Lattice discretisation of the Green-Schwarz superstring and AdS/CFT

Philipp Töpfer



Humboldt-Universität zu Berlin Emmy Noether Research Group Introduction — 2 | 13

Outline

Study the AdS/CFT duality numerically at hand of the cusp anomaly function from a string theory perspective

Introduction ______ 2 | 13

Outline

Study the AdS/CFT duality numerically at hand of the cusp anomaly function from a string theory perspective

- ▶ Motivation
- ► Cusp anomaly function
- ► AdS/CFT duality
- ► String theory framework
- Numerical approach

Motivation

- ▶ Good predictions on the real world via QCD
- ▶ Fundamental aim: particle masses as functions of parameters

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 - ▶ measure scaling dimensions of local operators and Wilson loops

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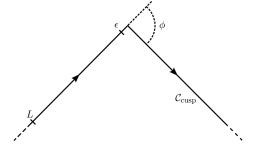
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- ► Twist two operators: scaling dimension $\Delta_S = 2 + S + \gamma_S$

Wilson loop with cusp

$$\langle \mathcal{W}[\mathcal{C}_{\mathrm{cusp}}] \rangle \sim e^{-\frac{f(g)}{2}|\phi| \ln \frac{L}{\epsilon}}$$

$$\gamma_S \simeq f(g) \ln S$$



 \Leftrightarrow

Wilson loop with cusp

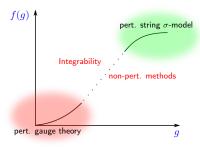
anomalous dimension γ_S

$$\langle \mathcal{W}[\mathcal{C}_{\text{cusp}}] \rangle \sim e^{-\frac{f(g)}{2}|\phi| \ln \frac{L}{\epsilon}}$$

$$\gamma_S \simeq f(g) \ln S$$

via cusp anomaly function

 \Leftrightarrow



Here
$$g\equiv rac{\sqrt{\lambda}}{4\pi}$$

$$\lambda=g_{\mathrm{YM}}^2N=rac{R^4}{lpha'^2}$$

AdS/CFT duality

conformal field theory (CFT)

↑ dynamically equivalent to

string theory on background containing Anti-de Sitter (AdS) space as a factor

Most symmetric setting similar to QCD:

$${\cal N}={f 4}$$
 SYM in 4d, $(g_{
m YM},N)$



† dynamically equivalent to

 \downarrow

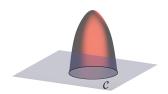
Type IIB superstring theory on $AdS_5 \times S^5$, (g_s, R)



${\cal N}=4$ SYM

Wilson loop

$$\langle \mathcal{W}[\mathcal{C}] \rangle = \frac{1}{N} \operatorname{Tr} \mathcal{P} e^{\oint (iA_{\mu}\dot{x}^{\mu} + \phi_{i}\dot{y}^{i}) ds}$$



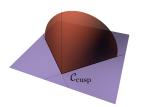
Type IIB strings

$$Z_{\text{string}}[\mathcal{C}] = \int \mathcal{D}X \mathcal{D}h \, e^{-S_{\text{string}}(X,h)} \sim e^{-A(\mathcal{C})}$$

Cusp anomaly of $\mathcal{N}=4$ SYM from string theory

vev of light-like cusped Wilson loop

$$\langle \mathcal{W}[\mathcal{C}_{\mathrm{cusp}}] \rangle \sim e^{\frac{f(g)}{2}|\phi| \ln \frac{L}{\epsilon}}$$



↑ AdS/CFT

$$Z_{\text{cusp}} = \int \mathcal{D}\delta X \, \mathcal{D}\delta \Psi \, e^{-S_{\text{IIB}}(X_{\text{cl}} + \delta X, \delta \Psi)} = e^{-\frac{f(g)}{2} \frac{V_2}{4}}$$

String partition function with cusp vacuum $(V_2 = \int \mathrm{d}t \mathrm{d}s)$

Strings on the lattice and AdS/CFT



Green-Schwarz string in AdS light-cone gauge

▶ Sigma-model in $AdS_5 \times S^5$ with RR flux $S = g \int d\tau \, d\sigma \, \left[G_{\mu\nu} \partial X^\mu \partial X^\nu + \overline{\Theta} \Gamma(D+F_5) \Theta \partial X + \ldots \right]$ [Metsaev Tseytlin 1998]

► Fix κ -symmetry and apply bosonic light-cone gauge \Rightarrow action at most quartic in complex Grassmann fields θ^i, η^i

Green-Schwarz string in null-cusp background

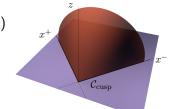
▶ In Poincaré patch

$$\mathrm{d}s^2_{AdS_5} = \tfrac{\mathrm{d}z^2 + \mathrm{d}x^+ \, \mathrm{d}x^- + \, \mathrm{d}x^* \, \mathrm{d}x}{z^2}, \quad x^\pm = x^3 \pm x^0, \quad x = x^1 + ix^2$$
 classical solution $\left(\tau, \sigma \in (0, \infty)\right)$: surface

$$z = \sqrt{\frac{\tau}{\sigma}}, \quad x^+ = \tau, \quad x^- = -\frac{1}{2\sigma}$$

bounded by a null cusp $(AdS_5 \text{ boundary at } 0 = z^2 = -2x^+x^-)$

• expand around classical solution + $(\tau, \sigma) \rightarrow (t, s) = (\ln \tau, \ln \sigma)$ $S_{\text{cusp}} = g \int dt ds \mathcal{L}_{\text{cusp}}$



Green-Schwarz string in null-cusp background

linearising quartic fermion contributions

$$\mathcal{L}_{\text{cusp}} = \left| \partial_t x + \frac{m}{2} x \right|^2 + \frac{1}{z^4} \left| \partial_s x - \frac{m}{2} x \right|^2 + \left(\partial_t z^M + \frac{m}{2} z^M \right)^2 + \frac{1}{z^4} \left(\partial_s z^M - \frac{m}{2} z^M \right)^2 + \phi^2 + (\phi_I)^2 + \Psi^T \mathcal{O}_F \Psi$$

- ▶ 8 bosonic coordinates: $x, x^*, z^M \ (M = 1, ..., 6), \ z = \sqrt{z_M z^M}$
- ▶ 17 auxiliary fields $\phi, \phi_I \ (I = 1, ..., 16)$
- ▶ 8 fermionic variables $\Psi \equiv (\theta^i, \theta_i, \eta^i, \eta_i)$, $\theta^i = \theta_i^{\dagger}$, $\eta^i = \eta_i^{\dagger}$, (i = 1, ..., 4)

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$$\mathcal{O}_{\mathrm{F}} = \begin{pmatrix} 0 & i\mathbb{1}_{4}\partial_{t} & -i\rho^{M}\left(\partial_{s} + \frac{m}{2}\right)\frac{z^{M}}{z^{3}} & 0 \\ i\mathbb{1}_{4}\partial_{t} & 0 & 0 & -i\rho^{\dagger}_{M}\left(\partial_{s} + \frac{m}{2}\right)\frac{z^{M}}{z^{3}} \\ i\frac{z^{M}}{z^{3}}\rho^{M}\left(\partial_{s} - \frac{m}{2}\right) & 0 & 2\frac{z^{M}}{z^{4}}\rho^{M}\left(\partial_{s}x - m\frac{x}{2}\right) & i\mathbb{1}_{4}\partial_{t} - A^{T} \\ 0 & i\frac{z^{M}}{z^{3}}\rho^{\dagger}_{M}\left(\partial_{s} - \frac{m}{2}\right) & i\mathbb{1}_{4}\partial_{t} + A & -2\frac{z^{M}}{z^{4}}\rho^{\dagger}_{M}\left(\partial_{s}x^{*} - m\frac{x}{2}^{*}\right) \end{pmatrix}$$

$$A = -\frac{\sqrt{6}}{z}\phi \mathbb{1}_4 + \frac{1}{z}\widetilde{\phi} + \frac{1}{z^3}\rho_N^*\widetilde{\phi}^T\rho^L z^N z^L + i\frac{z^N}{z^2}\rho^{MN}\partial_t z^M$$

Strings on the lattice and AdS/CFT



Numerical approach

requires calculating expectation values

$$Z_{\text{cusp}} = \int \mathcal{D}\delta X \mathcal{D}\delta \Psi \ e^{-S_{\text{cusp}}} = e^{-\frac{f(g)}{2}\frac{V_2}{4}}$$

Lattice simulations

$$\langle A \rangle = \frac{1}{Z} \int \mathcal{D}\phi \, A[\phi] e^{-S[\phi]}, \quad Z = \int \mathcal{D}\phi \, e^{-S[\phi]}$$

discretise the worldsheet with const. lattice spacing \boldsymbol{a}

$$\Lambda = \{(n_0, n_1) | n_\alpha = 0, \dots, (N_\alpha - 1)\}$$
 so that

$$\xi^{\alpha} = (\tau, \sigma) \equiv (an_0, an_1) \equiv an$$

▶ PI measure: discretise fields
$$\phi \to \phi(n)$$

$$\mathcal{D}\phi \to \prod_n \mathrm{d}\phi(n)$$

▶ Operators:
$$\partial_{\alpha}\phi(n) \rightarrow \frac{1}{a}[\phi(n+\widehat{\alpha}) - \phi(n)] \Rightarrow S \rightarrow S_{\text{disc}}$$

$$\Rightarrow Z_{\rm disc} = \int \prod d\phi(n) e^{-S_{\rm disc}[\phi]} \sim$$

multidimensional integral treatable via MC techniques

Strings on the lattice and AdS/CFT

