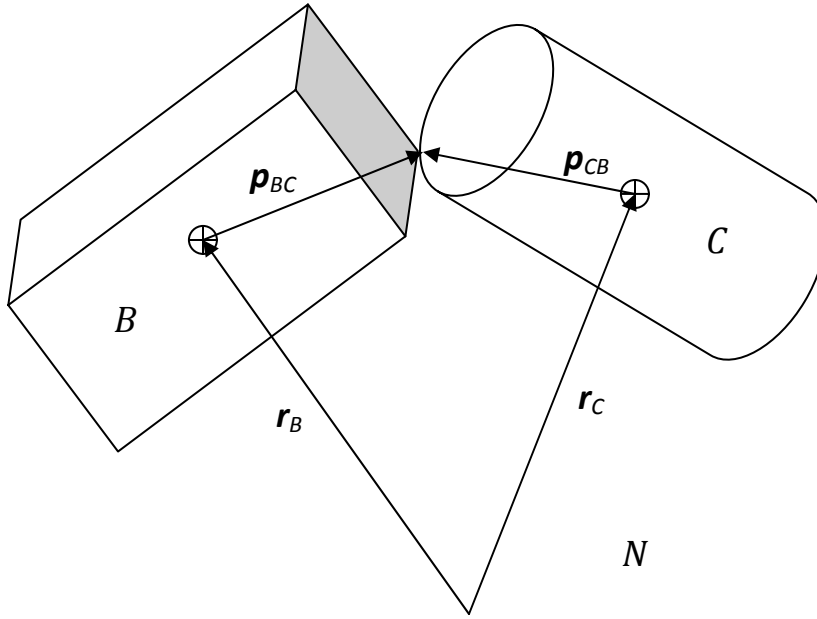


Ball joint (spherical joint) constraint expression: a point fixed on body B must be at the same location in inertial space as a point fixed on body C :



In vector notation, this is $\mathbf{r}_B + \mathbf{p}_{BC} = \mathbf{r}_C + \mathbf{p}_{CB}$ or $(\mathbf{r}_B + \mathbf{p}_{BC}) - (\mathbf{r}_C + \mathbf{p}_{CB}) = 0$. For the Udwadia-Kalaba form of the constraint, we need to take two derivatives of this expression. The first derivative (in the N frame) is:

$$\begin{aligned} \frac{{}^N d}{dt} ((\mathbf{r}_B + \mathbf{p}_{BC}) - (\mathbf{r}_C + \mathbf{p}_{CB})) &= 0 \\ \frac{{}^N d}{dt} (\mathbf{r}_B) + \frac{{}^B d}{dt} \mathbf{p}_{BC} + \boldsymbol{\omega}^{B/N} \times \mathbf{p}_{BC} - \frac{{}^N d}{dt} \mathbf{r}_C - \frac{{}^C d}{dt} \mathbf{p}_{CB} - \boldsymbol{\omega}^{C/N} \times \mathbf{p}_{CB} &= 0 \end{aligned}$$

Since \mathbf{p}_{BC} and \mathbf{p}_{CB} are fixed in the frames B and C , respectively,

$$\frac{{}^N d}{dt} (\mathbf{r}_B) - \frac{{}^N d}{dt} \mathbf{r}_C + \boldsymbol{\omega}^{B/N} \times \mathbf{p}_{BC} - \boldsymbol{\omega}^{C/N} \times \mathbf{p}_{CB} = 0$$

The second derivative (in the N frame) is:

$$\begin{aligned} \frac{{}^N d^2}{dt^2} \mathbf{r}_B - \frac{{}^N d^2}{dt^2} \mathbf{r}_C + \frac{{}^N d}{dt} (\boldsymbol{\omega}^{B/N} \times \mathbf{p}_{BC}) - \frac{{}^N d}{dt} (\boldsymbol{\omega}^{C/N} \times \mathbf{p}_{CB}) &= 0 \\ \frac{{}^N d^2}{dt^2} \mathbf{r}_B - \frac{{}^N d^2}{dt^2} \mathbf{r}_C + \frac{{}^N d}{dt} (\boldsymbol{\omega}^{B/N}) \times \mathbf{p}_{BC} + \boldsymbol{\omega}^{B/N} \times (\boldsymbol{\omega}^{B/N} \times \mathbf{p}_{BC}) - \frac{{}^N d}{dt} (\boldsymbol{\omega}^{C/N}) \times \mathbf{p}_{CB} - \boldsymbol{\omega}^{C/N} \times (\boldsymbol{\omega}^{C/N} \times \mathbf{p}_{CB}) &= 0 \end{aligned}$$

Now we need to get things in the form $A\ddot{\mathbf{x}} = \mathbf{b}$. The QuIRK state vector consists of the center-of-mass position for each body and the quaternion rotation from the body to inertial frame.

Exploiting the relations $\frac{{}^N d}{dt} (\boldsymbol{\omega}^{B/N}) = \frac{{}^B d}{dt} (\boldsymbol{\omega}^{B/N})$ and $\mathbf{a} \times \mathbf{b} = -\mathbf{b} \times \mathbf{a}$,

$$\frac{{}^N d^2}{dt^2} \mathbf{r}_B - \frac{{}^N d^2}{dt^2} \mathbf{r}_C - \mathbf{p}_{BC} \times \frac{{}^B d}{dt} (\boldsymbol{\omega}^{B/N}) + \mathbf{p}_{CB} \times \frac{{}^C d}{dt} (\boldsymbol{\omega}^{C/N}) + \boldsymbol{\omega}^{B/N} \times (\boldsymbol{\omega}^{B/N} \times \mathbf{p}_{BC}) - \boldsymbol{\omega}^{C/N} \times (\boldsymbol{\omega}^{C/N} \times \mathbf{p}_{CB}) = 0$$

Now we cast every vector onto the N basis vectors:

$${}^N \ddot{\mathbf{r}}_B - {}^N \ddot{\mathbf{r}}_C - ({}^N Q^B {}^B \mathbf{p}_{BC})^\times {}^N Q^B {}^B \dot{\boldsymbol{\omega}}^{B/N} + ({}^N Q^C {}^C \mathbf{p}_{CB})^\times {}^N Q^C {}^C \dot{\boldsymbol{\omega}}^{C/N} + \left(({}^N Q^B {}^B \boldsymbol{\omega}^{B/N})^\times \right)^2 {}^N Q^B {}^B \mathbf{p}_{BC} - \left(({}^N Q^C {}^C \boldsymbol{\omega}^{C/N})^\times \right)^2 {}^N Q^C {}^C \mathbf{p}_{CB} = 0$$

A superscript $^\times$ represents the skew-symmetric cross-product matrix form of a vector, and ${}^N Q^B$ is the direction-cosine matrix that transforms a vector from the B basis to the N basis coordinates.

Now we insert the identities ${}^B \boldsymbol{\omega}^{B/N} = T_B \dot{\mathbf{q}}_B$ and ${}^C \boldsymbol{\omega}^{C/N} = T_C \dot{\mathbf{q}}_C$ along with their derivatives:

$${}^B \boldsymbol{\omega}^{B/N} = \frac{2[q_4 I - q_{123}^\times - q_{123}]}{T_B} \dot{\mathbf{q}}_B$$

$${}^N \ddot{\mathbf{r}}_B - {}^N \ddot{\mathbf{r}}_C - ({}^N Q^B {}^B \mathbf{p}_{BC})^\times {}^N Q^B (\dot{T}_B \dot{\mathbf{q}}_B + T_B \ddot{\mathbf{q}}_B) + ({}^N Q^C {}^C \mathbf{p}_{CB})^\times {}^N Q^C (\dot{T}_C \dot{\mathbf{q}}_C + T_C \ddot{\mathbf{q}}_C) + \left(({}^N Q^B T_B \dot{\mathbf{q}}_B)^\times \right)^2 {}^N Q^B {}^B \mathbf{p}_{BC} - \left(({}^N Q^C T_C \dot{\mathbf{q}}_C)^\times \right)^2 {}^N Q^C {}^C \mathbf{p}_{CB} = 0$$

Finally, we rearrange terms so that those involving the second derivatives of states are isolated on one side of the expression:

$$\begin{aligned} & {}^N \ddot{\mathbf{r}}_B - {}^N \ddot{\mathbf{r}}_C - ({}^N Q^B {}^B \mathbf{p}_{BC})^\times {}^N Q^B T_B \ddot{\mathbf{q}}_B + ({}^N Q^C {}^C \mathbf{p}_{CB})^\times {}^N Q^C T_C \ddot{\mathbf{q}}_C \\ & = ({}^N Q^B {}^B \mathbf{p}_{BC})^\times {}^N Q^B \dot{T}_B \dot{\mathbf{q}}_B - ({}^N Q^C {}^C \mathbf{p}_{CB})^\times {}^N Q^C \dot{T}_C \dot{\mathbf{q}}_C - \left(({}^N Q^B T_B \dot{\mathbf{q}}_B)^\times \right)^2 {}^N Q^B {}^B \mathbf{p}_{BC} \\ & + \left(({}^N Q^C T_C \dot{\mathbf{q}}_C)^\times \right)^2 {}^N Q^C {}^C \mathbf{p}_{CB} \\ & \left[\begin{array}{cc} I & -I \end{array} \begin{array}{cc} -({}^N Q^B {}^B \mathbf{p}_{BC})^\times {}^N Q^B T_B & ({}^N Q^C {}^C \mathbf{p}_{CB})^\times {}^N Q^C T_C \end{array} \right] \begin{bmatrix} {}^N \ddot{\mathbf{r}}_B \\ {}^N \ddot{\mathbf{r}}_C \\ \ddot{\mathbf{q}}_B \\ \ddot{\mathbf{q}}_C \end{bmatrix} \\ & = \left[({}^N Q^B {}^B \mathbf{p}_{BC})^\times {}^N Q^B \dot{T}_B \dot{\mathbf{q}}_B - ({}^N Q^C {}^C \mathbf{p}_{CB})^\times {}^N Q^C \dot{T}_C \dot{\mathbf{q}}_C - \left(({}^N Q^B T_B \dot{\mathbf{q}}_B)^\times \right)^2 {}^N Q^B {}^B \mathbf{p}_{BC} \right. \\ & \left. + \left(({}^N Q^C T_C \dot{\mathbf{q}}_C)^\times \right)^2 {}^N Q^C {}^C \mathbf{p}_{CB} \right] \end{aligned}$$

Now we have the 3×14 and 3×1 matrices:

$$\begin{aligned} A &= \begin{bmatrix} I & -I & -({}^N Q^B {}^B \mathbf{p}_{BC})^\times {}^N Q^B T_B & ({}^N Q^C {}^C \mathbf{p}_{CB})^\times {}^N Q^C T_C \end{bmatrix} \\ b &= \begin{bmatrix} ({}^N Q^B {}^B \mathbf{p}_{BC})^\times {}^N Q^B \dot{T}_B \dot{\mathbf{q}}_B - ({}^N Q^C {}^C \mathbf{p}_{CB})^\times {}^N Q^C \dot{T}_C \dot{\mathbf{q}}_C - \left(({}^N Q^B T_B \dot{\mathbf{q}}_B)^\times \right)^2 {}^N Q^B {}^B \mathbf{p}_{BC} \\ + \left(({}^N Q^C T_C \dot{\mathbf{q}}_C)^\times \right)^2 {}^N Q^C {}^C \mathbf{p}_{CB} \end{bmatrix} \end{aligned}$$