A Compact ENDF (ACE) Format Specification

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Todo list

| Don't forget the 2.0.1 | Header | Opening | | | | | | | | | | 8 |
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| This needs to be done | <u>)</u> | | | | | | | | | | | 13 |

1 Introduction

The ACE format consists of two *types* and many *classes* of data. The data are kept in an ACE Table. The term ACE Table and ACE file are often used interchangeably.

1.1 Types of ACE-Formatted Data

There are two types of ACE-formatted data; simply called Type 1 and Type 2.

- **Type 1** Standard formatted tables. These tables contain ASCII text and are machine independent; they are readable on every machine.
- Type 2 Standard unformatted tables. These tables are binary and can be generated from the Type 1 files. They are more compact and faster to read than the Type 1 ACE Tables but are machine/platform dependent; they are not readable on every machine.

Traditionally Type 2 ACE files were more commonly used because they were smaller in size and faster to read. However due to the fact that they are not portable across machines and platforms they have fallen out of fashion.

1.2 Classes of ACE-Formatted Data

There are many classes of ACE-formatted data:

- 1. continuous-energy neutron (see Section 4),
- 2. discrete-reaction neutron,
- 3. neutron dosimetry (see Section 5),
- 4. $S(\alpha, \beta)$ thermal (see Section 6),
- 5. continuous-energy photoatomic (see Section 7),
- 6. continuous-energy electron interaction,
- 7. continuous-energy photonuclear interaction,
- 8. multigroup-energy neutron, and
- 9. multigroup-energy photoatomic.

Each of these classes of data are described later in this document.

An ACE Table is an entity that contains evaluation-dependent data about one of the many classes of data for a specific material—an target isotope, isomer, or element. For a given ZAID, the data contained on a Type 1 and Type 2 tables are identical. Simulations run with one type of data should produce identical results as those run with the other type of data.

1.3 ACE Libraries

A collection of ACE data tables that derive from a single set of evaluation files are typically grouped together in a "library"—not to be confused from the evaluation library from which they derive. Multiple ACE data tables can concatenated into the same logical file on the computer, although this has fallen somewhat out of fashion due to the large amount of data on each ACE table derived from modern evaluation files. Applications

that use ACE-formatted data should produce the same results regardless of whether the tables are contained in one logical file on the computer or spread across many.

2 ACE Tables

An ACE Table consists of a Header followed by an array (XSS) containing the actual data. The Header and XSS array are the same regardless of whether the ACE Table is Type 1 or Type 2. Each line in a Type 1 ACE Table is 80 characters or less.

2.1 ACE Header

The first section of an ACE Table is the Header. The ACE Header contains metadata¹ about the ACE Table. The Header consists of four parts:

- 1. Opening,
- 2. IZAW array,
- 3. NXS array, and
- 4. JXS array.

An example of an ACE Table Header (from ¹H in the ENDf71x library) is given in Figure 1 with each part highlighted a different color.

| 1 | 1001.80c | 0.999167 | 2.5301 | E-08 12/17/ | /12 | | | |
|----|------------|------------|--------|--------------|------|------------|-----------|---------|
| 2 | H1 ENDF71x | (jlconlin) | Ref. | see jlconlin | (ref | 09/10/2012 | 10:00:53) | mat 125 |
| 3 | 0 | 0. | 0 | 0. | 0 | 0. | 0 | 0. |
| 4 | 0 | 0. | 0 | 0. | 0 | 0. | 0 | 0. |
| 5 | 0 | 0. | 0 | 0. | 0 | 0. | 0 | 0. |
| 6 | 0 | 0. | 0 | 0. | 0 | 0. | 0 | 0. |
| 7 | 17969 | 1001 | 590 | 3 | 0 | 1 | 1 | 0 |
| 8 | 0 | 1 | 1 | 0 | 0 | 0 | 0 | 0 |
| 9 | 1 | 0 | 2951 | 2954 | 2957 | 2960 | 2963 | 4352 |
| 10 | 4353 | 5644 | 5644 | 5644 | 6234 | 6235 | 6236 | 6244 |
| 11 | 6245 | 6245 | 6246 | 16721 | 0 | 16722 | 0 | 0 |
| 12 | 0 | 0 | 0 | 0 | 0 | 16723 | 16724 | 16725 |

Figure 1: Header example. The (Legacy) Opening (lines 1–2) is in red, the IZAW array (lines 3–6) is in blue, the NXS array (lines 7–8) is in teal, and the JXS array (lines 9–12) is in violet.

2.1.0.0.1 Legacy Header Opening There are two slightly different formats for the Header Opening. The most common one found is called here the Legacy Opening and is the one demonstrated in the Header example in Figure 1.

The Legacy Opening consists of several variables given over two 80-character lines. The variables and the Fortran format for reading the variable is given in Table 1

2.1.0.0.2 2.0.1 Header Opening

Don't forget the 2.0.1 Header Opening

¹data about the data

| Line | Variable | Format | Description |
|------|----------|------------|--|
| 1 | HZ | A10 | ZAID (see Section 3.1) |
| 1 | AW | E12.0 | Atomic weight ratio |
| 1 | TZ | E12.0 | Temperature |
| 1 | _ | 1 X | (blank space) |
| 1 | HD | A10 | Processing date |
| 2 2 | HK HM | A70 A10 | Descriptive string 10-character material identifier |

Table 1: Variables in the Legacy Opening part of the ACE Header.

| Line | Variable | Format | Description |
|-------------|----------|------------|-----------------------------------|
| 1 | VERS | A10 | Version format string |
| 1 | HZ | A24 | SZAID (see Section 3.2) |
| 1 | SRC | A24 | Evaluation source |
| 2 | AW | E12.0 | Atomic weight ratio |
| 2 | TZ | E12.0 | Temperature |
| 2 | _ | 1 X | (blank space) |
| 2 | HD | A10 | Processing date |
| 2 | N | I10 | Number of comment lines to follow |
| 3 - (N + 2) | | A70 | N comment lines |

Table 2: Variables in the 2.0.1 Opening part of the ACE Header.

There is a limitation to the number of unique ZA IDs for a given ZA; 100 different IDs, in fact, for each class of ACE Table. To overcome this limitation, a new Header Opening[1] was developed in 2012 and updated a few years later to correct some errors.

```
2.0.1 1001.800nc ENDF/B-VIII.0-B1

0.999167 2.5301e-08 2018-05-02 2

1001.00c 0.999167 2.5301E-08 05/02/18

H1 Lib80x (jlconlin) Ref. see jlconlin (ref 01/29/2018 07:54) mat 125
```

Figure 2: Header Opening example. The Legacy Opening is shown in blue while the 2.0.1 Opening consists of the red and the blue portions.

Note that a Legacy Header Opening can be contained in the comment section of the 2.0.1 Header Opening. This was designed explicitly to allow backwards compatibility while application codes were modified to be able to handle. An example of this is shown in Figure 2. Codes that cannot read the 2.0.1 Header can be told (typically via an xsdir[1] entry) to start reading the ACE Table several lines after the beginning of the

2.0.1 Header.

Following the Opening of the Header are three arrays, IZAW, NXS, and JXS respectively. They are each described below. Immediately following the JXS array is the XSSarray.

2.1.1 IZAW Array

The IZAW array follows on the lines immediately following the Header. It consists of 16 pairs of ZA's (IZ) and atomic weight ratios (AW). The IZ entries are still needed for $S(\alpha, \beta)$ Tables to indicate for which isotope(s) the scattering data are appropriate.

The 16 pairs of numbers are spread over 4 lines. The Fortran format for reading/writing the numbers on one line is: 4(I7,F11.0).

2.1.2 NXS Array

The NXS array comes on the 2 lines after the IZAW array. The NXS array has 16 integer elements; 8 on each line. The Fortran format for reading/writing the numbers on each line is: 819. The first element of the NXS array indicates how many numbers are in the XSS array. The remainder of the NXS array elements (usually) indicate how many of different pieces of data there is.

2.1.3 JXS Array

The JXS array comes on the 4 lines after the NXS array. The JXS array has 32 integer elements; 8 on each line. The Fortran format for reading/writing the numbers on each line is: 819. The JXS array contains indices to the XSS array where difference pieces of data begins.

The specific definition of the elements of the NXS and JXS arrays are dependent on the class of data in the Table and are defined in the section of this document that describes each class of data.² Note that not all elements of the arrays are (currently) being used, allowing for future expansion.

2.2 The XSS Array

After the ACE Header comes the XSS array. It is typically *very* large with hundreds of thousands of elements. It is broken up into blocks with the blocks being dependent on the class of data that is contained in the table. The description and definition of each of these blocks can be found in the descriptions later in this document.

The data is written with 4 floating-point numbers on each 80-character line. All data in the XSS array can be read using the Fortran format: 4E20.0 for each line.

²See, for example, Table 3 and Table 4.

```
2.0.0
              1001.710nc
                                       ENDFB-VII.1
  0.999167 2.5301E-08 12/17/12
                                    3
_3 \hspace{-0.5mm} \mid The next two lines are the first two lines of 'old-style' ACE.
    1001.80c
                 0.999167 2.5301E-08
                                        12/17/12
  H1 ENDF71x (jlconlin) Ref. see jlconlin (ref 09/10/2012 10:00:53)
                                                                             mat 125
  1.000000000E-11
                      1.03125000000E-11
                                           1.06250000000E-11
                                                                1.09375000000E-11
  1.1250000000E-11
                       1.15625000000E-11
                                           1.18750000000E-11
                                                                1.21875000000E-11
  1.2500000000E-11
                       1.28125000000E-11
                                           1.31250000000E-11
                                                                1.34375000000E-11
  1.3750000000E-11
                       1.43750000000E-11
                                           1.5000000000E-11
                                                                1.5625000000E-11
10
  1.6250000000E-11
                       1.6875000000E-11
                                           1.7500000000E-11
                                                                1.8125000000E-11
  1.8750000000E-11
                       1.9375000000E-11
                                           2.0000000000E-11
                                                                2.09375000000E-11
  2.18750000000E-11
                       2.28125000000E-11
                                           2.3750000000E-11
                                                                2.46875000000E-11
  2.56250000000E-11
                       2.65625000000E-11
                                           2.7500000000E-11
                                                                2.84375000000E-11
  2.9375000000E-11
                       3.03125000000E-11
                                           3.1250000000E-11
                                                                3.21875000000E-11
  3.3125000000E-11
                       3.40625000000E-11
                                           3.5000000000E-11
                                                                3.59375000000E-11
```

Figure 3: ACE Header with beginning of XSS array for ¹H from the ENDF71x library. Note this uses the 2.0.1 Header with backwards compatibility with the Legacy Header.

3 Unique ACE Table Identifier

This needs to be done.

Each ACE Table needs to have an identifier to uniquely distinguish the data that is contained in the Table.

3.1 **ZAID**

3.2 SZAID

With the introduction of the 2.0.1 ACE Header, the identifier was modified to better specify the metastable state of the material as well as expand the available space for identifiers.

The new identifier is referred to as a SZAID³.

³pronounced "ess-ZAID"

4 Continuous-Energy and Discrete Neutron Transport Tables

The format of individual blocks found on neutron transport tables is identical for continuousenergy and discrete-reaction ACE Tables; the format for both are described in this section. The blocks of data are:

- 1. **ESZ Block**—contains the main energy grid for the Table and the total, absorption, and elastic cross sections as well as the average heating numbers. The ESZ Block block always exists. See Section 4.3.1.
- 2. NU Block—contains prompt, delayed and/or total $\overline{\nu}$ as a function of incident neutron energy. The NU Block exists only for fissionable isotopes; that is, if JXS(2) \neq 0. See Section 4.3.2.
- 3. MTR Block—contains a list of ENDF MT numbers for all neutron reactions other than elastic scattering. The MTR Block exists for all isotopes that have reactions other than elastic scattering; that is, all isotopes with NXS(4) \neq 0. See Section 4.3.3.
- 4. LQR Block—contains a list of kinematic Q-values for all neutron reactions other than elastic scattering. The LTR Block exists if NXS(4) \neq 0. See Section 4.3.4.
- 5. TYR Block—contains information about the type of reaction for all neutron reactions other than elastic scattering. Information for each reaction includes the number of secondary neutrons and whether secondary neutron angular distributions are in the laboratory or center-of-masssystem. The TYR Block exists if NXS(4) \neq 0. See Section 4.3.5.
- 6. LSIG Block—contains a list of cross section locators for all neutron reacitons other than elastic scattering. The LSIG Block exists if NXS(4) \neq 0. See Section 4.3.6
- 7. SIG Block—contains cross sections for all reactions other than elastic scattering. The SIG Block exists if NXS(4) \neq 0. See Section 4.3.7.
- 8. LAND Block—contains a list of angular-distribution locators for all reactions producing secondary neutrons. The LAND Block always exists. See Section 4.3.8.
- 9. AND Block—contains list angular distributions for all reactions producing secondary neutrons. The AND Block always exists. See Section 4.3.9.
- 10. LDLW Block—contains a list of energy distributions for all reactions producing secondary neutrons except for elastic scattering. The LDLW Block exists if NXS(5)≠ 0. See Section 4.3.10.
- 11. DLW Block—contains energy distributions for all reactions producing secondary neutrons except for elastic scattering. The DLW Block exists if NXS(5) \neq 0. See Section 4.3.11.
- 12. **GPD Block**—contains the total photon production cross section tabulated on the ESZ energy grid and a $30 \times$ matrix of secondary photon energies. The GPD Block exists only for those older evaluations that provide coupled neutron/photon information; that is, if $JXS(12) \neq 0$. See Section 4.3.12.
- 13. MTRP Block—contains a list of MT numbers for all photon production reactions. The term "photon production reaction" is used for any information describing a specific neutron-in, photon-out reaction. The MTR Block exists if NXS(6) \neq 6. See

- Section 4.3.3.
- 14. LSIGP Block—contains a list of cross section locators for all photon production reactions. The LSIGP Block exists if NXS(6) \neq 0. See Section 4.3.6.
- 15. SIGP Block—contains cross sections for all photon production reactions. The SIGP Block exists if NXS(6) \neq 0. See Section 4.3.13.
- 16. LANDP Block—contains a list of angular-distribution locators for all photon production reactions. The LANDP Block exist if NXS(6) \neq 0. See Section 4.3.14
- 17. ANDP Block—contains photon angular distributions for all photon production reactions. The ANDP Block exists if NXS(6) \neq 0. See Section 4.3.15.
- 18. LDLWP Block—contains a list of energy-distribution locators for all photon production reactions. The LDLWP Block exists if NXS(6) ≠ 0. See Section 4.3.10.
- 19. DLWP Block—contains photon energy distributions for all photon production reactions. The DLWP Block exists if NXS(6) \neq 0. See Section 4.3.11.
- 20. YP Block—contains a list of MT identifiers of neutron reaction cross sections required as photon production yield multipliers. The YP Block exists if NXS(6) \neq 0. See Section 4.3.16.
- 21. FIS Block—contains the total fission cross section tabulated on the ESZ energy grid. The FIS Block exists if $JXS(21) \neq 0$. See Section 4.3.17.
- 22. UNR Block—contains the unresolved resonance range probability tables. The UNR Block exists if $JXS(23) \neq 0$. See Section 4.3.18.

4.1 NXS Array

Table 3: ${\sf NXS}$ array element definitions for ${\sf NXS}$ ACE Table.

| Element | Name | Description |
|---------|------|--|
| 1 | _ | Length of second block of data (XSS array) |
| 2 | ZA | 1000 * Z + A |
| 3 | NES | Number of energies |
| 4 | NTR | Number of reactions excluding elastic scattering |
| 5 | NR | Number of reactions having secondary neutrons excluding elastic scattering |
| 6 | NTRP | Number of photon production reactions |
| | | |
| 8 | NPCR | Number of delayed neutron precurser families |
| 9 | S | Excited state † |
| 10 | Z | Atomic number [†] |
| 11 | Α | Atomic mass number [†] |
| | | |
| 14 | | $\mathrm{Reserved}^{\ddagger}$ |
| 15 | | $\mathrm{Reserved}^{\ddagger}$ |
| 16 | | $\mathrm{Reserved}^{\ddagger}$ |
| | | |

 $^{^\}dagger$ These values were introduced with the new 2.0.0 Header [1].

[‡] These entries are reserved for the use of transport codes (i.e., MCNP).

4.2 JXS Array

Table 4: ${\sf JXS}$ array element definitions for ${\sf JXS}$ ACE Table.

| Element | Name | Location Description |
|---------|-------|--|
| 1 | ESZ | Energy table |
| 2 | NU | Fission ν data |
| 3 | MTR | MT array |
| 4 | LQR | Q-value array |
| 5 | TYR | Reaction type array |
| 6 | LSIG | Table of cross section locators |
| 7 | SIG | Cross sections |
| 8 | LAND | Table of angular distribution locators |
| 9 | AND | Angular distributions |
| 10 | LDLW | Table of energy distribution locators |
| 11 | DLW | Energy distributions |
| 12 | GPD | Photon production data |
| 13 | MTRP | Photon production MT array |
| 14 | LSIGP | Table of photon production cross section locators |
| 15 | SIGP | Photon production cross sections |
| 16 | LANDP | Table of photon production angular distribution locators |
| 17 | ANDP | Photon production angular distributions |
| 18 | LDLWP | Table of photon production energy distribution locators |
| 19 | DLWP | Photon production energy distributions |
| 20 | YP | Table of yield multipliers |
| 21 | FIS | Total fission cross section |
| 22 | END | Last word of this table |
| 23 | LUNR | Probability tables |
| 24 | DNU | Delayed $\overline{\nu}$ data |
| 25 | BDD | Basic delayed data (λ 's, probabilities) |
| 26 | DNEDL | Table of energy distribution locators |
| 27 | DNED | Energy distributions |
| 32 | | |

4.3 Format of Individual Data Blocks

4.3.1 ESZ Block

The format of the ${\sf ESZ}$ Block is given in Table 5.

Table 5: ESZ Block.

| Location in XSS | Parameter | Description |
|------------------|----------------------------------|--|
| S_{ESZ} | $E(l), l = 1, \dots, N_E$ | Energies |
| $S_{ESZ} + N_E$ | $\sigma_t(l), l = 1, \dots, N_E$ | Total cross section |
| $S_{ESZ}+2N_E$ | $\sigma_a(l), l=1,\ldots,N_E$ | Total neutron disappearance cross section [†] |
| $S_{ESZ} + 3N_E$ | $\sigma_{el}(l), l=1,\ldots,N_E$ | Elastic cross section |
| $S_{ESZ}+4N_E$ | $H_{ave}(l), l = 1, \dots, N_E$ | Average Heating numbers |

[†] The disappearance cross section is defined in [2, Appendix B] as MT101

Note: S_{ESZ} is index of the XSS array where the ESZ Block starts, JXS(1), and N_E is the number of energy energy points, NXS(3).

4.3.2 NU Block

There are three possibilities for the NU Block:

- 1. No NU Block. (JXS(2)=0)
- 2. Either prompt or total $\overline{\nu}$ is given (but not both). (JXS(2) > 0) The NU array begins at location XSS(KNU) where KNU=JXS(2).
- 3. Both prompt and total $\overline{\nu}$ are given. (JXS(2) > 0) The prompt NU array begins at XSS(KNU) where KNU=JXS(2); the total NU array begins at XSS(KNU) where KNU = JXS(2) + ABS(XSS(JXS(2)))+1

In addition to prompt and/or total $\overline{\nu}$, delayed $\overline{\nu}$ is given when JXS(24) > 0. Delayed $\overline{\nu}$ array begins at XSS(KNU) where KNU=JXS(24). Delayed $\overline{\nu}$ must be given in tabulated as described below in Table 7.

The format of the NU Block has two forms (if it exists); polynomial (see Table 6) and tabulated (see Table 7). The format is specified by the LNU flag located in the XSS array at index KNU where KNU is defined above.

Table 6: NU Block—Polynomial function form.

| Location in XSS | Parameter | Description |
|--|---|--|
| $\begin{array}{c} KNU \\ KNU{+1} \\ KNU{+2} \end{array}$ | $egin{aligned} LNU = 1 \ N_C \ C(l), l = 1, \dots, N_C \end{aligned}$ | Polynomial function flag Number of coefficients Coefficients |

When using the polynomial function form of the NU array, $\bar{\nu}$ is reconstructed as

$$\overline{\nu}(E) = \sum_{l=1}^{N_C} C(l) E^{l-1},$$
(1)

where the energy, E, is given in MeV.

Table 7: NU Block—Tabulated form.

| Location in XSS | Parameter | Description |
|----------------------------|-------------------------------------|--|
| KNU | LNU=2 | Tabulated data flag |
| $KNU{+}1$ | N_R | Number of interpolation regions |
| $KNU{+}2$ | $NBT(l), l = 1, \dots, N_R$ | ENDF interpolation parameters |
| $KNU {+} 2 {+} N_R$ | $INT(l), l = 1, \dots, N_R$ | ENDF interpolation scheme [†] |
| $KNU {+} 2{+}2N_R$ | N_E | Number of energies |
| $KNU{+}3{+}2N_R$ | $E(l), l = 1, \dots, N_E$ | Tabulated energy points |
| $KNU {+} 3 {+} 2N_R + N_E$ | $\overline{\nu}(l), l=1,\ldots,N_E$ | Tabulated $\overline{\nu}$ values |

[†] If $N_R = 0$, NBT and INT are omitted and linear-linear interpolation is assumed.

If delayed neutron data exist (when JXS(24)>0), the precursor distribution format is given as in Table 8. The decay constant for the first group DEC_1 is given at XSS(JXS(25)). The precursor distribution immediately follows as described in Table 8. The indices (locators) of the XSS array where each precursor distribution begins (S_{DNU}) can found using the format described in Section 4.3.10 and Section 4.3.11, where LED=JXS(26) and NMT=NXS(8).

Table 8: Delayed $\overline{\nu}$ precursor distribution...

| Location in XSS | Parameter | Description |
|----------------------------|-----------------------------|--|
| S_{DNU} | DEC_i | Decay constant for the <i>i</i> -th group |
| $S_{DNU}{+}1$ | N_R | Number of interpolation regions |
| $S_{KNU}{+}2$ | $NBT(l), l = 1, \dots, N_R$ | ENDF interpolation parameters [†] |
| $S_{KNU}{+}2{+}N_R$ | $INT(l), l = 1, \dots, N_R$ | ENDF interpolation scheme |
| $S_{DNU}{+}2{+}2N_R$ | N_E | Number of energies |
| $S_{DNU}{+}3{+}2N_R$ | $E(l), l = 1, \dots, N_E$ | Tabulated energy points |
| $S_{DNU} + 3 + 2N_R + N_E$ | $P(l), l = 1, \dots, N_E$ | Corresponding probabilities |

[†] If $N_R = 0$, NBT and INT are omitted and linear-linear interpolation is assumed.

Note: S_{DNU} is the index of the XSS array where the delayed $\overline{\nu}$ precursor distribution begins; the first one is at $S_{DNU}=JXS(25)$.

4.3.3 MTR & MTRP Blocks

The format of the MTR Block (for incident neutron reactions) and MTRP Block (for photon production reactions) is given in Table 10. The starting index depends on whether it is the MTR Block or MTRP Block and are given in Table 9.

| Block | LMT | NMT |
|-------|---------|--------|
| MTR | JXS(3) | NXS(4) |
| MTRP | JXS(13) | NXS(6) |

Table 9: LMT and NMT values for the MTR Block and MTR Block.

Table 10: MTR & MTRP Block.

| Location in XSS | Parameter | Description |
|---|---------------|---|
| LMT LMT+1 | MT_1 MT_2 | First ENDF Reaction available Second ENDF Reaction available |
| $ \begin{array}{c} \dots \\ LMT + NMT + 1 \end{array} $ | MT_{NMT} | Last ENDF reaction available |

For the MTR Block, $MT_1, ..., MT_{NMT}$ are standard ENDF MTnumbers; that is, MT=16=(n, 2n); MT=17=(n, 3n); etc. For a complete listing of MT numbers, see [2, Appendix B].

For the MTRP Block, the MT numbers are somewhat arbitrary. To understand the scheme used for numbering the photon production MTs, it is necessary to realize that in the ENDF format, more than one photon can be produced by a particular neutron reaction that is itself specified by a single MT. Each of these photons is produced with an individual energy-dependent cross section. For example, MT102 (radiatiive capture) might be responsible for 40 photons, each with its own cross section, angular distribution, and energy distribution. We need 40 photon MTs to represent the data; the MTs are numbered 1002001, 1002002, ..., 1002040. Therefore, if ENDF MT N is responsible for M photons, we shall number the photon MTs 1000*N+1, 1000*N+2, ..., 1000*N+M.

4.3.4 LQR Block

The format of the LQR Block, containing the reaction-specific Q-values, is given in Table 11. The index at the start of the LQR Block, S_{LQR} =JXS(4). The number of reactions, NMT, is the same through the ACE Table, NMT=NXS(4).

Table 11: LQR Block.

| Location in XSS | Parameter | Description |
|-------------------------|--|--|
| $S_{LQR} \ S_{LQR} + 1$ | $egin{array}{c} Q_1 \ Q_2 \end{array}$ | Q -value for reaction MT_1 Q -value for reaction MT_2 |
| $S_{LQR} + NMT - 1$ | Q_{NMT} | Q -value for reaction MT_NMT |

4.3.5 TYR Block

The format of the TYR Block is given in Table 12. The index at the start of the TYR Block, $S_{\text{TYR}} = JXS(5)$. The number of reactions, NMT, is the same through the ACE Table, NMT=NXS(4).

Table 12: TYR Block.

| Location in XSS | Parameter | Description |
|-------------------------|----------------|--|
| $S_{TYR} \ S_{TYR}{+1}$ | $TY_1 \\ TY_2$ | Neutron release for reaction MT_1 Neutron release for reaction MT_2 |
| $S_{TYR} + NMT - 1$ | TY_{NMT} | Neutron release for reaction MT_{NMT} |

The possible values of TY are ± 1 , ± 2 , ± 3 , ± 4 , ± 19 , 0, and integers greater than 100 in absolute value; the sign indicates the system for scattering: negative=center-of-mass, positive=Lab. Thus if $\mathsf{TY}_i = +3$, three neutrons are released for reaction MT_i and the data on the cross section tables used to determine the exiting neutrons' angles are given in the Lab frame of reference. $\mathsf{TY} = 19$ indicates fission. The number of secondary neutrons released is determined from the fission $\overline{\nu}$ data found in the NU Block. $\mathsf{TY}_i = 0$ indicates absorption (ENDF reactions $\mathsf{MT} > 100$); no neutrons are released. $\|\mathsf{TY}_i\| > 100$ signifies reactions other than fission that have energy-dependent neutron multiplicities. The number of secondary neutrons released is determined from the yield data found in the DLW Block. The MT_i s are given in the MTR Block.

4.3.6 LSIG& LSIGP Blocks

The LSIG Block and LSIGP Block give the locators for cross section array for each reaction MT. A locator is a *relative* index in the XSS array where some piece of data. In this case, the data is the cross section values. The format of the LSIG Block (for incident neutron cross sections) and LSIGP Block (for photon production cross sections) is given in Table 13. The format for the incident neutron cross section arrays is given in Section 4.3.7. The format for the photon production cross sections is given in Section 4.3.13.

All locators are relative to JXS(7) for the LSIG Block or JXS(15) for the LSIGP Block. That is, LXS=JXS(7) for the LSIG Block and LXS=JXS(15) for the LSIGP Block. So the actual cross section data begins at the index LOCA+LXS. The MTs are given in the MTR Block and the MTRP Block for the LSIG Block and the LSIGP Block respectively. LOCA_i must be monotonically increasing.

Table 13: LSIG & LSIGP Block.

| Location in XSS | Parameter | Description |
|-----------------|---------------------|--|
| LXS LXS+1 | - | Location of cross sections for reaction MT_1 Location of cross sections for reaction MT_2 |
| LXS+NMT-1 | LOCA _{NMT} | Location of cross sections for reaction MT_NMT |

4.3.7 SIG Block

The SIG Block contains the incident neutron cross section data. (The photon production cross section is in the SIGP Block.) The format of the SIG Block is given in Table 14. The cross section data begins at the index specified by the locator from the LSIG Block; the format for which is given in Table 15.

Table 14: SIG Block.

| Location in XSS | Description |
|--|--|
| LXS+LOCA ₁ -1 LXS+LOCA ₂ -1 | Cross section array for reaction MT_1 Cross section array for reaction MT_2 |
| LXS+LOCA _{NMT} -1 | Cross section array for reaction MT_NMT |

Note: The number of cross section arrays NMT=NXS(4).

The LOCA_i values are given in the LSIG Block and are all relative to JXS(7). The energy grid index IE_i corresponds to the first energy in the grid at which a cross section is given. The MT_i s are defined in the MTR Block.

Table 15: Cross section array for the i-th reaction..

| Location in XSS | Parameter | Description |
|---------------------------|--|--|
| $LXS + LOCA_{i\text{-}}1$ | IE_i | Energy grid index for reaction MT_i |
| $LXS + LOCA_i$ | $N_{E,i}$ | Number of consecutive entries for MT_i |
| $LXS + LOCA_i{+}1$ | $\sigma_i[E(l)]$ for $l = IE_i, \dots, IE_i + N_{E,i} - 1$ | Cross section for reaction MT_i |

Note: The energy grid, E(l) is given in the ESZ Block.

4.3.8 LAND Block

The LAND Block contains locators for the angular distributions for all reactions producing secondary neutrons. The LAND Block always exists and begins at $S_{\text{LAND}}=JXS(8)$. All locators (LOCB) are relative JXS(9); that is, the angular distribution begins at $JXS(9)+LOCB_i$. The LOCB_i locators must be monotonically increasing. The format of the LAND Block is given in Table 16.

Table 16: LAND Block.

| Location in XSS | Parameter | Description |
|------------------|---------------------|---|
| S_{LAND} | $LOCB_1 = 1$ | Location of angular distribution data for elastic scattering reaction Location of angular distribution data for re- |
| $S_{LAND} + 1$ | $LOCB_2$ | action MT_1 |
| $S_{LAND} + NMT$ | LOCB _{NMT} | Location of angular distribution data for reaction MT_NMT |

Note: S_{LAND} =JXS(8) and NMT=NXS(5) is the number of reactions (excluding elastic scattering).

4.3.9 AND Block

The AND Block contains angular distribution data for all reactions that produce secondary neutrons. The format of the AND Block is given in Table 17. The angular distribution data begins at the index specified by the locator LOCB from the LAND Block. If $LOCB_i=0$ (given in the LAND Block), no angular distribution data are given for reaction i and isotropic scattering is assumed in either the Lab or center-of-mass system. The choice of Lab or center-of-mass system depends upon the value for reaction i in the TYR Block. If $LOCB_i=-1$ no angular distribution data are given for reaction i in the AND Block. The angular distribution data are specified through law=44 in the DLW Block.

Table 17: AND Block.

| Location in XSS | Description |
|--|---|
| $\begin{array}{c} \mathtt{JXS(9)} + \mathtt{LOCB_1-1} \\ \mathtt{JXS(9)} + \mathtt{LOCB_2-1} \\ \mathtt{JXS(9)} + \mathtt{LOCB_{NMT}-1} \end{array}$ | Angular distribution array for elastic scattering Angular distribution array for reaction MT_1 Angular distribution array for reaction MT_{NMT} |

Note: The format for the angular distribution of the *i*-th array is given in Table 18.

Table 18: Angular distribution array for the i-th reaction.

| Location in XSS | Parameter | Description |
|-----------------------------|-----------|--|
| JXS(9)+LOCB _i -1 | N_E | Number of energies at which angular distributions are tabulated. |

Table 18: Angular distribution array for the *i*-th reaction (continued)

| Location in XSS | Parameter | Description |
|---|--|---|
| $\begin{array}{c} \texttt{JXS(9)} \! + \! \texttt{LOCB}_{\mathrm{i}} \\ \texttt{JXS(9)} \! + \! \texttt{LOCB}_{\mathrm{i}} \! + \! N_E \end{array}$ | $E(l), l = 1,, N_E$ $L_C(l), l = 1,, N_E$ | Energy grid Location of tables associated with $E(l)$ |

The angular distribution arrays (Table 18) contains additional locators, L_C ; the sign of these locators is a flag:

- if $L_C(l) > 0$, then $L_C(l)$ points to a 32 equiprobable bin distribution (see Table 19);
- if $L_C(l) < 0$, then $L_C(l)$ points to a tabulated angular distribution (see Table 20);
- if $L_C(l) = 0$, then distribution is isotropic and no further data is needed.

Table 19: Format for the 32 equiprobable bin distribution.

| Location in XSS | Parameter | Description |
|----------------------------------|-----------------------------|---|
| $\texttt{JXS(9)} + L_C(l) - 1$ | $P(1,K)$ $K = 1, \dots, 33$ | 32 equiprobable cosine bins for scattering at energy $E(l)$. |

Table 20: Format for the tabulated angular distribution..

| Location in XSS | Parameter | Description |
|-----------------|--|--------------------------------------|
| $LDAT_l + 1$ | JJ | Interpolation flag [†] |
| $LDAT_l + 2$ | N_P | Number of points in the distribution |
| $LDAT_l + 3$ | $CS_{\mathrm{out}}(j), j=1,\ldots,N_P$ | Cosine scattering angular grid |
| $LDAT_l + 4$ | $PDF(j), j = 1, \dots, N_P$ | Probability density function |
| $LDAT_l + 5$ | $CDF(j), j = 1, \ldots, N_P$ | Cumulative density function |

[†] 0 histogram interpolation,

Note: LDAT_l = JXS(9) + $|L_C(l)| - 1$

4.3.10 LDLW & LDLWP Blocks

The LDLW Block and LDLW Block give the locators for the energy distribution for every reaction that produces secondary neutrons or secondary photons (respectively). The format of the LDLW Block (for secondary neutrons) and LDLW Block (for secondary photons) is given in Table 22. The locators for the delayed neutron precursors (see Section 4.3.2) also use the same format. The format for the distribution arrays is given in Section 4.3.11.

¹ linear-linear interpolation

The LDLW Block exists if NXS(5) $\neq 0$ while the LDLWP Block exists if NXS(6) $\neq 0$. The starting index, LED, depends on what data is being read; the starting values and the number of locators, NMT, are given in Table 21.

| Block | LED | NMT |
|------------------|---------|--------|
| LDLW | JXS(10) | NXS(5) |
| LDLWP | JXS(18) | NXS(6) |
| delayed neutrons | JXS(26) | NXS(8) |

Table 21: LED and NMT values for the LDLW Block and LDLWP Block.

Table 22: LDLW Block.

| Location in XSS | Parameter | Description |
|-----------------|---------------------|--|
| LED | $LOCC_1$ | Location of energy distribution data for reaction MT_1 or group 1 (if delayed neutron) |
| $LED{+}1$ | $LOCC_2$ | Location of energy distribution data for reaction MT_2 or group 2 (if delayed neutron) |
| • • • | | Location of energy distribution data for reac- |
| LED+NMT-1 | LOCC _{NMT} | tion MT_{NMT} or group NMT (if delayed neutron) |

Note: The LOCC_i must be monotonically increasing.

All locators point to data *relative* to JED (see Section 4.3.11) in the XSS array. The MT values are given in the MTR Block for LDLW Block or MTRP Block for LDLWP Block.

4.3.11 DLW & DLWP Blocks

The DLW Block contains secondary energy distributions for all reactions producing secondary neutrons—expect for elastic scattering. The DLWP Block contains secondary energy distribution for all photon-producing reactions. Both the DLW Block and DLWP Block have the same format. The energy distributions are given starting with a locator, LOCC, which were given in the LDLW Block and LDLWP Block. The locators are relative to the JED parameter. The value for JED and NMT (the number of reactions) is dependent on whether it is the DLW Block or DLWP Block. These values are given in Table 23.

| Block | JED | NMT |
|------------------|---------|--------|
| DLW | JXS(11) | NXS(5) |
| DLWP | JXS(19) | NXS(6) |
| delayed neutrons | JXS(27) | NXS(8) |

Table 23: JED and NMT for the DLW Block and DLW Block.

Table 24: DLW Block.

| Location in XSS | Description |
|---|--|
| $\begin{array}{c} JED + LOCC_1\text{-}1 \\ JED + LOCC_2\text{-}1 \end{array}$ | Energy distribution array for reaction MT_1 Energy distribution array for reaction MT_2 |
| JED+LOCC _{NMT} -1 | Energy distribution array for reaction MT_NMT |

The *i*-th array has the form shown in

Table 25: Format for the secondary energy distribution..

| Location in XSS | Parameter | Description |
|---------------------------------------|-----------------------------|--|
| JED+LOCC _i -1 | LNW_1 | Location of next law. [†] |
| $JED{+}LOCC_i$ | LAW_1 | Name of this law |
| $JED{+}LOCC_i{+}1$ | $IDAT_1$ | Location of data for this law relative to $JED + LOCC_i - 1$ |
| $JED {+} LOCC_i {+} 2$ | N_R | Number of interpolation regions to define law applicability regime |
| $JED + LOCC_i + 3$ | $NBT(l), l = 1, \dots, N_R$ | ENDF interpolation parameters |
| $JED {+} LOCC_i {+} 3 {+} N_R$ | $INT(l), l = 1, \dots, N_R$ | ENDF interpolation scheme [‡] |
| $JED {+} LOCC_i {+} 3 {+} 2N_R$ | N_E | Number of energies |
| $JED{+}LOCC_i{+}4{+}2N_R$ | $E(l), l = 1, \dots, N_E$ | Tabulated energy points |
| $JED {+} LOCC_i {+} 4 {+} 2N_R + N_E$ | $P(l), l = 1, \dots, N_E$ | Probability of law validity* |
| $JED + IDAT_1 - 1$ | $LDAT(l), l, \dots, L$ | Law data for LAW_1 . |
| $JED + LNW_1 - 1$ | LNW_2 | Location of next law |
| $JED + LNW_1$ | LAW_2 | Name of this law |
| JED + LNW + 1 | $IDAT_2$ | Location of data for this law relative to JED |
| • • • | | |

[†] If $\mathsf{LNW}_i = 0$ then LAW_1 is used regardless of other circumstances.

The format for the law data depends on the law. The length, L, of the law data array, LDAT, is determined from parameters with LDAT. The various LDAT arrays and their formats are given in the following tables. Laws 2 (Table 27) and 4 (Table 28) are used to describe spectra of secondary photons from neutron collisions. All laws—except for Law 2—are used to describe the spectra of scattered neutrons.

In the following tables, we provide relative locations of data in the LDAT array rather

[‡] If $N_R = 0$, NBT and INT are omitted and linear-linear interpolation is assumed.

^{*} If the particle energy E < E(1), then P(E) = P(1). If $E > E(N_E)$, then $P(E) = P(N_E)$. If more than one law is given, then LAW₁ is used only if $\xi < P(E)$ where ξ is a random number between 0 and 1.

than the absolute locations in the XSS array. Table 25 defines the starting location of the LDAT array within the XSS array.

4.3.11.1 LAW=1—Tabular Equiprobable Energy Bins

Table 26: LAW=1 (From ENDF Law 1).

| Location | Parameter | Description |
|-----------------------|--|--|
| LDAT(1) | N_R | Number of interpolation regions between tables of E_{out} |
| LDAT(2) | $NBT(l), l = 1, \dots, N_R$ | ENDF interpolation parameters |
| $LDAT(2{+}N_R)$ | $INT(l), l = 1, \dots, N_R$ | ENDF interpolation scheme † |
| $LDAT(2{+}2N_R)$ | N_E | Number of incident energies tabulated |
| $LDAT(3{+}2N_R)$ | $E_{\mathrm{in}}(l), l=1,\ldots,N_E$ | List of incident energies for which E_{out} is tabulated |
| $LDAT(3{+}2N_R+N_E)$ | NET | Number of outgoing energies in each E_{out} table |
| $LDAT(4{+}2*N_R+N_E)$ | $E_{\mathrm{out}_1}(l), l=1,\ldots,NET$ | $E_{\mathrm{out}} \mathrm{tables}^{\ddagger}$ |
| | $E_{\mathrm{out}_2}(l), l = 1, \dots, NET$ | |
| | | |
| | $E_{\operatorname{out}_{N_E}}(l), l = 1, \dots, NET$ | |

[†] If $N_R = 0$, NBT and INT are omitted and linear-linear interpolation is assumed.

4.3.11.2 LAW=2—Discrete Photon Energy

Table 27: LAW=2—Discrete Photon Energy.

| Location | Parameter | Description |
|----------|-----------|--|
| LDAT(1) | LP | Indicator of whether the photon is a primary or non-primary photon |
| LDAT(2) | EG | Photon energy or binding energy |

Note: If LP=0 or LP=1, the photon energy is EG. If LP=2, the photon energy is

$$\mathsf{EG} + \left(\frac{\mathsf{AWR}}{\mathsf{AWR} + 1}\right) E_N$$

where AWR is the atomic weight ratio and E_N is the incident neutron energy.

[‡] E_{out} tables consist of NET boundaries of NET-1 equally likely energy intervals. Linear-linear interpolation is used between intervals.

4.3.11.3 LAW=3—Level Scattering

$$\mathsf{LDAT}(1) = \left(\frac{A+1}{A}\right)|Q| \tag{2}$$

$$LDAT(2) = \left(\frac{A}{A+1}\right)^2 \tag{3}$$

$$E_{\text{out.}}^{\text{CM}} = \text{LDAT}(2) * (E - \text{LDAT}(1)) \tag{4}$$

where

 $E_{\text{out}}^{\text{CM}} = \text{outgoing center-of-mass energy}$

E = incident energy

A = atomic weight ratio

Q = Q-value

The outgoing neutron energy in the laboratory system is:

$$E_{\text{out}}^{\text{LAB}} = E_{\text{out}}^{\text{CM}} + \left\{ E + 2\mu_{\text{CM}}(A+1)(E \cdot E_{\text{out}}^{\text{CM}})^{1/2} \right\} / (A+1)^2$$
 (5)

where $\mu_{\rm CM}$ is the cosine of the center-of-mass scattering angle

4.3.11.4 LAW=4—Continuous Tabular Distribution

Table 28: LAW=4 (From ENDF-6 LAW=1).

| Location | Parameter | Description |
|---|---|---|
| | N_R $NBT(l), l = 1, \dots, N_R$ $INT(l), l = 1, \dots, N_R$ | Interpolation scheme between tables of $E_{\rm out}$ ENDF interpolation parameters ENDF interpolation scheme [†] |
| $LDAT(2{+}2N_R)$ | N_E | Number of energies at which distributions are tabulated |
| $\begin{array}{c} LDAT(3{+}2N_R) \\ LDAT(3{+}2N_R + N_E) \end{array}$ | $E(l), l = 1, \dots, N_E$ $L(l), l = 1, \dots, N_E$ | Incident neutron energies Locations of distributions [‡] |

[†] If $N_R = 0$, NBT and INT are omitted and linear-linear interpolation is assumed.

The data associated with each incident neutron energy begins at the location L(l). The format for the data is given in Table 29, where for E(1) let $K=3+2N_R+2N_E$.

[‡] Relative to JXS(11) (neutron reactions), JXS(19) (photon-producing reactions), or JXS(27) (delayed neutrons).

Table 29: Secondary energy distribution for each incident energy in LAW=4...

| Location | Parameter | Description | |
|---|--|--------------------------------------|--|
| | Data : | for E(1) | |
| LDAT(K) | INTT' | Interpolation parameter | |
| $LDAT(K{+}1)$ | N_p | Number of points in the distribution | |
| $LDAT(K{+}2)$ | $E_{\mathrm{out}}(l), l = 1, \dots, N_p$ | outgoing energy grid | |
| $LDAT(K{+}2+N_p)$ | $PDF(l), l = 1, \dots, N_p$ | Probability Density Function | |
| $LDAT(K{+}2 + 2N_p)$ | $CDF(l), l = 1, \dots, N_p$ | Cumulative Density Function | |
| Data for $E(2)$ —same format for $E(1)$ | | | |
| Data for $\mathbf{E}(\mathbf{N_E})$ —same format for $E(1)$ | | | |

The first element in the data is INTT' or the interpolation parameter, which is a combination of two other parameters:

- 1. the number of discrete photon lines, N_D , and
- 2. the interpolation scheme for the subsequent data, INTT.

INTT has two valid values:

INTT=1 histogram distribution, and

INTT=2 linear-linear distribution.

If the value of INTT' > 10, then

$$INTT' = 10N_D + INTT$$

where INTT is the interpolation scheme and the first N_D values of N_p points describe discrete photon lines. The remaining $(N_p - N_D)$ values describe a continuous distribution. In this way, the distribution may be discrete, continuous, or a discrete distribution superimposed upon a continuous background.

4.3.11.5 LAW=5—General Evaporation Spectrum

Table 30: LAW=5 (From ENDF-6, MF=5, LF=5).

| Location | Parameter | Description |
|--|---|--|
| $\begin{array}{c} LDAT(1) \\ LDAT(2) \\ LDAT(2{+}N_R) \end{array}$ | $\begin{aligned} N_R \\ NBT(l), l &= 1, \dots, N_R \\ INT(l), l &= 1, \dots, N_R \end{aligned}$ | Interpolation scheme between T 's |
| $\frac{LDAT(2{+}2N_R)}{LDAT(3{+}2N_R)}$ | $N_E \\ E(l), l = 1, \dots, N_E$ | Number of incident energies tabulated Incident energy table |

Table 30: LAW=5 (From ENDF-6, MF=5, LF=5) (continued)

| Location | Parameter | Description |
|-------------------------|-----------------------------|--|
| $LDAT(3{+}2N_R+N_E)$ | $\theta(l), l=1,\ldots,N_E$ | Effective temperature tabulated on incident energies |
| $LDAT(3{+}2N_R + 2N_E)$ | NET | Number of X 's tabulated |
| $LDAT(4{+}2N_R + 2N_E)$ | $X(l), l=1,\ldots,NET$ | Equiprobable bins |

$$E_{\text{out}} = X(\xi)\theta(E) \tag{6}$$

where:

 $X(\xi)$ is a randomly sampled table of X's;

 $\theta(E)$ is the effective temperature tabulated on incident energy; and

E is the incident energy.

4.3.11.6 LAW=7—Simple Maxwellian Fission Spectrum

Table 31: LAW=7 (From ENDF-6, MF=5, LF=7).

| Location | Parameter | Description |
|--|---|--|
| $\begin{array}{c} LDAT(1) \\ LDAT(2) \\ LDAT(2{+}N_R) \end{array}$ | N_R NBT $(l), l = 1, \dots, N_R$ INT $(l), l = 1, \dots, N_R$ | Interpolation scheme between T 's |
| $\frac{LDAT(2{+}2N_R)}{LDAT(3{+}2N_R)}$ | $N_E \\ E(l), l = 1, \dots, N_E$ | Number of incident energies tabulated Incident energy table |
| $LDAT(3{+}2N_R+N_E)$ | $\theta(l), l = 1, \dots, N_E$ | Effective temperature tabulated on incident energies |
| $LDAT(3 + 2N_R + 2N_E)$ | U | Restriction energy |

The outgoing energy, E_{out} , can be calculated as

$$f(E \to E_{\rm out}) = \frac{\sqrt{E_{\rm out}}}{I} e^{-E_{\rm out}/\theta(E)}$$
 (7)

where:

I is the normalization constant

$$I = \theta^{3/2} \frac{\sqrt{\pi}}{2} \operatorname{erf} \left(\sqrt{(E-U)/\theta} \right) - \sqrt{(E-U)/\theta} \ e^{-(E-U)/\theta}, \tag{8}$$

 θ is tabulated as a function of incident energy, E; and

U is a constant introduced to define the proper upper limit for the final particle energy such that $0 \le E_{\text{out}} \le (E - U)$.

4.3.11.7 LAW=9—Evaporation Spectrum

Table 32: LAW=9 (From ENDF-6, MF=5, LF=9).

| Location | Parameter | Description |
|---|---|---|
| $\begin{array}{c} LDAT(1) \\ LDAT(2) \\ LDAT(2{+}N_R) \end{array}$ | N_R NBT $(l), l = 1, \dots, N_R$ INT $(l), l = 1, \dots, N_R$ | Interpolation scheme between T 's |
| $\begin{array}{c} LDAT(2{+}2N_R) \\ LDAT(3{+}2N_R) \end{array}$ | N_E $E(l), l = 1, \dots, N_E$ | Number of incident energies tabulated Incident energy table |
| $\begin{aligned} LDAT(3 + 2N_R + N_E) \\ LDAT(3 + 2N_R + 2N_E) \end{aligned}$ | $\theta(l), l = 1, \dots, N_E$ U | Effective temperature tabulated on incident energies Restriction energy |

The outgoing energy, $E_{\rm out}$, can be calculated as

$$f(E \to E_{\text{out}}) = \frac{\sqrt{E_{\text{out}}}}{I} e^{-E_{\text{out}}/\theta(E)}$$
(9)

where:

I is the normalization constant

$$I = \theta^2 \left[1 - e^{-(E-U)/\theta} \left(1 + \frac{E-U}{\theta} \right) \right], \tag{10}$$

 θ is tabulated as a function of incident energy, E; and

U is a constant introduced to define the proper upper limit for the final particle energy such that $0 \le E_{\text{out}} \le (E - U)$.

Note: Equation 9 is the same as Equation 7; just the definitions of I in Equation 8 and Equation 10 are different.

4.3.11.8 LAW=11—Energy Dependent Watt Spectrum

Table 33: LAW=11 (From ENDF-6, MF=5, LF=11).

| Location | Parameter | Description |
|--|---|---|
| $\begin{array}{c} LDAT(1) \\ LDAT(2) \\ LDAT(2{+}N_{R_a}) \end{array}$ | N_{R_a} $NBT_a(l), l = 1, \dots, N_{R_a}$ $INT_a(l), l = 1, \dots, N_{R_a}$ | Interpolation scheme between a 's |
| $LDAT(2{+}2N_{R_a})$ | N_{E_a} | Number of incident energies tabulated for $a(E_{\rm in})$ table |

Table 33: LAW=11 (From ENDF-6, MF=5, LF=11) (continued)

| Location | Parameter | Description |
|--|---|---|
| $\frac{LDAT(3{+}2N_{R_a})}{LDAT(3{+}2N_{R_a}+N_{E_a})}$ | $E_a(l), l = 1, \dots, N_{E_a}$ $a(l), l = 1, \dots, N_{E_a}$ | Incident energy table Tabulated a 's |
| $\begin{split} & \text{let L} = 3 + 2\left(N_{R_a} + N_{E_a}\right) \\ & \text{LDAT(L)} \\ & \text{LDAT(L+1)} \\ & \text{LDAT(L+1} + N_{R_b}) \end{split}$ | N_{R_b} $NBT_b(l), l = 1, \dots, N_{R_b}$ $INT_b(l), l = 1, \dots, N_{R_b}$ | Interpolation scheme between b 's |
| $LDAT(L{+}1{+}2N_{R_b})$ | N_{E_b} | Number of incident energies tabulated for $b(E_{in})$ table |
| $\begin{array}{l} LDAT(L+2+2N_{R_b}) \\ LDAT(L+2+2N_{R_b}+N_{E_b}) \\ LDAT(L+2+2N_{R_b}+2N_{E_b} \end{array}$ | $E_b(l), l = 1, \dots, N_{E_b}$ $b(l), l = 1, \dots, N_{E_b}$ U | Incident energy table Tabulated b's Rejection energy |

The outgoing energy, $E_{\rm out}$, can be calculated as

$$f(E \to E_{\rm out}) = \frac{e^{-E_{\rm out}/a}}{I} \sinh\left(\sqrt{bE_{\rm out}}\right)$$
 (11)

where:

I is the normalization constant

$$I = \frac{1}{2} \sqrt{\frac{\pi a^3 b}{4}} e^{(ab/4)} \left[\operatorname{erf} \left(\sqrt{\frac{E - U}{a}} - \sqrt{\frac{ab}{4}} \right) + \operatorname{erf} \left(\sqrt{\frac{E - U}{a}} + \sqrt{\frac{ab}{4}} \right) \right] - a e^{-(E - U)/a} \sinh \sqrt{b(E - U)}; \quad (12)$$

a and b are tabulated energy-dependent parameters; and

U is a constant introduced to define the proper upper limit for the final particle energy such that $0 \le E_{\text{out}} \le (E - U)$.

4.3.11.9 LAW=22—Tabular Linear Functions of Incident Energy Out

Table 34: LAW=22 (From UK Law 2).

| Location | Parameter | Description |
|--|---|---------------------------------------|
| $\begin{array}{c} LDAT(1) \\ LDAT(2) \\ LDAT(2{+}N_R) \end{array}$ | N_R $NBT(l), l = 1, \dots, N_R$ $INT(l), l = 1, \dots, N_R$ | Interpolation parameters |
| $LDAT(2{+}2N_R)$ | N_E | Number of incident energies tabulated |

Table 34: LAW=22 (From UK Law 2) (continued)

| Location | Parameter | Description |
|---|---|--|
| $\frac{LDAT(3{+}2N_R)}{LDAT(3{+}2N_R+N_E)}$ | $E_{\mathrm{in}}(l), l = 1, \dots, N_E$ $LOCE(l), l = 1, \dots, N_E$ | Tabulated incident energies for E_{out} tables Locators of E_{out} tables |
| Data for $E_{\rm in}(1)$ Let L LDAT(L) LDAT(L+1) LDAT(L+1+NF ₁) LDAT(L+1+2NF ₁) | $= 3 + 2N_R + 2N_E$: NF_1 $P_{1k}, k = 1, \dots, NF_1$ $T_{1k}, k = 1, \dots, NF_1$ $C_{1k}, k = 1, \dots, NF_1$ | |
| Data for $E_{\rm in}(2)$: | | |

Tables of P_{ik}, C_{ik} , and T_{ik} are given at a number of incident energies, E_{in} . If

$$E_{\rm in}(l) \le E < E_{\rm in}(l+1) \tag{13}$$

then the secondary neutron energy is:

$$E_{\text{out}} = C_{ik} \left(E - T_{ik} \right), \tag{14}$$

where k is chosen according to

$$\sum_{j=1}^{k} P_{ij} < \xi \le \sum_{k=1}^{k+1} P_{ij} \tag{15}$$

for a given random number, $\xi \in [0, 1)$.

4.3.11.10 LAW=24

Table 35: LAW=24 (From UK Law 6).

| Location | Parameter | Description |
|--|---|--|
| $\begin{array}{c} LDAT(1) \\ LDAT(2) \\ LDAT(2{+}N_R) \end{array}$ | N_R $NBT(l), l = 1, \dots, N_R$ $INT(l), l = 1, \dots, N_R$ | Interpolation scheme between T 's |
| $\overline{LDAT(2{+}2N_R)}$ | N_E | Number of incident energies tabulated |
| $LDAT(3{+}2N_R)$ | $E_{\mathrm{in}}(l), l=1,\ldots,N_E$ | List of incident energies for which T is tabulated |
| $LDAT(3{+}2N_R+N_E)$ | NET | Number of outgoing values in each table |

Table 35: LAW=24 (From UK Law 6) (continued)

| Location | Parameter | Description |
|---------------------------------|--|--|
| $\overline{LDAT(4{+}2N_R+N_E)}$ | $T_1(l), l=1,\ldots, NET$ $T_2(l), l=1,\ldots, NET$ \ldots $T_{N_E}(l), l=1,\ldots, NET$ | Tables have NET boundaries with NET-1 equally likely intervals. Linear-linear interpolation is used between intervals. |

The outgoing energy, $E_{\rm out}$ can be calculated as:

$$E_{\text{out}} = T_k(l) * E \tag{16}$$

where:

 $T_k(l)$ is sampled from the tables and

E is the incident energy.

4.3.11.11 LAW=44—Kalbach-87 Formalism

Table 36: LAW=44 (From ENDF-6 MF=6 LAW=1, LANG=2).

| Location | Parameter | Description |
|---|---|--|
| | N_R $NBT(l), l = 1, \dots, N_R$ $INT(l), l = 1, \dots, N_R$ | Interpolation scheme between tables of E_{out} ENDF interpolation parameters ENDF interpolation scheme [†] |
| $LDAT(2+2N_R)$ | N_E | Number of energies at which distributions are tabulated |
| $\begin{array}{c} LDAT(3{+}2N_R) \\ LDAT(3{+}2N_R + N_E) \end{array}$ | $E(l), l = 1, \dots, N_E$ $L(l), l = 1, \dots, N_E$ | Incident neutron energies Locations of distributions [‡] |

[†] If $N_R = 0$, NBT and INT are omitted and linear-linear interpolation is assumed.

The data associated with each incident neutron energy begins at the location L(l). The format for the data is given in Table 37, where for E(1) let $K=3+2N_R+2N_E$.

Table 37: Secondary energy distribution for each incident energy in LAW=44.

| Location | Parameter | Description | | | | |
|--------------------------|-----------|-------------|------|---|--|--|
| Data for $\mathbf{E}(1)$ | | | | | | |
| | | | 0 1: | 7 | | |

[‡] Relative to JXS(11) (neutron reactions), JXS(19) (photon-producing reactions), or JXS(27) (delayed neutrons).

Table 37: Secondary energy distribution for each incident energy in LAW=44 (continued)

| Location | Parameter | Description | |
|--|--|--------------------------------------|--|
| LDAT(K) | INTT' | Interpolation parameter | |
| $LDAT(K{+}1)$ | N_p | Number of points in the distribution | |
| $LDAT(K{+}2)$ | $E_{\mathrm{out}}(l), l = 1, \dots, N_p$ | outgoing energy grid | |
| $LDAT(K{+}2+N_p)$ | $PDF(l), l = 1, \dots, N_p$ | Probability Density Function | |
| $LDAT(K{+}2 + 2N_p)$ | $CDF(l), l = 1, \dots, N_p$ | Cumulative Density Function | |
| $LDAT(K{+}2 + 3N_p)$ | $R(l), l = 1, \dots, N_p$ | Precompound fraction r | |
| $LDAT(K{+}2 + 4N_p)$ | $A(l), l = 1, \dots, N_p$ | Angular distribution slope value a | |
| Data for $\mathbf{E}(2)$ —same format for $E(1)$ | | | |
| ••• | | | |

Data for $E(N_E)$ —same format for E(1)

The first element in the data is INTT' or the interpolation parameter, which is a combination of two other parameters:

- 1. the number of discrete photon lines, N_D , and
- 2. the interpolation scheme for the subsequent data, INTT.

INTT has two valid values:

INTT=1 histogram distribution, and

INTT=2 linear-linear distribution.

If the value of INTT' > 10, then

$$INTT' = 10N_D + INTT$$

where INTT is the interpolation scheme and the first N_D values of N_p points describe discrete photon lines. The remaining $(N_p - N_D)$ values describe a continuous distribution. In this way, the distribution may be discrete, continuous, or a discrete distribution superimposed upon a continuous background.

The angular distributions for neutrons are then sampled from:

$$p(\mu, E_{\rm in}, E_{\rm out}) = \frac{1}{2} \frac{a}{\sinh(a)} \left[\cosh(a\mu) + r \sinh(a\mu) \right]. \tag{17}$$

4.3.11.12 LAW=61—Like LAW=44, but tabular angular distribution instead of Kalbach-87

Table 38: LAW=61.

| Location | Parameter | Description |
|----------|-----------|---------------------------------|
| LDAT(1) | N_R | Number of interpolation regions |

Table 38: LAW=61 (continued)

| Location | Parameter | Description |
|---|--|---|
| $\frac{LDAT(2)}{LDAT(2{+}N_R)}$ | $\begin{aligned} NBT(l), l &= 1, \dots, N_R \\ INT(l), l &= 1, \dots, N_R \end{aligned}$ | ENDF interpolation parameters ENDF interpolation scheme [†] |
| $LDAT(2{+}2N_R)$ | N_E | Number of energies at which distributions are tabulated |
| $\begin{array}{c} LDAT(3{+}2N_R) \\ LDAT(3{+}2N_R + N_E) \end{array}$ | $E(l), l = 1, \dots, N_E$ $L(l), l = 1, \dots, N_E$ | Incident neutron energies Locations of distributions ‡ |

[†] If $N_R = 0$, NBT and INT are omitted and linear-linear interpolation is assumed.

The data associated with each incident neutron energy begins at the location L(l). The format for the data is given in Table 39, where for E(1) let $K=3+2N_R+2N_E$.

Table 39: Secondary energy distribution for each incident energy in LAW=61.

| Location | Parameter | Description | | |
|--|--|--|--|--|
| Data for E(1) | | | | |
| LDAT(K) | INTT' | Interpolation parameter | | |
| $LDAT(K{+}1)$ | N_p | Number of points in the distribution | | |
| $LDAT(K{+}2)$ | $E_{\mathrm{out}}(l), l = 1, \dots, N_p$ | outgoing energy grid | | |
| $LDAT(K{+}2+N_p)$ | $PDF(l), l = 1, \dots, N_p$ | Probability Density Function | | |
| $LDAT(K{+}2 + 2N_p)$ | $CDF(l), l = 1, \dots, N_p$ | Cumulative Density Function | | |
| $LDAT(K{+}2+3N_p)$ | $LC(l), l = 1, \dots, N_p$ | Location of tables associated with incident energies $E(l)$. See Table 40 | | |
| Data for E(2) —same format for $E(1)$ | | | | |
| \mathbf{Data} for $\mathbf{E}(\mathbf{N_E})$ —same format for $E(1)$ | | | | |

If the value of INTT' > 10, then

$$INTT' = 10N_D + INTT$$

where INTT is the interpolation scheme and the first N_D values of N_p points describe discrete photon lines. The remaining $(N_p - N_D)$ values describe a continuous distribution. In this way, the distribution may be discrete, continuous, or a discrete distribution superimposed upon a continuous background.

The J-th array for the tabular angular distribution has the form shown in Table 40. For the angular distribution, the locators L are relative to JXS(11) for neutron reactions

[‡] Relative to JXS(11) (neutron reactions), JXS(19) (photon-producing reactions), or JXS(27) (delayed neutrons).

or JXS(19) for photon-producing reactions. Thus,

$$\begin{split} \mathsf{L} &= \mathsf{JXS}(\mathsf{11}) + |\mathsf{LC}(J)| - 1 \text{ (for neutron reactions)}, \\ \mathsf{L} &= \mathsf{JXS}(\mathsf{19}) + |\mathsf{LC}(J)| - 1 \text{ (for photon-producing reactions)}. \end{split}$$

Table 40: Angular distribution for LAW=61.

| Location | Parameter | Description |
|----------------------|--|--------------------------------------|
| LDAT(L+1) | JJ | Interpolation flag |
| $LDAT(L{+}2)$ | N_P | Number of points in the distribution |
| $LDAT(L{+}3)$ | $CS_{\mathrm{out}}(j), j=1,\ldots,N_P$ | Cosine scattering angular grid |
| $LDAT(L{+}3{+}N_P)$ | $PDF(j), j = 1, \dots, N_P$ | Probability density function |
| $LDAT(L{+}3{+}2N_P)$ | $CDF(j), j = 1, \dots, N_P$ | Cumulative density function |

4.3.11.13 LAW=66—N-body phase space distribution

Table 41: LAW=66 (From ENDF-6 MF=6 LAW=6).

| Location | Parameter | Description |
|----------|-----------|--|
| LDAT(1) | | Number of bodies in the phase space |
| LDAT(2) | A_P | Total mass ratio for the NPSX particles. |

The outgoing energy is

$$E_{\text{out}} = T(\xi)E_i^{\text{max}} \tag{18}$$

where

$$E_i^{\text{max}} = \frac{A_p - 1}{A_p} \left(\frac{A}{A+1} E_{\text{in}} + Q \right) \tag{19}$$

and $T(\xi)$ is sampled from:

$$P_i(\mu, E_{\rm in}, T) = C_n \sqrt{T} \left(E_i^{\rm max} - T \right)^{3n/2 - 4}$$
 (20)

4.3.11.14 LAW=67—Laboratory Angle-Energy Law

Table 42: LAW=67 (From ENDF-6 MF=6 LAW=7).

| Location | Parameter | Description |
|---|---|--|
| $\begin{array}{c} LDAT(1) \\ LDAT(2) \\ LDAT(2{+}N_R) \end{array}$ | N_R NBT $(l), l = 1, \dots, N_R$ INT $(l), l = 1, \dots, N_R$ | Number of interpolation regions ENDF interpolation parameters ENDF interpolation scheme [†] |
| $LDAT(2{+}2N_R)$ | N_E | Number of energies at which distributions are tabulated |
| $\begin{array}{l} LDAT(3{+}2N_R) \\ LDAT(3{+}2N_R + N_E) \end{array}$ | $E(l), l = 1, \dots, N_E$ $L(l), l = 1, \dots, N_E$ | Incident neutron energies Locations of distributions ‡ |

[†] If $N_R = 0$, NBT and INT are omitted and linear-linear interpolation is assumed.

The data associated with each distribution begins at location L(l). The format for the data is given in Table 43, where for E(1) let $K = 3 + 2N_R + 2N_e$.

Table 43: Angular distribution for LAW=67.

| Location | Parameter | Description |
|----------------------|-----------------------------|---|
| LDAT(K) LDAT(K+1) | INTMU NMU | Interpolation scheme [†] Number of secondary cosines |
| $LDAT(K{+}2)$ | $XMU(l), l = 1, \dots, NMU$ | Secondary cosines |
| $LDAT(K{+}2{+}NMU)$ | $LMU(l), l = 1, \dots, NMU$ | Locations of data for each secondary cosine. See Table 44 |

[†] INTMU=1 histogram distribution, INTMU=2 linear-linear distribution.

The format for the secondary energy distribution (for each cosine bin, XMU) is given in Table 44. For the energy distribution, the locators, LMU, are relative to JXS(11) or JXS(19). Thus,

$$L_l = JXS(11) + LMU(l)$$
 (for neutron reactions),
 $L_l = JXS(19) + LMU(l)$ (for photon-producing reactions).

[‡] Relative to JXS(11) (neutron reactions), JXS(19) (photon-producing reactions), or JXS(27) (delayed neutrons).

Table 44: Secondary energy distribution for each cosine bin in LAW=67.

| Location | Parameter | Description |
|-------------------------|--|---|
| $LDAT(L_l)$ | INTEP | Interpolation parameter between secondary energies [†] |
| $LDAT(L_l+1)$ | NPEP | Number of secondary energies |
| $LDAT(L_l+2)$ | $E_P(l), l=1,\ldots,NPEP$ | Secondary energy grid |
| $LDAT(L_l + 2 + NPEP)$ | $\mathrm{PDF}(l), l = 1, \dots, NPEP$ | Probability density function |
| $LDAT(L_l + 2 + 2NPEP)$ | $\mathrm{CDF}(l), l = 1, \ldots, NPEP$ | Cumulative density function |

[†] INTEP=1 histogram distribution,

INTEP=2 linear-linear distribution.

4.3.11.15 Energy-Dependent Neutron Yields

There are additional numbers to be found for neutrons in the DLW Block and DLWP Block. For those reactions with entries in the TYR Block that are greater than 100 in absolute value, there must be neutron yields, Y(E) provided as a function of neutron energy. The neutron yields are handled similarly to the average number of neutrons per fission, $\nu(E)$ that is given for the fission reactions. These yields are a part of the coupled energy-angle distributions given in File 6 of ENDF-6 data.

The *i*-th array has the form given in Table 45, where $KY = JED + |TY_i| - 101$.

Table 45: Energy-Dependent Neutron Yields.

| Location | Parameter | Description |
|--------------------------|-----------------------------|--|
| KY | N_R | Number of interpolation regions |
| $KY{+}1$ | $NBT(l), l = 1, \dots, N_R$ | ENDF interpolation parameters |
| $KY{+}1{+}N_R$ | $INT(l), l = 1, \dots, N_R$ | ENDF interpolation scheme [†] |
| $KY{+}1{+}2N_R$ | N_E | Number of energies |
| $KY{+}2{+}2N_R$ | $E(l), l = 1, \dots, N_E$ | Tabular energy points |
| $KY {+} 2 {+} N_R + N_E$ | $Y(l), l = 1, \dots, N_E$ | Corresponding energy-dependent yields |

[†] If $N_R = 0$, NBT and INT are omitted and linear-linear interpolation is assumed.

4.3.12 GPD Block

The GPD Block contains the *total* photon production cross section, tabulated on the energy grid given in the ESZ Block, the size of which is given by NXS(3). The GPD Block only exists if $JXS(12) \neq 0$.

There are 30 groups for the incident neutron energies, the boundaries of which are shown in Table 46. For each incident neutron energy group, the outgoing photon energies are discretized into 20 equiprobable energy groups, thus creating a 30×20 matrix. The

outgoing energies are given in the GPD Block as shown in Table 47. Note that this matrix is only used for older tables that do not provide expanded photon production data.

Table 46: Discrete neutron energy boundaries.

| Group # | Upper Boundary (MeV) | Group # | Upper Boundary (MeV) |
|---------|-------------------------|---------|-------------------------|
| 1 | 1.39×10^{-10} | 16 | 0.184 |
| 2 | 1.52×10^{-7} | 17 | 0.303 |
| 3 | 4.14×10^{-7} | 18 | 0.500 |
| 4 | 1.13×10^{-6} | 19 | 0.823 |
| 5 | 3.06×10^{-6} | 20 | 1.353 |
| 6 | 8.32×10^{-6} | 21 | 1.738 |
| 7 | 2.26×10^{-5} | 22 | 2.232 |
| 8 | 6.14×10^{-5} | 23 | 2.865 |
| 9 | 1.67×10^{-4} | 24 | 3.68 |
| 10 | 4.54×10^{-4} | 25 | 6.07 |
| 11 | 1.235×10^{-3} | 26 | 7.79 |
| 12 | 3.35×10^{-3} | 27 | 10.0 |
| 13 | 9.23×10^{-3} | 28 | 12.0 |
| 14 | 2.48×10^{-2} | 29 | 13.5 |
| 15 | 6.76×10^{-2} | 30 | 15.0 |

The format of this Block is given in Table 47. The XSS array index at the start of the GPD Block, S_{GPD} =JXS(12).

Table 47: GPD Block.

| Location in XSS | Parameter | Description |
|---|--------------------------------------|--|
| S_{GPD} | $\sigma_{\gamma}(l), l=1,\ldots,NES$ | Total photon production cross section |
| $S_{GPD} + NES$ | $E_1(K), K = 1, 20$ | 20 equiprobable outgoing photon energies for incident neutron $E < E_N(2)$ |
| $S_{GPD} + NES + 20$ | $E_2(K), K = 1, 20$ | 20 equiprobable outgoing photon energies for incident neutron $E_N(2) \le E < E_N(3)$ |
| $S_{GPD} + NES + (\mathrm{i}\text{-}1)^*20$ | $E_i(K), K = 1, 20$ | 20 equiprobable outgoing photon energies for incident neutron $E_N(i) \leq E < E_N(i+1)$ |
| $S_{\sf GPD} + {\sf NES} + (30-1)^*20$ | $E_N(K), K = 1, 20$ | 20 equiprobable outgoing photon energies for incident neutron $E \geq E_N(30)$ |

4.3.13 SIGP Block

The SIGP Block contains the photon production cross section data. The format of the SIGP Block is given in Table 48. The cross section data begins at the index specified by the locator, LOCA_i, given in the LSIG Block (see Section 4.3.6). The MTs are defined in the MTRP Block (see Section 4.3.3). All indices to the XSS array are *relative* to JXS(15).

Table 48: SIGP Block.

| Location in XSS | Parameter | Description |
|---|-----------------------|--|
| $\begin{array}{c} JXS(15) \!+\! LOCA_1\text{-}1 \\ JXS(15) \!+\! LOCA_2\text{-}1 \end{array}$ | $MFTYPE_1$ $MFTYPE_2$ | Cross section array for reaction MT_1 Cross section array for reaction MT_2 |
| JXS(15)+LOCA _{NMT} -1 | MFTYPE _{NMT} | Cross section array for reaction MT_NMT |

Note: The number of photon production cross section arrays NMT=NXS(6).

The format of the i-th cross section array has two possible forms depending on the first number in the array, MFTYPE.

1. If MFTYPE=12 or MFTYPE=16, yield data taken from ENDF File 12 or 6, respectively (see Table 49). With this format, the photon production cross section can be constructed using Equation 21:

$$\sigma_{\gamma,i}(E) = Y(E) * \sigma_{\mathsf{MTMULT}}(E). \tag{21}$$

2. If MFTYPE=13, cross section data from ENDF File 13 (see Table 50).

Table 49: Photon production array if MFTYPE=12 or 16.

| Location in XSS | Parameter | Description |
|--|-----------------------------|--|
| $JXS(15) + LOCA_{i} - 1$ | MFTYPE | 12 or 16 |
| ${\tt JXS(15)}\!+\!{\tt LOCA}_i$ | MTMULT | Neutron MT whose cross section should multiply the yield |
| $\mathtt{JXS(15)}\!+\!LOCA_{\mathrm{i}}\!+\!1$ | N_R | Number of interpolation regions |
| $\mathtt{JXS}(\mathtt{15})\!+\!LOCA_{\mathrm{i}}\!+\!2$ | $NBT(l), l = 1, \dots, N_R$ | ENDF interpolation parameters [†] |
| $\begin{array}{c} {\tt JXS(15)\!+\!LOCA_i+2}\\ +N_R\end{array}$ | $INT(l), l = 1, \dots, N_R$ | ENDF interpolation scheme |
| $\begin{array}{c} \texttt{JXS(15)} \!+\! \texttt{LOCA}_{\text{i}} \!+\! 2 \\ +\! 2*N_R \end{array}$ | N_E | Number of energies at which the yield is tabulate |
| $\begin{array}{c} \texttt{JXS(15)} \!+\! \texttt{LOCA}_{\text{i}} \!+\! 3 \\ +\! 2*N_R \end{array}$ | $E(l), l = 1, \dots, N_E$ | Energies |
| $\begin{array}{c} \texttt{JXS(15)} \!+\! \texttt{LOCA}_{\text{i}} \!+\! 3 \\ +2*N_R+N_E \end{array}$ | $Y(l), l = 1, \ldots, N_E$ | Yields |

Continued on next page

Table 49: Photon production array if MFTYPE=12 or 16 (continued)

| Location in XSS | Parameter | Description | |
|-----------------|-----------|-------------|--|
|-----------------|-----------|-------------|--|

[†] If $N_R = 0$, NBT and INT are omitted and linear-linear interpolation is used.

Table 50: Photon production cross section array if MFTYPE=13.

| Location in XSS | Parameter | Description |
|---|--|--|
| $ \overline{ JXS(15) \! + \! LOCA_{i} \! - \! 1 } $ | MFTYPE | 13 |
| ${\tt JXS(15)\!+\!LOCA}_i$ | IE | Energy grid index |
| $\mathtt{JXS(15)}\!+\!LOCA_{i}\!+\!1$ | N_E | Number of consecutive entries |
| $\mathtt{JXS(15)}\!+\!LOCA_{i}\!+\!2$ | $\sigma_{\gamma,i}[E(K)],$ $K = IE, \dots, IE + N_E - 1$ | Photon production cross sections for reaction MT_i |

Note: The MT_i s are defined in the MTRP Block.

4.3.14 LANDP Block

The LANDP Block gives locator information for angular distribution arrays for photon production reactions and exists if NXS(6) \neq 0. All locators (LOCB) in the LANDP Block are relative to JXS(17); that is, the angular distribution arrays begin at JXS(17) + LOCB_i. The number of photon-producing reactions is NMT=NXS(6). The LOCB_i must be monotonically increasing. The MTs are defined in the MTRP Block (see Section 4.3.3). The format of the LANDP Block is given in Table 51.

Table 51: LANDP Block.

| Location in XSS | Parameter | Description |
|------------------------|---------------------|---|
| JXS(16) | $LOCB_1 {=} 1$ | Location of angular distribution data for reaction MT_1 |
| $\mathtt{JXS}(16) + 1$ | $LOCB_2$ | Location of angular distribution data for reaction MT_2 |
| JXS(16)+NMT-1 | LOCB _{NMT} | Location of angular distribution data for reaction MT_NMT |

Note: The LOCB_i must be monotonically increasing. The format for the angular distribution of the i-th reaction is given in Table 53.

4.3.15 ANDP Block

The ANDP Block contains angular distribution data for all photon-producing reactions and exists if NXS(6) \neq 0. The format of the ANDP Block is given in Table 52; the format of each angular distribution array is given in Table 53. The angular distribution data begins at the index specified by the locator, LOCB, from the LANDP Block; if LOCB_i = 0, there are no angular distribution data given for reaction i and isotropic scattering is assumed in the Lab system.

Table 52: ANDP.

| Location in XSS | Description |
|--|--|
| $\begin{array}{c} \mathtt{JXS(17)} \! + \! \mathtt{LOCB}_1 \text{-} 1 \\ \mathtt{JXS(17)} \! + \! \mathtt{LOCB}_2 \end{array}$ | Angular distribution array for reaction MT_1 Angular distribution array for reaction MT_2 |
| JXS(17)+LOCB _{NMT} -1 | Angular distribution array for reaction MT_{NMT} |

Note: NMT=NXS(6) is the number of photon-producing reactions.

Table 53: Angular distribution array for the *i*-th photon-producing reaction.

| Location in XSS | Parameter | Description |
|---|--------------------------------|--|
| JXS(17)+LOCB _i -1 | N_E | Number of energies at which angular distributions are tabulated. |
| $\mathtt{JXS(17)}\!+\!LOCB_i$ | $E(l), l = 1, N_E$ | Energy grid |
| $\texttt{JXS(17)} \!+\! \texttt{LOCB}_{\text{i}} \!+\! N_E$ | $L_C(l), l = 1, \dots, N_E$ | Location of tables associated with $E(l)^{\dagger}$ |
| ${	t JXS(17)} + L_C(1) - 1$ | $P_1(K), K = 1, \dots, 33$ | 32 equiprobable cosine bins for scattering at energy $E(1)$ |
| ${	t JXS(17)} + L_C(2) - 1$ | $P_2(K), K = 1, \dots, 33$ | 32 equiprobable cosine bins for scattering at energy $E(2)$ |
| | | |
| $\mathtt{JXS(17)} + L_C(N_E) - 1$ | $P_{N_E}(K), K = 1, \dots, 33$ | 32 equiprobable cosine bins for scattering at energy $E(N_E)$ |

[†] All values of $L_C(l)$ are relative to JXS(17). If $L_C(l) = 0$, no table is given for energy E(l) and scattering is assumed to be isotropic in the Lab system.

4.3.16 YP Block

The YP Block contains a list of MT identifiers of neutron cross sections that are used as yield multipliers in Equation 21 to calculate the photon production cross sections and are referenced by the MTMULT parameter in Table 49. The YP Block exists if NXS(6) \neq 0. The format of the YP Block is given in Table 54.

Table 54: YP Block.

| Location in XSS | Parameter | Description |
|----------------------|---|---|
| JXS(20) JXS(20)+1 | $\begin{array}{l} NYP \\ MTY(l), l = 1, \dots, NYP \end{array}$ | Number of neutron MTs to follow Neutron MTs. |

4.3.17 FIS Block

The FIS Block contains the total fission cross section. The FIS Block exists if JXS(21) $\neq 0$, but is generally not provided; the total fission cross section is redundant as the total fission cross section is the summation of first-, second-, third-, and fourth-chance fission (MT=19, 20, 21, and 38);

$$\sigma_{f,t}(E) = \sigma_{(n,f)} + \sigma_{(n,nf)} + \sigma_{(n,2nf)} + \sigma_{(n,3nf)}.$$
 (22)

The format of the FIS Block is given in Table 55.

Table 55: FIS Block.

| Location in XSS | Parameter | Description |
|-----------------------------------|---|--|
| JXS(21) JXS(21)+1 JXS(21)+2 | IE N_E $\sigma_f[E(l)], K = IE, \ldots, IE + N_E - 1$ | Energy grid index Number of consecutive entries Total fission cross sections |

Note: The energy E(l) is given in the ESZ Block.

4.3.18 UNR Block

The UNR Block contains the unresolved resonance range probability tables. It exists if $JXS(23) \neq 0$ and begins at location JXS(23) in XSS. The UNR Block has several flags that have special meaning:

ILF The ILF flag is the inelastic competition flag.

 $\mathsf{ILF} < 0$ The inelastic cross section is zero within the entire unresolved energy range.

 $\mathsf{ILF} > 0$ The value of ILF is a special MT number whose tabulation is the sum of the inelastic levels.

ILF = 0 The sum of the contribution of the inelastic reactions will be made using a balance relationship involving the smooth cross sections.

An exception to this scheme is typically made when there is only one inelastic level within the unresolved energy range, because the flag can then just be set to its MT number and the special tabulation is not needed.

IOA The IOA is the other absorption flag for determining the contribution of "other

absorptions" (no neutron out or destruction reactions).

- **IOA** < 0 The "other absorption" cross section is zero within the entire unresolved resonance range.
- $\mathsf{IOA} > 0$ The value of IOA is a special MT number whose tabulation is the sum of the "other absorption" reactions.
- IOA = 0 The sum of the contribution of the "other absorption" reactions will be made using a balanced relationship involving the smooth cross sections.

An exception to this scheme is typically made when there is only one "other absorption" reaction within the unresolved energy range, because the flag can then just be set to its MT number and the special tabulation is not needed.

IFF The IFF is the factors flag.

 $\mathsf{IFF} = 0$ The tabulations in the probability tables are cross sections.

IFF = 1 The tabulations in the probability tables are factors that must be multiplied by the corresponding "smooth" cross sections to obtain the actual cross sections.

The format of the UNR Block is given in Table 57. The P(i,j,k) values, where

- $-i = 1, \ldots, N,$
- $j = 1, \ldots, 6,$
- $-k=1,\ldots,M,$

are what make up the probability tables. The argument j has special meaning depending on its value as shown in Table 56.

Table 56: Possible values for the j argument.

| j | Description |
|---|-----------------------------------|
| 1 | cumulative probability |
| 2 | total cross section/factor |
| 3 | elastic cross section/factor |
| 4 | fission cross section/factor |
| 5 | (n,γ) cross section/factor |
| 6 | neutron heating number/factor |

Table 57: UNR Block.

| Location in XSS | Parameter | Description |
|------------------------|-----------|---|
| JXS(23) | N | Number of incident energies where there is a probability table. |
| ${\tt JXS(23)}{+}1$ | M | Length of probability table. |
| $\mathtt{JXS}(23) + 2$ | INT | Interpolation parameter between tables. [†] |
| ${\tt JXS(23)}{+}3$ | ILF | Inelastic competition flag. |
| JXS(23)+4 | IOA | Other absorption flag. |

Continued on next page

Table 57: UNR Block (continued)

| Location in XSS | Parameter | Description |
|---|--|--|
| $\begin{array}{c} {\tt JXS(23)+5} \\ {\tt JXS(23)+6} \\ {\tt JXS(23)+6+} N \end{array}$ | $\begin{aligned} & IFF \\ & E(i), i = 1, \dots, N \\ & P(i, j, k) \end{aligned}$ | Factors flag. Incident energies. Probability tables. |

 $^{^{\}dagger}$ 2 linear-linear interpolation,

The ordering of the probability table entries, P(i, j, k) is given in Table 58, which begins at PTABLE = JXS(23) + 6 + N.

Table 58: Order of probability table elements P(i, j, k).

| Location in XSS | Parameter | Description |
|-----------------------|-------------------------|--|
| PTABLE | CDF_1 | Cumulative probabilities for energy $i = 1$ |
| $PTABLE {+} M$ | $\sigma_{t,1}$ | Total cross section/factors for energy $i = 1$ |
| $PTABLE {+} 2M$ | $\sigma_{s,1}$ | Elastic cross section/factors for energy $i = 1$ |
| $PTABLE{+}3M$ | $\sigma_{f,1}$ | Fission cross section/factors for energy $i = 1$ |
| $PTABLE{+}4M$ | $\sigma_{(n,\gamma),1}$ | (n, γ) cross section/factors for energy $i = 1$ |
| $PTABLE{+}5M$ | H_1 | Heating number/factors for energy $i = 1$ |
| PTABLE $+(i-1)*6M$ | CDF_i | Cumulative probabilities for energy i |
| PTABLE $+(N-1)*6M+5M$ | H_N | Heating numbers/factors for energy $i=N$ |

[†] The cumulative probabilities are monotonically increasing from an implied (but not included) lower value of zero to the upper value of P(i, 1, k = M) = 1.0.

⁵ log-log interpolation

| Element | Name | Description |
|---------|------|--|
| 1 | | Length of second block of data (XSS array) |
| 2 | ZA | 1000 * Z + A |
| 3 | | |
| 4 | NTR | Number of reactions |
| | | |
| 16 | | |

Table 59: NXS array element definitions for neutron dosimetry ACE Table.

| Element | Name | Location Description |
|---------|------|---------------------------------|
| 1 | LONE | First word of table |
| 2 | | |
| 3 | MTR | MT array |
| | | |
| 6 | LSIG | Table of cross section locators |
| 7 | SIGD | Cross sections |
| | | |
| 22 | END | Last word of this table |
| | | |
| 32 | | |

Table 60: JXS array element definitions for neutron dosimtry ACE Table.

5 Neutron Dosimetry

- 5.1 NXS Array
- 5.2 JXS Array
- 6 Thermal Scattering $S(\alpha, \beta)$
- 6.1 NXS Array
- 6.2 JXS Array

7 Continuous-Energy Photon

7.1 NXS Array

| Element | Name | Description |
|---------|-------|--|
| 1 | _ | Length of second block of data (XSS array) |
| 2 | IDPNI | Inelastic scattering mode |
| 3 | NIL | Inelastic dimensioning parameter |
| 4 | NIEB | Number of inelastic exiting energies |
| 5 | IDPNC | Elastic scattering mode |
| 6 | NCL | Elastic dimensioning parameter |
| 7 | IFENG | Secondary energy mode |
| | | |
| 16 | | |

Table 61: NXS array element definitions for thermal scattering ACE Table.

| Element | Name | Location Description |
|---------|------|--------------------------------------|
| 1 | ITIE | Inelastic energy table |
| 2 | ITIX | Inelastic cross sections |
| 3 | ITXE | Inelastic energy/angle distributions |
| 4 | ITCX | Elastic cross sections |
| 5 | ITCA | Elastic angular distributions |
| | | |
| 32 | | |

Table 62: JXS array element definitions for thermal scattering ACE Table.

Table 63: ${\sf NXS}$ array element definitions for ${\sf NXS}$ ACE Table.

| Element | Name | Description |
|---------|------|---|
| 1 2 | | Length of second block of data (XSS array) Atomic number |
| 3 | NES | Number of energies |
| 4 | NFLC | Length of the flourescence data divided by 4 |
| 5 | NSH | Number of electron shells |
| | | |
| 16 | | |

7.2 JXS Array

Table 64: JXS array element definitions for JXS ACE Table.

| Element | Name | Location Description |
|---------|-------|---|
| 1 | ESZG | Energy table |
| 2 | JINC | Incoherent form factors |
| 3 | JCOH | Coherent form factors |
| 4 | JFLO | Fluorescence data |
| 5 | LHNM | Heating numbers |
| 6 | LNEPS | Number of electrons per shell |
| 7 | LBEPS | Binding energy per shell |
| 8 | LPIPS | Probability of interaction per shell |
| 9 | LSWD | Array of offsets to the shell-wise data |
| 10 | SWD | Shell-wise data in PDF and CDF form |
| | | |
| 32 | _ | |

References

- [1] Jeremy Lloyd Conlin, Forrest B. Brown, A. C. Kahler, M. Beth Lee, D. Kent Parsons, and Morgan C. White. "Updating the Format of ACE Data Tables". In: *Transactions of the American Nuclear Society*. Ed. by Julie Rule. Vol. 107. LA-UR-12-22033. San Diego, California: American Nuclear Society, Nov. 2012, pp. 631–633 (cit. on pp. 10, 16).
- [2] A. Trkov, M. Herman, and D. A. Brown. ENDF-6 Formats Manual: Data Formats and Procedures for the Evaluated Nuclear Data Files. Tech. rep. BNL-90365-2009 Rev.2. National Nuclear Data Center, Brookhaven National Laboratory, Dec. 2011 (cit. on pp. 18, 20).