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MEng Title

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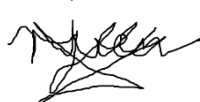
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Abstract

English

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Nomenclature

Ocean Vessel Dynamics

X, Y, Z	Coordinates of force vector decomposed in the body-fixed frame(surge, sway and heave forces)
K, M, N	Coordinates of moment vector decomposed in the body-fixed frame(roll, pitch and yaw moment)
u, v, w	Coordinates of linear velocity vector decomposed in the body-fixed frame(surge, sway and heave velocities)
p, q, r	Coordinates of angular velocity vector decomposed in the body-fixed frame(roll, pitch and yaw angular velocities)
x, y, z	Coordinates of position vector decomposed in the body-fixed frame(surge, sway and heave positions)
ϕ, θ, ψ	Coordinates of Euler angle vector decomposed in the body-fixed frame(roll, pitch and yaw Euler angles)

Acronyms and abbreviations

SNAME	Society of Naval Architects and Engineers
CG	Center of gravity of the vessel

Chapter 1

Introduction

1.1. Background

1.2. Problem Statement

1.3. Summary of Work

1.4. Scope

1.5. Format of Report

Chapter 2

Literature Review

Chapter 3

Modeling of Ocean Vessels

This chapter models a standard ocean vessel in six degrees of freedom. It also introduces the definitions associated with movement in each direction of freedom. The chapter also take into account the forces and moments generated by hydrodynamics and restoration of an ocean vessel. The chapter continues to model the environmental disturbances experience by a semi-submerged ocean vessel. The environmental disturbances are wind, waves and ocean currents.

3.1. Standard Ocean Vessel Notation

An ocean vessels are modelled in six degrees of freedom, requiring six independent coordinates to determine its position and orientation. The first three coordinates corresponding to position (x, y, z) and their first time derivatives, translation motion along the x -, y -, and z -axes. The last three coordinates (ϕ, θ, ψ) and their first time derivatives describing orientation and rotational motion [1]. Figure 3.1 illustrates the motion variables of an ocean vessel with the six independent coordinates.

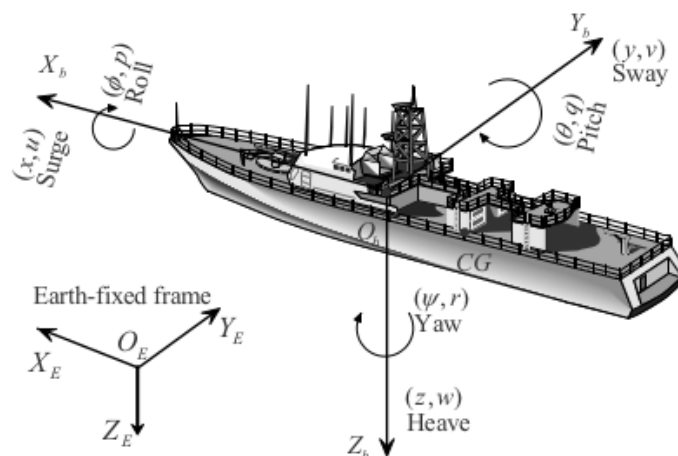


Figure 3.1: Motion variables for an ocean vessel

The SNAME(Society of Naval Architects and Marine Engineers) established the notation for the six different motion components as *surge*, *sway*, *heave*, *roll*, *pitch* and *yaw*. Table A.1 summarizes the SNAME notation for ocean vessels.

Two reference models are used to determine the equations of motion, namely the inertial to earth frame $O_e X_e Y_e Z_e$ that may be displaced to overlap with the vessel's fixed coordinates in some initial condition and the body-fixed frame $O_b X_b Y_b Z_b$, illustrated in Figure 3.1. The most common used position for the body-fixed frame results in symmetry about the $O_b X_b Z_b$ -plane and approximate symmetry about the $O_b Y_b Z_b$. The body axes coincides with the axes of inertia and are usually defines as follows: $O_b X_b$ is the longitudinal axis, $O_b Y_b$ is the transverse axis and $O_b Z_b$ is the normal axis. Below are the vectors used to describe the general motion of an ocean vessel:

$$\mathbf{n} = [\mathbf{n}_1 \mathbf{n}_2]^T \quad (3.1)$$

$$\mathbf{v} = [\mathbf{v}_1 \mathbf{v}_2]^T \quad (3.2)$$

$$\tau = [\tau_1 \tau_2]^T \quad (3.3)$$

$$\mathbf{n}_1 = [x \ y \ z]^T \quad (3.4) \quad \mathbf{n}_2 = [\phi \ \theta \ \psi]^T \quad (3.5)$$

$$\mathbf{v}_1 = [u \ v \ w]^T \quad (3.6) \quad \mathbf{v}_2 = [p \ q \ r]^T \quad (3.7)$$

$$\tau_1 = [X \ Y \ Z]^T \quad (3.8) \quad \tau_2 = [K \ M \ N]^T \quad (3.9)$$

where \mathbf{n} denotes the position and orientation vector with coordinates in the earth fixed frame, \mathbf{v} denotes the linear and angular velocity vector with coordinates in the body-fixed frame and τ denotes the forces and moments acting on the vessel in the body-fixed frame. The vessel dynamics are divided into two parts known as *kinematics* and *kinetics*.

3.2. Kinematics

Kinematics looks at the motion of the vessel without directly considering the forces affecting the motion. The first time derivative of the position vectors \mathbf{n}_1 and \mathbf{n}_2 is related to the linear velocity vector \mathbf{v}_1 and \mathbf{v}_2 via the following transformations,

$$\dot{\mathbf{n}}_1 = \mathbf{J}_1(\mathbf{n}_2) \mathbf{v}_1 \quad (3.10)$$

$$\dot{\mathbf{n}}_2 = \mathbf{J}_2(\mathbf{n}_2) \mathbf{v}_2 \quad (3.11)$$

where $\mathbf{J}_1(\mathbf{n}_2)$ and $\mathbf{J}_2(\mathbf{n}_2)$ are transformation matrices, which is related through the functions of the Euler angles: roll(ϕ), pitch(θ) and yaw(ψ). The \mathbf{J}_1 transformation matrix is given by

$$\mathbf{J}_1(\mathbf{n}_2) = \begin{bmatrix} \cos(\psi) \cos(\theta) & -\sin(\psi) \cos(\theta) + \sin(\phi) \sin(\theta) \cos(\psi) & \sin(\psi) \sin(\phi) + \sin(\theta) \cos(\psi) \cos(\phi) \\ \sin(\psi) \cos(\theta) & \cos(\psi) \cos(\phi) + \sin(\phi) \sin(\theta) \sin(\psi) & -\cos(\psi) \sin(\phi) + \sin(\theta) \sin(\psi) \cos(\phi) \\ -\sin(\theta) & \sin(\phi) \cos(\theta) & \cos(\phi) \cos(\theta) \end{bmatrix} \quad (3.12)$$

and the transformation matrix \mathbf{J}_2 is given by,

$$\mathbf{J}_2(\mathbf{n}_2) = \begin{bmatrix} 1 & -\sin(\phi) \tan(\theta) & \cos(\phi) \tan(\theta) \\ 0 & \cos(\phi) & -\sin(\phi) \\ 0 & \sin(\phi)/\cos(\theta) & \cos(\phi)/\cos(\theta) \end{bmatrix} \quad (3.13)$$

When $\theta = \pi/2$, the transformation matrix $\mathbf{J}_2(\mathbf{n}_2)$ becomes singular, however this is unlikely to happen when practically testing an ocean vessel, because of the metacentric restoring forces. Combining Equation 3.12 and Equation 3.13 results in the kinematics of an ocean vessel.

$$\begin{bmatrix} \dot{\mathbf{n}}_1 \\ \dot{\mathbf{n}}_2 \end{bmatrix} = \begin{bmatrix} \mathbf{J}_1(\mathbf{n}_2) & 0_{3 \times 3} \\ 0_{3 \times 3} & \mathbf{J}_2(\mathbf{n}_2) \end{bmatrix} \begin{bmatrix} \mathbf{v}_1 \\ \mathbf{v}_2 \end{bmatrix} = \dot{\mathbf{n}} = \mathbf{J}(\mathbf{n})\mathbf{v} \quad (3.14)$$

3.3. Kinetics

The Newton-Euler formulation [2] defines the balancing forces and moments for a rigid body with a mass of m as follows,

$$\mathbf{f}_{Ob} = \mathbf{m}[\dot{\mathbf{v}}_{Ob}^E + \dot{\mathbf{w}}_{Ob}^E \times \mathbf{r}_{Ob} + \mathbf{w}_{Ob}^E \times \mathbf{v}_{Ob} + \mathbf{w}_{Ob}^E \times (\mathbf{w}_{Ob}^E \times \mathbf{r}_{Ob})] \quad (3.15)$$

$$\mathbf{m}_{Ob} = \mathbf{I}_o \mathbf{w}_{Ob}^E + \dot{\mathbf{w}}_{Ob}^E \times \mathbf{I}_o \mathbf{w}_{Ob}^E + m \mathbf{r}_{Ob} \times (\dot{\mathbf{v}}_{Ob} + \mathbf{w}_{Ob}^E \times \mathbf{v}_{Ob}) \quad (3.16)$$

where \mathbf{f}_{Ob} is the balancing forces, \mathbf{m}_{Ob} the balancing moments and \mathbf{I}_o is the inertia matrix about O_b defined as,

$$\mathbf{I}_o = \begin{bmatrix} I_x & -I_{xy} & -I_{xz} \\ -I_{yx} & I_y & -I_{yz} \\ -I_{zx} & -I_{zy} & I_z \end{bmatrix} \quad (3.17)$$

I_x, I_y and I_z are the moments of inertia about the $O_b X_b$, $O_b Y_b$ and $O_b Z_b$ axes. $I_{xy} = I_{yx}$, $I_{xz} = I_{zx}$ and $I_{yz} = I_{zy}$ are the products of inertia. These quantities are defined as

$$I_x = \int_V (y^2 + z^2) \rho_m dV \quad (3.18) \quad I_{xy} = \int_V xy \rho_m dV \quad (3.19)$$

$$I_y = \int_V (x^2 + z^2) \rho_m dV \quad (3.20) \quad I_{xz} = \int_V xz \rho_m dV \quad (3.21)$$

$$I_z = \int_V (x^2 + y^2) \rho_m dV \quad (3.22) \quad I_{zy} = \int_V zy \rho_m dV \quad (3.23)$$

where ρ_m are the mass density and V the volume of the rigid body. By substituting the definitions defined in Table A.2 into Equations 3.15 and 3.16, results in the equation below,

$$\mathbf{M}_{RB} \dot{\mathbf{v}} + \mathbf{C}_{RB}(\mathbf{v})\mathbf{v} = \tau_{RB} \quad (3.24)$$

where $\mathbf{v} = [u \ v \ w \ p \ q \ r]^T$ is the generalized velocity vector decomposed in the body-fixed frame and $\boldsymbol{\tau}_{\mathbf{RB}} = [X \ Y \ Z \ K \ M \ N]^T$ is the generalized vector of external forces and moments. The rigid body system inertia matrix $\mathbf{M}_{\mathbf{RB}}$ and the rigid body Coriolis and centripetal matrix $\mathbf{C}_{\mathbf{RB}}$ is defined in Equation A.1 and A.2. The generalized external force and moment vector, $\boldsymbol{\tau}_{\mathbf{RB}}$, is a sum of the hydrodynamic force and moment vector $\boldsymbol{\tau}_{\mathbf{H}}$, external disturbance force and moment vector $\boldsymbol{\tau}_{\mathbf{E}}$ and propulsion force and moment vector $\boldsymbol{\tau}$.

3.4. Hydrodynamic Forces and Moments

Hydrodynamic forces and moments can be defined as the forces and moments on a ocean body when the body is forced to oscillate with the wave excitation and no wave are incident on the body. As shown in [3], the hydrodynamic forces and moments acting on a rigid body can be assumed to be linearly superimposed. The forces and moments can be subdivided into three components,

1. Added mass due to the inertia of the surrounding fluid
2. Radiation-induced potential damping due to the energy carried away by the generated surface waves
3. Restoring forces due to Archimedian forces

The hydrodynamic forces and moments vector $\boldsymbol{\tau}_{\mathbf{H}}$ is expressed in the equation below,

$$\boldsymbol{\tau}_{\mathbf{H}} = -\mathbf{M}_{\mathbf{A}}\dot{\mathbf{v}} - \mathbf{C}_{\mathbf{A}}(\mathbf{v})\mathbf{v} - \mathbf{D}(\mathbf{v})\mathbf{v} - \mathbf{g}(\mathbf{n}) \quad (3.25)$$

where $\mathbf{M}_{\mathbf{A}}$ is the added mass matrix, $\mathbf{C}_{\mathbf{A}}(\mathbf{v})$ is the hydrodynamic Coriolis and centripetal matrix, $\mathbf{D}(\mathbf{v})$ is the damping matrix and $\mathbf{g}(\mathbf{n})$ is the position and orientation depending vector of restoring forces and moments. The added mass $\mathbf{M}_{\mathbf{A}}$ is given below,

$$\mathbf{M}_{\mathbf{A}} = \begin{bmatrix} X_{\dot{u}} & X_{\dot{v}} & X_{\dot{w}} & X_{\dot{p}} & X_{\dot{q}} & X_{\dot{r}} \\ Y_{\dot{u}} & Y_{\dot{v}} & Y_{\dot{w}} & Y_{\dot{p}} & Y_{\dot{q}} & Y_{\dot{r}} \\ Z_{\dot{u}} & Z_{\dot{v}} & Z_{\dot{w}} & Z_{\dot{p}} & Z_{\dot{q}} & Z_{\dot{r}} \\ K_{\dot{u}} & K_{\dot{v}} & K_{\dot{w}} & K_{\dot{p}} & K_{\dot{q}} & K_{\dot{r}} \\ M_{\dot{u}} & M_{\dot{v}} & M_{\dot{w}} & M_{\dot{p}} & M_{\dot{q}} & M_{\dot{r}} \\ N_{\dot{u}} & N_{\dot{v}} & N_{\dot{w}} & N_{\dot{p}} & N_{\dot{q}} & N_{\dot{r}} \end{bmatrix} \quad (3.26)$$

The hydrodynamic Coriolis and centripetal matrix is given below,

$$\mathbf{C}_{\mathbf{A}}(\mathbf{v}) = \begin{bmatrix} 0 & 0 & 0 & 0 & -a_3 & a_2 \\ 0 & 0 & 0 & a_3 & 0 & -a_1 \\ 0 & 0 & 0 & -a_2 & a_1 & 0 \\ 0 & -a_3 & a_2 & 0 & -b_3 & b_2 \\ a_3 & 0 & -a_1 & b_3 & 0 & -b_1 \\ -a_2 & a_1 & 0 & -b_2 & b_1 & 0 \end{bmatrix} \quad (3.27)$$

where a_1 , a_2 , a_3 , b_1 , b_2 and b_3 are defined in Equations A.3, A.4, A.5, A.6, A.7 and A.8.

The general hydrodynamic damping experienced by ocean vessels is the potential damping, skin friction, wave drift damping and damping due to vortex shedding. The hydrodynamic damping can be expressed in a general form as below,

$$\mathbf{D}(\mathbf{v}) = \mathbf{D} + \mathbf{D}_n(\mathbf{v}) \quad (3.28)$$

where the linear damping matrix \mathbf{D} is given below,

$$\mathbf{D} = - \begin{bmatrix} X_u & X_v & X_w & X_p & X_q & X_r \\ Y_u & Y_v & Y_w & Y_p & Y_q & Y_r \\ Z_u & Z_v & Z_w & Z_p & Z_q & Z_r \\ K_u & K_v & K_w & K_p & K_q & K_r \\ M_u & M_v & M_w & M_p & M_q & M_r \\ N_u & N_v & N_w & N_p & N_q & N_r \end{bmatrix} \quad (3.29)$$

3.5. Restoring Forces and Moments

The restoring forces and moments are described by the symbol $\mathbf{g}(\mathbf{n})$. If ∇ is the volume of fluid displaced by the ocean vessel. The acceleration of gravity, g and the water density ρ . The submerged weight of the body and buoyancy forces are defined by

$$W = mg \quad (3.30)$$

$$B = \rho g \nabla \quad (3.31)$$

With the above definition for body and buoyancy forces, the restoring force and moment vector $\mathbf{g}(\mathbf{n})$ is due to gravity and buoyancy forces and is given by

$$\mathbf{g}(\mathbf{n}) = \begin{bmatrix} (W - B)\sin(\theta) \\ -(W - B)\cos(\theta)\sin(\phi) \\ -(W - B)\cos(\theta)\cos(\phi) \\ -(y_g W - y_b B)\cos(\theta)\cos(\phi) + (z_g W - z_b B)\cos(\theta)\sin(\phi) \\ (z_g W - z_b B)\sin(\theta) + (x_g W - x_b B)\cos(\theta)\cos(\phi) \\ -(x_g W - x_b B)\cos(\theta)\sin(\phi) - (y_g W - y_b B)\sin(\theta) \end{bmatrix} \quad (3.32)$$

where (x_b, y_b, z_b) denote coordinates of the center of buoyancy.

3.6. Environmental Disturbances

The forces and moments induced by the environmental disturbances is defined by the vector $\tau_{\mathbf{E}}$ and includes ocean currents, waves(wind generated) and wind.

$$\tau_{\mathbf{E}} = \tau_{\mathbf{E}}^{\text{cu}} + \tau_{\mathbf{E}}^{\text{wa}} + \tau_{\mathbf{E}}^{\text{wi}} \quad (3.33)$$

where $\tau_{\mathbf{E}}^{\text{cu}}$, $\tau_{\mathbf{E}}^{\text{wa}}$ and $\tau_{\mathbf{E}}^{\text{wi}}$ are vectors of forces and moments induced by ocean currents, waves and wind.

3.6.1. Current-induced Forces and Moments

The current induced forces and moments vector $\tau_{\mathbf{E}}^{\text{cu}}$ is given by

$$\tau_{\mathbf{E}}^{\text{cu}} = (\mathbf{M}_{\text{RB}} + \mathbf{M}_{\text{A}})\dot{\mathbf{v}}_{\text{c}} + \mathbf{C}(\mathbf{v}_{\text{r}})\mathbf{v}_{\text{r}} - \mathbf{C}(\mathbf{v})\mathbf{v} + \mathbf{D}(\mathbf{v}_{\text{r}})\mathbf{v}_{\text{r}} - \mathbf{D}(\mathbf{v})\mathbf{v} \quad (3.34)$$

where $\mathbf{v}_{\text{r}} = \mathbf{v} - \mathbf{v}_{\text{c}}$ and $\mathbf{v}_{\text{c}} = [u_{\text{c}}, v_{\text{c}}, w_{\text{c}}, 0, 0, 0]^T$ is a vector irrotational body-fixed current velocities. Take the earth-fixed velocity vector denoted by $[u_{\text{c}}^E, v_{\text{c}}^E, w_{\text{c}}^E]^T$, then the body-fixed components $[u_{\text{c}}, v_{\text{c}}, w_{\text{c}}]^T$ can be calculated by

$$\begin{bmatrix} u_{\text{c}} \\ v_{\text{c}} \\ w_{\text{c}} \end{bmatrix} = \mathbf{J}_{\mathbf{I}}^T(\mathbf{n}_2) \begin{bmatrix} u_{\text{c}}^E \\ v_{\text{c}}^E \\ w_{\text{c}}^E \end{bmatrix} \quad (3.35)$$

3.6.2. Wave-induced Forces and Moments

The vector $\tau_{\mathbf{E}}^{\text{wa}}$ of the wave-induced forces and moments is given by

$$\tau_{\mathbf{E}}^{\text{wa}} = \begin{bmatrix} \sum_{i=1}^N \rho g B L T \cos(\beta) s_i(t) \\ \sum_{i=1}^N \rho g B L T \sin(\beta) s_i(t) \\ 0 \\ 0 \\ 0 \\ \sum_{i=1}^N \frac{1}{24} \rho g B L (L^2 - B^2) \sin(2\beta) s_i^2(t) \end{bmatrix} \quad (3.36)$$

where β is the vessel's heading(encounter) angle, illustrated in Figure , ρ is the water density, L is the length of the vessel, B is the breadth of the vessel and T is the draft of the vessel. Ignoring the higher-order terms of the wave amplitude, the wave slope $s_i(t)$ for the wave component i is defined by

$$s_i(t) = A_i \frac{2\pi}{\lambda_i} \sin(\omega_{ei}t + \phi_i) \quad (3.37)$$

where A_i is the wave amplitude, λ_i is the wave length, ω_{ei} is the encounter frequency and ϕ_i is a random phase uniformly distributed and constant with time $[0, 2\pi)$ corresponding to the wave component i .

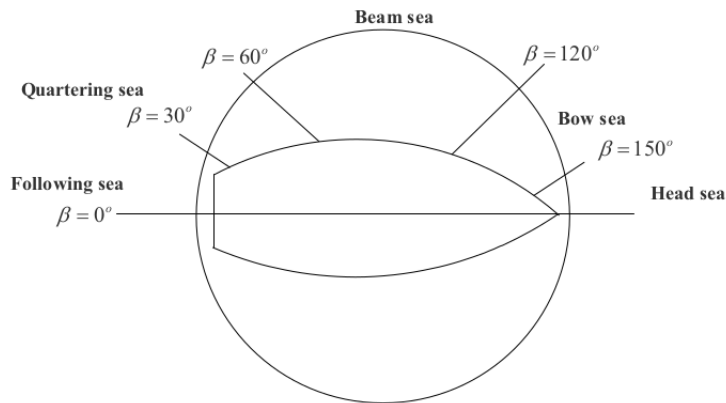


Figure 3.2: Ocean vessel's heading angle

3.6.3. Wind-induced Forces and Moments

When the ocean vessel is at rest the vector τ_E^{wi} of the wind induced forces and moments is given by

$$\tau_E^{wi} = \begin{bmatrix} C_X(\gamma_\omega)A_{F\omega} \\ C_Y(\gamma_\omega)A_{L\omega} \\ C_Z(\gamma_\omega)A_{F\omega} \\ C_K(\gamma_\omega)A_{L\omega}H_{L\omega} \\ C_M(\gamma_\omega)A_{F\omega}H_{F\omega} \\ C_N(\gamma_\omega)A_{L\omega}L_{oa} \end{bmatrix} \quad (3.38)$$

where V_ω is the wind speed, ρ_a is the air density, $A_{F\omega}$ is the frontal projected area, $A_{L\omega}$ is the lateral projected area, $H_{F=\omega}$ is the centroid of $A_{F\omega}$ above the water line, $H_{L\omega}$ is the centroid of $A_{L\omega}$ above the water line, L_{oa} is the over all length of the vessel, γ_ω is the angle of relative wind of the vessel bow, illustrated in Figure 3.3 and is given by

$$\gamma_\omega = \psi - \beta_\omega - \pi \quad (3.39)$$

where β_ω being the wind direction. All the wind coefficients(look-up tables) $C_X(\gamma_\omega)A_{F\omega}$, $C_Y(\gamma_\omega)A_{L\omega}$, $C_Z(\gamma_\omega)A_{F\omega}$, $C_K(\gamma_\omega)A_{L\omega}H_{L\omega}$, $C_M(\gamma_\omega)A_{F\omega}H_{F\omega}$ and $C_N(\gamma_\omega)A_{L\omega}L_{oa}$ are computed numerically or by experiments in a wind tunnel as shown in [4].

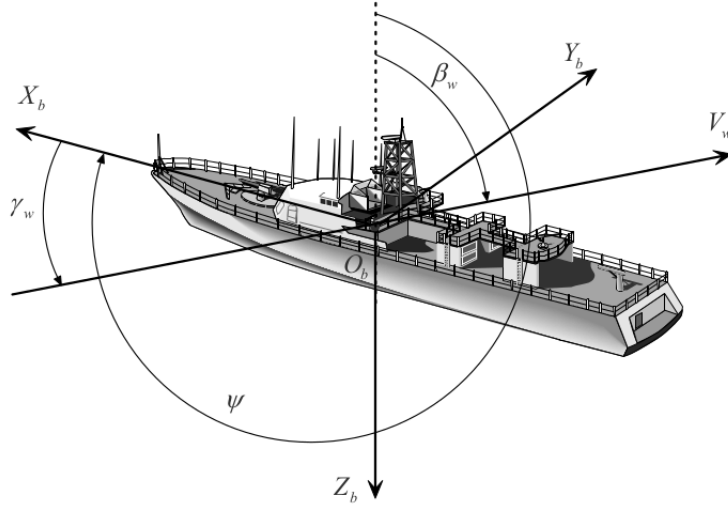


Figure 3.3: Wind angle on vessel

When the vessel is moving the vector τ_E^{wi} is given by

$$\tau_E^{wi} = \begin{bmatrix} C_X(\gamma_{r\omega})A_{F\omega} \\ C_Y(\gamma_{r\omega})A_{L\omega} \\ C_Z(\gamma_{r\omega})A_{F\omega} \\ C_K(\gamma_{r\omega})A_{L\omega}H_{L\omega} \\ C_M(\gamma_{r\omega})A_{F\omega}H_{F\omega} \\ C_N(\gamma_{r\omega})A_{L\omega}L_{oa} \end{bmatrix} \quad (3.40)$$

where

$$V_{r\omega} = \sqrt{u_{r\omega}^2 + v_{r\omega}^2} \quad (3.41)$$

$$\gamma_{r\omega} = -\arctan2(v_{r\omega}, u_{r\omega}) \quad (3.42)$$

with

$$u_{r\omega} = u - V_{\omega} \cos(\beta_{\omega} - \psi) \quad (3.43)$$

$$v_{r\omega} = v - V_{\omega} \sin(\beta_{\omega} - \psi) \quad (3.44)$$

3.7. Summary

The combined six degrees of freedom equations of motion is shown below:

$$\dot{n} = J(n)v \quad (3.45)$$

$$M\dot{v} = -C(v)v - D(v)v - g(n) + \tau + \tau_E \quad (3.46)$$

where

$$M = M_{RB} + M_A \quad (3.47)$$

$$C(v) = C_{RB}(v) + C_A(v) \quad (3.48)$$

$\mathbf{J}(\mathbf{n})$, equation 3.12 and 3.13, is the transformation matrix which translate \mathbf{v}_1 and \mathbf{v}_2 through the functions of the Euler angles to $\dot{\mathbf{n}}_1$ and $\dot{\mathbf{n}}_2$. $\mathbf{C}(\mathbf{v})$ is the linear combination of the rigid body Coriolis and centripetal matrix $\mathbf{C}_{RB}(\mathbf{v})$, equation A.2 and the hydrodynamic Coriolis and centripetal matrix $\mathbf{C}_A(\mathbf{v})$, equation 3.27. $\mathbf{D}(\mathbf{v})$, equation 3.29, is the hydrodynamic damping and $\mathbf{g}(\mathbf{n})$, equation 3.28, is the restoring forces and moments, equation 3.32. The propulsion forces and moments is modelled by τ and the environmental disturbances by τ_E .

3.8. Simulation Results

Chapter 4

Control Properties of Ocean Vessels

In this chapter, control properties of ocean vessels is presented. Practical implementation and constraints are also introduced. A review of previous work on control of ocean vessels is also discussed.

4.1. Controllability Properties

4.1.1. Acceleration Constraints

The number(m_c) of independent control inputs is smaller than the number(n_c) of degrees of freedom to be controlled for a ocean vessel model. As such you remove all zero elements in τ and denote the resulting vector by τ_a . Thus, if $\tau_a \in \mathbb{R}^{m_c}$ and $\mathbf{n} \in \mathbb{R}^{n_c}$, then $m_c \geq n_c$.

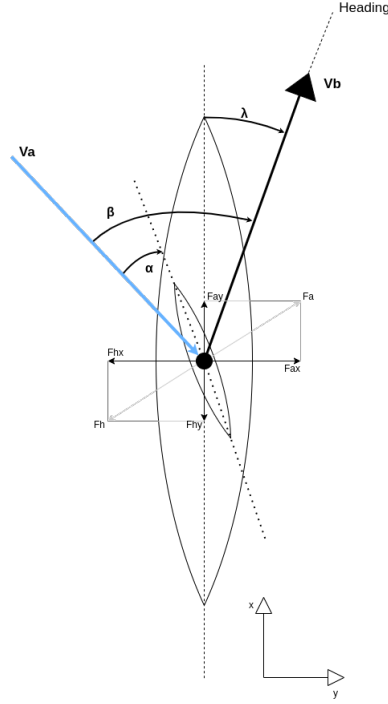
Chapter 5

Modeling of a Fixed-Wing Sail, Keel and Rudder

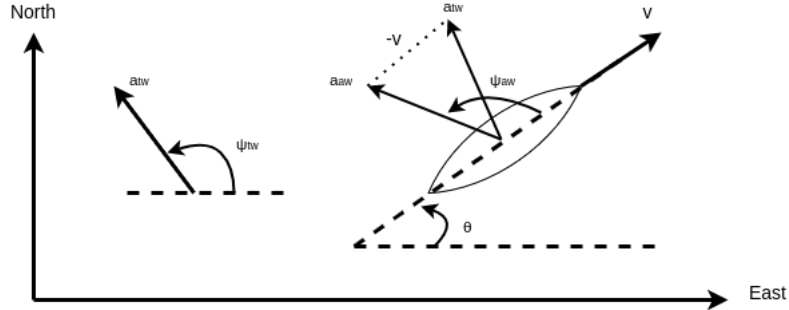
This chapter the modeling of a fixed-wing sail is considered. Traditional sails consist of a mainsail and a jib [5]. The sail considered in the chapter is a fixed-wing sail that is fully autonomous. The sail takes inspiration from a free rotating fixed wing sail [6] and fixed-wing sail [7]. The chapter models the forces experienced by adding a fixed-wing sail to the model described in Chapter 3. Also discussed in this chapter is the modelling of forces, moments and constraints caused by a rudder and keel.

5.1. Fixed-wing Sail

A simplified model, all the acting forces are assumed to be acting in the horizontal plane, are shown in Figure 5.1. The aerodynamic force has both a side force component F_{ax} and a forward force component F_{ay} . These forces need to be balanced by the corresponding hydrodynamic side force F_{Hx} and the resistance F_{Hy} . In order to generate the hydrodynamic side force, the hull needs to have a velocity relative to the water and an angle of attack, the leeway angle λ relative to the water flow field. The rudder of the boat is used to generate side force with the aim of controlling the longitudinal position of the sailboat.

**Figure 5.1:** Top view of sailboat

Wind is both described in the n-frame and in the b-frame corresponding to true wind(tw) and apparent wind(aw) respectively. The definition of true wind and apparent wind is illustrated in Figure 5.2.

**Figure 5.2:** Definition of true and apparent wind

The apparent wind in Cartesian relative to the direction of the boat, i.e., the first coordinate corresponding to the heading of the boat can be calculated from true wind by

$$\mathbf{W}_{c,aw} = \begin{bmatrix} a_{tw} \cos(\psi_{tw} - \theta) - v \\ a_{tw} \sin(\psi_{tw} - \theta) \end{bmatrix} \quad (5.1)$$

The corresponding polar coordinates are thus

$$\mathbf{W}_{p,aw} = \begin{bmatrix} a_{aw} \\ \psi_{aw} \end{bmatrix} = \begin{bmatrix} |\mathbf{W}_{c,aw}| \\ \text{atan2}(\mathbf{W}_{c,aw}) \end{bmatrix} \quad (5.2)$$

The angle of attack, α_s , on the sail is determined by the direction of the apparent wind and the angle of the sail, $\alpha_s = \psi_{aw} - \lambda$. The force on the sail is this given by,

$$F_s = -La_{aw}\sin(\lambda - \psi_{aw}) \quad (5.3)$$

where L is the lift force and D is the drag force and are defined by:

$$L = \frac{1}{2}\rho A v_a^2 C_L \quad (5.4)$$

$$D = \frac{1}{2}\rho A v_a^2 C_D \quad (5.5)$$

where ρ is the air density, A is the plane area of the foil, v_a is the apparent wind velocity defined in equation, C_L and C_D are the lift and drag coefficients respectively. These coefficients depends on the winds angle of attack, α and the foils shape. In [8] it is shown that the coefficients can be approximated as,

$$C_L = k_1 \sin(2\alpha) \quad (5.6)$$

$$C_D = k_1 (1 - \cos(2\alpha)) \quad (5.7)$$

Chapter 6

Stability Analysis

Chapter 7

Platform Development

Chapter 8

Control Techniques for a Sail and Rudder

Chapter 9

Model Simulation

Chapter 10

Results

Chapter 11

Conclusion

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Appendix A

Additional Modelling Information

A.1. Notation and Vector Definitions

Table A.1: SNAME Notation for ocean vessels

Degree of freedom		Force and moment	Linear and angular velocity	Position and Euler angles
1	Surge	X	u	x
2	Sway	Y	v	y
3	Heave	Z	w	z
4	Roll	K	p	ϕ
5	Pitch	M	q	θ
6	Yaw	N	r	ψ

Table A.2: Rigid body motion vectors

Vector	Components	Definition
\mathbf{f}_{Ob}	$[X \ Y \ Z]^T$	force decomposed in the body-fixed frame
\mathbf{m}_{Ob}	$[K \ M \ N]^T$	moment decomposed in the body-fixed frame
\mathbf{v}_{Ob}	$[u \ v \ w]^T$	linear velocity decomposed in the body-fixed frame
\mathbf{w}_{Ob}^E	$[p \ q \ r]^T$	angular velocity of the body-fixed relative to the earth-fixed frame
\mathbf{r}_{Ob}	$[x_g \ y_g \ z_g]^T$	vector from O_b to CG decomposed in the body-fixed frame

A.2. Modeling Equations

$$\mathbf{M}_{\mathbf{RB}} = \begin{bmatrix} m & 0 & 0 & mz_g & mz_g & -my_g \\ 0 & m & 0 & 0 & 0 & mx_g \\ 0 & 0 & m & -mx_g & -mx_g & 0 \\ 0 & -mz_g & -my_g & I_x & -I_{xy} & -I_{xz} \\ mz_g & 0 & -mx_g & -I_{xy} & I_y & -I_{yz} \\ -my_g & mx_g & 0 & -I_{xz} & -I_{zy} & I_z \end{bmatrix} \quad (\text{A.1})$$

$$\mathbf{C}_{\mathbf{RB}}(\mathbf{v}) = \begin{bmatrix} 0 & 0 & 0 & m(y_g q + z_g r) & -m(x_g q - w) & -m(x_g r + v) \\ 0 & 0 & 0 & -m(y_g p + w) & m(z_g r + x_g p) & -m(y_g r - u) \\ 0 & 0 & 0 & -m(z_g p - v) & -m(z_g q + u) & m(x_g p + y_g q) \\ -m(y_g q + z_g r) & m(y_g p + w) & m(y_g p - v) & 0 & -I_{yz} q - I_{xz} q + I_z r & I_{yz} r + I_{xy} p - I_y q \\ m(x_g p - w) & -m(z_g r - x_g p) & m(z_g q + u) & I_{yz} q + I_{xz} p - I_z r & 0 & -I_{xz} r - I_{xy} q + I_x p \\ m(x_g r + v) & m(y_g r - u) & -m(x_g p + y_g q) & -I_{yz} r - I_{xy} p + I_y q & I_{xz} r + I_{xy} q - I_x p & 0 \end{bmatrix} \quad (\text{A.2})$$

$$a_1 = X_{\dot{u}}u + X_{\dot{v}}v + X_{\dot{w}}w + X_{\dot{p}}p + X_{\dot{q}}q + X_{\dot{r}}r \quad (\text{A.3})$$

$$a_2 = Y_{\dot{u}}u + Y_{\dot{v}}v + Y_{\dot{w}}w + Y_{\dot{p}}p + Y_{\dot{q}}q + Y_{\dot{r}}r \quad (\text{A.4})$$

$$a_3 = Z_{\dot{u}}u + Z_{\dot{v}}v + Z_{\dot{w}}w + Z_{\dot{p}}p + Z_{\dot{q}}q + Z_{\dot{r}}r \quad (\text{A.5})$$

$$b_1 = K_{\dot{u}}u + K_{\dot{v}}v + K_{\dot{w}}w + K_{\dot{p}}p + K_{\dot{q}}q + K_{\dot{r}}r \quad (\text{A.6})$$

$$b_2 = M_{\dot{u}}u + M_{\dot{v}}v + M_{\dot{w}}w + M_{\dot{p}}p + M_{\dot{q}}q + M_{\dot{r}}r \quad (\text{A.7})$$

$$b_3 = N_{\dot{u}}u + N_{\dot{v}}v + N_{\dot{w}}w + N_{\dot{p}}p + N_{\dot{q}}q + N_{\dot{r}}r \quad (\text{A.8})$$