

FFT: sin和cos的FFT

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$\delta(x)$ 的定义式

本文用到的FFT对为:

$$F(\omega) = \int_{-\infty}^{+\infty} f(t) e^{-i\omega t} dt$$

对 $f(t) = \sin \omega_0 t$,

有: $F(\omega) = \int_{-\infty}^{+\infty} \sin \omega_0 t e^{-i\omega t} dt$

$$= \int_{-\infty}^{+\infty} \frac{e^{i\omega_0 t} - e^{-i\omega_0 t}}{2j} e^{-i\omega t} dt$$

$$= \frac{1}{2j} \int_{-\infty}^{+\infty} e^{it(\omega_0 - \omega)} dt + \frac{-1}{2j} \int_{-\infty}^{+\infty} e^{it(-\omega_0 - \omega)} dt$$

用到了 δ 的定义式
Dirac function

$$= -\pi j [\delta(\omega_0 - \omega) - \delta(\omega_0 + \omega)]$$

振幅谱: $C(\omega) = \sqrt{R[\frac{F(\omega)}{\pi}]^2 + I[\frac{F(\omega)}{\pi}]^2}$, 当 ω 取 $\omega_0, -\omega_0$ 时均为1.

相位谱: $\varphi(\omega) = \arctan \frac{-I[F(\omega)]}{R[F(\omega)]}$. 当取 ω_0 时: $\arctan(\infty) = \frac{\pi}{2}$, $1 \cdot \cos(\omega_0 t - \frac{\pi}{2})$

~~当取 $-\omega_0$ 时: $\arctan(-\infty) = -\frac{\pi}{2}$~~

$$= \sin \omega_0 t$$

注意: 这里 ω 只能取正.

还厚了!!
原.

类似有 $\cos \omega_0 t$ 的, 则

"相位谱里无负频"和代码
周期延拓的FFT不同.

$$\begin{cases} \cos \omega_0 t \xleftrightarrow{FT} \pi [\delta(\omega - \omega_0) + \delta(\omega + \omega_0)] \\ \sin \omega_0 t \xleftrightarrow{FT} -j\pi [\delta(\omega - \omega_0) - \delta(\omega + \omega_0)] \end{cases}$$

δ(x) 定义

数学上有若干种:

1. $\delta(\vec{r}) = \frac{1}{-4\pi} \nabla^2 \frac{1}{r}$, $\delta(\vec{r}) = \delta(x)\delta(y)\delta(z)$

2. $\delta(x) = \lim_{\epsilon \rightarrow 0} \frac{1}{\pi} \frac{\epsilon}{x^2 + \epsilon^2}$, 适合数值计算

3. $\delta(x) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ikx} dk$, 可用于求 $\sin \omega t$, $e^{i\omega t}$ 的 FFT.

1. $r = \sqrt{x^2 + y^2 + z^2}$

$\nabla^2(\frac{1}{r}) = \frac{1}{r} \frac{\partial}{\partial r}$ $r \neq 0$ 时:

$\nabla^2(\frac{1}{r}) = \frac{1}{r} \frac{\partial^2}{\partial r^2} (r \cdot \frac{1}{r}) = 0$

$\int \nabla^2(\frac{1}{r}) dV = \int \nabla \cdot (\nabla \frac{1}{r}) dV = \oint (\nabla \frac{1}{r}) \cdot d\vec{S} = -\oint \frac{1}{r^2} \frac{\vec{r}}{r} \cdot d\vec{S} = -4\pi.$

得: $\nabla^2(\frac{1}{r}) = -4\pi \delta(\vec{r})$

2. $\int_{-\infty}^{+\infty} \lim_{\epsilon \rightarrow 0} \frac{1}{\pi} \frac{\epsilon}{x^2 + \epsilon^2} dx = \frac{1}{\pi} \int_{-\infty}^{+\infty} \lim_{\epsilon \rightarrow 0} \frac{1}{(\frac{x}{\epsilon})^2 + 1} d(\frac{x}{\epsilon}) = \frac{1}{\pi} \lim_{\epsilon \rightarrow 0} \left[\tan^{-1} \frac{x}{\epsilon} \right]_{-\infty}^{+\infty} = 1$

$= \frac{1}{\pi} \lim_{\epsilon \rightarrow 0} \tan^{-1} \frac{x}{\epsilon} \Big|_{-\infty}^{+\infty} = 1$

3. $\frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ikx} dk = \frac{1}{2\pi} \lim_{\eta \rightarrow 0^+} \left(\int_0^{+\infty} e^{ik(x+i\eta)} dk + \int_{-\infty}^0 e^{ik(x-i\eta)} dk \right)$

$= \lim_{\eta \rightarrow 0^+} \frac{1}{\pi} \frac{\eta}{x^2 + \eta^2} = \delta(x)$

↑
上文已证.

↑
证明合理, 见背面.

$$\int_{0^-}^{0^+} e^{ik(x+i\eta)} dk = 0. \text{ 显然成立.}$$

$$\lim_{\eta \rightarrow 0^+} \lim_{M \rightarrow \infty} \int_0^M e^{ik(x+i\eta)} dk = \text{无穷大. } x+i\eta \rightarrow x.$$

$$\text{若 } M\eta \rightarrow 0, \text{ 即 } \eta \text{ 足够小, 则 } \int_0^M e^{ik(x+i\eta)} dk \rightarrow \int_0^M e^{ikx} dk$$

但这证明也有不严谨处, 比如, $M\eta \rightarrow 0$ 为什么成立, 也可能趋于 ∞ .

$$x+i- = 2i \cdot \frac{1}{2} \cdot \frac{1}{x-i} = 2i \left(\frac{1}{x} + \frac{1}{x^2} + \dots \right) = \frac{2i}{x} + \frac{2i}{x^2} + \dots$$

$$\frac{1}{x+i-} = \frac{1}{x-i} + \frac{1}{x^2} + \dots$$

$$\left[\frac{1}{x+i-} - \frac{1}{x-i} \right] \frac{1}{x} = \left[\frac{1}{x^2} + \frac{1}{x^3} + \dots \right] \frac{1}{x} = \frac{1}{x^3} + \frac{1}{x^4} + \dots$$

$$1 = \int_{-\infty}^{\infty} \frac{1}{x^2} dx = \dots$$

$$\left(\frac{1}{x^2} - \frac{1}{x^2} \right) \frac{1}{x} = \frac{1}{x^3} - \frac{1}{x^3} = 0$$

$$\frac{1}{x^2} = \frac{1}{x^2} + \frac{1}{x^2} = \frac{2}{x^2}$$

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