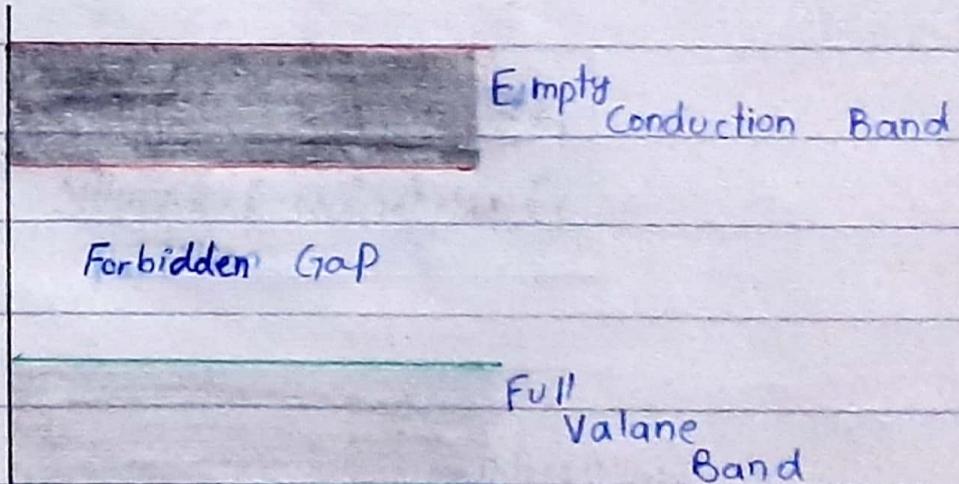
Energy Band:



Insulators:

In insulators the electrons are tightly bound in their Valance shell and they are not free to move. In these materials the Valance band is completely filled and the Conduction band is Empty. So there is a Large Energy Gap between the conduction band and Valance band. And due to this difference they cannot move from Valance band to the conduction band.

Energy Band:



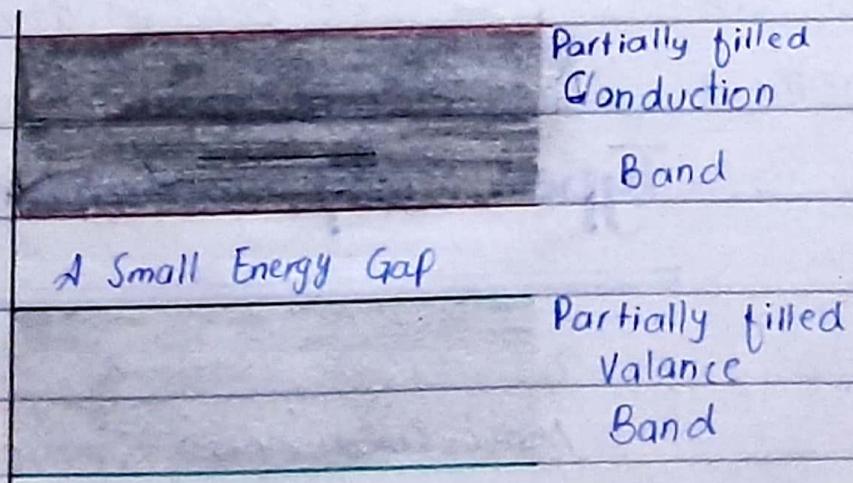
: A Large Difference between Conduction Band and Valance Band.

Semi-Conductors:

They are those materials which have Partially Filled Valence band and Partially filled Conduction band. A very narrow energy gap between the two bands exists and normally order of 1eV.

These materials are not conductors nor insulators but it depends on condition that they behave like conductors or insulators.

Energy Gap:



: A Small Gap between Valence & Conduction Gap

A small energy gap means by providing a small amount of energy the electrons can jump from valence band to conduction band easily.

Now there is a question that why we use semi-conductors in our mostly experiments instead of conductors while conductors conducts electricity easily on low temperature?

In such cases, we use semi-conductors because we can control the conductivity of semi-conductors.

But how we control the conductivity of semi-conductors? we can control the conductivity of semi-conductors by a method called Dopping

Types Of Semi-Conductors

There are two types of semi-conductors which are:

Extrinsic Semiconductors:

The impure semiconductor (in which impurity is added) is called extrinsic semi-conductors.

Intrinsic Semiconductors:

A semiconductor in its pure form is called intrinsic Semiconductors.

The electrical behaviour of semiconductor purely depends on the purity of semiconductors.

Example:

→ Silicon

→ Germanium

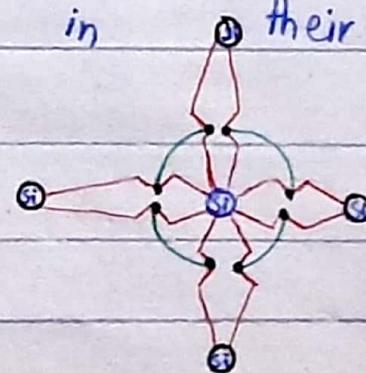
Dopping :

The process of adding selected impurity in pure semiconductors is called doping.

→ The impurity atoms are added in 1 to 10^6 ratio.

→ Due to addition of impurity always extrinsic semiconductors generate.

→ Since semiconductors belong to the 4th group of periodic Table. Means they have 4 electrons in their valance shell and form four bonds with their neighbour atoms. This effectively allocates 8 electrons to the outermost orbit of each atom which is a stable state. Due to these covalent bonds the electrons are bound in their respective shells.



N-Type Material: (Previous Fig: is its example).

When a silicon material is dopped with pentavalent impurity i.e. atom from the fifth group like

⇒ Arsenic

⇒ Antimony

⇒ Phosphorus etc

four electrons of the impurity atom make four covalent bonds with the silicon atoms and the fifth electron remain unbounded.

This become the free electron in the crystal. This type of dopped or extrinsic semiconductors is called N-Type semiconductor.

Donor The impurity atom is called donor atom because it donates a free electron. This electron is thermally excited into the conduction band.

P-Type Semiconductor:

When a silicon crystal is doped with a Trivalent impurity. i.e.

⇒ Aluminum

⇒ Boron

⇒ Indium

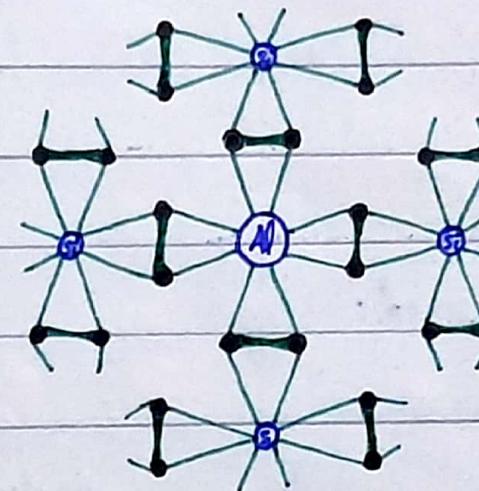
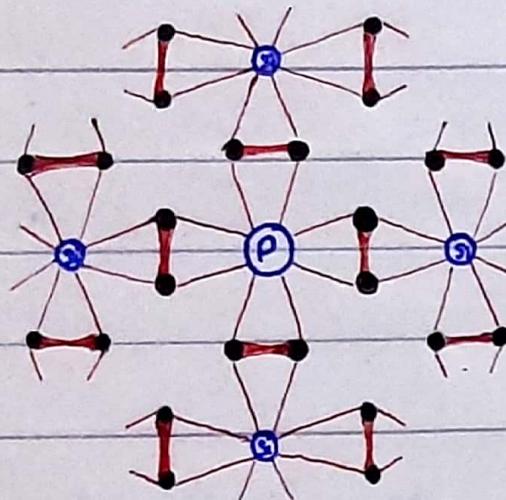
⇒ Gallium

Three Valence electrons of the impurity atom three covalent bonds with the silicon atoms, there is a missing electron to form the forth covalent bond with the forth neighbor silicon atom.

In other words, there is a Vacancy in that bond. This Vacancy is called the hole and is considered to have a positive charge. This Vacancy can accommodate an electron. This type of semiconductor is called

P-type Semiconductor.

Acceptor: The impurity atom in this case is called the acceptor impurity because it can accept an electron from a nearby silicon atom thus by creating a hole in the Valence band.



States:

Electrons in single isolated atom have discrete energies and are bound to their nuclei. In solid, when large number of atoms, say N come close enough to each other to form solid, each energy level splits up into N sub levels, called states. This action is due to the action of the forces exerted due to other atoms in solid state. These energy state are discrete but so close to each other they look like continuous Energy Band.

Bands

Now let's discuss about bands. Basically there are three bands of Materials.

- Valence Band
- Conduction Band
- Forbidden Band

Valence Band:

The electrons in the outermost shell are called Valence electrons and the energy band formed by these electrons is called Valence Band.

This band may be either completely filled or partially filled but cannot be empty.

Conduction Band:

The band above the valence band is called conduction band. In this band electrons move freely and conduct electric current. The electrons in this band are called conduction electrons or free electrons. This band may be completely filled or partially filled.

Forbidden Energy Band:

There is a range of energy states which cannot be occupied by electrons. These are called forbidden energy states and its range is called forbidden energy gap.

Gauss's Law Applications

Sphere:

A Sphere is a three-dimensional geometrical object that is analogue to two dimensional circle.

A Sphere is a set of points that are all at the same distance r from a given point in three-dimensional space. That given point is the centre of the sphere, and r is the sphere's radius.

Volume

$$\frac{4}{3}\pi r^3$$

Surface Area

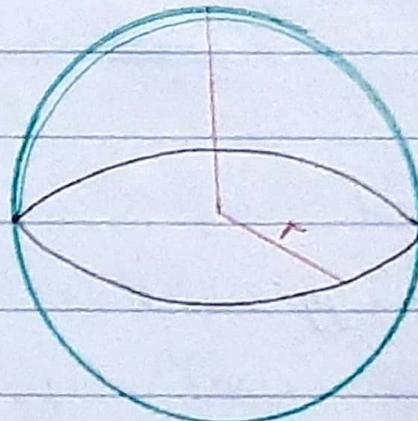
$$4\pi r^2$$

Euler char

$$2$$

Symmetry Group

$$O(3)$$



When we study a sphere in physics, we study about two types of sphere:

- Hollow Sphere

- Solid Sphere

And Solid Sphere is classified into two more classes which are

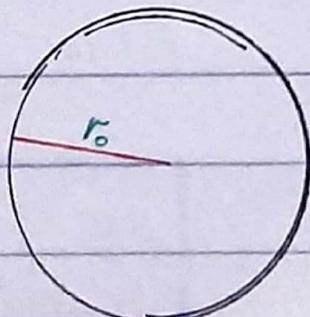
- Conductive Sphere

- Non-Conductive Sphere

Hollow Sphere:

Let take a hollow sphere. Since the sphere is hollow means there is no charge lies at centre of sphere. when we will apply electric Field around it all charges will be distributed on the surface of the sphere. As:

This is a hollow sphere of radius 'r' when we apply electric field around it all charges will come



On the surface of Sphere. To calculate the Electric Field, we imagine a Gaussian surface of radius r outside the sphere.

The total charge enclosed by the surface is Q . In order to calculate the Electric Field the electric flux we have

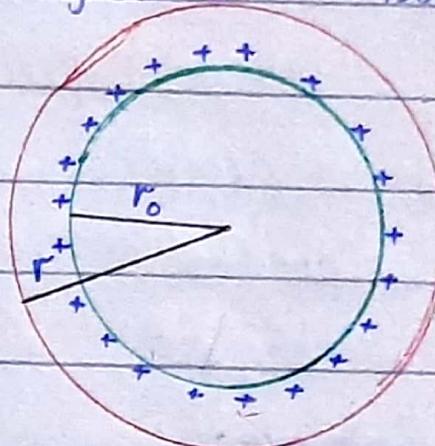
$$E \cdot dA = \frac{1}{\epsilon_0} Q_{en}$$

$$E(A) = \frac{1}{\epsilon_0} (Q)$$

$$\therefore A = 4\pi r^2$$

$$E(4\pi r^2) = \frac{Q}{\epsilon_0}$$

$$E = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2}$$



We have the above electric field if the Radius of Gaussian surface is greater than the Radius of sphere ($r > r_0$)

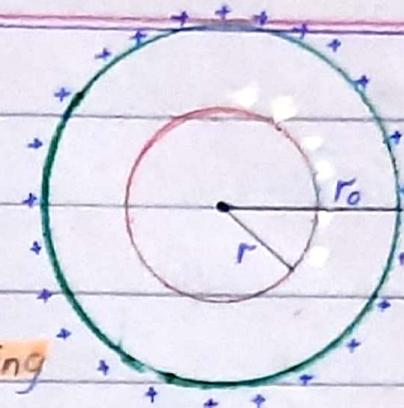
NOTE

When a Sphere is Gaussian Surface

close to Gaussian Surface, if a point is at the center of the sphere, then the electric field is zero.

If $r < r_0$

It means if we consider a Gaussian surface which has smaller radius than the radius of sphere. Means we are considering the surface inside sphere.



If we do this and now want to calculate the electric field of sphere, we have

$$E \cdot dA = \frac{1}{\epsilon_0} Q_{\text{en}}$$

$\therefore Q = 0$ (inside the hollow sphere)

So

$$\oint \vec{E} \cdot d\vec{A} = \frac{1}{\epsilon_0} Q_{\text{en}}$$

$$E(A) = \frac{1}{\epsilon_0} (0)$$

$$E(4\pi r^2) = 0$$

$$E = 0$$

Due to absence of charge the electric field inside the sphere is zero.

→ It is a point to be noted that there is no difference between Hollow Sphere and Solid Conductive Sphere. In both terms charge store at outside the sphere. Means, inside the sphere the charge is zero for both hollow and solid conductive sphere.

Non-conductive Solid Sphere

In case of Non-conductive solid sphere, if we want to calculate the electric field of solid sphere, first of all we must be know that either the charge is uniformly distributed or Non-uniformly distributed.

Uniformly Charge Distribution

Let's consider a Non-conductive solid sphere in which charges are distributed uniformly

distributed. In order to calculate the electric field of the sphere we study three cases which are given below:

Case-I:

In Case-I consider a gaussian surface outside the sphere. Means

$$r > r_0$$

Hence the charge enclosed by the gaussian surface is Q_{en} . As we know,

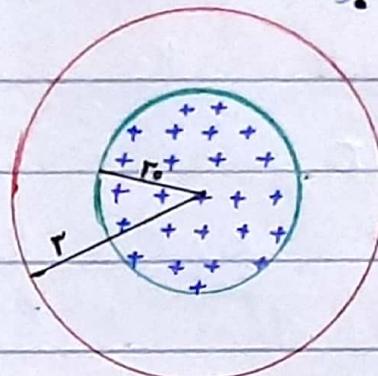
$$\oint \vec{E} \cdot d\vec{A} = \frac{1}{\epsilon_0} Q_{en}$$

$$E(A) = \frac{1}{\epsilon_0} Q_{en}$$

$$\therefore A = 4\pi r^2$$

$$E(4\pi r^2) = \frac{1}{\epsilon_0} Q_{en}$$

$$E = \frac{1}{4\pi\epsilon_0} \frac{Q_{en}}{r^2}$$



In the above equation we can see that the all factors are constant except r^2 . Mean electric field intensity can be changed by changing the value of r^2 . Here

$$E \propto \frac{1}{r^2}$$

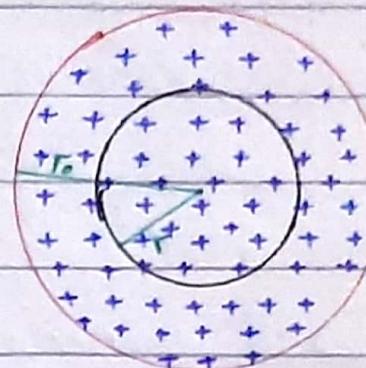
Case-II

In this case consider a gaussian surface inside the sphere. Means

$$r < r_0$$

Hence, we know that

$$\oint \vec{E} \cdot d\vec{A} = \frac{1}{\epsilon_0} Q_{\text{en}} \quad \dots \dots \quad (i)$$



Here To calculate the electric field first of all we have to calculate the value of Q_{en} .

As we know that the charge Density of the sphere can be given as

$$\rho_E = \frac{Q}{V}$$

As we know that

$$\text{Volume of sphere} = \frac{4}{3} \pi r^3$$

so we have

$$\rho_E = \frac{Q_{\text{en}}}{\frac{4}{3} \pi r^3}$$

• r (radius of Sphere)

• Q (Total charge enclosed by Sphere.)

This is the charge density for the whole sphere.

If we want to calculate the charge density of Gaussian surface. Since the charge is uniformly distributed so we have:

$$\rho_E = \frac{Q_{\text{en}}}{\frac{4}{3}\pi r^3}$$

r (radius of Gaussian Surface)

Q_{en} (charge enclosed by Gaussian Surface)

Since charge is uniformly distributed so we can write

Density of sphere = Density of Gaussian Surface

$$\frac{Q_{\text{en}}}{\frac{4}{3}\pi r_0^3} = \frac{Q_{\text{en}}}{\frac{4}{3}\pi r^3}$$

$$Q_{\text{en}} = \frac{Q \times r^3}{r_0^3}$$

Put this value in eq. i

$$\oint \vec{E} \cdot d\vec{A} = \frac{1}{\epsilon_0} Q_{\text{en}}$$

$$E(A) = \frac{1}{\epsilon_0} \frac{Q}{r^3}$$

$$\therefore A = 4\pi r^2$$

$$E(4\pi r^2) = \frac{1}{\epsilon_0} \frac{r^2}{r^3} Q$$

$$E = \frac{1}{4\pi\epsilon_0} \frac{r}{r^3} Q$$

In this case we can see that all the factors are constant except Radius of Gaussian surface 'r'. Mean Electric Field depends on 'r'

$$E \propto r$$

Case-III

In this case we will use the above case, as we said the Electric Field depends upon radius of Gaussian surface.

$$E \propto r$$

98 we started to increase the radius of gaussian surface and at one stage it becomes equal to r_0 . Then the electric field can be.

$$E = \frac{1}{4\pi\epsilon_0} \frac{r}{r_0^3} Q$$

As $r = r_0$

so

$$E = \frac{1}{4\pi\epsilon_0} \frac{\left(\frac{r_0}{r}\right) Q}{r_0^3}$$

$$E = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2 r_0^2}$$

This is the electric Field.

Graph:

We can also express these cases in term of graph. As, we have seen in case # 01

$$E \propto \frac{1}{r^2}$$

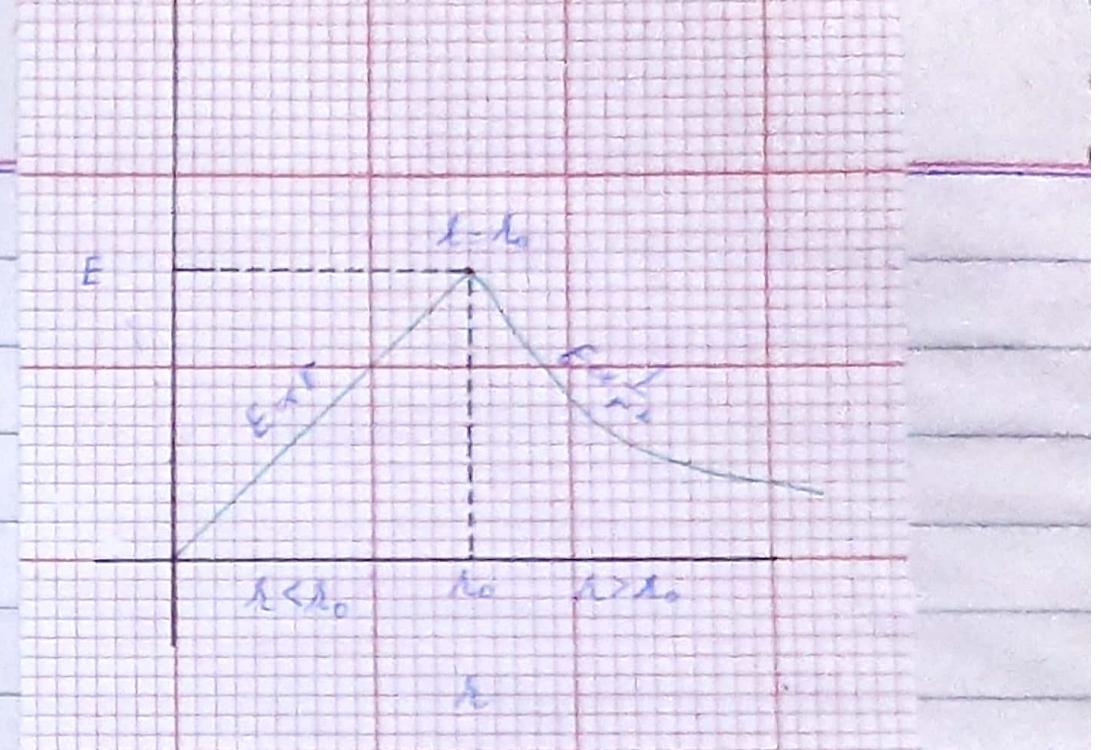
And in Case # 02

$E \propto r$

And in Case # 03

$r = r_0$

So there graph can be



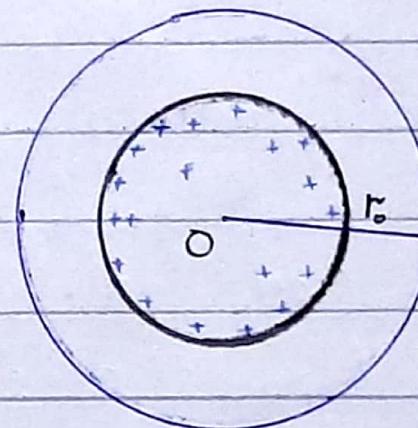
Non-Uniformly Charge Distribution

We have talked about the Non-conductive Solid Sphere and have learnt that in these spheres charges may be distributed uniformly or non-uniformly. Now we are going to talk about Non-conductive Solid sphere in which charges are distributed Non-uniformly.

Consider a non-conductive solid sphere having the Radius of r_0 . The charge density of the sphere can be given as:

$$\rho = \alpha r^2$$

The above equation tells that if the radius of the sphere increases then the charge density of the sphere also increases and vice-versa.



If the radius of the sphere is r_0 and the charge enclosed by the sphere is Q . In order to calculate the Electric Field Intensity of the sphere, we will consider a gaussian surface. Here if we consider the gaussian surface outside the sphere, it means all

the charges are enclosed in a sphere. Means Result will be same as result was for the conductive Solid or Hollow sphere. ($\lambda > r_0$)

But Here charge is non-uniformly distributed, so we have to divide the sphere into small patches. Make sure, Each patch of sphere has a thickness of dr .

From above equation

$$S_E = \alpha r^2$$

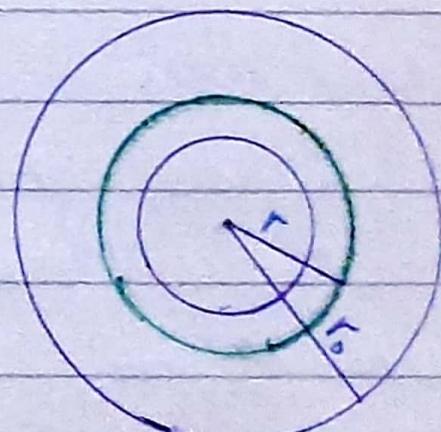
In order to find the Electric Field Intensity we have to find the value of

- α
- E

Here from the Figure:

We have

$$-S_E = \frac{Q}{V}$$



$$Q = \rho_E V$$

By Applying integration

$$Q = \int_0^R \rho_E dV$$

we have

$$dV = 4\pi r^2 dr$$

$$\rho_E = \alpha r^2$$

so

$$Q = \int_0^{R_0} \alpha r^2 4\pi r^2 dr$$

$$= \int_0^{R_0} \alpha 4\pi r^4 dr$$

$$= \alpha 4\pi \int_0^{R_0} r^4 dr$$

$$Q = \alpha 4\pi \left[\frac{r^5}{5} \right]_0^{R_0}$$

$$Q = \alpha 4\pi \left\{ \frac{R_0^5}{5} - \frac{0}{5} \right\}$$

$$Q = 4\pi \alpha \frac{R_0^5}{5}$$

$$\alpha = \frac{5Q}{4\pi R_0^5}$$

Since, the charge is non-uniformly distributed but depends on the Radius of sphere **patched** for **Each patch**.

Now we are going to calculate the total charge enclosed by the gaussian surface.
I have

$$Q_{\text{en}} = \int_0^R \rho_E dV$$

Similarly Now we have

$$Q_{en} = \frac{4\pi (5Q)}{4\pi r_0^2} \frac{r_0^5}{8}$$

$$Q_{en} = \int_0^r \alpha r^2 4\pi r^2 dr$$

$$= \int_0^r 4\pi \alpha r^4 dr$$

$$= 4\pi \alpha \int_0^r r^4 dr$$

$$= 4\pi \alpha \left[\frac{r^5}{5} \right]_0^r$$

$$= 4\pi \alpha \left\{ \frac{r^5}{5} - \frac{0}{5} \right\}$$

$$Q_{en} = 4\pi \alpha \frac{r^5}{5}$$

$$\therefore \alpha = \frac{5Q}{4\pi r_0^5}$$

$$Q_{en} = \frac{r_0^5}{r_0^5} Q$$

According to Gauss's Law

$$\oint \mathbf{E} \cdot d\mathbf{A} = \frac{1}{\epsilon_0} Q_{en}$$

$$\because \mathbf{E} \cdot \mathbf{A} = \frac{1}{\epsilon_0} Q_{en}$$

$$\because A = 4\pi r^2$$

$$\therefore Q_{en} = \frac{r_0^5}{r_0^5} Q$$

By putting

$$E (4\pi r^2) = \frac{1}{\epsilon_0} \left(\frac{r_0^5}{r_0^5} \right) Q$$

$$E = \frac{1}{\epsilon_0} \left(\frac{r_0^5}{r_0^5} \right) Q \times \frac{1}{4\pi r^2}$$

$$E = \frac{1}{4\pi\epsilon_0} \frac{\lambda^3 Q}{R^5}$$

This is the total electric charge intensity of the Non-Conductive Solid sphere in which charges are non-uniformly distributed.

Some other Applications

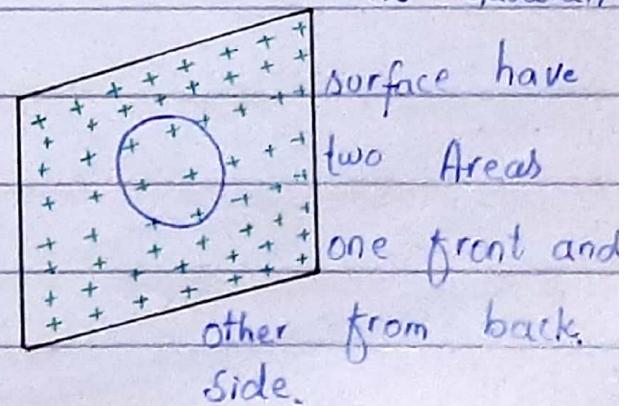
If we have a surface area on which charge distribution is uniform. In order to calculate the electric field intensity we have the formula:

$$\int EdA = \frac{1}{\epsilon_0} Q_{en}$$

$$2EA = \frac{1}{\epsilon_0} (\sigma A)$$

$$E = \frac{\sigma}{2\epsilon_0}$$

1

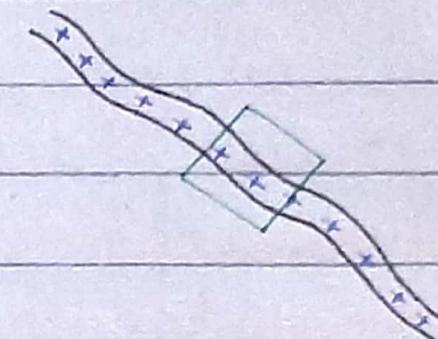


2: If we have a non-uniformly shaped object but charges are distributed uniformly. By imagine a gaussian surface we can calculate the Electric Field Intensity as:

$$\int E \cdot dA = \frac{1}{\epsilon_0} Q$$

$$EA = \frac{1}{\epsilon_0} \sigma A$$

$$E = \frac{\sigma}{\epsilon_0}$$



CHAPTER 24

Capacitance, Dielectrics, Electrical Energy Storage

Capacitor:

A Capacitor is a device that can store electric charges.

Normally, A capacitor consists of two conducting objects (usually plates or sheets) placed near each other but not touching. Capacitors are widely used in electronic circuits.

Capacitance:

The ability of a capacitor to store charges is called Capacitance of a capacitor.

- It is represented by C .
- It can be given as:

$$Q = CV$$

$$C = \frac{Q}{V}$$

- Its SI unit is Coulomb per Volt called Farad.

Potential Gradient

The quantity $\frac{\Delta V}{d}$ is the rate of potential change of maximum value of potential with distance and termed as Potential Gradient.

$$\Delta V = -Ed$$

- Negative sign shows that the direction of E is

along the decreasing potential. (Move from a point of low potential to a point of high potential)

Determination Of Capacitance

usually we talk

about three types of capacitors which are

- Parallel plate capacitors
- Spherical plate capacitors
- Cylindrical capacitors
- **Parallel Plate Capacitors**

Consider two parallel plates one of which is positively charged while other is negatively charged. In order to calculate the capacitance of these parallel plates.

a b

a

a non-conductive material is placed between the these two capacitor plates.

According to Gauss's Law we have.

$$\oint \mathbf{E} \cdot d\mathbf{A} = \frac{1}{\epsilon_0} Q_{\text{en}}$$

$$V_b - V_a = - \int_a^b \mathbf{E} \cdot d\mathbf{l}$$

$$\Delta V = - \int_a^b \mathbf{E} \cdot d\mathbf{l} \cos\theta$$

$$\theta = 180^\circ \text{ Meas } \cos\theta = \cos(180^\circ) = -1$$

∴ W Potential Gradient (v/m)

NOTE:

/* +ve charges are present in large amount but -ve charges

$$\Delta V = - \int_a^b \mathbf{E} \cdot d\mathbf{l} (-1)$$

are in small amount. That's the reason we apply Limit from Min R_a

to Max R_a (R_a-R_b)*/

$$\Delta V = \int_a^b \mathbf{E} \cdot d\mathbf{l} \quad (l \text{ is distance between two plates})$$

we have

$$\Delta V = Ed$$

As we know

$$EA = \frac{1}{\epsilon_0} Q$$

$$E = \frac{Q}{\epsilon_0 A}$$

Putting in above Equation

$$\Delta V = \left(\frac{Q}{\epsilon_0 A} \right) d$$

$$Q = \epsilon_0 A$$

$$\Delta V = \frac{Q}{d}$$

$$\therefore \frac{Q}{V} = C$$

$$C = \frac{\epsilon_0 A}{d}$$

This is the capacitance of parallel plate capacitors.

Cylindrical Capacitor

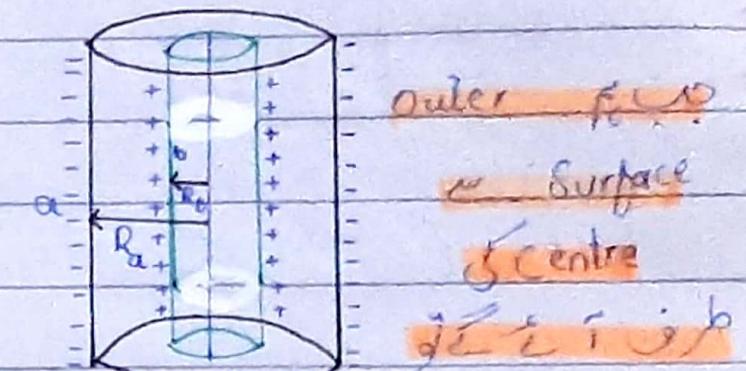
Consider a cylindrical capacitor, in this way the internal capacitor is +vely charged and external cylindrical capacitor is -vely charged. The Radius of +ve charged capacitor is R_b and of -ve charged capacitor is R_a . In order to calculate the capacitance of this cylindrical capacitor is:

Use Along Curved Surface of Electric Field & Cylinder to Capacitor $V_b - V_a$ (Gauss's Law)

$$\oint \mathbf{E} \cdot d\mathbf{A} = \frac{1}{\epsilon_0} Q_{\text{en}} \dots (i)$$

And Potential Gradient

$$V_b - V_a = - \int_{R_a}^{R_b} E \cdot d\mathbf{r} \dots (ii)$$



Now using (i)

$$\oint \mathbf{E} \cdot d\mathbf{A} = \frac{1}{\epsilon_0} Q_{\text{en}}$$

$$EA = \frac{1}{\epsilon_0} Q_{\text{en}}$$

$$V_b - V_a = \frac{Q_{\text{en}}}{\epsilon_0} E d$$

$$EA = \frac{1}{\epsilon_0} Q_{\text{en}}$$

$$\therefore \text{Surface Area} = 2\pi r d$$

$$E(2\pi r d) = \frac{1}{\epsilon_0} Q_{\text{en}}$$

$$E = \frac{1}{2\pi \epsilon_0 r d} Q_{\text{en}} \dots (iii)$$

Putting equation (iii) in (ii)

$$\Delta V = - \int_{R_a}^{R_b} \frac{1}{2\pi \epsilon_0 r d} Q_{\text{en}} \cdot dr$$

$$= - \frac{Q_{\text{en}}}{2\pi \epsilon_0 d} \int_{R_a}^{R_b} \frac{dr}{r}$$

$$\therefore \int \frac{dx}{x} = \ln x$$

$$= -\frac{Q}{2\pi l \epsilon_0} \left| \ln r \right|_{R_a}^{R_b}$$

$$= -\frac{Q}{2\pi l \epsilon_0} \{ \ln(R_b) - \ln(R_a) \}$$

$$= -\frac{Q_{en}}{2\pi l \epsilon_0} \ln(R_b) + \frac{Q_{en}}{2\pi l \epsilon_0} \ln(R_a)$$

$$= \frac{Q_{en}}{2\pi l \epsilon_0} \{ \ln(R_a) - \ln(R_b) \}$$

$$\therefore \ln A - \ln B = \ln A/B$$

$$\delta V = \frac{Q}{2\pi l \epsilon_0} \frac{\ln(R_a)}{R_b}$$

$$\frac{Q}{\delta V} = \frac{2\pi l \epsilon_0}{\ln(R_a/R_b)}$$

→ This is the capacitance of cylindrical capacitor.

$$C = \frac{2\pi l \epsilon_0}{\ln(R_a/R_b)}$$

Spherical Capacitor

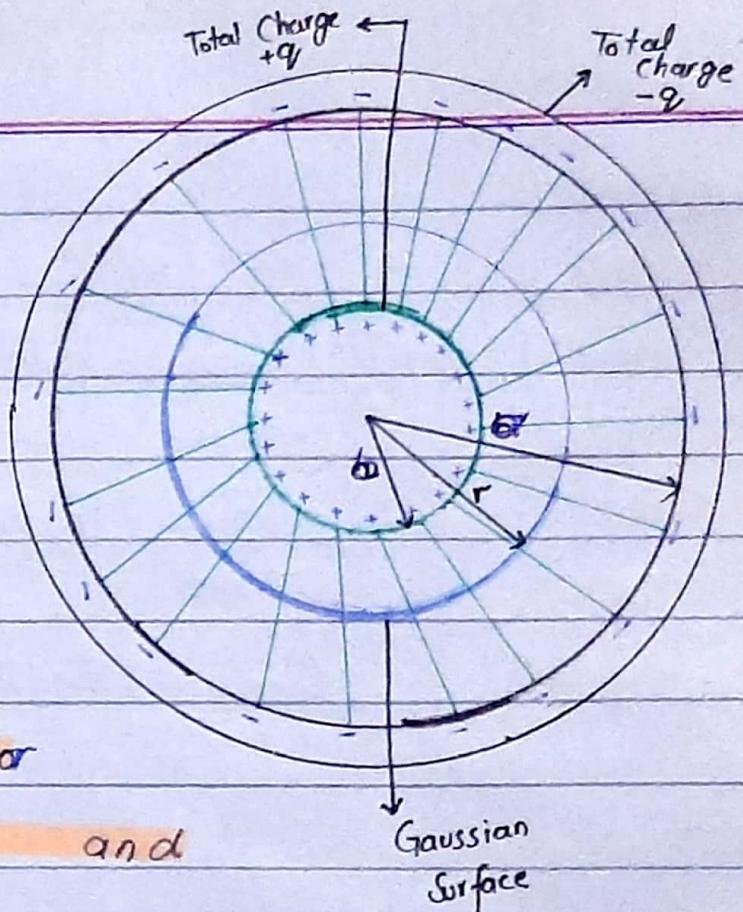
Consider a spherical plate capacitor in which similarly like above capacitors negative and positively charged capacitors are given. To calculate the capacitance of the capacitor we will use Gauss's Law and potential Gradient.

We have

$$\oint \mathbf{E} \cdot d\mathbf{A} = \frac{1}{\epsilon_0} Q_{\text{en}}$$

$$EA = \frac{1}{\epsilon_0} Q$$

$$\therefore A = 4\pi r^2$$



$$E(4\pi r^2) = \frac{1}{\epsilon_0} Q$$

$$E = \frac{1}{4\pi\epsilon_0 r^2} Q$$

Now Potential Gradient

$$V_b - V_a = - \int_{R_a}^{R_b} E \cdot dr$$

$$= - \frac{1}{4\pi\epsilon_0} \frac{Q}{r} \Big|_{R_a}^{R_b}$$

$$= - \frac{1}{4\pi\epsilon_0} \frac{Q}{r} \Big|_{R_a}^{R_b}$$

↪ Multiply & (-) & (-) \leftarrow
Apply $\mathop{\lim}_{r \rightarrow 0}$

Putting Value of E

$$V = - \int_{R_a}^{R_b} \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} dr$$

$$= \frac{1}{4\pi\epsilon_0} Q \left\{ \frac{1}{R_b} - \frac{1}{R_a} \right\}$$

$$V = \frac{1}{4\pi\epsilon_0} Q \left\{ \frac{R_a - R_b}{R_a R_b} \right\}$$

$$= - \frac{1}{4\pi\epsilon_0} \frac{Q}{R_a} \int_{R_a}^{R_b} \frac{dr}{r^2}$$

$$\frac{Q}{V} = \frac{4\pi\epsilon_0 R_a R_b}{R_a - R_b}$$

$$= - \frac{1}{4\pi\epsilon_0} \frac{Q}{R_a} \int_{R_a}^{R_b} r^{-2} dr$$

$$C = \frac{4\pi\epsilon_0 R_a R_b}{R_a - R_b}$$

$$= - \frac{1}{4\pi\epsilon_0} \frac{Q}{R_a} \Big|_{-2+1}^{R_b}$$

→ This is the Capacitance
of Spherical plate Capacitors.

CHAPTER 27

MAGNETISM

Magnetic Field:

The space that a region of magnetism around a magnet where the effect described its as magnetism, such as between the reflection and of the compass needle, can be detected is called Magnetic Field.

Magnetic Line of forces:

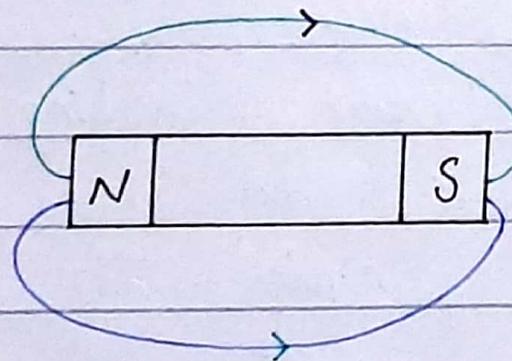
The path of an isolated

North magnetic pole in a magnetic field represents a magnetic line of force.

A single pole of a Magnet cannot be obtained.

↳ Impossible to separate North pole from South pole.

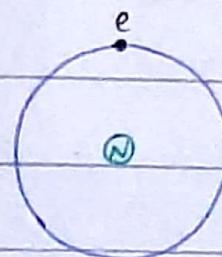
→ The direction of Magnetic Field lines is from North pole to South pole. As:



How does Magnetic Field Produce?

Magnetic Field is produced due to change in Flux.

If we have an electron in a orbit, with a nucleour N. as:



Here we can see that the nucleous is attracting the electron towards it self and electron is also repelling by outer electrons, so it will not stay at a constant place, it start revolving, Due to which Flux is produced and due to revolution electron is changing it's position, means change in flux is occurring, As a result, a Magnetic Field Produced.

لوب

J_s is clockwise & $\text{I} \leftarrow$ is current, Even J_s Material \leftarrow
 $\text{E} \leftarrow$ Anti clockwise & $\text{I} \rightarrow$

*P*HYSICS FOR *Scientists*

q *Engineers*

(DOUGLAS C. Giancoli)

AFTER Mid TERM

CHAPTER # 15

WAVES

What is wave?

The disturbance which travel from one point to another point with or without any medium is called a wave.

Types Of Wave

As we know that this world is full of waves. But two main types of wave are given below:

- ◆ Mechanical Waves
- ◆ Electromagnetic Waves

ELECTROMAGNETIC WAVES

Electromagnetic waves are those waves which required no any medium for their propagation.

Example:

The examples of electromagnetic waves are

- ◆ Visible Light
- ◆ Radio waves
- ◆ Television Signals

• why do they don't need any Medium for propagation?

Electromagnetic waves don't need any medium for their propagation because in these waves Electric Field and Magnetic Field both are perpendicular to each other, as a result, the resultant Field is zero.

NOTE : -

The Speed of Mechanical waves is very less as compared to the Electromagnetic waves

M ECHANICAL W AVES

Mechanical Waves are those waves which required a Medium for their propagation.

Example:

The examples of Mechanical waves are given below:

- Throw a pebble in water, water elements disturb.
- Sound waves
- Jump Rope waves.

NOTE:

A Mechanical wave is a wave that is not capable of transmitting its energy through a vacuum.

Types Of Mechanical Waves

There are two types of Mechanical waves.

- ◆ Longitudinal Waves
- ◆ Transverse waves
- ◆ Circular Transverse Waves

Longitudinal waves:

The waves in which Particles of Medium moves parallel to the wave motion or to extent of wave.

Example:

An example of longitudinal wave is compression moving along a slinky.

Transverse Waves:

The wave in which particles of medium move perpendicular to the motion of wave are called Transverse Waves.

Examples:

- ◆ Radio waves
- ◆ Sound waves
- ◆ Water waves

How does the above waves are Transverse waves?

The above waves are called Transverse waves. Let's check through an example of water waves.

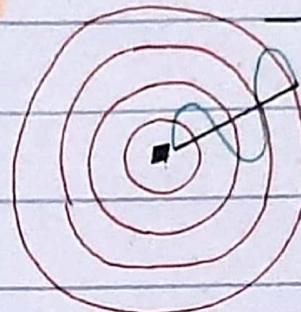
When we threw an pebble into water and see that waves produced, and how

these waves are transverse waves. - Motion of wave

Because, when pebble produce waves \curvearrowright Motion of Particles

we can see the angle between particles of Medium and wave motion is 90° Means both are

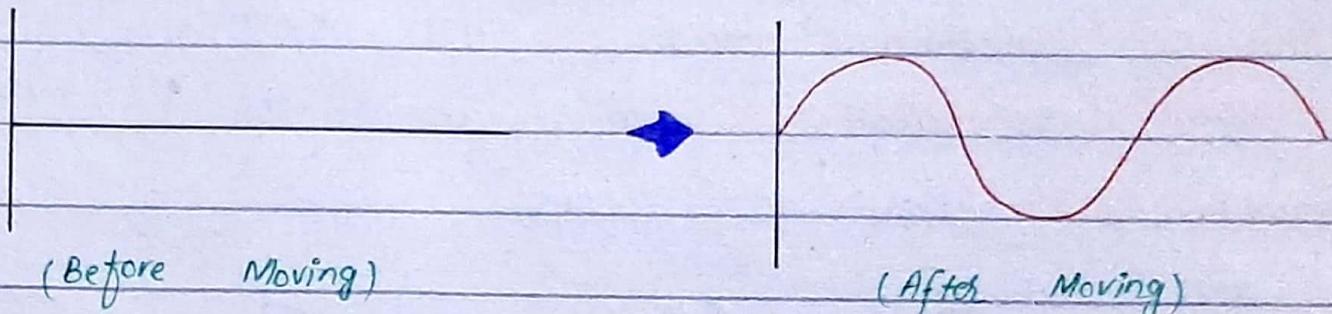
perpendicular. As, shown in figure:



THE WAVE EQUATION

Many types of waves satisfy an important general equation thus it is the equivalent to Newton's second law of motion for particles. This equation of motion for a wave is called the wave equation.

Let's take a rope and apply force on it in order to create a wave in the string. As

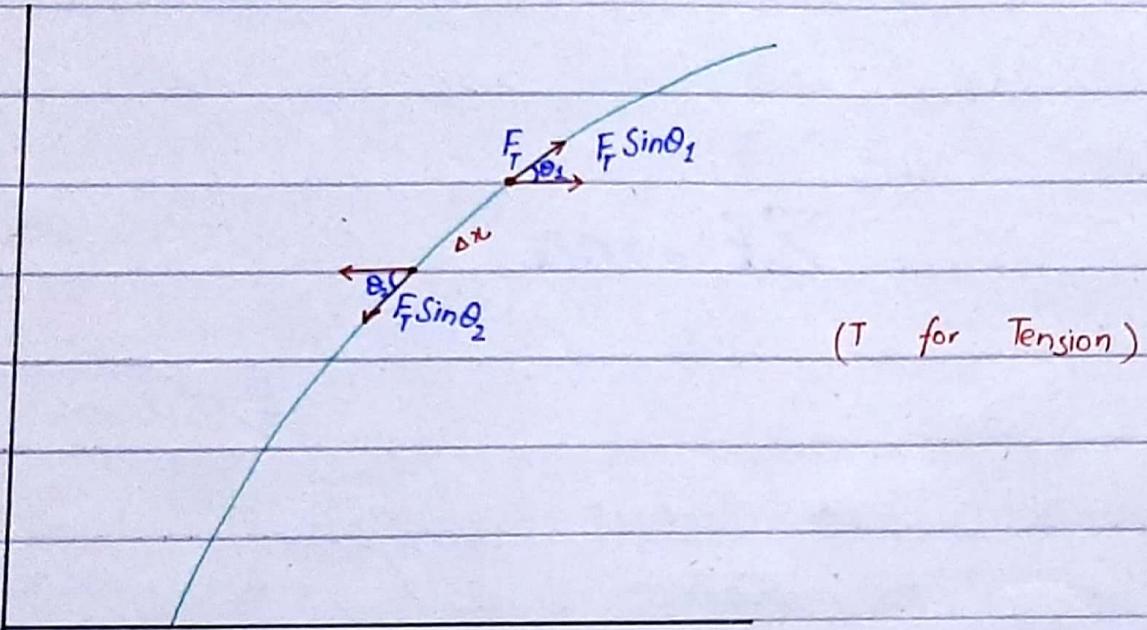


Here, we see that when we have applied force on the string, an acceleration has been produced in the direction of motion of string. And we know that rope has some mass also. So According To Newton's 2nd Law of Motion.

$$F = ma \text{ --- (i)}$$

In order to calculate the wave equation, let's take a small part of part of wave and resolve it into its components.

Figure:



(T for Tension)

We assume the amplitude of wave is small as compared to the wavelength, so that each point on the string can be assumed to move only vertically, and the Tension in the string F_T does not vary during vibration. Hence By Applying Newton's 2nd Law of Motion, we get the equation.

$$F_y = ma_y \dots (ii)$$

As, we know at a time, not only one force is working on y-axis. So, we can write the above equation as:

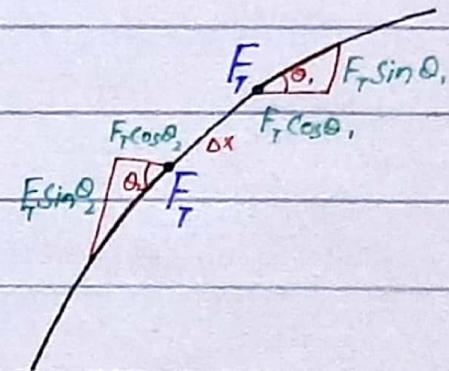
$$\sum F_y = ma_y \dots (iii)$$

From the figure,

we can write the above equation in this manner

$$F_T \sin \theta_2 - F_T \sin \theta_1 = ma_y \dots (iv)$$

Since we know that, we are considering a small patch of string and this patch of string will also have some mass. So we can find the mass of this special patch from Linear mass density which can be given as:



$$\mu = \frac{m}{l}$$

Since from figure, we have

$$l = \Delta x \quad (\text{The difference b/w Two forces})$$

so

$$\mu = \frac{m}{\Delta x}$$

and

$$m = \mu \Delta x \dots \dots (v)$$

By putting the equation (v) in eq. (iv)

$$F_1 \sin \theta_2 - F_2 \sin \theta_1 = \mu \Delta x a_y \dots \dots (vi)$$

Since we know that the acceleration is the double derivative of displacement. As. By taking first derivative of displacement we get velocity and by taking second derivative the velocity turns into acceleration.

And we have the acceleration

$$a_y = \frac{\partial v}{\partial t^2} \Rightarrow (v = \frac{s}{t})$$

$$\Rightarrow (a = \frac{v}{t})$$

Since the motion is only vertically and we use the partial derivative notation because the displacement D is a function of both x and t .

Now

Put the value of acceleration in equation

(vi)

$$F_T \sin\theta_2 - F_T \sin\theta_1 = \mu \alpha x \frac{\partial^2 D}{\partial t^2} \dots \dots \text{(vii)}$$

As the angles θ_1 and θ_2 are assumed small. So $\sin\theta \approx \tan\theta$ and $\tan\theta$ is equal to the slope of string at each point.

$$\sin\theta \approx \tan\theta \approx \frac{\partial D}{\partial x} = \text{Slope}$$

By using equation (vii)

$$F_T \sin\theta_2 - F_T \sin\theta_1 = \mu \alpha x \frac{\partial^2 D}{\partial t^2}$$

$$F_T (\sin\theta_2 - \sin\theta_1) = \mu \alpha x \frac{\partial^2 D}{\partial t^2}$$

From above relation $\sin\theta \approx \tan\theta$ so

$$F_T (\tan\theta_2 - \tan\theta_1) = \mu \Delta x \frac{\partial^2 D}{\partial t^2}$$

Since

$\tan\theta = \text{Slope} = S$ (so above equation becomes)

$$F_T (S_2 - S_1) = \mu \Delta x \frac{\partial^2 D}{\partial t^2}$$

$$F_T (\Delta S) = \mu \Delta x \frac{\partial^2 D}{\partial t^2}$$

$$F_T \frac{\partial S}{\partial x} = \mu \frac{\partial^2 D}{\partial t^2} \dots \text{viii}$$

Since

$$S = \frac{\partial D}{\partial x} \rightarrow S = \frac{\partial D}{\partial x} \rightarrow \text{if } \Delta x \approx 0$$

$$S = \frac{\partial D}{\partial x} \text{ (also a function of time)}$$

Put this in equation (viii)

$$F_T \frac{\partial}{\partial x} \left(\frac{\partial D}{\partial x} \right) = \mu \frac{\partial^2 D}{\partial t^2}$$

$$\frac{F_I}{\mu} \frac{\partial^2 D}{\partial x^2} = \frac{\partial^2 D}{\partial t^2} \dots \dots \text{(ix)}$$

As we know velocity of transverse wave

$$v = \sqrt{\frac{F_I}{\mu}}$$

Taking square

$$v^2 = \frac{F_I}{\mu} \dots \dots \text{(x)}$$

Put in eq. (ix)

$$v^2 = \frac{\partial^2 D}{\partial x^2} = \frac{\partial^2 D}{\partial t^2}$$

$$\frac{\partial^2 D}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 D}{\partial t^2}$$

This is called one-dimensional wave Equation. It can describe not only small amplitude wave on a stretched string, but also small amplitude longitudinal waves (such as sound waves) in gases, liquid, and elastic solid.

Noise & Music

Noise:

The sound which has unpleasant effect on air and has no periodicity is called Noise.

→ Noise is produced by horns of vehicles, crackers, machines etc.

Music:

The sound which has pleasant effect on air and has periodicity is called Music.

→ Music is produced by musical instruments like piano, Guitar, Flute etc.

Note:

It is not important that every musical instrument cannot produce noise.

Periodicity

Periodicity means that if elements are arranged in an order, the properties of elements would repeat some period.

Periodicity Makes your sound Noise or Music.

If the Periodicity is controlled, it will produce Music and if the Periodicity is not controlled then Noise produces.

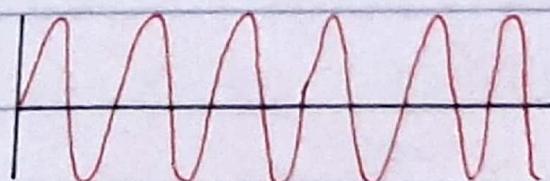
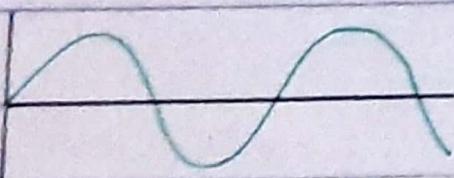
PITCH:

{Position of a single sound in the complete range of Sound}

The pitch of sounds is determined by the frequency of vibration of the sound waves that produce them.

We can differentiate between different sounds on the basis of pitch.

- Sounds having high-frequency produced high-pitched noise.
- Sound waves having low-frequency produce low-pitched sounds.



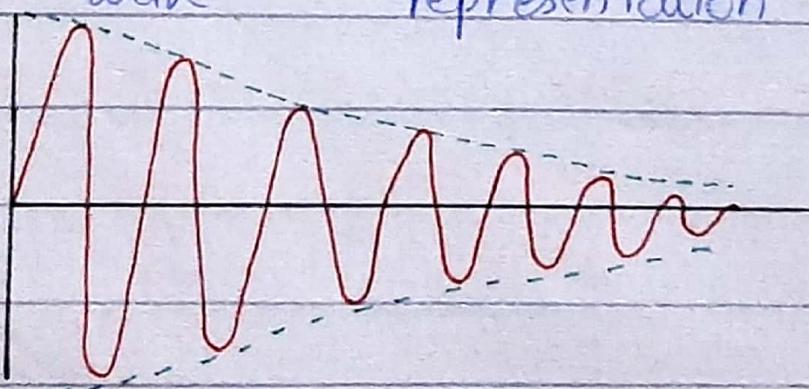
Frequency:

A frequency is how frequent sound waves pass a certain point in a second.

Harmonic Oscillator:

It is called Harmonic. Because the solution of Newton's second Law (a second order differential equation that determines the motion of the object) are sines and cosines of time with the particular frequency.

Its wave representation is given below:



Intensity:

It is the Power delivered per unit area.

Intensity of wave is proportional to the square of amplitude of the wave.

Usually, we measure the intensity in Bell or decibell.

For Sound Waves:

$$I_0 = 10^6$$

$$I = 10^{12}$$

$$\log \frac{I}{I_0} = \log \frac{10^{12}}{10^6} = \log 10^{12-6} = \log 10^6 = "6 \text{dB}"$$

$$\text{Intensity of sound waves} = \frac{E}{A \cdot t}$$

Standing/Stationary Waves

The waves travelling in same medium but coming from opposite direction super impose.

Example:

Waves produced by stringed musical instruments.

Node:

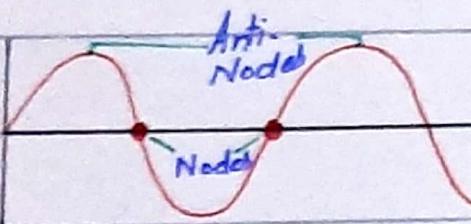
A node is a point along a standing wave where the wave has minimum amplitude.

Anti-Node:

An anti-node is a point along a standing wave where the wave has maximum amplitude.

Stationary Waves In Air Column

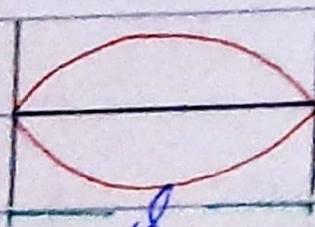
Graphically Representation:



Consider we have a wave, and let's check the relation of frequency with wavelength and increasing the number of cycles.

First Half Cycle:

Let Consider the half cycle of wave



As we know

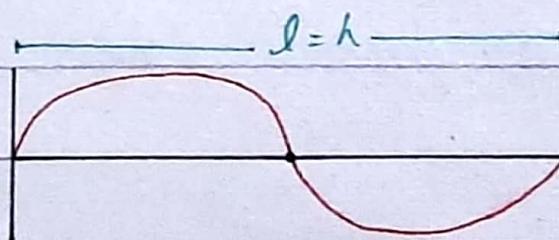
$$V = f \lambda$$

$$f = \frac{V}{\lambda}$$

$$f = \frac{V}{2\lambda}$$

For 1-Complete Cycle:

Now let's check for one complete cycle.



As we know the wavelength is the distance from the crest of one wave to crest of next wave.

Now we have

$$\Rightarrow \lambda = \lambda$$

$$f_2 = \frac{v}{\lambda_2}$$

$$f_2 = \frac{v}{\lambda} \quad (\text{crossing 1 - ing by 2})$$

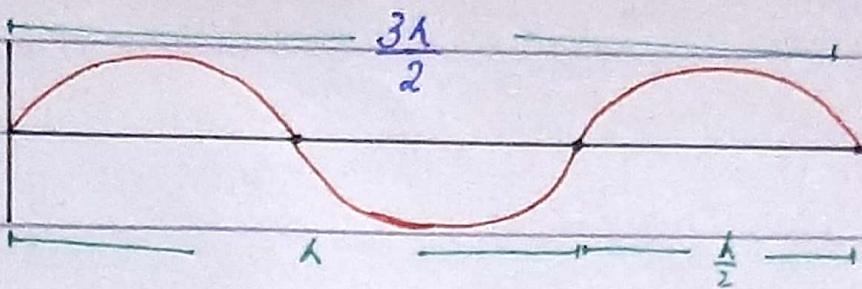
$$f_2 = 2 \left(\frac{v}{2\lambda} \right)$$

$$\therefore \frac{v}{2\lambda} = f_1$$

$$f_2 = 2 f_1$$

Now For 1.5 Cycles:

Now If we consider one complete and one half cycle. And we have length of wave $\frac{3\lambda}{2}$. As:



$$\Rightarrow l = \frac{3\lambda}{2}$$

Now

$$\Rightarrow \lambda_3 = \frac{2l}{3}$$

$$f_3 = \frac{v}{\lambda_3}$$

$$f_3 = \frac{v}{2l/3}$$

$$f_3 = 3 \left(\frac{v}{2l} \right) \quad \therefore \frac{v}{2l} = f_1$$

$$f_3 = 3 f_1$$

For n -cycles:

From above relations we can say that, if we have n -number of cycles then, we can find a relation such as:

$$f_n = n f_1$$

From Relation we can say that

As the wavelength decreases frequency increases.

TRANSVERSE WAVE OF FREQUENCY

The frequency of a transverse wave is found by determining the number of high points (crests) that go by a point in one

Second.

A wave with high frequency has many crests that pass by the point.

FORMULA

$$V = \sqrt{\frac{F_T}{\mu}}$$

As

$$V = f\lambda$$

$$f = \frac{V}{\lambda}$$

$$f \propto V \quad \text{and} \quad f \propto \frac{1}{\lambda}$$

Distance of strings from string & Musical Instruments

Height of string from position of string & length of

Speed of light is the speed of strings, it is the frequency of sound.

$$v = \sqrt{\frac{F}{\mu}}$$

String at different positions, same string at different positions, different strings at different positions, the periodicity of sounds is called Music. Noise is non-periodic.

IMP. Points:

From the above topic we have learnt that

- Loudness of a sound depends on intensity.
 - Higher Intensity the sound will be louder.
 - Lower Intensity the sound will be lower.

◆ The Frequency Produce Sounds.

◆ If we combine different frequencies sounds it has different effect on human's ear.

• If there is a Periodicity in sound and having pleasant effect on human's ear then it will called Music.

• If there is no Periodicity in sound and having unpleasant effect on human's ear then it will called Noise.

Since we know that frequency can be given as:

$$f = \frac{v}{\lambda}$$

In Case of Transverse Waves the velocity can be given as:

$$v = \sqrt{\frac{F_T}{\mu}}$$

- v depends on μ .
 - For each sound the value of μ is different.
 - Due to different velocities the frequency of sound changes and new sounds produce. Because $f \propto v$
- ⇒ Higher the frequency, higher will be velocity.
- ⇒ Lower the velocity, lower will be frequency.
- Music has everlasting effect on human's ear.