$\begin{array}{c|c} \square \underline{PHYS115} \ \square \underline{PHYS121} \ \boxtimes \underline{PHYS123} \\ \square \underline{PHYS116} \ \square \underline{PHYS122} \ \square \underline{PHYS124} \end{array}$

Lab Cover Letter

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		ur TA) See your TA for detailed feedback. you need to improve this aspect of your work.
Pape	er Subtotals (points)	() Discussion & Conclusions (6)
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	Clarity of Presentation Format	Suggestions to reduce errors () Paper Total (60 points)
()	Abstract (4) Quantity or principle How measurement was made	(30 points for CME or EPF) () Notebook (10 points)
	Numerical Results Conclusion	Format (proper style, following directions) Apparatus (brief description of equipment, including sketches)
()	Intro & Theory (9) Basic principle	Data (including computer file names and manually recorded data) Experimental Technique (describing your
	Main equations to be used Apparatus What will be plotted Fitting parameters related	procedures; stating & justifying uncerts.) Analysis (results and errors)
— ()	Exp. Procedures (15)	() Worksheet(s)/Fill-in-the-Blank- Report (30 points) if applicable
	Description Stating and justifying uncertainties Data Record Quality of Lab Work	() Adjustments – late submissions, improper procedures, etc. – or bonus points for exceptional work.
()	Analysis & Error Analysis (20) Discussion Equations & Calculations	() Total Grade
	Presentation inc. Graphs, Tables Results Reported & Reasonable Underlined items addressed	Graded by(TA's initial)

Rotational Kinetic Energy Lab

Wolf S. Mermelstein

October 31, 2023

${\bf Abstract}$

 ${\rm content...}$

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1 Introduction

1.1 Moment of Inertia

Moment of inertia, often denoted by I, is a function of the specific geometry and mass distribution of an object. Moment of inertia is implicitly relative to the axis of rotation. Where R is the distance from the axis of rotation and M is the mass of the object, for a point mass the moment of inertia, I, is given by

$$I = MR^2 \tag{1}$$

The entire moment of inertia can be computed by thinking of a given object as a collection of tiny masses. As the masses' volumes shrink down to some small volume with some proportionately small mass, the can then be eventually said to be the differential dM. Integrating all the small point masses across the entire object implies that the entire moment of inertia of an object, I_{tot} , is given by

$$I = \sum_{k=1}^{\infty} \frac{m}{k} R^2 \tag{2}$$

$$= \int R^2 dM \tag{3}$$

This is conceptually helpful in understanding moment of inertia for arbitrary shapes, but is not practically useful for non simple (i.e. circles, squares, collections of discreet point masses) shapes, such as the mass-loaded, spoked wheel that we used in our experiment. As a result of such, it is often helpful to actually measure moment of inertia instead of attempting to compute it.

1.2 Conservation of Energy

The translational kinetic energy of an object in motion with mass M moving at speed v is given to be

$$K_T = \frac{1}{2}Mv^2 \tag{4}$$

Since we know that

$$\frac{\theta}{2\pi} = \frac{s}{2\pi R}$$

$$s = \theta R \tag{5}$$

and

$$\frac{d}{dt}s = \frac{d}{dt}\theta R$$

$$v = \omega r \tag{6}$$

We can then derive from equation 4 that the rotational kinetic energy, K_R , is

$$K_R = \frac{1}{2}M(\omega R)^2$$

$$= \frac{1}{2}(MR^2)\omega^2$$

$$= \frac{1}{2}I\omega^2$$
 (7)

where I is defined to be the moment of inertia about the axis of rotation.

For the mass in figure 2 descends downwards due to gravity, it begins to lose its gravitational potential energy, U_W . The total energy of the system is internally conserved, however a small amount of energy is lost due to friction. So, where ΔU_W is the change in the gravitational potential energy of the counterweight, K_T is the translational kinetic energy of the counterweight, and K_R is the rotational kinetic energy of the counterweight, we state that

$$\Delta U_W + K_T + K_R = W_f \tag{8}$$

Which, using equations 7 and 4, implies that

$$\Delta U_W + (\frac{1}{2}Mv^2) + (\frac{1}{2}I\omega^2) = W_f$$
 (9)

 ΔU_W should be negative, and K_T & K_R positive because the mass is falling, and, thus, losing gravitational kinetic energy, whilst simultaneously proportionately gaining kinetic energy.

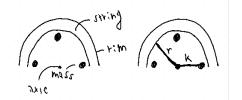
Using the fact that gravitational potential energy for an object at height h of mass m in an environment where gravity can be approximated to g is given to be

$$U_G = (M \cdot g \cdot h) \tag{10}$$

Plugging this in, and renaming h be y, we get the final equation

$$-(M \cdot g \cdot y) + (\frac{1}{2}Mv^2) + (\frac{1}{2}I\omega^2) = W_f$$
 (11)

Figure 1: Visual representation of k and r



1.3 Working Equation

Since for our specific experiment we used paperclips attached to the counterweight to cancel out friction, we can instead rewrite equation 11 to be

$$-(M \cdot g \cdot y) + (\frac{1}{2}Mv^2) + (\frac{1}{2}I\omega^2) = 0$$
 (12)

Carefully noting that we have discluded the energy of the moving paperclips, as it is negligible in comparison to the other energies of the system. And, to further simplify things, we will define y to be vertically positive, so as to make the equation into

$$-(M \cdot g \cdot y) + (\frac{1}{2}Mv^2) + (\frac{1}{2}I\omega^2) = 0$$
 (13)

Then, using relationship 6, we plug in $\frac{v}{r}$ for ω , resulting in the equation

$$-(M \cdot g \cdot y) + (\frac{1}{2}Mv^2) + (\frac{1}{2}I(\frac{v}{r})^2) = 0$$
$$= \frac{1}{2} \cdot v^2 \cdot (M + \frac{I}{r^2})$$
(14)

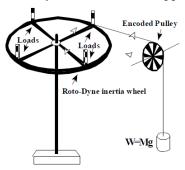
Or, as will be used more frequently for our computations, the equivalent equation in the form

$$gy = \frac{1}{2}(1 + \frac{I}{Mr^2}) \cdot v^2 \tag{15}$$

2 Procedure

Before conducting our experiment, we took measurements of various parts of our setup. First we obtained d from the lab manual, twice the distance from the axle of the wheel to the string, and then we measured k, the distance from the axle of the wheel to the masses. We determined r to be half the diameter. For r we used the provided uncertainty of ± 0.002 m, whereas for k we measured very carefully and chose the uncertainty to be 0.001m

Figure 2: Roto-Dyne Inertia Wheel Apparatus [2]



$$d = 0.200 \pm 0.002 \text{ m}$$

 $r = 0.100 \pm 0.002 \text{ m}$
 $k = 0.073 \pm 0.001 \text{ m}$

We used a counterweight with a given mass of 0.06kg to provide a torque to spin our Roto-Dyne wheel, as can be seen in figure 2. To account for friction, we incrementally added paperclips to the bottom of the counterweight. We continued to add paperclips up until the mass would fall at a constant speed to counteract the force of friction, using $Log-ger\ Pro^{\mathbb{T}M}$ software and an encoded pulley to monitor acceleration and velocity. Let M_c be the mass of the counterweight and M_p be the mass of the paperclips opposing friction, not used in our computations but still important to the experimental design.

$$M_p = 0.0015 \pm 0.0001 \text{ kg}$$
 $M_c = 0.06 \text{ kg}$

Also, we were provided with the mass of the Roto-Dyne wheel, M_R and the mass loads, M_L .

$$M_R = 1.5 \text{ kg}$$

 $M_L = .225 \pm .002 \text{ kg}$

For the encoded pulley that we used to measure velocity and length of unrolled string it was given that the gaps between intervals of measurement, Δs , was

$$\Delta s = 0.015 \text{ m}$$

- 3 Results
- 4 Analysis
- 5 Conclusion

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