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## Lab Cover Letter

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		Sig. figs.		Discussion of pos. errors
		Units		Suggestions to reduce errors
		Clarity of Presentation		66
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			( )	Paper Total (60 points)
(	)	Abstract (4)		(30 points for CME or EPF)
		Quantity or principle	( )	Notebook (10 points)
		How measurement was made		Format (proper style, following directions)
		Numerical Results		Apparatus (brief description of equipment,
		Conclusion		including sketches)
				Data (including computer file names and
(	)	Intro & Theory (9)		manually recorded data)
		Basic principle		Experimental Technique (describing your
		Main equations to be used		procedures; stating & justifying uncerts.)
		Apparatus		Analysis (results and errors)
		What will be plotted		,
		Fitting parameters related	( )	Workshoot(s)/Fill in the Plank
			( )	Worksheet(s)/Fill-in-the-Blank-
(	)	Exp. Procedures (15)	Repor	t (30 points) if applicable
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		Stating and justifying uncertainties	( )	Adjustments – late submissions,
		Data Record	( )	improper procedures, etc. – or bonus points
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## Rotational Kinetic Energy Lab

## Wolf S. Mermelstein

## October 31, 2023

## ${\bf Abstract}$

 ${\rm content...}$ 

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#### Introduction 1

#### Moment of Inertia 1.1

Moment of inertia, often denoted by I, is a function of the specific geometry and mass distribution of an object. Moment of inertia is implicitly relative to the axis of rotation. Where R is the distance from the axis of rotation and M is the mass of the object, for a point mass the moment of inertia, I, is given by

$$I = MR^2 \tag{1}$$

The entire moment of inertia can be computed by thinking of a given object as a collection of tiny masses. As the masses' volumes shrink down to some small volume with some proportionately small mass, the can then be eventually said to be the differential dM. Integrating all the small point masses across the entire object implies that the entire moment of inertia of an object,  $I_{tot}$ , is given by

$$I = \sum_{k=1}^{\infty} \frac{m}{k} R^2 \tag{2}$$

$$= \int R^2 dM \tag{3}$$

This is conceptually helpful in understanding moment of inertia for arbitrary shapes, but is not practically useful for non simple (i.e. circles, squares, collections of discreet point masses) shapes, f the counterweight, we state that such as the mass-loaded, spoked wheel that we used in our experiment. As a result of such, it is often helpful to actually measure moment of inertia instead of attempting to compute it.

#### 1.2Conservation of Energy

The translational kinetic energy of an object in motion with mass M moving at speed v is given to be

$$K_T = \frac{1}{2}Mv^2 \tag{4}$$

Since we know that

$$\frac{\theta}{2\pi} = \frac{s}{2\pi R}$$

$$s = \theta R \tag{5}$$

and

$$\frac{d}{dt}s = \frac{d}{dt}\theta R$$

$$v = \omega r \tag{6}$$

We can then derive from equation 4 that the rotational kinetic energy,  $K_R$ , is

$$K_R = \frac{1}{2}M(\omega R)^2$$

$$= \frac{1}{2}(MR^2)\omega^2$$

$$= \frac{1}{2}I\omega^2$$
 (7)

where I is defined to be the moment of inertia about the axis of rotation.

For the mass in figure ?? descends downwards due to gravity, it begins to lose its gravitational potential energy,  $U_W$ . The total energy of the system is internally conserved, however a small amount of energy is lost due to friction. So, where  $\Delta U_W$  is the change in the gravitational potential energy of the counterweight,  $K_T$ is the translational kinetic energy of the counterweight, and  $K_R$  is the rotational kinetic energy

$$\Delta U_W + K_T + K_R = W_f \tag{8}$$

Which, using equations 7 and 4, implies that

$$\Delta U_W + (\frac{1}{2}Mv^2) + (\frac{1}{2}I\omega^2) = W_f$$
 (9)

 $\Delta U_W$  should be negative, and  $K_T \& K_R$  positive because the mass is falling, and, thus, losing gravitational kinetic energy, whilst simultaneously proportionately gaining kinetic energy.

Using the fact that gravitational potential energy for an object at height h of mass m in an environment where gravity can be approximated to g is given to be

$$U_G = (M \cdot g \cdot h) \tag{10}$$

Plugging this in, and renaming h be y, we get the final equation

$$-(M \cdot g \cdot y) + (\frac{1}{2}Mv^2) + (\frac{1}{2}I\omega^2) = W_f \quad (11)$$

### 1.3 Working Equation

Since for our specific experiment we used paperclips attached to the counterweight to cancel out friction, we can instead rewrite equation 11 to be

$$-(M \cdot g \cdot y) + (\frac{1}{2}Mv^2) + (\frac{1}{2}I\omega^2) = 0$$
 (12)

Carefully noting that we have discluded the energy of the moving paperclips, as it is negligible in comparison to the other energies of the system. And, to further simplify things, we will define y to be vertically positive, so as to make the equation into

$$-(M \cdot g \cdot y) + (\frac{1}{2}Mv^2) + (\frac{1}{2}I\omega^2) = 0 \qquad (13)$$

Then, using relationship 6, we plug in  $\frac{v}{r}$  for  $\omega$ , resulting in the equation

$$\begin{split} -(M\cdot g\cdot y) + (\frac{1}{2}Mv^2) + (\frac{1}{2}I(\frac{v}{r})^2) &= 0 \\ &= \frac{1}{2}\cdot v^2\cdot (M \\ &\qquad (14) \end{split}$$

Or, as will be used later for our computations, the equivalent equation in the form

$$gy = \frac{1}{2}(1 + \frac{I}{Mr^2}) \cdot v^2 \tag{15}$$

v is a value that is determined by our Logger  $Pro^{\mathsf{TM}}$ software. It is computed with an advanced proprietary algorithm, but is similar to

$$v \approx \frac{\Delta s}{\Delta T} \tag{16}$$

Figure 1: Visual representation of k and r

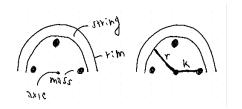
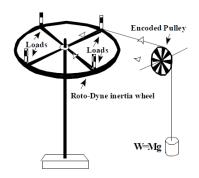


Figure 2: Roto-Dyne wheelInertia Wheel Apparatus [2]



## 2 Procedure

### 2.1 Taking Measurements

Before conducting our experiment, we took measurements of various parts of our setup. First we obtained r from the lab manual, the distance from the axle of the wheel to the string, and then we measured k, the distance from the axle of the wheel to the masses. Our measurement for k was originally 0.073, but we shifted it to 0.173 because the original value did not make result of misreading the ruler. For r we used the provided uncertainty of  $\pm 0.002$ m, whereas for k we measured very carefully and chose the uncertainty to be 0.01m

$$r = 0.200 \pm 0.002 \text{ m}$$
 (17)

$$k = 0.173 \pm 0.01 \text{ m}$$
 (18)

We used a counterweight with a given mass of 0.06kg to provide a torque to spin our Roto-Dyne wheel, as can be seen in figure ??. To account for friction, we incrementally added paperclips to the bottom of the counterweight. We continued to add paperclips up until the mass would fall at a constant speed to counteract the force of friction, using  $Logger\ Pro^{\mathbb{M}}$  software and an encoded pulley to monitor acceleration and velocity. Let  $M_c$  be the mass of the counterweight and  $M_p$  be the mass of the paperclips opposing friction, not used in our computations but still important to the experimental design.

$$M_p = 0.0015 \pm 0.0001 \text{ kg}$$
 (19)

$$M_c = 0.06 \text{ kg}$$
 (20)

Also, we were provided with the mass of the Roto-Dyne wheel,  $M_R$  and the mass loads,  $M_L$ , of which there were 4.

$$M_R = 1.5 \text{ kg}$$
 (21)

$$M_L = 0.225 \pm 0.002 \text{ kg}$$
 (22)

For the encoded pulley that we used to measure velocity and length of unrolled string it was given that the gaps between intervals of measurement,  $\Delta s$ , was

$$\Delta s = 0.015 \text{ m}$$

## 2.2 Estimating Moment of Inertia

Before measuring the moment of inertia, we decided to make a rough approximation. To do this, we used two different common models, the moment of inertia of a disc,  $\frac{1}{2}Mr^2$ , and the moment of inertia of a ring,  $Mr^2$ . These computations will not be accurate since the actual Rote-Dyne disc is neither a perfect disc nor a hoop. For our rough estimates, we took the radius to be r, and the mass to be

$$M_{tot} = M_R + 4M_L$$
$$= 2.4 \text{ kg}$$
(23)

For the disc estimate, we got

$$I_{\rm disc} = \frac{1}{2}Mr^2$$
  
= 0.048 kg m<sup>2</sup> (24)

For the hoop estimate, we got

$$I_{\text{hoop}} = Mr^2$$
  
= 0.096 kg m<sup>2</sup> (25)

To determine an overall estimate,  $I_{\rm est}$ , we averaged  $I_{\rm disc}$  and  $I_{\rm hoop}$ , and then set the uncertainty to be that average. Assuming that the wheel is a hoop is an under-estimate, and assuming that the wheel is a disc is an over-estimate, so the actual moment of inertia should be somewhere between the two.

$$I_{\text{est}} = \frac{I_{\text{disc}} + I_{\text{hoop}}}{2}$$
  
= 0.072 ± 0.072 kg m<sup>2</sup> (26)

#### 2.3 Monte Carlo Simulation

Before actually measuring the moment of inertia, we performed a Monte Carlo Simulation using our estimated I so that we could compare a graph of data of the estimated value to a graph of data for the actual value later on.

To perform this simulation, we began by creating a new  $Origin\ Pro^{\mathbb{T}}$  document, and arranged a table including a column for  $\Delta s$ , the overall displacement (vertical distance it has fallen) of the string, and  $\Delta T_0$ , the time elapsed since dropping the counterweight to arrive at that overall displacement. The values for  $\Delta s$  were computed using the equation

$$y_i = i\Delta s \tag{27}$$

Which is simply stating that the total displacement of the rope is equal to the amount of a single displacement, a notch in the encoded pulley, times the number of increments, which is the total number of notches that passed by the laser at a given time point. To compute the  $\Delta T_0$  from the  $\Delta s$  we derived the following equation from equations 16 and 15.

First, we manipulated equation 15 by solving for v.

$$gy = \frac{1}{2}(1 + \frac{I}{Mr^2}) \cdot v^2$$

$$v^2 = \frac{\frac{1}{2}(1 + \frac{I}{Mr^2})}{gy}$$

$$v = \sqrt{\frac{\frac{1}{2}(1 + \frac{I}{Mr^2})}{qy}}$$
(28)

Next we solved equation 16 for t. For our use case here, we will allow  $\Delta T$  to be  $\Delta T_0$ .

$$v \approx \frac{\Delta s}{\Delta T}$$

$$\Delta T = \frac{\Delta s}{v}$$

$$\Delta T_0 = \frac{\Delta s}{v}$$
(29)

Then, we solved for  $\Delta T_0$  by plugging in v from equation 29 and  $\Delta s$  into equation 29.

$$\Delta T_0 = \frac{\Delta s}{v}$$

$$\Delta T_0 = \frac{\Delta s}{\sqrt{\frac{\frac{1}{2}(1 + \frac{I}{Mr^2})}{gy}}}$$

$$\Delta T_0^2 = \frac{\Delta s^2}{\frac{\frac{1}{2}(1 + \frac{I}{Mr^2})}{gy}}$$

$$\Delta T_0^2 = \frac{\Delta s^2(1 + \frac{I}{Mr^2})}{2gy}$$

$$\Delta T_0 = \Delta S \sqrt{\frac{(1 + \frac{I}{Mr^2})}{2gy}}$$
(30)

With a simple script we had  $Origin\ Pro^{\mathbb{M}}$  apply this equation to each row, utilizing that respective row's  $\Delta s$  value. Now, since this data is purely based on an estimated moment of inertia value, we applied Monte Carlo randomization. To do this, we shifted each  $\Delta T0_i$  by some  $\Delta T0_Gi$  obtained from a Gaussian distribution G with the given mean 0 and  $\sigma$  equal to the estimated uncertainty for  $\Delta t$ ,  $\delta_{\Delta t}$ , which was given to be 0.0002s, resulting with a column with values for  $\delta t_R$ . This can be expressed formulaically

$$\delta T_R = \Delta T 0 + \Delta t \cdot \text{grnd}() \tag{31}$$

After running this computation, for the first three values of  $\Delta T_r$  we got

Table 1: Samples of random Monte Carlo data generation

trial $\#$	s
1	0.15409
2	0.10925
3	0.08891
1	0.15381
2	0.10883
3	0.08919
1	0.01540
2	0.10891
3	0.08891
1	0.15385
2	0.10918
3	0.08894

## 3 Results

### 3.1 Plotting Data

For the actual results, we measured the velocity, acceleration, time (duration), and position of the falling counterweight mass using the encoded pulley and  $Logger\ Pro^{\mathbb{T}}$ . We then imported the data into  $Origin\ Pro^{\mathbb{T}}$  for analysis and to help us compile plots. To keep our data consistent, we decided to trim the first three rows and all rows after row 44. Data beyond that in either direction was problematic because of us abruptly setting up and stopping the counterweight from hitting the floor.

To visualize our data, we plotted  $v^2$  against displacement for both the system with and without load masses, and did the same for our simulation. For our  $v^2$  column in  $Origin\ Pro^{\mathbb{N}}$  we used

$$v = \frac{\Delta s}{\Delta T}$$

$$v^2 = (\frac{\Delta s}{\Delta T})^2$$

$$v^2 = (\frac{0.015}{\Delta T_i})^2$$
(32)

To include error bars in our  $Origin \ Pro^{\mathbb{M}}$  plot, we applied the derivative method row-wise to

equation 32, using equation 34. We obtained the value

$$\delta_{\Delta T} = 0.0002s \tag{33}$$

from our lab manual, which is a consequence of the intrinsic inaccuracy of  $Logger\ Pro^{TM}$  and our recording hardware. Solving for  $\delta_{v^2}$ , we got

$$\delta_{v^2} = \left| \frac{\partial}{\partial \Delta T_i} \left( \left( \frac{0.015}{\Delta T_i} \right)^2 \right) \cdot \delta_{\Delta t} \right|$$

$$= \left| 0.015^2 \cdot \frac{\partial}{\partial \Delta T_i} (\Delta T^2) \cdot \delta_{\Delta t} \right|$$

$$= \left| 0.000225 \cdot -2\Delta T^{-3} \cdot \delta_{\Delta t} \right|$$

$$= \frac{.00000009}{\Delta T^3}$$
(34)

Then we used  $Origin Pro^{TM}$  to compute a line of best fit, of which the regression and slope has been superposed onto our plot figures, figure??, ??, and ??.

These aforementioned plots can be found in the appendix, section 6. Additionally, tables ?? and ?? contain the actual datatables with raw Logger  $Pro^{\mathsf{TM}}$ data used for generating the plots.

#### $\mathbf{4}$ ${f Analysis}$

**Objective** For the analysis of our data, our objective was to compute  $I_E$ , the measured moment of inertia for just the mass loads on the wheel. In order to do this, we took many intermediate steps to prevent mistakes and measurements. This included a Monte Carlo simulation to verify our computation for  $I_1$ , the moment of the Roto-Dyne wheel including the mass loads, and an careful estimate of  $I_E$  using the definition of moment of inertia.

#### Computing $I_1$ (w/masses) 4.1

The first step towards accomplishing this goal was to determine whether our estimate for the moment of inertia of the Roto-Dyne wheelwas reasonably close to our measured moment value. In order to make this determination we first worked been discussed earlier in our procedure. The two to derive a formula for I, the measured value of the moment of inertia. To start, we note the

equation for the slope of the Roto-Dyne wheelwith the load masses,  $B_1$  and  $B_2$ , the slope without load masses, obtained through  $Origin Pro^{TM}$ 's provided linear fit line.

Using the elementary equation for some slope, m, with a linear function plotted on an x-y graph,

$$m = \frac{\Delta y}{\Delta x} \tag{35}$$

We determined that for our vertical axis  $v^2$ and horizontal axis y,

$$B_1 = \frac{v^2}{y} B_2 = \frac{v^2}{y} (36)$$

First, we solved equation 15 for  $v^2$  as a function of y

$$gy = \frac{1}{2}(1 + \frac{I}{Mr^2}) \cdot v^2 \tag{15}$$

$$v^2 = \frac{2gy}{1 + \frac{I}{Mr^2}} \tag{37}$$

Then we plugged this value for  $v^2$  into equation 36, finally solving for  $I_1$ 

$$B_{1} = \frac{\frac{2gy}{1 + \frac{I_{1}}{Mr^{2}}}}{y}$$

$$= \frac{\frac{2gy}{1 + \frac{I_{1}}{Mr^{2}}}}{y} \cdot \frac{1 + \frac{I_{1}}{Mr^{2}}}{1 + \frac{I_{1}}{Mr^{2}}}$$

$$= \frac{2g}{(1 + \frac{I_{1}}{Mr^{2}})}$$

$$By \frac{I_{1}}{Mr^{2}} = 2g - B$$

$$I_{1} = Mr^{2} \cdot \frac{2g - B}{B_{1}}$$

$$I_{1} = Mr^{2} \cdot (\frac{2g}{B_{1}} - 1)$$
(38)

**Solving for**  $I_1$  We then plugged in the values to solve for  $I_1$  for the system including the load masses. Most of the values have already novel values are  $B_1$  and its uncertainty, which will henceforth be denoted as  $\delta_B$ , provided by Origin  $Pro^{TM}$ 's linear regression model, and q,

the generally accepted approximation of nearearth gravity. The more interesting situation, the situation where there are load masses in the Roto-Dyne wheel, is the one we made an estimate and did a Monte Carlo simulation for. We will compute the system without masses for the sake of completion later.

$$r = 0.200 \pm 0.002 \text{ m}$$
 (17)

$$g = 9.81 \pm 0.01 \text{ ms}^2 \tag{39}$$

$$M = 0.0600 \pm 0.0001 \text{ kg}$$
 (20)

$$B_1 = 0.68204 \pm 0.00168 \,\frac{\mathrm{m}}{\mathrm{s}^2} \tag{40}$$

So,

$$I_1 = (.06)(.2)^2 \cdot (\frac{2(9.81)}{0.68204} - 1)$$
  
 $\approx 0.06663 \text{ kg m}^2$  (41)

Computing Error To compute the error, we applied the derivative method to equation 38, only factoring the error on  $I_1$  due to B as strictly requested by our lab manual. It is, however, worth noting that we believe that the uncertainties due to M and r could be significant, especially given that the r term is squared. These uncertainties will, however, be taken into account for computation of  $I_E$  later.

$$\delta_{I} = \left| \frac{\partial}{\partial B} (Mr^{2} \cdot (\frac{2g}{B} - 1)) \cdot \delta_{B} \right| 
= \left| \frac{\partial}{\partial B} (Mr^{2} \cdot (2gB^{-1} - 1)) \cdot \delta_{B} \right| 
= 2Mgr^{2} \delta_{B} \cdot \left| \frac{\partial}{\partial B} (B^{-1} - 1) \cdot \delta_{B} \right| 
= 2Mgr^{2} \delta_{B} \cdot \left| \frac{-1}{B^{2}} \right| 
= 2\delta_{B} \frac{Mgr^{2}}{B^{2}}$$
(42)

Plugging in values, we get

$$\delta_I = 2\delta_B \frac{Mgr^2}{B^2} \approx 0.00017 \text{ kg m}^2$$
 (43)

So, we can succinctly state

$$I_1 = 0.06663 \pm 0.00017 \text{ kg m}^2$$
 (44)

Looking back at our predicted value for  $I_1$  from section 2.2, we chose an estimate of I, denoting it  $I_{\text{est}}$ . Putting the previously estimated  $I_1$ , " $I_{1,\text{est}}$ " inline with the measured  $I_1$ , " $I_1$ ", we have

$$I_{1,\text{est}} = 0.072 \pm 0.072 \text{ kg m}^2$$
 (26)

$$I_1 = 0.06663 \pm 0.00017 \text{ kg m}^2$$
 (44)

It is obvious that the actual value falls within one to two uncertainty intervals of the estimated value. The more important question, however, is whether the estimated value falls within one to two uncertainty intervals of the actual value. Two uncertainty intervals is

$$\begin{aligned} \delta_{I_1} \cdot 2 \\ &= 0.00017 \cdot 2 \\ &= 0.00034 \end{aligned}$$

So, mathematically, we are checking

$$0.07182 \notin$$

$$\notin [0.07182 - 0.00034, 0.07182 + 0.00034]$$

$$\notin [0.07148, 0.07216]$$
(45)

Which is false, so our measured indeed value does not fall within one to two uncertainty intervals from our estimate. To see how off we are, a percentage value also can be computed,  $\Delta I_{\rm skew}$ 

$$\Delta I_{\text{skew}} = 100 \cdot \left(1 - \frac{I_1}{I_{1,\text{est}}}\right)$$

$$\approx 100 \cdot \left(1 - \frac{0.06663}{0.072}\right)$$

$$= 100 \cdot \left(1 - 0.92541\right)$$

$$= 7.45\%$$
(46)

## 4.2 Computing $I_2$ (w/o masses)

**Solving for**  $I_2$  Next we compute the value of  $I_2$ , the moment of inertia of the Roto-Dyne

wheelwithout mass loads, using the same equations for  $I_1$ , equation 38 and for  $\delta_{I_2}$ , equation 42.

For the system without masses,  $Origin Pro^{\mathbb{M}}$  produced the linear regression value

$$B_2 = 1.26494 \pm 0.00889 \ \frac{\mathrm{m}}{\mathrm{s}^2} \tag{47}$$

$$I_2 = Mr^2 \cdot (\frac{2g}{B} - 1) \pm 2\delta_B \frac{Mgr^2}{B^2}$$

$$= (0.0600)(0.200)^2 \cdot (\frac{2(9.81)}{(1.26494)} - 1)$$

$$\pm 2(0.00889) \frac{(0.0600)(9.81)(0.200)^2}{(1.26494)^2}$$

Which evaluates to

$$I_2 = 0.03482 \pm 0.00026 \text{ kg m}^2$$
 (48)

for the massless system. It makes intuitive sense that the moment of inertia is lesser. Inertia is a measure of all the infinite infinitesimally small point masses of an object. There are less masses in the system without mass loads, so the moment of inertia indeed should be notably lesser.

## 4.3 Computing $I_E$ (mass loads)

We are now very close to being able to find  $I_E$ , the moment of just the mass loads. Knowing the experimental moment of inertia for the system with the mass loads and the system without the mass loads, we can compute  $I_E$  by subtracting the moment of inertia of the system with the mass loads by the system without the mass loads.

$$I_E = I_1 - I_2$$
 (49)  
= 0.06663 - 0.03482  
= 0.03181 kg m<sup>2</sup> (50)

We will base the uncertainty of this measurement on the literal, complete equation, subbing I into equation 49 for equation 38.

$$I_{E} = I_{1} - I_{2}$$

$$= Mr^{2} \cdot (\frac{2g}{B_{1}} - 1) - Mr^{2} \cdot (\frac{2g}{B_{2}} - 1)$$

$$= Mr^{2} ((\frac{2g}{B_{1}} - 1) - (\frac{2g}{B_{2}} - 1))$$

$$= Mr^{2} \frac{2g(B_{2} - B_{1})}{B_{1}B_{2}}$$

$$(51)$$

Writing the equation like this does not help us solve for  $I_E$ , in fact if anything it makes it more convoluted. Rather, the purpose of this literal expression is to allow us to compute an uncertainty based on the values that it is based on. We believe that  $\delta_{I_E,M}$  and  $\delta_{I_E,g}$  are negligible in comparison to  $\delta_{I_E,r}$ ,  $\delta_{I_E,B_1}$ , and  $\delta_{I_E,B_2}$ , since, besides the fact that the actual value for  $\delta_M$  is 0.0001kg, r is squared which makes its uncertainty even more significant. To start, we recall

$$B_1 = 0.68204 \pm 0.00168 \frac{\text{m}}{\text{s}^2}$$
 (40)

$$B_2 = 1.26494 \pm 0.00889 \frac{\text{m}}{\text{s}^2}$$
 (47)

$$r = 0.200 \pm 0.002 \text{ m}$$
 (17)

$$M = 0.0600 \pm 0.0001 \text{ kg}$$
 (20)

$$g = 9.81 \pm 0.01 \frac{m}{s^2} \tag{39}$$

To compute the overall uncertainty  $\delta_{I_E}$ , we will first compute  $\delta_{I_E,r}$  using the derivative method, and  $\delta_{I_E,B_1}$ , and  $\delta_{I_E,B_2}$  using the computational method, and then will add all the uncertainties in quadrature.

$$\delta_{I_E,r} = \left| \frac{\partial}{\partial r} (Mr^2 \frac{2g(B_2 - B_1)}{B_1 B_2}) \cdot \delta_r \right|$$

$$= \left| 2rM \frac{2g(B_2 - B_1)}{B_1 B_2} \cdot \delta_r \right|$$

$$= 0.00063 \text{ kg m}^2$$
(52)

$$\delta_{I_E,B_2} = |Mr^2 \frac{2g(B_2 - B_1)}{B_1 B_2} - Mr^2 \frac{2g((B_2 + \delta_{B_2}) - B_1)}{B_1 (B_2 + \delta_{B_2})}|$$

$$\approx 0.00026 \text{ kg m}^2$$
(53)

$$\delta_{I_E,B_1} = |Mr^2 \frac{2g(B_2 - B_1)}{B_1 B_2} - Mr^2 \frac{2g((B_1 + \delta_{B_1}) - B_1)}{(B_1 + \delta_{B_1}) B_2}|$$

$$\approx 0.00017 \text{ kg m}^2$$
(54)

So, to find the overall uncertainty we add in quadrature, getting

$$\begin{split} \delta_{I_E} &= \\ &= \sqrt{\delta_{I_E,r}^2 + \delta_{I_E,B_1}^2 + \delta_{I_E,B_2}^2} \\ &= \sqrt{0.00063^2 + 0.00026^2 + 0.00017^2} \\ &= 0.00070 \text{ kg m}^2 \end{split} \tag{56}$$

We can then succinctly state that

$$I_E = 0.03181 \pm 0.00070 \text{ kg m}^2$$
 (57)

### 4.3.1 Estimating $I_E$

Now, like how we came up with an approximate value for the Roto-Dyne wheel with masses and loads by approximating it to be a disc and a ring and averaging the two, we wanted to come up with a estimate in similar fashion for  $I_E$ . To do this, we treated the four loads  $M_L$ , of  $0.225\pm0.002$ kg each, as point masses for computing an estimate for  $I_E$ . To do this, we applied equation 1 to all 4 point masses. Note that we use k here instead of r since the masses are at distance k from the center of the Roto-Dyne wheel.

$$I_{\text{est},E} = \sum Mk^2$$

$$= 4(M_L k^2)$$

$$= 4 \cdot 0.225 \cdot 0.173^2$$

$$= 0.02693 \text{kg m}^2$$
(58)

Computing  $\delta_{I_{\mathrm{est},E}}$  is fairly trivial, requiring us to propagate  $\delta_k$  and  $\delta_{M_L}$ 's effect on  $\delta_{I_{\mathrm{est},E}}$  and then sum in quadrature. For this we used the derivative method.

$$\delta_{I_{\text{est},E},k} = \left| \frac{\partial}{\partial k} (4(M_L k^2)) \cdot \delta_k \right|$$

$$= |8M_L k \delta_k|$$

$$= |8 \cdot 0.225 \cdot 0.173 \cdot 0.01|$$

$$= 0.00311$$
(59)

$$\delta_{I_{\text{est},E},M_L}$$

$$= \left| \frac{\partial}{\partial M_L} (4M_L k^2 \delta_{M_L}) \right|$$

$$= \left| 4k^2 \delta_{M_L} \right|$$

$$= \left| 4 \cdot 0.173^2 \cdot 0.002 \right|$$

$$= 0.00023 \text{kg m}^2$$
(60)

Adding the uncertainties  $\delta_{I_{{\rm est},E},M_L}$  and  $\delta_{I_{{\rm est},E},r}$  in quadrature we get

$$\delta_{I_{\text{est},E}} = \sqrt{\delta_{I_{\text{est},E},r}^2 + \delta_{I_{\text{est},E},M_L}^2}$$

$$= \sqrt{0.00311^2 + 0.00023^2}$$

$$= 0.00311 \text{kg m}^2$$
(61)

Succinctly, we can thus state  $I_{est,E}$  to be

$$I_{\text{est},E} = 0.02693 \pm 0.00311 \text{ kg m}^2$$
 (62)

Now we can compare this to our value for  $I_E$  to see if it is reasonable and agrees two within two uncertainty intervals. We recall that our value for  $I_E$  is

$$I_E = 0.03181 \pm 0.00070 \text{ kg m}^2$$
 (57)

Two uncertainty intervals from  $I_E$  is

$$\delta_{I_E} \cdot 2$$
=0.00311 · 2
=0.00622 (63)

So, upon inspection of the mathematical statement  $% \left( 1\right) =\left( 1\right) \left( 1\right)$ 

 $0.02693 \in [0.03181 - 0.00622, 0.03181 + 0.00622]$   $0.02693 \in [0.02559, 0.03803]$  (64)

we determine that indeed our estimated uncertainty lies within two uncertainty intervals. This provides us assurance that our value for  $I_E$  is indeed reasonable.

## 5 Conclusion

## Acknowledgments

I would like to thank Christopher Richner and Lily Kagy, CWRU Department of Physics, for their help in obtaining the experimental data, collaborating on preparation of the figures, and checking calculations. Additionally, I would like to thank Olivia Green, CWRU Department of Physics, for helping facilitate our lab.

## References

- [1] Resnick Halliday and Walker. Fundamentals of Physics. Addison-Wesley Professional, 6 edition, 2023.
- [2] D. Schultz. General Physics I: Mechanics Lab Manual. CWRU Bookstore, Spring 2004.

## 6 Appendix

Figure 3: Monte Carlo Simulation of Rotational Kinetic Energy Experiment plot

# Monte Carlo Simulation of Rotational Kinetic Energy Experiment

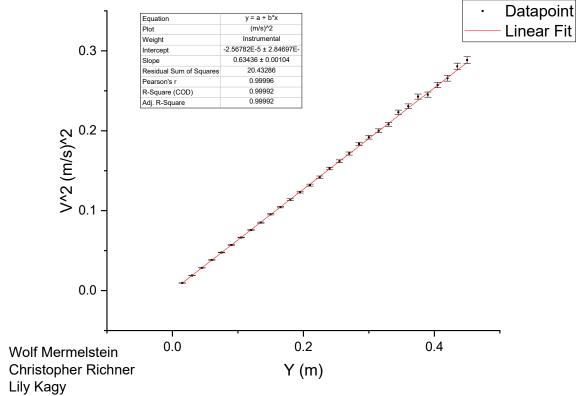


Figure 4: Masses  $v^2$  vs sDist plot

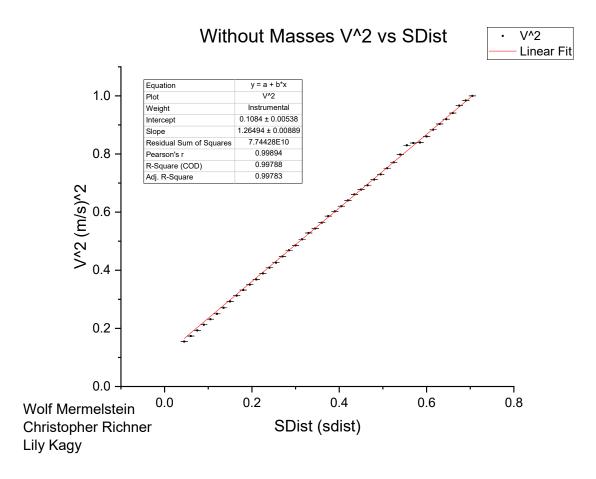


Figure 5: With Masses  $v^2$  vs s Dist plot

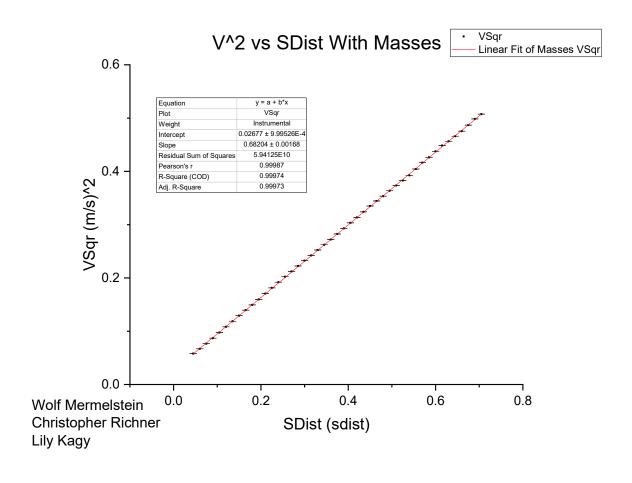


Table 2: Data for falling counterweight without mass loads in place

Time	STime	SDist	SVel	SAccel
t	stime	sdist	svel	saccel
s	s	m	m/s	$m/s\hat{2}$
0.2003832	0.3409468	0.045	0.240637469	0.2974361002
0.2421024	0.4010488	0.06	0.2587705469	0.3059740303
0.2727832	0.4571542	0.075	0.2768823709	0.3396616489
0.3099096	0.5096392	0.09	0.2950082819	0.3510465914
0.3375832	0.5590336	0.105	0.3125067287	0.3574728706
0.371984	0.6057826	0.12	0.3289729322	0.3469787101
0.3977972	0.6503324	0.135	0.3445762028	0.3535079613
0.4301136	0.6929332	0.15	0.3594867431	0.3465042602
0.4541972	0.7338524	0.165	0.3734799691	0.3374400019
0.4841948	0.7733118	0.18	0.3866063477	0.3278706007
0.506816	0.8114928	0.195	0.3997294453	0.3595444818
0.5350836	0.8484054	0.21	0.4130797945	0.3638046842
0.5563028	0.8841552	0.225	0.4258449397	0.3503333039
0.5829836	0.9188832	0.24	0.4381633461	0.3590888604
0.6031832	0.9526498	0.255	0.4497965385	0.3299465981
0.6285816	0.9856	0.27	0.4607194593	0.3330491259
0.6478832	1.0177834	0.285	0.4715773455	0.3417015924
0.6720832	1.0492332	0.3	0.4820423673	0.3238048217
0.6905832	1.0800322	0.315	0.492183323	0.3347201767
0.7137832	1.1101992	0.33	0.5023979233	0.3424834134
0.7315832	1.1397582	0.345	0.5121068853	0.3144374509
0.7539216	1.1687902	0.36	0.5216114938	0.3403302933
0.7710904	1.1972824	0.375	0.5315253346	0.355568289
0.792702	1.2252416	0.39	0.5413072074	0.3441565057
0.8092096	1.2527124	0.405	0.5507273167	0.3416705778
0.8302836	1.2797228	0.42	0.559880448	0.3360780896
0.8462136	1.3063022	0.435	0.5692255573	0.3671062769
0.8665272	1.3324336	0.45	0.5787253594	0.359973184
0.8820832	1.358147	0.465	0.5869743702	0.2816386499
0.9017832	1.383547	0.48	0.5946036744	0.3190939627
0.9168836	1.4086054	0.495	0.602973345	0.3489192067
0.9359832	1.4333056	0.51	0.6110845183	0.3078506391
0.9506832	1.4577022	0.525	0.6186489591	0.3122719003
0.9693164	1.481802	0.54	0.6266088147	0.3483025149
0.983612	1.5055832	0.555	0.6359267169	0.4353334838
1.0018836	1.5289834	0.57	0.6452122961	0.3582989806
1.0158992	1.5520834	0.585	0.652922225	0.3092271113
1.0336832	1.5749334	0.6	0.661597301	0.4500795393
1.0474092	1.5974336	0.615	0.669666613	0.2671862678
1.0647832	1.6197336	0.63	0.6756628186	0.27059002
1.0782148	1.6418362	0.645	0.6826017458	0.3572933224
1.0952812	1.663686	0.66	0.6896131877	0.2844921273
1.1083832	1.6853406	0.675	0.69766979	0.4596086524
1.1251172	1.7066906	0.69	0.7060389196	0.32438475
1.1379832	1.7278332	0.705	0.7123894853	0.2763517435

Table 3: Data for falling counterweight without mass loads in place

Time	STime	SDist	SVel	SAccel
t	$_{ m stime}$	sdist	svel	saccel
s	S	m	m/s	$m/s\hat{2}$
0.0942928	0.1784214	0.045	0.3931293508	0.6331416361
0.118666	0.2154714	0.06	0.4167682347	0.6429114748
0.13679	0.2505154	0.075	0.4395279097	0.6560084549
0.1593652	0.2838154	0.09	0.4611842509	0.6446726969
0.1763652	0.3156328	0.105	0.4813227422	0.6212064342
0.1974776	0.3461942	0.12	0.5002408825	0.6168350372
0.2134652	0.375645	0.135	0.5204425281	0.7550560866
0.2334652	0.403888	0.15	0.5411476689	0.7111579009
0.2485668	0.4311196	0.165	0.5590171297	0.6012483335
0.2675656	0.457576	0.18	0.5759222213	0.6767102446
0.2819648	0.4832344	0.195	0.5919181069	0.5701239701
0.3000652	0.5082738	0.21	0.6071150987	0.6437223597
0.3138984	0.5326654	0.225	0.6238645636	0.7296573993
0.3312004	0.5563802	0.24	0.6392812271	0.5705149372
0.3445048	0.5796034	0.255	0.6528166664	0.5951675979
0.361188	0.602345	0.27	0.668869986	0.8166345302
0.3740644	0.624472	0.285	0.6847467198	0.618420718
0.390102	0.6461654	0.3	0.6968292431	0.4955147032
0.4024884	0.6675292	0.315	0.7117831029	0.9044103837
0.417674	0.6883282	0.33	0.7267542228	0.53518959
0.4292784	0.7088134	0.345	0.737173249	0.4820351523
0.4445652	0.7290278	0.36	0.7508551143	0.8716399856
0.4560652	0.7487786	0.375	0.7655336953	0.6147383898
0.4705868	0.768221	0.39	0.7764021555	0.5032779268
0.4818648	0.7874214	0.405	0.7876149561	0.664697796
0.495882	0.8063156	0.42	0.8001386413	0.6609667181
0.5067812	0.8249194	0.435	0.8130346045	0.72541275
0.5206656	0.8432194	0.45	0.8231695853	0.3822354223
0.5311832	0.8613652	0.465	0.8321622217	0.6089180588
0.5446652	0.8792732	0.48	0.8437710657	0.6875800459
0.554982	0.8969236	0.495	0.8544283269	0.520014252
0.5680952	0.9143864	0.51	0.8665462466	0.8678410374
0.5782656	0.931549	0.525	0.8783503608	0.5077214323
0.5911116	0.948543	0.54	0.8933737653	1.260361948
0.6010652	0.965139	0.555	0.9110416678	0.8688140546
0.6135784	0.9814764	0.57	0.915255149	-0.3530060006
0.6232648	0.9979174	0.585	0.9163732773	0.4890230607
0.6353656	1.0142154	0.6	0.9275233967	0.8792576405
0.6447648	1.0302654	0.615	0.9399816885	0.673177482
0.6569652	1.046133	0.63	0.9503827681	0.6378058508
0.6663652	1.0618334	0.645	0.9590269548	0.4633363769
0.6780932	1.0774156	0.66	0.9699853551	0.9431916229
0.6870656	1.0927652	0.675	0.9830297701	0.7564507123
0.6985632	1.1079356	0.69	0.9919658955	0.4216494522
0.7075648	1.1230088	0.705	0.999952488	0.6380581699