

A (kind of) practical review on global stability approches

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Informal Workshop "D&D", 10 juillet 2018

Linear global stability : basic principles, and a few useful tricks

Qu'es aquò??

Three ideas to really speed up your linear stability computations!

StabFem : a software that may save your life !

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Demonstration for the wake of a cylinder

List of test-cases currently available

Nonlinear global stability approaches : status and future

Weakly nonlinear approach

Self-consistent approach

Harmonic-balance approach

Qu'es aquò ??

Instability problems are ubiquitous in fluid mechanics

Numerical resolution of such problems resorts to a specific class of numerical methods, which replace time-stepping of the full equations by assumptions about temporal dependence (modal expansion, amplitude equations, etc...)

These methods are complementary to direct numerical resolution methods (i.e. time-stepping).

We refer to *global stability* when the geometry requires resolution in 2 (or 3) spatial dimensions.

(as opposed to *local stability* which takes advantage of invariance directions (parallel flow, etc...) to bring the problem to spatially 1D).

Base flow

We look for a steady base-flow $(\mathbf{u}_b; p_b)$ satisfying the steady Navier-Stokes equations, i.e. $NS(\mathbf{u}_b, p_b) = 0$.

Suppose that we have a 'guess' for the base flow $[\mathbf{u}_b^g, p_b^g]$ which almost satisfies the equations. We look for a better approximation under the form

$$[\mathbf{u}_b, p_b] = [\mathbf{u}_b^g, p_b^g] + [\delta\mathbf{u}_b, \delta p_b] = 0. \quad (1)$$

Injecting into the Navier-Stokes equation lead to

$$NS(\mathbf{u}_b^g, p_b^g) + NSL_{\mathbf{u}_b^g}(\delta\mathbf{u}_b, \delta p_b)$$

Where NSL is the linearised Navier-Stokes operator.

\Rightarrow matricial problem with the form $A \cdot \delta X = Y$. The procedure of Newton iteration is to solve iteratively this set of equations up to convergence.

Linear stability

$$\mathbf{u} = \mathbf{u}_b + \epsilon \hat{\mathbf{u}} e^{\lambda t} \quad (2)$$

The eigenmodes is governed by the linear problem

$$\lambda \hat{\mathbf{u}} = NSL_{\mathbf{u}_b}(\hat{\mathbf{u}}, \hat{p})$$

After discretization we end up with an eigenvalue problem with the matricial form

$$\lambda B \hat{X} = A \hat{X} \quad (3)$$

Iterative method : single-mode shift-invert iteration

$$X^n = (A - \lambda_{shift} B)^{-1} B X^{n-1}$$

Generalization : Arnoldi

Adjoint problem

Define a scalar product :

$$\langle \phi_1, \phi_2 \rangle = \int_{\Omega} \overline{\phi_1} \cdot \phi_2 \, d\Omega$$

We can first define the *adjoint linearised Navier-Stokes operator* NSL^\dagger defined by the property :

$$\begin{aligned} \forall(\mathbf{u}, p; \mathbf{u}^\dagger, p^\dagger), \quad & \langle NSL_{\mathbf{U}}^\dagger(\mathbf{u}^\dagger, p^\dagger), \mathbf{u} \rangle + \langle \nabla \cdot \mathbf{u}^\dagger, p \rangle \\ &= \langle \mathbf{u}^\dagger, NSL_{\mathbf{U}}(\mathbf{u}, p) \rangle + \langle p^\dagger, \nabla \cdot \mathbf{u} \rangle. \end{aligned} \quad (4)$$

We can then define the adjoint eigenmodes as the solutions to the eigenvalue problem

$$\forall(\mathbf{u}, p), \quad \lambda^\dagger \langle \hat{\mathbf{u}}^\dagger, \mathbf{u} \rangle = \langle NSL_{\mathbf{U}}^\dagger(\hat{\mathbf{v}}, \hat{p}^\dagger), \mathbf{u} \rangle + \langle \nabla \cdot \hat{\mathbf{u}}^\dagger, p \rangle \quad (5)$$

Matricial form :

$$\overline{\lambda}^\dagger B \hat{X}^\dagger = A^T \hat{X}^\dagger. \quad (6)$$

Adjoint mode and structural sensitivity

Significance of the adjoint mode :
(optimal perturbation)

The adjoint eigenmode also allows us to introduce the so-called *structural sensitivity tensor* that is defined as

$$\mathbf{S}(\mathbf{x}) = \frac{||\hat{\mathbf{u}}^\dagger|| \ ||\hat{\mathbf{u}}||}{\langle \hat{\mathbf{u}}^\dagger, \hat{\mathbf{u}} \rangle}, \quad (7)$$

which has become popular in the recent years.

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 - ▶ Compressible case : to handle non-reflective boundary conditions (example : cylinder wake ; with Javier Sierra)

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- ▶ 😞 Syntax may be touchy and debugging sometimes awkward...

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Computation chain : freefem solvers / shell scripts / postprocessing with tecplot/gnuplot...

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- ▶ => Necessity of a set of "drivers" in a high-level language to monitor computations and draw the results in "command-line" or "script" mode.
- ▶ Philosophy (objective) : one work (one paper) = 1 unique program to generate all results and produce all figures. (cf. Basilisk...)

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- ▶ Maintained on Github

<https://github.com/erbafeidavid/StabFem>

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- Etage 1 : Solveurs FreeFem++ "briques de base".
Un solveur par "classe de problèmes" (2D incompressible, 2D compressible, Axisymétrique incompressible...) et par "type de calcul" (calcul d'un champ de base, stabilité linéaire, ...)

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Remarques : les programmes FreeFem doivent pouvoir être utilisés directement en dehors du driver StabFem, notamment pour faciliter le développement/débuggage...)

Les contributeurs "utilisateurs" (ex. étudiant M1/M2) ne travaillent qu'à l'étage 3 et ne devraient travailler que sur ces 3 fichiers.

Les contributeurs "développeurs" travaillent aux étages inférieurs.

Format d'échange des données

- ▶ Format de fichier d'échange ".ff2m", généré par FreeFem++ et relu par Matlab

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1  ### Data generated by Freefem++ ;  
2  Temperature  
3  Format :  
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Ligne 4 : "TypeField1 NameField1 TypeField2 NameField2... "
TypeField can be "real" (scalar data), "real.N" (vectorial data), "P1" (data associated to mesh), ...

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- Le fichier est lu et importé sous forme d'une *structure matlab*
- Illustration : cas "EXAMPLE_Lshape"

First step : Generation of a mesh and "guess" base flow

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bf=SF_Init('Mesh_Cylinder.edp', [-40 80 40]);
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What the `SF_Init` driver does :

- ▶ Runs the relevant FreeFem++ program `Mesh_Cylinder.edp` with the corresponding parameters (here size of the domain),

This program generates the following output files : `mesh.msh` (mesh data),

`mesh.ff2m` (mesh information), `SF_Init.ff2m` (auxiliary information),

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- ▶ Reads all the output files,
- ▶ Returns a matlab "structure" object containing all the data needed for post-processing and subsequent usage.

Computation of a Base flow : principle

We look for a steady base-flow $(\mathbf{u}_b; p_b)$ satisfying the steady Navier-Stokes equations, i.e. $NS(\mathbf{u}_b, p_b) = 0$.

Suppose that we have a 'guess' for the base flow $[\mathbf{u}_b^g, p_b^g]$ which almost satisfies the equations. We look for a better approximation under the form

$$[\mathbf{u}_b, p_b] = [\mathbf{u}_b^g, p_b^g] + [\delta\mathbf{u}_b, \delta p_b] = 0. \quad (8)$$

Injecting into the Navier-Stokes equation lead to

$$NS(\mathbf{u}_b^g, p_b^g) + NSL_{\mathbf{u}_b^g}(\delta\mathbf{u}_b, \delta p_b)$$

Where NSL is the linearised Navier-Stokes operator.

\Rightarrow matricial problem with the form $A \cdot \delta X = Y$. The procedure of Newton iteration is to solve iteratively this set of equations up to convergence.

Computation of a Base flow : implementation

```
bf=SF_BaseFlow(bf,'Re',10);
```

What the `SF_Init` driver does :

- ▶ Copies the previous base flow into file `BaseFlow_guess.txt` which will be read by Freefem++,

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Mesh adaptation

Linear stability

$$\mathbf{u} = \mathbf{u}_b + \epsilon \hat{\mathbf{u}} e^{\lambda t} \quad (9)$$

The eigenmodes is governed by the linear problem

$$\lambda \hat{\mathbf{u}} = NSL_{\mathbf{u}_b}(\hat{\mathbf{u}}, \hat{p})$$

After discretization we end up with an eigenvalue problem with the matricial form

$$\lambda B \hat{X} = A \hat{X} \quad (10)$$

Iterative method : single-mode shift-invert iteration

$$X^n = (A - \lambda_{shift} B)^{-1} B X^{n-1}$$

Generalization : Arnoldi

Eigenvalue computation : implementation

```
SF_Stability(bf,'shift',0.04) + 0.74i,'nev',1,' type ' , 'D' ) ;
```

What the `SF_Stability` driver does :

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Eigenvalue computation : implementation

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- ▶ Copies the base flow into file `BaseFlow.txt` which will be needed by Freefem++,
- ▶ Runs the FreeFem++ solver (here `Stab_2D.edp`) with the corresponding parameters (shift, number of eigenvalues, direct eigenmode),

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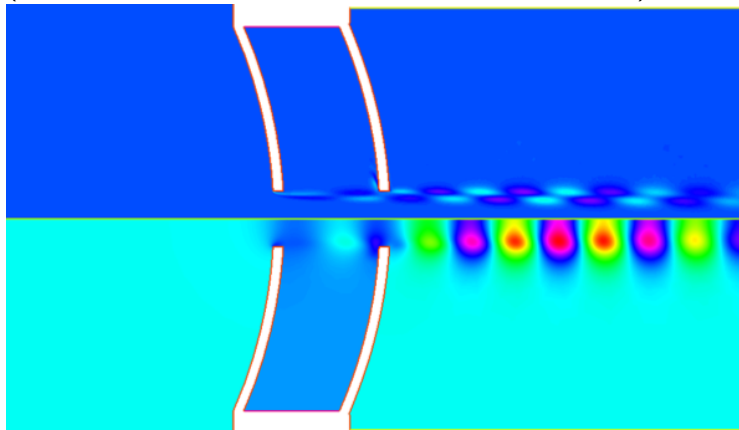
StabFem : list of test-cases currently available (or under development...)

- ▶ CYLINDER -> 2D incompressible, objet fixe
- ▶ CYLINDER_VIV -> 2D incompressible, objet mobile
(Stage Diogo Ferrera-Sabino)
- ▶ IMPACTINGJET -> 2D incompressible, 3D stability.
(with David LoJacono)
- ▶ CYLINDER_Compressible -> 2D compressible
(Javier Serra, Vincenzo Citro...)
- ▶ BIRDCALL -> 2D-axisymmetric, incompressible or "augmented incompressible"
(with R. Longobardi, V. Citro....)
- ▶ POROUS_DISK -> 2D-axisymmetric, with porous object
(stage Adrien Rouvière)
- ▶ LiquidBridges -> 2D axi, with deformable free surface
(stage Nabil Achour)
- ▶ ROTATING_POLYGONS
(with Jérôme Mougel...)
- ▶

=> Illustration dans le cas CYLINDER

Whistling jets : axisymmetric flow through a two-hole configuration

(with Raffaele Longobari, Vincenzo Citro & others...)



2D around a spring-mounted cylinder...

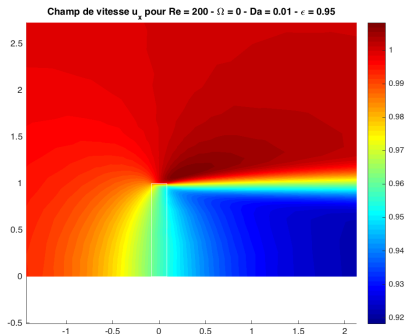
(with Diogo Ferreira Sabino & Olivier Marquet)

2D flow around a compressible cylinder...

(in fast progress with Javier Sierra...)

Flow around (and through) a porous (& rotating) disk...

(with Adrien Rouvière & D. Lo Jacono)



Liquid bridges...

Reference : Chireux et al., Phys. Fluids, 2015.

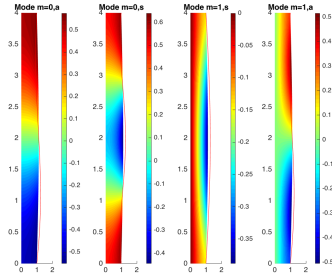
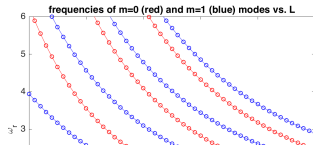
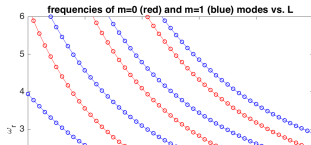


Figure – Oscillation modes of a liquid bridge of aspect ratio $L/R = 4$ and reduced volume $V = \dots$



rotating polygons...

Reference : Mougel et al., JFM 2018

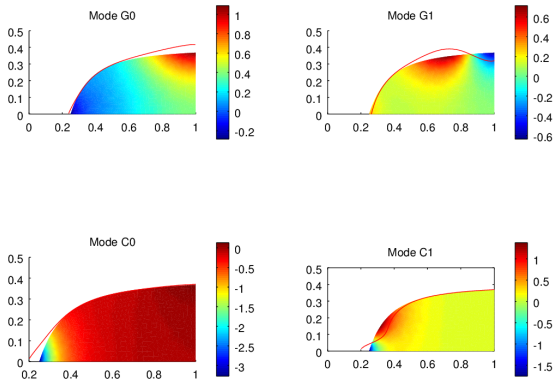
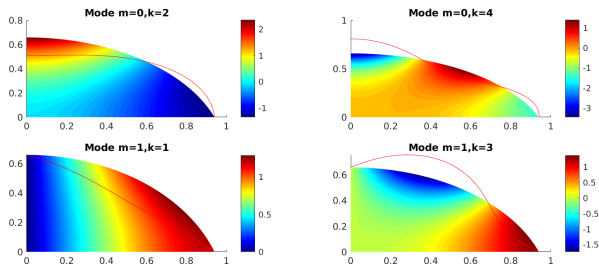


Figure – Oscillation modes of a potential vortex for $a = H/R = 0.3$ and $m = 3$ (figure 5, 6 of Mougel et al.).

Sessile drops...

With Nabil Achour (& Paul Bonnefis)



- Linear oscillations modes ("pined" or "fixed angle" conditions)
- Next steps : investigate contact-line dynamics with WNL methods (cf. Viola, Brun & Gallaire, 2018)

Linear global stability : basic principles, and a few useful tricks

Qu'es aquò ??

Three ideas to really speed up your linear stability computations !

StabFem : a software that may save your life !

Qu'es aquò ???

Demonstration for the wake of a cylinder

List of test-cases currently available

Nonlinear global stability approaches : status and future

Weakly nonlinear approach

Self-consistent approach

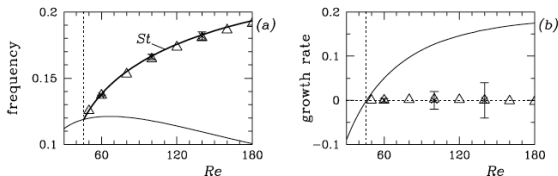
Harmonic-balance approach

Nonlinear global stability approaches : review

The linear stability approach approach is the right tool to predict instability threshold (Re_c) and the shedding frequency at threshold ($St_c = \omega_c/2\pi$).

But for $Re > Re_c$ it badly predicts the frequency of the limit cycle.

D. BARKLEY: CYLINDER MEAN FLOW



It has been remarked that stability analysis of the *mean flow* obtained by time-averaging the limit cycle gives better predictions (Barkley, Leontini,...).

=> Objective of nonlinear stability approaches : provide rational approach to describe the nonlinear oscillation cycle, and provide amplitude equations to describe the transients.

Weakly nonlinear approach (Sipp & Lebedev, 2007)

Starting point : weakly non-linear expansion, with multiple scale method.

$$\epsilon = \frac{1}{Re_c} - \frac{1}{Re}; \quad \tau = \epsilon^2 t$$

$$\begin{aligned} \mathbf{u} &= \mathbf{u}_{bc} + \epsilon \left[A_{wnl}(\tau) \hat{\mathbf{u}} e^{i\omega_c t} + c.c. \right] \\ &+ \epsilon^2 \left[\mathbf{u}_\epsilon + |A_{wnl}|^2 \mathbf{u}_{2,0} + \left(A_{wnl}^2 \mathbf{u}_{2,2} e^{2i\omega_c t} + c.c. \right) \right] + \mathcal{O}(\epsilon^3) \end{aligned} \quad (11)$$

Resolution at order 2 :

$$\mathcal{LN}S_{\mathbf{u}_{bc}}(\mathbf{u}_\epsilon) - 2\nabla \cdot \mathbf{D}(\mathbf{u}_{bc}) = 0, \quad (12)$$

$$\mathcal{LN}S_{\mathbf{u}_{bc}}(\mathbf{u}_{2,0}) = \mathcal{C}(\hat{\mathbf{u}}, \bar{\hat{\mathbf{u}}}), \quad (13)$$

$$\mathcal{LN}S_{\mathbf{u}_{bc}}(\mathbf{u}_{2,2}) - 2i\omega_c \mathbf{u}_{2,2} = \frac{1}{2} \mathcal{C}(\hat{\mathbf{u}}, \hat{\mathbf{u}}). \quad (14)$$

Compatibility conditions at order 3 :

$$\frac{\partial A_{wnl}}{\partial \tau} = \Lambda A_{wnl} - (\nu_0 + \nu_2) |A_{wnl}|^2 A_{wnl}, \quad (15)$$

$$\Lambda = - \frac{\langle \hat{\mathbf{u}}^\dagger, (\mathcal{C}(\mathbf{u}_\epsilon, \hat{\mathbf{u}}) + 2\nabla \cdot \mathbf{D}(\hat{\mathbf{u}})) \rangle}{\langle \hat{\mathbf{u}}^\dagger, \hat{\mathbf{u}} \rangle}, \quad (16)$$

$$\nu_0 = \frac{\langle \hat{\mathbf{u}}^\dagger, \mathcal{C}(\mathbf{u}_{20}, \hat{\mathbf{u}}) \rangle}{\langle \hat{\mathbf{u}}^\dagger, \hat{\mathbf{u}} \rangle}, \quad \nu_2 = \frac{\langle \hat{\mathbf{u}}^\dagger, \mathcal{C}(\mathbf{u}_{22}, \bar{\hat{\mathbf{u}}}) \rangle}{\langle \hat{\mathbf{u}}^\dagger, \hat{\mathbf{u}} \rangle}, \quad (17)$$

Self-Consistent approach (Mantic-Lugo, Arratia & Gallaire, 2014)

Starting point : Pseudo-eigenmode decomposition

$$\mathbf{u} = \mathbf{u}_m + A_{sc} \left[\tilde{\mathbf{u}}_1 e^{\sigma_{sc} t + i \omega_{sc} t} + \overline{\tilde{\mathbf{u}}_1} e^{\sigma_{sc} t - i \omega_{sc} t} \right], \quad \left(\|\tilde{\mathbf{u}}_1\| = 1/\sqrt{2} \right) \quad (18)$$

where A_{sc} is an amplitude parameter, and $\lambda_{sc} = \sigma_{sc} + i \omega_{sc}$ is a pseudo-eigenvalue which depends upon the parameter A_{sc} .

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=> SC-model equations

$$\mathcal{N}S(\mathbf{u}_m) - A_{sc}^2 \mathcal{C}(\tilde{\mathbf{u}}_1, \overline{\tilde{\mathbf{u}}_1}) = 0, \quad (19a)$$

$$(\sigma_{sc} + i \omega_{sc}) \tilde{\mathbf{u}}_1 = \mathcal{L} \mathcal{N}S_{\mathbf{u}_m}(\tilde{\mathbf{u}}_1). \quad (19b)$$

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=> Amplitude equation :

$$\frac{\partial A_{sc}}{\partial t} = \sigma_{sc}(A_{sc})A_{sc}$$

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Resolution method of (Mantic-Lugo 2014) : double iterative loop

- Inner loop :

Fix A_{sc} , iteratively solve (eigenvalue problem + calculation of mean flow) up to convergence for $(\sigma_{sc} + i\omega_{sc})$ as function of A_{sc} .

- Outer loop :

iterate over A_{sc} to reach $\sigma_{sc} = 0$

Self-Consistent approach : a direct resolution method

Let's forget about transients, and look directly for a description of the saturated cycle :

$$\mathbf{u} = \mathbf{u}_m + \mathbf{u}_{1,c} \cos(\omega t) + \mathbf{u}_{1,s} \sin(\omega t), \quad (20)$$

where $\mathbf{u}_{1,c}$ and $\mathbf{u}_{1,s}$ are two *real* fields
and ω is the (real) oscillation frequency of the limit cycle.

=> Equations :

$$\mathcal{NS}(\mathbf{u}_m) = \frac{\mathcal{C}(\mathbf{u}_{1,c}, \mathbf{u}_{1,c}) + \mathcal{C}(\mathbf{u}_{1,s}, \mathbf{u}_{1,s})}{4}, \quad (21a)$$

$$\omega \mathbf{u}_{1,s} = \mathcal{LN}_{\mathbf{u}_m}(\mathbf{u}_{1,c}), \quad (21b)$$

$$-\omega \mathbf{u}_{1,c} = \mathcal{LN}_{\mathbf{u}_m}(\mathbf{u}_{1,s}). \quad (21c)$$

We need an extra scalar equation to fix the phase of the cycle, e.g :

$$\Im\{F_y(\mathbf{u}_1)\} = 0. \quad (22)$$

=> Direct resolution with Newton method !

Remarks :

- We need a good guess : use the WNL to produce it !
- Computation can be optimized using preconditioning (Ask Olivier...)

Harmonic-Balance

We start with the following expansion :

$$\mathbf{u} = \mathbf{u}_m + \mathbf{u}_{1,c} \cos(\omega t) + \mathbf{u}_{1,s} \sin(\omega t) + \mathbf{u}_{2,c} \cos(2\omega t) + \mathbf{u}_{2,s} \sin(2\omega t), \quad (23)$$

arriving to a system of equations :

$$\mathcal{NS}(\mathbf{u}_m) = \frac{\mathcal{C}(\mathbf{u}_{1,c}, \mathbf{u}_{1,c}) + \mathcal{C}(\mathbf{u}_{1,s}, \mathbf{u}_{1,s}) + \mathcal{C}(\mathbf{u}_{2,c}, \mathbf{u}_{2,c}) + \mathcal{C}(\mathbf{u}_{2,s}, \mathbf{u}_{2,s})}{4}, \quad (24a)$$

$$\omega \mathbf{u}_{1,s} = \mathcal{L}_{\mathbf{u}_m}(\mathbf{u}_{1,c}) - \frac{1}{2} \left(\mathcal{C}(\mathbf{u}_{1,c}, \mathbf{u}_{2,c}) + \mathcal{C}(\mathbf{u}_{1,s}, \mathbf{u}_{2,s}) \right), \quad (24b)$$

$$-\omega \mathbf{u}_{1,c} = \mathcal{L}_{\mathbf{u}_m}(\mathbf{u}_{1,s}) - \frac{1}{2} \left(\mathcal{C}(\mathbf{u}_{1,c}, \mathbf{u}_{2,s}) - \mathcal{C}(\mathbf{u}_{1,s}, \mathbf{u}_{2,c}) \right), \quad (24c)$$

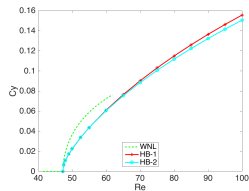
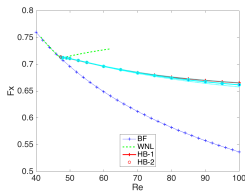
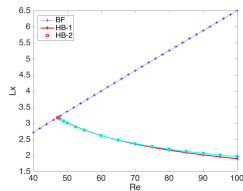
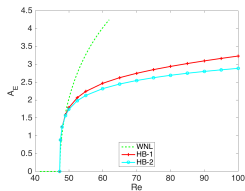
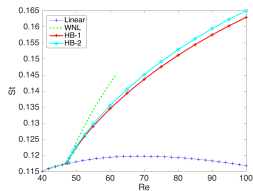
$$2\omega \mathbf{u}_{2,s} = \mathcal{L}_{\mathbf{u}_m}(\mathbf{u}_{2,c}) - \frac{1}{4} \left(\mathcal{C}(\mathbf{u}_{1,c}, \mathbf{u}_{1,c}) - \mathcal{C}(\mathbf{u}_{1,s}, \mathbf{u}_{1,s}) \right), \quad (24d)$$

$$-2\omega \mathbf{u}_{2,c} = \mathcal{L}_{\mathbf{u}_m}(\mathbf{u}_{2,s}) - \frac{1}{2} \mathcal{C}(\mathbf{u}_{1,s}, \mathbf{u}_{1,c}), \quad (24e)$$

$$\Im\{F_y(\mathbf{u}_1)\} = 0. \quad (25)$$

=> Direct Newton resolution again

Harmonic-Balance : results for the cylinder !



Conclusions

The future of StabFem

Recent progress

- ▶ Multi-platform objective : MacOS OK ; Unix OK ; Windows 10 currently 50 % compatible.
main issues with windows : cp = copy,...
- ▶ Plotting options : recent intergration of "pdeplot2dff" from Markus "chloros" in place of pdeplot/pdetools .
other solutions for plotting : tecplot converter, vtk converter, ...
- ▶ Compatibility with Octave : currently 50 % compatible.
Main issues with octave : importdata, plotting (now solved), inputParser (now solved).
- ▶ Translation in Python ??

Besoins

- ▶ Maintaining a fully opensource (Matlab-Octave or Python ?) and fully multiplatform version (windows).
- ▶ Managing a list of test cases (non-regression tests, etc...)
- ▶ Help simplifying/rationalizing the programmation style.
- ▶ Gestion of errors / debugging / "verbosity" ...
- ▶ Upgrading to 3D / parallel computation ? (currently not priority)
- ▶ Support with github (/ gitlab ?)
- ▶ Documentation
Automatic generation from comments in programs ?
Doxygen ??