

Problem 1. *Griffiths 3.44*

First note that $\rho = 0$ everywhere except for $(0, 0, z)$ for $z \in [-a, a]$. The formula for potential is given by

$$V(r) = \frac{1}{4\pi\epsilon_0} \sum_{n=0}^{\infty} \frac{1}{r^{n+1}} \int (r')^n P_n \cos \theta \rho(r') d\tau'.$$

Substituting $r' = z \cos \theta$, and evaluating when $\theta = 0$ and $\theta = \frac{\pi}{2}$, we see:

$$\begin{aligned} V(r) &= \frac{1}{4\pi\epsilon_0} \sum_{n=0}^{\infty} \frac{1}{r^{n+1}} \frac{Q}{2a} \left[\int_0^a (r')^n P_n \cos \alpha \Big|_{\theta=0} dr' + \int_0^a (r')^n P_n \cos \alpha \Big|_{\theta=\pi} dr' \right] \\ &= \frac{1}{4\pi\epsilon_0} \sum_{n=0}^{\infty} \frac{1}{r^{n+1}} \frac{Q}{2a} \left[\int_0^a z^n P_n \cos \theta dz + \int_0^a z^n P_n \cos(\theta - \pi) dz \right] \quad (\text{making the substitution}) \\ &= \frac{1}{4\pi\epsilon_0} \sum_{n=0}^{\infty} \frac{1}{r^{n+1}} \frac{Q}{2a} [P_n \cos \theta + P_n \cos(\theta - \pi)] \int_0^a z^n dz \\ &= \frac{Q}{4\pi\epsilon_0} \sum_{n=0}^{\infty} \frac{a^n}{r^{n+1}} P_n \cos \theta \quad (\text{Since when } n \text{ is odd } P_n \cos \theta + P_n \cos(\theta - \pi) = 0) \end{aligned}$$

As Desired.

Problem 2. *Griffiths 3.47*

(a) We compute the average electric field as follows:

$$E_{\text{avg}} = \frac{1}{4/3\pi R^3} \int E(\mathbf{r}) d\tau' = \frac{1}{4/3\pi R^3} \int \frac{-q}{4\pi\epsilon_0} \frac{\hat{\mathbf{r}}}{r^2} d\tau' = \frac{1}{4/3\pi R^3} \cdot \frac{-q}{4\pi\epsilon_0} \int \frac{1}{r^2} \hat{\mathbf{r}} d\tau'.$$

Using formula 2.15 we compute the field of a sphere at as:

$$E(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\hat{\mathbf{r}}}{r^2} \rho(r') d\tau' = \frac{1}{4\pi\epsilon_0} \frac{1}{4/3\pi R^3} \int \frac{-q}{r^2} \hat{\mathbf{r}} d\tau'.$$

(b) Using Gauss' Law, we compute:

$$\oint \mathbf{E} \cdot d\mathbf{a} = \frac{q_{\text{enc}}}{\epsilon_0} \implies |\mathbf{E}| \cdot 4\pi r^2 = \frac{-q \cdot (4/3)\pi r^3}{\epsilon_0 \cdot 4/3\pi R^3} \implies \mathbf{E} = \frac{-qr}{4\pi\epsilon_0} \hat{\mathbf{r}} = \frac{-P}{4\pi\epsilon_0}.$$

(c) We can write an arbitrary charge distribution as $P = \sum q_i r'_i$. Since integration is linear by superposition the formula above should hold as well.

(d) If we place a charge q outside of the sphere, the average \mathbf{E} will be :

$$E_{\text{avg}} = \frac{1}{4/3\pi R^3} \int \mathbf{E} d\mathbf{a} = \frac{1}{4/3\pi R^3} \frac{-q}{r^2} \cdot \frac{4/3\pi R^3}{4\pi\epsilon_0} = \frac{q}{4\pi\epsilon_0 r^2} \hat{\mathbf{r}}.$$

We can use superposition to extend to any distribution of charges.

Problem 3. *Griffiths 3.56*

First note that $F_{\text{dip}}(r, \theta) = \frac{qp}{4\pi\epsilon_0 r^3} (2 \cos \theta \hat{r} + \sin \theta \hat{\theta})$. We assume that $\phi = 0$, and r is fixed. In spherical coordinates, we have that acceleration can be written as:

$$(\ddot{r} - r\dot{\theta}^2 - r\dot{\phi}^2 \sin^2 \theta) \hat{r} + (r\ddot{\theta} + 2\dot{r}\dot{\theta} - r\dot{\phi}^2 \sin \theta \cos \theta) \hat{\theta} + (r\ddot{\phi} + 2\dot{r}\dot{\phi} \cos \theta) \hat{\phi}.$$

Equating this to Force, and approximating $\sin \theta$ as θ , we get that the last term vanishes, and

$$\ddot{\theta} = \frac{q \cdot p}{4\pi\epsilon_0 r^4 m} \theta.$$

This is solved by a θ which is periodic in t .

Since r is fixed and ϕ is 0, we have that the particle traces out an arc.

Problem 4. *Griffiths 4.2*

We first compute the electric field due to the electron cloud. At a radius of r , the enclosed charge is:

$$Q_{\text{enc}} = \int_0^r \rho d\tau = \frac{4q}{3a^3} \int_0^r (r')^2 e^{-\frac{2r'}{a}} dr' = q \left(1 - e^{-\frac{2r}{a}} \left(1 + \frac{2r}{a} + \frac{2r^2}{a^2} \right) \right).$$

Therefore by Gauss' Law, we have that

$$E_{\text{in}} = \frac{1}{4\pi\epsilon_0 r^2} q \left(1 - e^{-\frac{2r}{a}} \left(1 + \frac{2r}{a} + \frac{2r^2}{a^2} \right) \right).$$

If we apply some small electric field E_{ext} , the nucleus will get shifted by some small quantity d so that $E_{\text{ext}} = E_{\text{in}}$. We Taylor expand $e^{-\frac{2r}{a}}$ as:

$$e^{-\frac{2r}{a}} = 1 - \frac{2d}{a} + \frac{2d^2}{a^2} - \frac{4d^3}{3a^3} + \dots$$

And so

$$1 - e^{-\frac{2d}{a}} \left(1 + \frac{2d}{a} + \frac{2d^2}{a^2} \right) = 1 - 1 - \frac{2d}{a} - \frac{2d^2}{a^2} + \frac{2d}{a} + \frac{4d^2}{a^2} - \frac{2d^2}{a^2} + \frac{4d^3}{3a^3} + \dots = \frac{4}{3} \frac{d^3}{a^3} + O\left(\frac{d^5}{a^5}\right).$$

Thus at d we have $E_{\text{in}} = E_{\text{out}}$ and so

$$E = \frac{1}{4\pi\epsilon_0} \cdot \frac{q}{d^2} \cdot \frac{4d^3}{3a^3} = \frac{qd}{3a^3\pi\epsilon_0} = \frac{P}{3a^3\pi\epsilon_0}.$$

Therefore $\alpha = 3a^3\pi\epsilon_0$.

Problem 5. *Griffiths 4.4*

We have that the electric field of q at distance r is

$$|E| = \frac{q}{4\pi\epsilon_0 r^2}.$$

We can place the charge on the x axis so that $\vec{p} = \alpha E = \frac{\alpha q}{4\pi\epsilon_0 r^2} \hat{x}$. We know that $E_{\text{dip}}(r) = \frac{1}{4\pi\epsilon_0 r^3} [3(\vec{p} \cdot \hat{r}) - p]$. Since $\hat{p} = \hat{x}$, we have that

$$3\vec{p} \cdot \hat{r} - p = \frac{3\alpha q}{4\pi\epsilon_0 r^3} \hat{x} - \frac{\alpha q}{4\pi\epsilon_0 r^3} \hat{x} = \frac{2\alpha q}{4\pi\epsilon_0 r^3} \hat{x}.$$

Therefore $F = qE = \frac{2\alpha q^2}{16\pi^2 \epsilon_0^2 r^3} \hat{x}$.

Problem 6. *Griffiths 4.6*

We first place an dipole at $-z$. The boundary condtions are the same, so we can instead compute the torque this dipole exerts on \mathbf{p} . We have that $\mathbf{E}_{\text{dip}}(r, \theta) = \frac{\mathbf{p}}{4\pi\epsilon_0 r^3} (2 \cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\boldsymbol{\theta}})$. At a distance of $2z$ this becomes:

$$\mathbf{E}_{\text{dip}} = \frac{\mathbf{p}}{4\pi\epsilon_0 (2z)^3} (2 \cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\boldsymbol{\theta}}).$$

We can also write the dipole as $\mathbf{p} = p \cos \theta \hat{\mathbf{r}} + p \sin \theta \hat{\boldsymbol{\theta}}$. We compute the torque as

$$\mathbf{N} = \mathbf{p} \times \mathbf{E} = \frac{p^2}{4\pi\epsilon_0 (2z)^3} \left[(\cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\boldsymbol{\theta}}) \times (2 \cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\boldsymbol{\theta}}) \right] = \frac{-p^2 \sin \theta \cos \theta}{4\pi\epsilon_0 (2z^3)} \hat{\boldsymbol{\phi}}.$$

If we allow the dipole to rotate, it will come to a rest when the torque is at a minimum. so at $\theta = \frac{\pi}{4}, \frac{5\pi}{4}$.