

P102, 12.

证明半无界问题

$$\begin{cases} u_{tt} - a^2 u_{xx} = f(u, t) & (x, t) \in (0, \infty) \times (0, \infty) \\ u|_{t=0} = \varphi(x), \quad u_t|_{t=0} = \psi(x) & 0 \leq x < \infty \\ u|_{x=0} = \mu(t). & t \geq 0. \end{cases}$$

解唯一.

Proof: 只需证

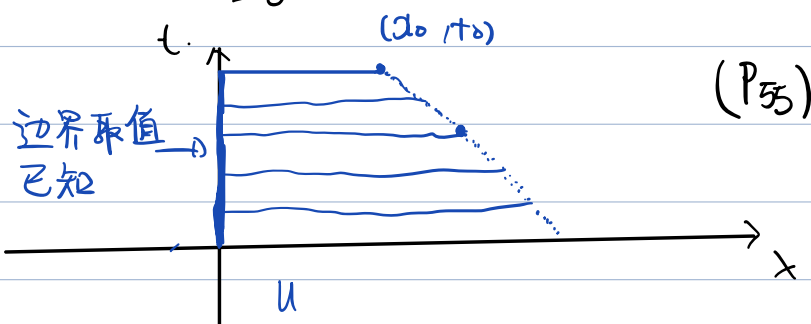
$$\begin{cases} u_{tt} - a^2 u_{xx} = 0 \\ u_{t=0} = u_t|_{t=0} = 0 \\ u|_{x=0} = 0 \end{cases}$$



只有零解.

任意给定 $(x_0, t_0) \in (0, +\infty) \times (0, +\infty)$. 定义能量

$$E(t) = \frac{1}{2} \int_0^{x_0 + a(t_0 - t)} u_t^2 + a^2 u_x^2 dx$$



① 对 t 求导

$$\begin{aligned} \frac{dE}{dt} &= -\frac{a}{2} (u_t^2 + a^2 u_x^2) \Big|_{x_0 + a(t_0 - t)} \\ &\quad + \int_0^{x_0 + a(t_0 - t)} (u_t u_{tt} + a^2 u_x u_{xt}) dx \\ &= -\frac{a}{2} (u_t^2 + a^2 u_x^2) \Big|_{x_0 + a(t_0 - t)}. \end{aligned}$$

$$+ \int_0^{x_0+a(t_0-t)} (u_t u_{tt} - a^2 u_t u_{xx}) dx$$

$$+ a^2 u_x u_t \Big|_0^{x_0+a(t_0-t)} \quad \textcircled{0}$$

$$u(0, \pm) = 0 \quad u_t(0, t) = \lim_{t \rightarrow 0^+} \frac{u(0, t) - u(0, 0)}{t} = 0$$

$$= \int_0^{x_0+a(t_0-t)} u_t (u_{tt} - a^2 u_{xx}) dx$$

$$+ \left(a^2 u_x u_t - \frac{a}{2} (u_t^2 + a^2 u_x^2) \right) \Big|_0^{x_0+a(t_0-t)}$$

$$[a^2 u_x u_t = \sqrt{a} u_t \cdot \sqrt{a} u_x \leq \frac{1}{2} (a u_t^2 + a^3 u_x^2)]$$

$$\leq 0.$$

$$\text{即 } \frac{dE}{dt} \leq 0.$$

② 对 $\frac{dE}{dt}$ 积分

$$0 \geq \int_0^\tau \frac{dE}{dt} dt = E(\tau) - E(0) \quad 0 \leq \tau \leq t_0$$

$$\text{即 } 0 \leq E(\tau) \leq E(0) = \frac{1}{2} \int_0^{x_0+a t_0} u_t^2(x, 0) + a^2 u_x^2(x, 0) dx$$

$$= 0.$$

$$\text{故 } E(\tau) = 0 \quad \forall \tau \in [0, t_0]$$

$$u_t = u_x = 0 \rightarrow u \text{ 在 } \text{阴影区域} \text{ 上是常值, 但 } u|_{x=0} = 0$$

故 $u = 0$ on 阴影区域 , 但 (x_0, t_0) 是任意的, 故我们可能在 $(0, \infty) \times [0, \infty)$ 上 $u = 0$.

类似于 (x_0, t_0) .

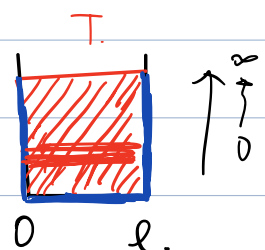


P83: 混合问题能量不能式

设 $u \in C^1(\bar{Q}_T) \cap C^2(Q_T)$. $Q_T = (0, l) \times (0, T)$.

是下述 PDE 的解

$$\begin{cases} u_{tt} - a^2 u_{xx} = f(x, t), & 0 < x < l, t > 0. \\ u(0, t) = u(l, t) = 0, & t > 0. \end{cases}$$



$$\begin{aligned} u(x, 0) &= \varphi(x) & 0 \leq x \leq l. \\ u_t(x, 0) &= \psi(x) & 0 \leq x \leq l. \end{aligned}$$

① 定义能量.

$$E(t) = \frac{1}{2} \int_0^l u_t^2 + a^2 u_x^2 dx \quad 0 \leq t \leq T.$$

② 求导

$$\begin{aligned} \frac{dE}{dt} &= \int_0^l u_t u_{tt} + a^2 u_x u_{xt} dx \\ &= \int_0^l \underline{u_t u_{tt}} - a^2 u_{xx} u_t dx + (a^2 u_x u_t) \Big|_0^l. \\ &= (a^2 u_x u_t) \Big|_0^l + \int_0^l u_t f(x, t) dx. \end{aligned}$$

③ 对 $\frac{dE}{dt}$ 积分.

$$\begin{aligned} E(T) - E(0) &= \int_0^T \int_0^l \boxed{u_t f(x, t)} dx dt + \int_0^T (a^2 u_x u_t) \Big|_0^l dt \\ &\leq \frac{1}{2} \int_0^T \int_0^l u_t^2 + a^2 u_x^2 dx dt + \frac{1}{2} \int_0^T \int_0^l f^2 dx dt \\ &\quad + \underbrace{\int_0^T a^2 u_x u_t \Big|_0^l dt}_{\sim 0}. \end{aligned}$$

$$u(0, t) = u(l, t) = 0.$$

$$\Omega(t) = \int_0^t E(\tau) d\tau.$$

$$\Rightarrow \frac{d\Omega}{dt}(T) \leq \underline{E(0) + \frac{1}{2} \int_0^T \int_0^l f^2 dx dt} + \Omega(T). \quad (\times)$$

③ 由 Gronwall inequality:

$$\begin{aligned} \frac{1}{2} \int_0^T \int_0^l u_t^2 + a^2 u_x^2 dx dt + \Omega(T) &\leq M (E(0) + \frac{1}{2} \int_0^T \int_0^l f^2 dx dt). \\ &= M \left(\frac{1}{2} \int_0^l \underline{u_t^2(x, 0) + a^2 u_x^2(x, 0)} dx + (\quad) \right) \\ &= \tilde{M} \left(\int_0^l \varphi^2(x) + a^2 (\varphi')^2 dx + \int_0^T \int_0^l f^2 dx dt \right). \end{aligned}$$

$$\begin{aligned} \sim \int_0^T \int_0^l u_t^2 + a^2 u_x^2 dx dt &\leq \\ \Omega(t) &M \left(\int_0^l \varphi^2 + a^2 (\varphi')^2 dx + \int_0^T \int_0^l f^2 dx dt \right) \end{aligned}$$

由(*) $\frac{dQ}{dt}(T) \leq M (\int_0^l \varphi^2 + a^2 (\varphi')^2 dx + \int_0^T \int_0^l f^2 dx dt).$

热方程能量不等式 (混合问题)

$Q_T = (0, l) \times [0, T]$, $u \in C_{x,t}^{1,0}(\bar{Q}_T) \cap C_{x,t}^{2,1}(Q_T)$. 是下述 PDE 的解:

$$\begin{cases} u_t - a^2 u_{xx} = f & (x,t) \in Q_T \\ u(x,0) = \varphi(x) & 0 \leq x \leq l. \\ u(0,t) = u(l,t) = 0. & 0 \leq t \leq T \end{cases}$$

Proof: 对 (*) 两侧同乘 u 并积分.

$$\int_0^T \int_0^l u_t u - a^2 u_{xx} u dx dt = \int_0^T \int_0^l f u dx dt.$$

$$\leq \frac{1}{2} \int_0^T \int_0^l f^2 dx dt + \frac{1}{2} \int_0^T \int_0^l u^2 dx dt.$$

$$\int_0^l u_{xx} u dx = - \int_0^l u_x^2 dx + \frac{u u_x}{2} \Big|_0^l (=0).$$

$u(0,t) = u(l,t) = 0$

$$= - \int_0^l u_x^2(x,t) dx$$

$$\int_0^T u_t u dt = \int_0^T u du = \frac{1}{2} u^2 \Big|_0^T.$$

$$= \frac{1}{2} (u^2(x,T) - u^2(x,0)).$$

$$= \frac{1}{2} (u^2(x,T) - \varphi^2(x))$$

$$= \int_0^l \int_0^T u_t u dt dx - a^2 \int_0^T \int_0^l u_{xx} u dx dt.$$

$$= \frac{1}{2} \int_0^l u^2(x,T) - \varphi^2(x) dx + a^2 \int_0^T \int_0^l u_x^2(x,t) dx dt.$$

$$\leq \frac{1}{2} \int_0^T \int_0^l f^2 dx dt + \frac{1}{2} \int_0^T \int_0^l u^2 dx dt.$$

$$\Rightarrow \int_0^l u^2(x,T) dx + a^2 \int_0^T \int_0^l u_x^2(x,t) dx dt \quad (\star)$$

$$\leq \underbrace{\int_0^T \int_0^L \varphi^2 dx}_{(H^1)} + \int_0^T \int_0^L f^2 dx dt + \int_0^T \int_0^L u^2 dx dt.$$

$$\Omega(T) = \int_0^T \int_0^L u^2 dx dt, \text{ 则 } \frac{d\Omega}{dt} = \int_0^L u^2(x, T) dx.$$

$$\frac{d}{dt}\Omega(T) \leq \frac{d}{dt}\Omega(T) + 2a^2 \int_0^T \int_0^L u_x^2(x, t) dx dt.$$

$$\leq \underline{F(T)} + \Omega(T).$$

由 Gronwall Inequality

$$\Omega(T) \leq MF(T)$$

代回 (*)

$$\boxed{\sup_{0 \leq t \leq T} \int_0^L u^2(x, t) dx} + a^2 \int_0^T \int_0^L u_x^2(x, t) dx dt \leq F(T) + \Omega(T) \leq (M+1) F(T).$$

$$= \tilde{M} \left(\int_0^L \varphi^2 dx + \int_0^T \int_0^L f^2 dx dt \right),$$

$$\sup_{0 \leq t \leq T} \int_0^L u^2(x, t) dx + 2a^2 \int_0^T \int_0^L u_x^2(x, t) dx dt \leq \tilde{M} \left(\int_0^L \varphi^2 dx + \int_0^T \int_0^L f^2 dx dt \right)$$

Remark: 若 $\varphi = 0, f = 0$. 则

$$\sup_{0 \leq t \leq T} \int_0^L u^2(x, t) dx = 0. \quad \underline{\int_0^T \int_0^L u_x^2(x, t) dx dt = 0}$$

$$\Rightarrow u = 0 \quad \forall (x, t) \in Q_T$$

P158 : 得到 $\frac{\partial u}{\partial t}$ 的一个上界估计.

$$u_t - a^2 u_{xx} = f \quad \text{两侧乘 } \frac{\partial u}{\partial t} \text{ 得到}$$

位势方程:

$$\begin{cases} -\Delta u + c(x)u = f(x) \star & x \in \Omega. \quad \star \\ \underline{u|_{\partial\Omega} = 0.} \end{cases}$$

$c(x) \geq c_0 > 0$, $u \in C^2(\Omega) \cap C^1(\bar{\Omega})$ 为 (4) 的解, 则有估计

$$\int_{\Omega} |\nabla u|^2 dx + \frac{c_0}{2} \int_{\Omega} |u|^2 dx \leq M \int_{\Omega} |f|^2 dx.$$

Proof: 在 $-\Delta u + c(x)u = f$ 两侧乘 u 积分

$$\int_{\Omega} -u \Delta u + c(x)u^2 dx = \int_{\Omega} f u dx$$

$$\begin{aligned} \int_{\Omega} u \Delta u dx &= - \int_{\Omega} |\nabla u|^2 dx + \int_{\partial \Omega} u \frac{\partial u}{\partial n} dS \\ &= - \int_{\Omega} |\nabla u|^2 dx. \end{aligned}$$

$$\begin{aligned} \int_{\Omega} -u \Delta u + c(x)u^2 dx \\ = \int_{\Omega} |\nabla u|^2 + c(x)u^2 dx \geq \int_{\Omega} |\nabla u|^2 dx + c_0 \int_{\Omega} u^2 dx. \end{aligned}$$

$$\Rightarrow \int_{\Omega} |\nabla u|^2 dx + c_0 \int_{\Omega} u^2 dx \leq \int_{\Omega} f u dx.$$

$$\begin{aligned} f u &= \varepsilon u \cdot \frac{1}{\varepsilon} f \quad (\varepsilon > 0) \\ &\leq \frac{1}{2} \varepsilon^2 u^2 + \frac{1}{2\varepsilon^2} f^2. \end{aligned}$$

$$\begin{aligned} \Rightarrow \int_{\Omega} |\nabla u|^2 dx + c_0 \int_{\Omega} u^2 dx &\leq \int_{\Omega} f u dx \\ &\leq \frac{1}{2} \varepsilon^2 \int_{\Omega} u^2 dx + \frac{1}{2\varepsilon^2} \int_{\Omega} f^2 dx \end{aligned}$$

$$\Rightarrow \int_{\Omega} |\nabla u|^2 dx + (c_0 - \frac{1}{2} \varepsilon^2) \int_{\Omega} u^2 dx \leq \frac{1}{2\varepsilon^2} \int_{\Omega} f^2 dx.$$

取 ε 满足 $c_0 - \frac{1}{2} \varepsilon^2 = \frac{c_0}{2} \iff \varepsilon^2 = c_0.$

$$\Rightarrow \int_{\Omega} |\nabla u|^2 dx + \frac{c_0}{2} \int_{\Omega} u^2 dx \leq \frac{1}{2c_0} \int_{\Omega} f^2 dx.$$

$$\begin{cases} -\Delta u = f \\ u|_{\partial \Omega} = 0 \end{cases} \quad \star$$

Lemma (Poincaré Inequality) For $u \in C_0^\infty(\Omega)$, \mathbb{R}^d .

$$\|u\|_{L^2(\Omega)} \leq 2d \|\nabla u\|_{L^2(\Omega)}$$

$$d = \text{diam}(\Omega) = \sup_{x, y \in \Omega} |x - y|.$$

$$\|u\|_{L^2(\Omega)} = \left(\int_{\Omega} |u|^2 dx \right)^{\frac{1}{2}}.$$

对 (★) 两侧乘 u 积分

$$\int_{\Omega} -\Delta u \cdot u \, dx = \int_{\Omega} f u \, dx \leq \frac{1}{2} \varepsilon^2 \int_{\Omega} u^2 dx + \frac{1}{2\varepsilon^2} \int_{\Omega} f^2 dx$$

$\int_{\Omega} |\nabla u|^2 dx$ }

$$\int_{\Omega} u^2 dx \leq 4d^2 \int_{\Omega} |\nabla u|^2 dx$$

$$\int_{\Omega} |\nabla u|^2 dx \leq \left(\frac{1}{2} \varepsilon^2 \cdot 4d^2 \right) \int_{\Omega} |\nabla u|^2 dx + \frac{1}{2\varepsilon^2} \int_{\Omega} f^2 dx$$

$$\text{取 } \frac{1}{2} \varepsilon^2 \cdot 4d^2 = \frac{1}{2}, \text{ 则 } \varepsilon = \frac{1}{2d}$$

$$\Rightarrow \int_{\Omega} |\nabla u|^2 dx \leq 4d^2 \int_{\Omega} f^2 dx$$

$$\subset \int_{\Omega} |u|^2 dx \leq$$

Remark: $f = 0 \Rightarrow u = 0$.

附注 3 (习题 13).

Fourier 变换.

$$\hat{f}(\xi) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} f(x) e^{-i\xi x} dx.$$

反演公式:

$f \in L^1(\mathbb{R}) \cap C^1(\mathbb{R})$, 则有

$$f(x) = \lim_{N \rightarrow \infty} \frac{1}{\sqrt{2\pi}} \int_{-N}^N \hat{f}(\lambda) e^{i\lambda x} d\lambda$$

$$\mathcal{F}: f \mapsto \hat{f}$$

$$\mathcal{F}^{-1}: \hat{f} \mapsto (\hat{f})^\vee$$

Fourier 变换性质.

① 线性性 $(\alpha f + \beta g)^\wedge = \alpha \hat{f} + \beta \hat{g}$

② 微商性:

$$\left(\frac{df}{dx}\right)^\wedge = i\lambda \hat{f}.$$

Proof:

$$\begin{aligned} \left(\frac{df}{dx}\right)^\wedge &= \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} \frac{df}{dx}(x) e^{-i\lambda x} dx = \frac{-1}{\sqrt{2\pi}} \int_{\mathbb{R}} f(x) \frac{d}{dx} e^{-i\lambda x} dx. \\ &= \frac{1}{\sqrt{2\pi}} i\lambda \int_{\mathbb{R}} f(x) e^{-i\lambda x} dx \\ &= i\lambda \hat{f} \end{aligned}$$

③ 卷积性质

$f, g \in L^1(\mathbb{R})$, 则

$$(f * g)^\wedge = \sqrt{2\pi} \hat{f} \hat{g}$$

其中 $f * g(x) = \int_{\mathbb{R}} f(x-t) g(t) dt.$

Proof: $(f * g)^\wedge(\lambda) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} f * g(x) e^{-i\lambda x} dx$

$$= \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} \left(\int_{\mathbb{R}} f(x-t) g(t) dt \right) e^{-i\lambda x} dx$$

$$= \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} \left(\int_{\mathbb{R}} f(x-t) g(t) dt \right) e^{-i\lambda(x-t)} \cdot e^{-i\lambda t} dx.$$

$$= \int_{\mathbb{R}} g(t) \left(\frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} f(x-t) e^{-i\lambda(x-t)} dx \right) e^{-i\lambda t} dt.$$

$$= \int_{\mathbb{R}} g(t) \hat{f}(\lambda) e^{-i\lambda t} dt.$$

$$= \sqrt{2\pi} \hat{f}(\lambda) \hat{g}(\lambda)$$

④ 伸缩性质

$f \in L^1(\mathbb{R}^n)$. 定义: ☆

$$\underline{f_t(x) = \frac{1}{t} f\left(\frac{x}{t}\right)}. \quad (t > 0)$$

$$\hat{f}_t(\lambda) = \hat{f}(\lambda t)$$

Proof: $\hat{f}_t(\lambda) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} \frac{1}{t} f\left(\frac{x}{t}\right) e^{-i\lambda x} dx$ $\frac{x}{t} \mapsto y$

$$= \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} f(y) e^{-i\lambda t y} dy.$$

$$= \hat{f}(\lambda t)$$

例:

$$f(x) = e^{-\frac{1}{2}x^2}. \quad \hat{f}(\lambda) = e^{-\frac{1}{2}\lambda^2}. \quad \text{Fourier 变换不动点}$$

Proof: $\hat{f}(\lambda) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-\frac{1}{2}x^2} e^{-i\lambda x} dx$

$$\begin{aligned} \frac{d}{d\lambda} \hat{f}(\lambda) &= \frac{1}{\sqrt{2\pi}} \frac{d}{d\lambda} \int_{\mathbb{R}} e^{-\frac{1}{2}x^2} e^{-i\lambda x} dx \\ &= \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-\frac{1}{2}x^2} -ix e^{-i\lambda x} dx \\ &= -i \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} x e^{-\frac{1}{2}x^2} e^{-i\lambda x} dx \\ &= i \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} \left(e^{-\frac{1}{2}x^2} \right)' e^{-i\lambda x} d\frac{1}{2}x^2 \\ &= i \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-i\lambda x} d e^{-\frac{1}{2}x^2} \\ &= -\lambda \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-\frac{1}{2}x^2} e^{-i\lambda x} dx \\ &= -\lambda \hat{f}(\lambda). \end{aligned}$$

$$\hat{f}(0) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-\frac{1}{2}x^2} dx$$

ODE: $\frac{d}{d\lambda} \hat{f}(\lambda) + \lambda \hat{f}(\lambda) = 0$

$$\hat{f}(0) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-\frac{1}{2}x^2} dx = 1 \quad \text{Gauss 积分}$$

$$g' + xg = 0 \quad \frac{d}{dx} g = -xg$$

$$\frac{dg}{g} = -x dx$$

$$d \log g = d(-\frac{1}{2}x^2)$$

$$g = C \exp(-\frac{1}{2}x^2)$$

$$\Rightarrow \hat{f}(\lambda) = C \exp(-\frac{1}{2}\lambda^2) = \exp(-\frac{1}{2}\lambda^2)$$

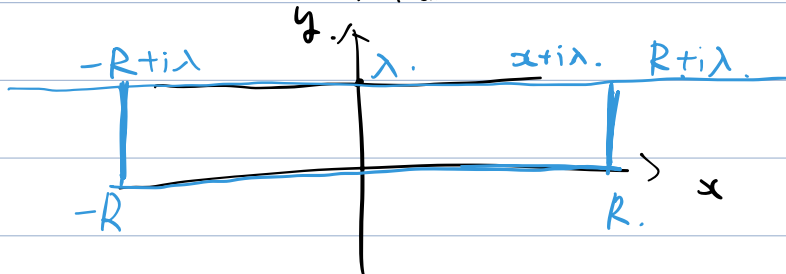
$$\hat{f}(\lambda) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-\frac{1}{2}x^2} e^{-i\lambda x} dx$$

$$= \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-\frac{1}{2}x^2 - i\lambda x} dx$$

$$\frac{1}{2}x^2 + i\lambda x = \frac{1}{2}(x+i\lambda)^2 + \frac{\lambda^2}{2}$$

$$= \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-\frac{1}{2}(x+i\lambda)^2} \cdot e^{-\frac{\lambda^2}{2}} dx$$

$$= e^{-\frac{\lambda^2}{2}} \cdot \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-\frac{1}{2}(x+i\lambda)^2} dx = 1$$



例: $f(x) = e^{-Ax^2}$

$$\frac{1}{2}x^2 \xrightarrow{x \mapsto \sqrt{\frac{A}{2}}y} Ay^2$$

$$f = ?$$

$$g(x) = e^{-\frac{1}{2}x^2}$$

$$\hat{g}(\lambda) = e^{-\frac{1}{2}\lambda^2}$$

$$\sqrt{2A} f(x) = \sqrt{2A} g(\sqrt{2A}x) \quad t = \sqrt{\frac{2}{A}}$$

$$\overset{||}{g_{\frac{1}{\sqrt{2A}}}}(x)$$

$$\text{例} \quad (g_{\frac{1}{\sqrt{2A}}})^\wedge(\lambda) = \hat{g}(\sqrt{\frac{1}{2A}}\lambda) = \exp(-\frac{1}{4A}\lambda^2)$$

↓

$$\Rightarrow \sqrt{2A} \hat{f}(\lambda) = \exp\left(-\frac{1}{4A}\lambda^2\right)$$

$$\Rightarrow \hat{f}(\lambda) = \frac{1}{\sqrt{2A}} \exp\left(-\frac{1}{4A}\lambda^2\right) \star$$

热方程:

$$\begin{cases} u_t - a^2 u_{xx} = f(x, t) \\ u(x, 0) = \varphi(x) \end{cases} \star$$

对 (\star) 关于 x 作 Fourier 变换 (t 看作常数).

$$\frac{\partial}{\partial t} \hat{u}(\lambda, t) = \frac{\partial}{\partial t} \hat{u}(\lambda, t)$$

$$\hat{u}_{xx}(\lambda, t) = (i\lambda)^2 \hat{u}(\lambda, t) = -\lambda^2 \hat{u}(\lambda, t).$$

$$\hat{f}(\lambda, t) \quad \hat{\varphi}(\lambda).$$

则 (\star) 变为:

$$\begin{cases} \frac{\partial}{\partial t} \hat{u}(\lambda, t) + a^2 \lambda^2 \hat{u}(\lambda, t) = \hat{f}(\lambda, t) \\ \hat{u}(\lambda, 0) = \hat{\varphi}(\lambda) \end{cases} \text{关于 } t \text{ 的 ODE.}$$

$$\frac{\partial}{\partial t} (\hat{u}(\lambda, t) e^{a^2 \lambda^2 t}) = \hat{f}(\lambda, t) e^{a^2 \lambda^2 t}$$

$$\hat{u}(\lambda, t) e^{a^2 \lambda^2 t} - \hat{\varphi}(\lambda) = \int_0^t \hat{f}(\lambda, \tau) e^{a^2 \lambda^2 \tau} d\tau.$$

$$\Rightarrow \hat{u}(\lambda, t) = \underbrace{e^{-a^2 \lambda^2 t}}_{\downarrow \mathcal{F}^{-1}} \hat{\varphi}(\lambda) + \int_0^t \hat{f}(\lambda, \tau) \underbrace{e^{-a^2 \lambda^2 (t-\tau)}}_{\uparrow \mathcal{F}} d\tau.$$

$$u(x, t).$$

$$\frac{1}{\sqrt{2\pi}} e^{-a^2 \lambda^2 (t-\tau)} = \hat{g}_{t-\tau}(\lambda)$$

$$\text{令 } \hat{g}_t(\lambda) = \frac{1}{\sqrt{2\pi}} e^{-a^2 \lambda^2 t}$$

$$\sqrt{2\pi} \hat{g}_t(\lambda) \cdot \hat{\varphi}(\lambda) = e^{-\alpha^2 \lambda^2} \hat{\varphi}(\lambda)$$

$$(g * \varphi)^\wedge(\lambda).$$

$$\Rightarrow \hat{u}(\lambda, t) = \sqrt{2\pi} \hat{g}_t(\lambda) \cdot \hat{\varphi}(\lambda) + \int_0^t \hat{f}(\lambda, \tau) \hat{g}_{t-\tau}(\lambda) \sqrt{2\pi} d\tau.$$

$$f(x, t) =: f_t(x)$$

$$\hat{f}(\lambda, t) = \hat{f}_t(\lambda).$$

$$= \sqrt{2\pi} \hat{g}_t(\lambda) \hat{\varphi}(\lambda) + \int_0^t \hat{f}_\tau(\lambda) \hat{g}_{t-\tau}(\lambda) \sqrt{2\pi} d\tau.$$

$$= (g * \varphi)^\wedge(\lambda) + \int_0^t (f_\tau * g_{t-\tau})^\wedge(\lambda) d\tau.$$

$$\Rightarrow u(x, t) = (g * \varphi)(x) + \int_0^t f_\tau * g_{t-\tau}(x) d\tau.$$

$$= \int_{\mathbb{R}^n} g_t(x-y) \varphi(y) dy$$

$$+ \int_0^t \left(\int_{\mathbb{R}^n} f_\tau(x-y) g_{t-\tau}(y) dy \right) d\tau$$

$$= \int_{\mathbb{R}^n} g_t(x-y) \varphi(y) dy + \int_0^t \left(\int_{\mathbb{R}^n} f(x-y, \tau) g_{t-\tau}(y) dy \right) d\tau.$$

$$g_t(x) \rightarrow \delta(x) \quad (t \rightarrow 0).$$