

# HURRICANE ISABEL (2003): NEW INSIGHTS INTO THE PHYSICS OF INTENSE STORMS. PART I

## Mean Vortex Structure and Maximum Intensity Estimates

BY MICHAEL T. MONTGOMERY, MICHAEL M. BELL, SIM D. ABERSON, AND MICHAEL L. BLACK

An unprecedented observational analysis of a category-5 hurricane suggests a mechanism by which high-entropy air inside the low-level eye can sustain the storm at an intensity above the currently formulated upper bound for tangential wind speed.

*“From a science point of view, I love how the atmosphere likes to throw curve balls when people think they understand ‘everything.’”*

—LANCE BOSART

**H**urricane Isabel was at or near category 5 on the Saffir–Simpson scale (surface wind speed  $> 67 \text{ m s}^{-1}$ ) from 12 to 14 September 2003. On 13 September, the strongest known horizontal wind ( $107 \text{ m s}^{-1}$ ) recorded in a hurricane was obtained by a National Center for Atmospheric Research (NCAR) GPS dropwindsonde at 1752 UTC (Aberson et al. 2006, hereafter Part II; Hock and Franklin 1999).<sup>1</sup> Multiple aircraft missions by the National Oceanic

and Atmospheric Administration (NOAA) and U.S. Air Force (USAF) Hurricane Hunters obtained a uniquely comprehensive observational dataset during this period. Multiple dropwindsondes, in situ flight level data, airborne Doppler radar data, and high-resolution satellite imagery provide an unprecedented view of the kinematic and thermodynamic structure of Isabel over 3 days. A synthesis of flight-level and dropwindsonde data during 12–14 September along with a detailed description of the dataset and

<sup>1</sup> A report of  $117 \text{ m s}^{-1}$  from a dropwindsonde released in Hurricane Katrina (2005) is unconfirmed because of an apparent loss of the raw data.

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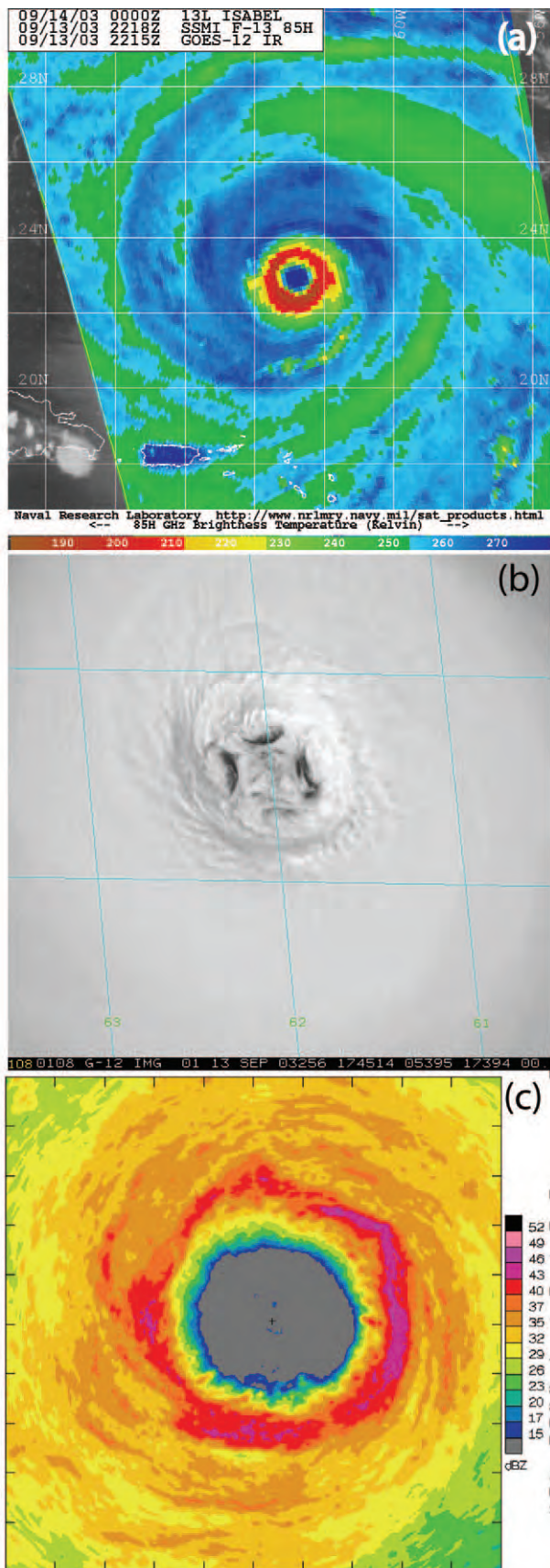
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analysis methodology is reported elsewhere (Bell and Montgomery 2006, manuscript submitted to *Mon. Wea. Rev.*); the focus of this two-part study is the analysis for 13 September.

A nearly circular core of deep eyewall convection at 2218 UTC 13 September is represented in Fig. 1a by the cold (red) 85-GHz brightness temperature. Nonaxisymmetric structures inside the eye with cloud bands connecting to the eyewall are also apparent in the low-level stratus clouds at 1745 UTC (Fig. 1b), along with a pentagon-shaped region of high reflectivity seen in a NOAA P-3 lower-fuselage radar image at approximately the same time (Fig. 1c). These observations, in conjunction with Geostationary Operational Environmental Satellite (GOES) super-rapid-scan animations (not shown), corroborate the existence of coherent mesovortices in the vicinity of the eyewall that have been predicted by high-resolution numerical simulations (Schubert et al. 1999, hereafter S99; Kossin and Schubert 2001; Persing and Montgomery 2003, hereafter PM03; Kossin and Schubert 2004; Braun et al. 2006) and liquid water laboratory experiments (Montgomery et al. 2002, and references therein). The mesovortices are believed to be the result of a combined barotropic–baroclinic instability associated with an annulus of high potential vorticity near and within the eyewall cloud that is generated by latent heating and vortex tube stretching (S99; Nolan and Montgomery 2002).

In addition to contributing to localized extreme winds, evidence suggests that eyewall mesovortices in the lower troposphere can actually enhance the storm-scale energetics. In simple models that neglect diabatic processes and/or the secondary circulation, eyewall wave and vortex structures exert a cyclonic torque on the eye, thereby spinning up the eye at the expense of the maximum tangential winds (Rotunno 1978; Smith 1980; Emanuel 1997; S99; Kossin and Schubert 2001). However, in conjunction with a secondary

**FIG. 1. Satellite and airborne radar imagery of Hurricane Isabel on 13 September. (a) Special Sensor Microwave Imager (SSM/I) 85-GHz brightness temperature at 2218 UTC (courtesy of NRL/Monterey); (b) visible image at 1745 UTC from GOES super-rapid-scan operations; and (c) radar reflectivity (dBZ) from lower-fuselage (5 cm) radar onboard NOAA aircraft while flying inside the eye at ~2 km altitude. The time of the radar image (1747 UTC) is 2 min after the visible image shown in (b). In (a), (b), and (c) the horizontal scales of the images are approximately 1400, 300, and 180 km, respectively. The mesovortices within the eye in (b) and the pentagonal shape of the high reflectivity in the eyewall in (c) are particularly striking features.**

circulation driven by latent heating, surface friction, and eddy processes, the breakdown of the eyewall and the corresponding (vorticity) mixing associated with the enstrophy cascade permits low-level inflowing air parcels to penetrate the highly impermeable eyewall (Rotunno 1984; S99). Parcels that reach the low-level eye and spend sufficient time inside it can increase their moist entropy above eyewall values through interaction with the ocean and because of the reduced surface pressure relative to that in the eyewall. From an energetics viewpoint, then, the net effect of the mean secondary and eddy circulation within the low-level eye is to transport and stir high-entropy air from the eye to the eyewall (S99; Braun 2002; PM03; Braun et al. 2006; Cram et al. 2006, hereafter C06; Eastin et al. 2005). In this way, the low-level eye can provide additional power to the hurricane engine relative to that obtained from the ocean underneath the eyewall, which is where current theory assumes all of the heat uptake occurs (cf. Emanuel 1997).

For the first time, observations suggest that this thermodynamic boost exceeds the weakening tendency associated with the breakdown of the eyewall. By utilizing this additional power, the hurricane can, in principle, obtain a “superintense” state with azimuthally averaged tangential winds that are faster than previously predicted upper bounds for the mean tangential wind at the swirling boundary layer top, as predicted by Emanuel and colleagues (Emanuel 1986, hereafter E86; Rotunno and Emanuel 1987, hereafter RE87; Emanuel 1988, 1991, 1995, 1997, hereafter E88, E91, E95, and E97, respectively; Bister and Emanuel 1998, hereafter BE98; Emanuel et al. 2004, hereafter E04). This theoretical upper bound, or maximum potential intensity (MPI), for given environmental conditions is referred to hereafter as E-MPI. The current work is believed to be the first study using sufficiently detailed in situ observations to test E-MPI theory and to suggest the existence of superintense storms in nature that were previously predicted using high-resolution computer simulations (PM03).

A detailed analysis of the dropwindsonde and flight-level data as they relate to the estimated azimuthal mean structure of Isabel on 13 September is provided herein. A record horizontal wind speed measured on the same day is presented in Part II, along with a discussion of similar observations from other intense storms. Together, these studies suggest new insight into the thermomechanics of intense hurricanes. The observations are consistent with high-resolution numerical simulations that produce superintense storms. An outline of the remaining

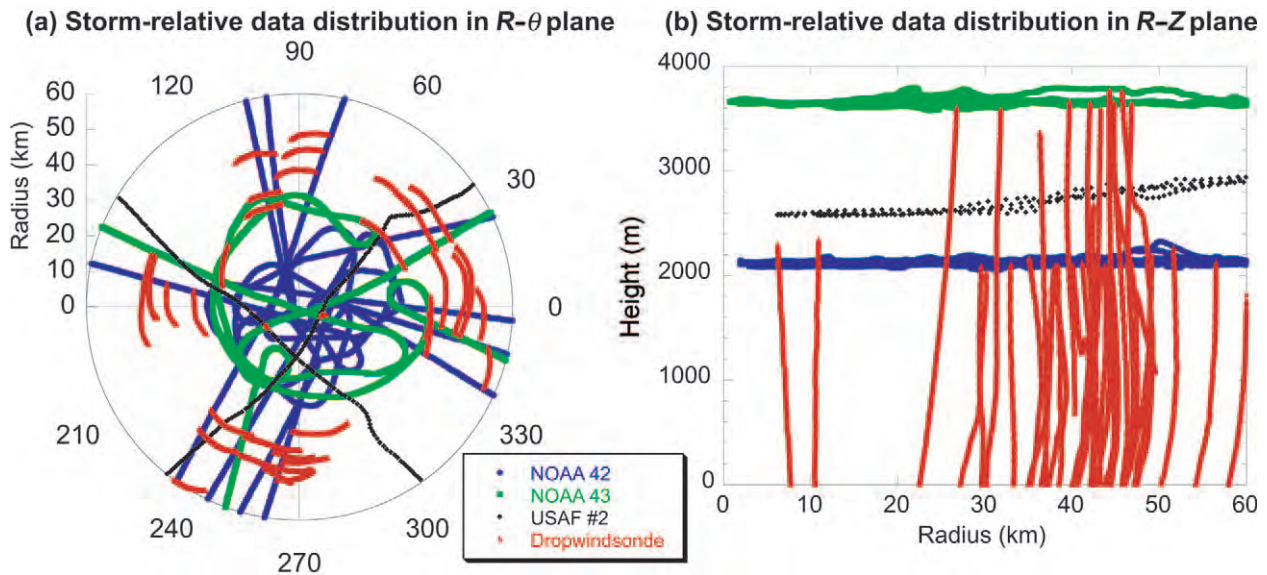
sections of Part I follows: “Estimated mean vortex structure” provides the methodology and analysis of the estimated mean inner-core structure of Isabel; “MPI estimates” presents an analysis of environmental conditions and a priori E-MPI estimates; “Conclusions” are given last.

## ESTIMATED MEAN VORTEX STRUCTURE.

**Data overview.** A two-dimensional analysis of the axisymmetric kinematic and thermodynamic structures of Isabel in the radial vertical plane was performed using in situ flight-level data and 35 eyewall and eye dropwindsondes released from three aircraft (two NOAA P-3s and a USAF C-130) flying at different altitudes. The sweeping dropwindsonde trajectories and multiple aircraft penetrations into the hurricane (Fig. 2a) show the approximately uniform azimuthal data coverage obtained from 1600 to 2300 UTC 13 September. The high radial and vertical resolutions of the ~30,700 data points (approximately half of which are from the dropwindsondes) recorded during this time period are depicted in Fig. 2b. Although the dense spatial resolution allows for high confidence in retrieved axisymmetric features in the eyewall region, greater uncertainty in the axisymmetric component exists inside the 30-km radius where the sampling was limited. In particular, the limited observations inside this radius preclude a meaningful separation of the kinematic and thermodynamic quantities into azimuthal mean (vortex) and asymmetric (eddy) components. Although these innermost data are included in the composite, the estimated axisymmetric structure is likely quantitatively inaccurate there due to large-amplitude mesovortex structures evident in the GOES super-rapid-scan animations and video recordings taken aboard the NOAA P-3 while flying inside the eye of Isabel.

**Compositing technique.** The dropwindsonde and in situ data were composited relative to a moving (~7 m s<sup>-1</sup>) storm center using the Barnes objective analysis technique (Barnes 1973; Koch et al. 1983). Although temporal evolution is not resolved by this method, the data and quality control parameters suggest that Isabel was near steady state during the time period of the analysis. Whereas the Barnes objective analysis has been used extensively with synoptic and satellite data, its application for inferring radius–height mean hurricane structure is believed to be novel. Different grid sizes were used in the vertical (250 m) and radial (2500 m) directions, with weighting functions allowing for the resolving of features on scales of greater than 1





**FIG. 2.** Dropwindsonde locations and trajectories and aircraft flight tracks relative to storm center from 1600 to 2300 UTC 13 Sep. Storm-relative data distribution in the (a) radius–azimuthal ( $R$ – $\theta$ ) plane and (b) radius–height ( $R$ – $Z$ ) plane, showing the NOAA P-3 (42 in blue, 43 in green) and USAF C-130 (in black) flight tracks, and dropwindsonde trajectories (in red). The dropwindsondes in (a) move cyclonically (counterclockwise).

and 10 km, respectively, to account for the relatively shallow observations indicated in Fig. 2. All data were treated as azimuthally independent, combining the primarily Eulerian aircraft data with the Lagrangian dropwindsonde data by their location in the radius–height plane. Storm-relative tangential and radial wind velocities were calculated by subtracting the storm motion at appropriate times derived from aircraft center fixes. Vertical velocity was measured at flight level by the aircraft inertial navigation system, and was also calculated along the dropwindsonde trajectory by removing the terminal fall speed from the measured dropwindsonde motion. The analyzed secondary circulation does not satisfy the continuity equation everywhere, because of uncertainty in the analysis of the axisymmetric radial wind arising from undersampling (e.g., near the 15-km radius from storm center; see Fig. 3), uncertainty in the precise location of the circulation center, and the estimated dropwindsonde terminal fall speed. Asymmetries not considered here may also contribute to the disagreement between measured vertical velocities and those calculated from mass continuity. The resulting transverse circulation components are nevertheless self-consistent and are believed to capture the zeroth-order features of the axisymmetric transverse circulation.

*Derived mean structure and evidence supporting the proposed “superintensity” mechanism.* The kinematic

and thermodynamic structure of Isabel derived from GPS dropwindsonde and flight-level data is displayed in the radius–height composites of storm-relative tangential and radial wind velocities (Fig. 3a) and equivalent potential temperature, specific absolute vertical angular momentum, and transverse secondary circulation vectors (Fig. 3b). The specific absolute vertical angular momentum (absolute circulation) is defined as  $M = rv + 1/2fr^2$ , where  $r$  is the radius from the vortex center,  $v$  is the storm-relative tangential wind, and  $f$  is the Coriolis parameter.

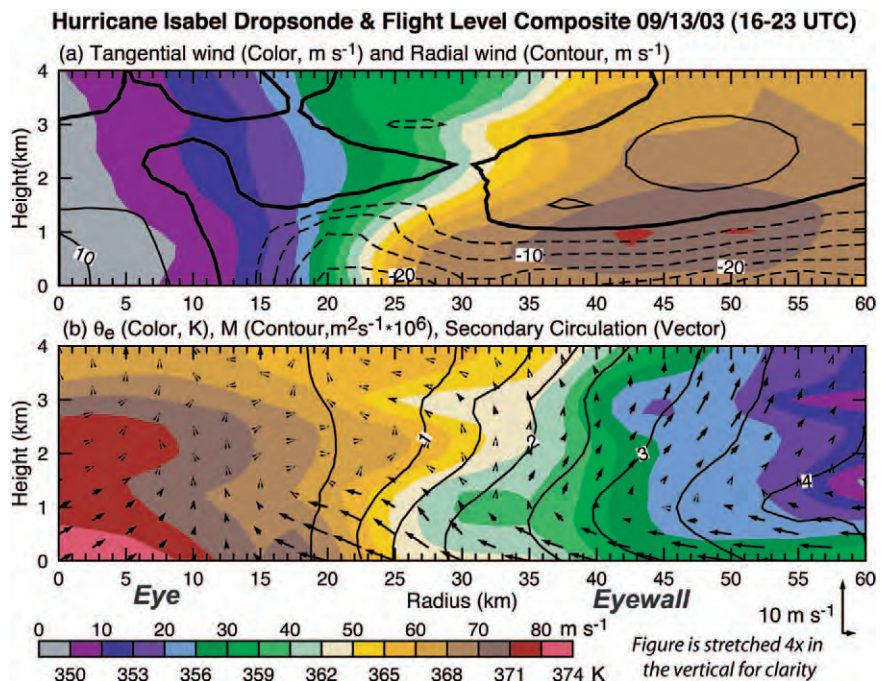
The maximum azimuthally averaged tangential wind speed ( $76 \text{ m s}^{-1}$ ) is located approximately 42-km radius from the center near 1000 m above the surface. Because this is an azimuthal average, it is likely comparable to a sustained wind at this level. The strong radial gradient of tangential wind along the inner edge of the eyewall is consistent with local Kelvin–Helmholtz (shear) instabilities and associated lateral mixing across the eyewall interface implied by the satellite and radar imagery in Figs. 1b and 1c (e.g., S99; Montgomery et al. 2002). An extreme wind speed maximum ( $107 \text{ m s}^{-1}$ ) that would likely have been recorded by an anemometer as a gust was observed by a dropwindsonde near 40-km radius at 1.4-km altitude (see Part II, for details). Near this radius, the tangential velocity was highly variable azimuthally, as would be expected with the presence of mesovortices in the eye–eyewall interface, as seen in satellite imagery. The lowest-level (0–250 m) radial

inflow of  $20 \text{ m s}^{-1}$  is located 25-km radius from the center, suggesting significant penetration of air from under the eyewall into the eye. Strong inflowing air that breaches the eyewall is modified by the sea-to-air latent heat flux inside the eye where widespread downdrafts that ordinarily tend to limit the enthalpy gain outside of the eyewall do not exist (e.g., E86; Eastin et al. 2005). Despite the rapid decrease of mean tangential wind speed inside the eye, the low pressure and significant inflow and convergence near the surface support nonzero and persistent sea-to-air latent heat flux within the eye, thus maintaining a reservoir of high- $\theta_e$  air that has been found in previous studies (LeeJoice 2000; Braun 2002; PM03; Eastin et al. 2005). High-entropy air returns to the eyewall in the outflow above 1 km, indicated by the outward bulge in  $\theta_e$  contours there (Fig. 3b). The small outward bulge in the absolute angular momentum near 35-km radius and 2-km altitude provides additional evidence of this exchange, consistent with an injection of this high- $\theta_e$  air into the eyewall.

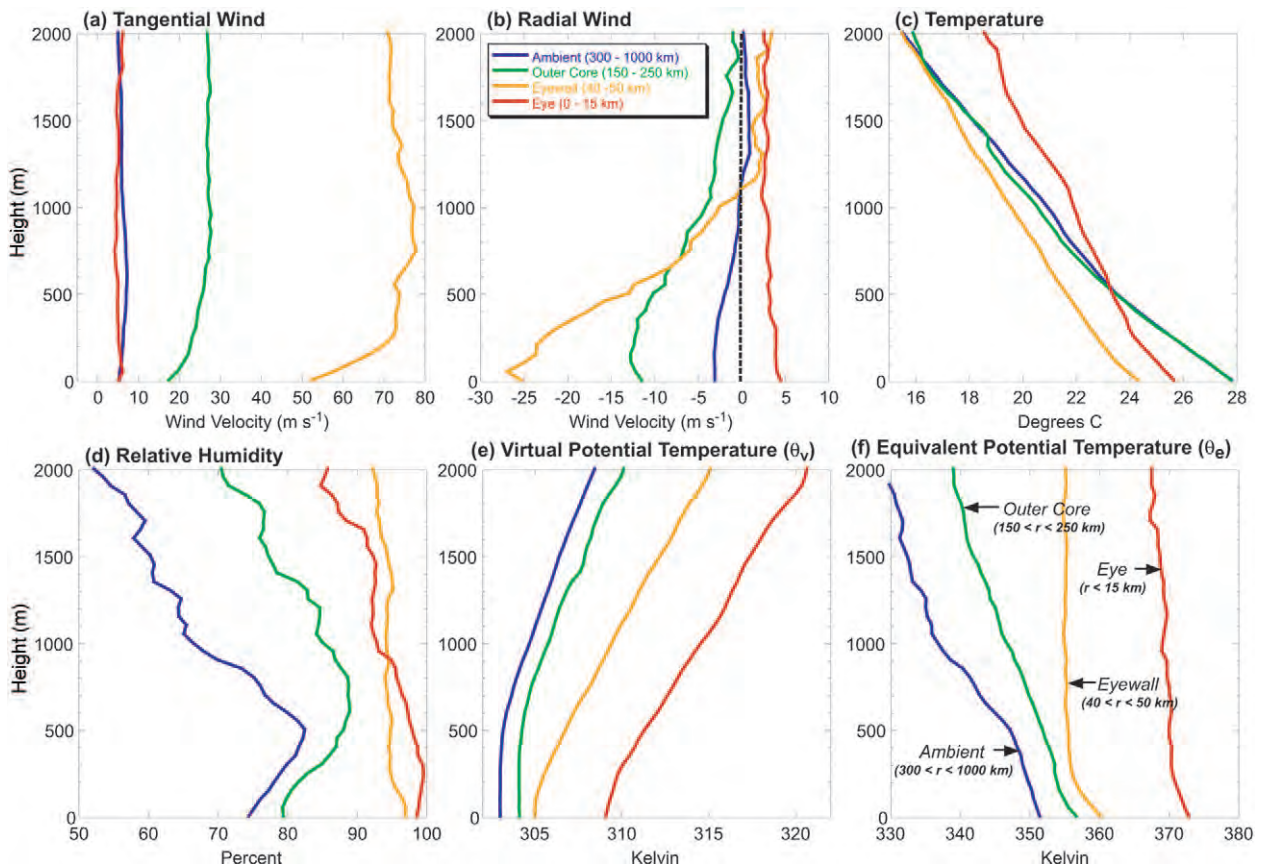
Continuous penetration of low-level air into the eye is believed to replenish the eye reservoir. The combined mean and asymmetric eddy flux of high- $\theta_e$  air from the low-level eye into the eyewall represents an additional power source to the hurricane heat engine. This “turboboost” of the engine invokes significant surface enthalpy fluxes well inside the radius at which E-MPI theory assumes they operate (PM03; cf. E95, E97, and BE98).

**Low-level structure.** The azimuthal mean low-level ( $0 < z < 2 \text{ km}$ ) kinematic and thermodynamic structures in the eye ( $r < 15 \text{ km}$ ) near the radius of maximum tangential wind (RMW;  $40 \text{ km} < r < 50 \text{ km}$ ), the outer core ( $r \sim 200 \text{ km}$ ), and the storm environment ( $300 \text{ km} < r < 1000 \text{ km}$ ) are shown in Fig. 4. Mean soundings with 50-m vertical resolution were constructed within these radial rings. The mean sounding near the RMW shows a shallow  $\theta_v$  mixed-

layer depth of only about 150 m, whereas the deep inflow layer extends up to 1.2 km. The maximum average tangential wind speed of  $\sim 76 \text{ m s}^{-1}$  near 1000 m is above the strong inflow. Whereas tangential and radial wind speeds of 27 and  $11 \text{ m s}^{-1}$ , respectively, are found in the outer core, only weak tangential wind and radial outflow exists inside the eye, and only weak inflow is present in the ambient environment. The relative humidity (RH) near and inside the RMW is greater than 90% in the lowest 2 km, indicating nearly saturated conditions even in the subcloud layer. The 0–50-m air temperature of  $\sim 24^\circ\text{C}$  near the RMW indicates a  $\sim 3.5 \text{ K}$  cooling from the outer core and environment. However, the low-level eye temperature is  $1.5^\circ\text{C}$  higher than near the RMW. The ambient average  $\theta_e$  in the mixed layer ( $\sim 500 \text{ m}$  deep) increases from 350 to  $354 \text{ K}$  in the outer core ( $\sim 200 \text{ km}$  radius), and then increases an additional 3 K as the air approaches the eyewall base at the RMW. An additional 14-K  $\theta_e$  increase is available to near-surface parcels that penetrate significantly into the eye and spend some time inside it. By accessing this energy surplus and transporting it to the eyewall, the mean tangential wind may increase beyond the currently formulated upper



**FIG. 3.** Radius–height azimuthal mean structure derived from GPS dropwindsonde and flight-level data from 1600 to 2300 UTC 13 Sep. (a) Storm-relative tangential wind (color) and radial wind (contour;  $\text{m s}^{-1}$ ); (b)  $\theta_e$  (color; K), absolute angular momentum (contour;  $\text{m}^2 \text{ s}^{-1} \times 10^6$ ), and transverse secondary circulation (vector). The origin (0, 0) denotes the eye center at the ocean surface.



**FIG. 4.** Low-level (0–2 km) mean soundings from dropwindsonde observations for the eye (0–15-km radius, red), nominal eyewall (40–50-km radius, orange), outer core (~200 km radius, green), and ambient environment (300–1000-km radius, blue) during 1600–2300 UTC 13 Sep. (a) Tangential wind ( $\text{m s}^{-1}$ ), (b) radial wind ( $\text{m s}^{-1}$ ; inflow/outflow transition indicated by dashed line), (c) temperature ( $^{\circ}\text{C}$ ), (d) RH (%), (e)  $\theta_v$  (K), and (f)  $\theta_e$  (K).

bounds. Whereas a detailed calculation of the residence time of air parcels in Isabel's eye is beyond the scope of this study, C06 showed that residence times in the eye for significant  $\theta_e$  gain in their numerical simulation of Hurricane Bonnie were commonly on the order of 40–60 min and were as low as 15 min. This suggests that the eye entropy reservoir may be replenished on even relatively short time scales.

**MPI ESTIMATES.** In E-MPI theory, the primary parameters needed to calculate the *a priori* MPI are 1) the sea surface temperature (SST) in the hurricane core, 2) the upper-level exhaust (outflow) temperature where parcels undergo radiational cooling to space, 3) the ambient RH near the sea surface, and 4) the sea-to-air exchange coefficients of enthalpy and angular momentum. These input parameters defined in E-MPI theory (E86; RE87; E95) were calculated using available data in Hurricane Isabel and its environment. Pertinent details of these calculations are summarized here.

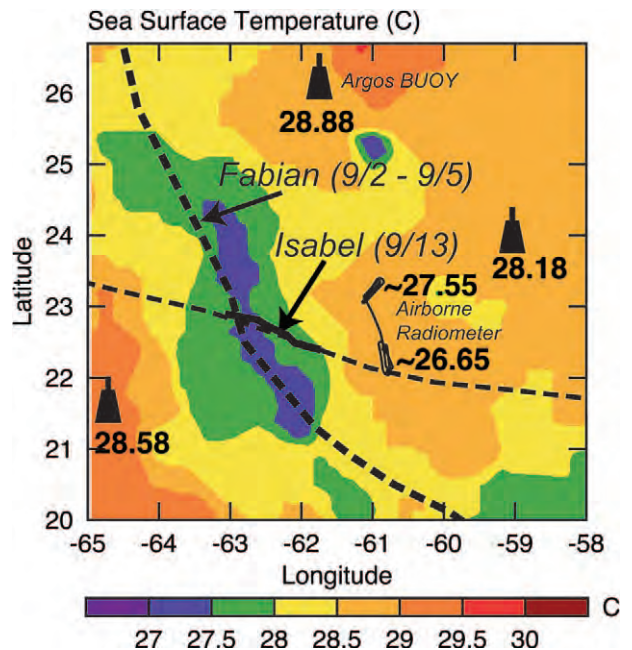
**Sea surface temperature.** The core SST (Fig. 5) was estimated from a combination of prestorm satellite imagery, Argos buoy data, and low-altitude aircraft radiometer measurements. Hurricane Fabian passed through this same area 10 days prior to Isabel, resulting in cooler water near the core region of Isabel than in the surrounding area during the time of the analysis. The P-3 aircraft radiometer measurements (Moss 1978) from NOAA 43, recorded at ~60 and ~120 m above sea level in the wake of Isabel, compare well with satellite SST estimates (McMillin and Crosby 1984; Brown and Minnett 1999; Legeckis and Zhu 1997) and suggest a  $27.5^{\circ}\text{C}$  effective SST near the storm core. Given the  $24^{\circ}\text{C}$  near-surface air temperature (Fig. 4c), a  $\sim 3^{\circ}\text{C}$  temperature deficit existed across the ocean–air interface, consistent with recent observations (Wroe and Barnes 2003; Cione et al. 2000).

For the *a priori* MPI estimates (see “MPI estimates: Azimuthal mean  $V_{\max}$  at boundary layer top”), a constant SST of  $27.5^{\circ}\text{C}$  is employed to represent an

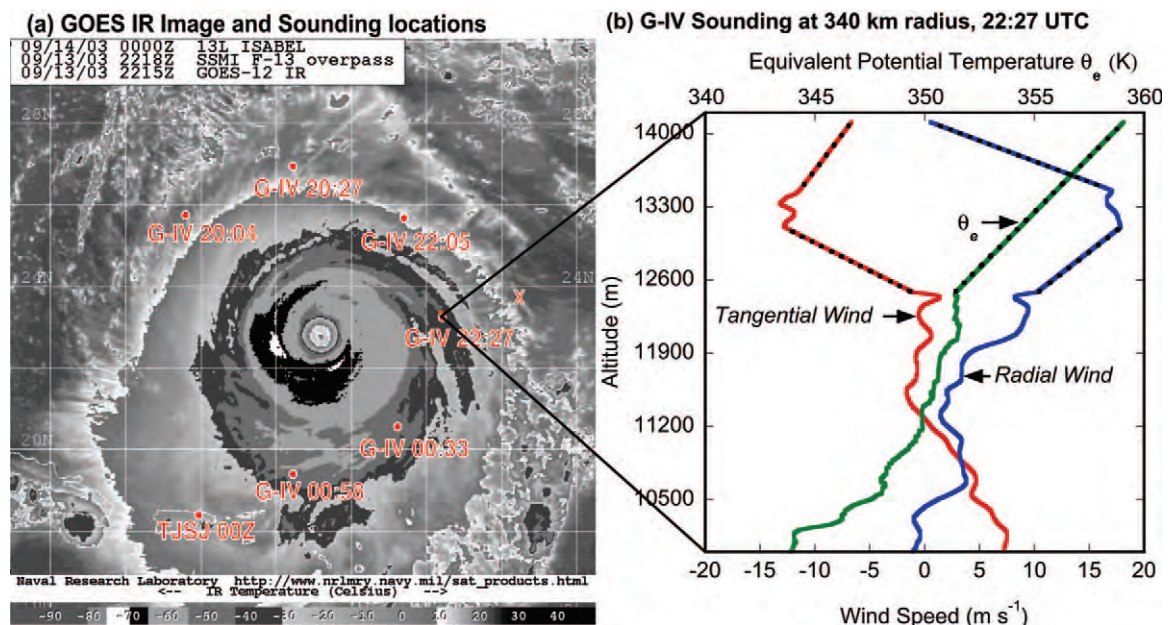


effective SST during the passage of Isabel over the wake of Fabian (from 1600 to 2300 UTC, see Fig. 5).<sup>2</sup> Although the airborne radiometer measured a reduction in SST by  $\sim 1\text{--}2^\circ\text{C}$  relative to satellite-derived SST estimates prior to Isabel's passage, this reduction is believed to be caused by Isabel and the corresponding shear-induced turbulent mixing from the cold thermocline (e.g., E04, and references therein). This negative feedback effect is not represented in E-MPI

<sup>2</sup> In E-MPI theory, the SST under the eyewall is critical for maintaining the maximum sustained tangential wind. Here, the boundary layer air parcels spiraling inward toward the eyewall from an outer-core region with an SST of  $\sim 28.5^\circ\text{C}$  are assumed to quickly lose their thermodynamic memory and to adjust rapidly to local ( $\sim 27.5^\circ\text{C}$ ) SST conditions on a boundary layer "eddy turnover" time scale  $\tau \sim H/w'$ , where  $H$  is the characteristic depth of the boundary layer and  $w'$  is a root-mean-square eddy vertical velocity. Letting  $H \sim 1\text{ km}$  and  $w' \sim 2\text{ m s}^{-1}$  gives  $\tau \sim 500\text{ s} \sim 8\text{ min}$ , a short time interval compared to that required for a boundary layer parcel to traverse horizontally inward from the outer core to the eyewall ( $L/U \sim 200\text{ km}/10\text{ m s}^{-1} \sim 5\text{ h}$ , where  $L$  represents a typical distance between the outer core and eyewall, and  $U$  is an upper bound to the average radial inflow within this region). Alternatively, an average ambient SST of  $\sim 28.5^\circ\text{C}$  only increases the primary intensity estimate of Fig. 7 by  $2\text{ m s}^{-1}$ .

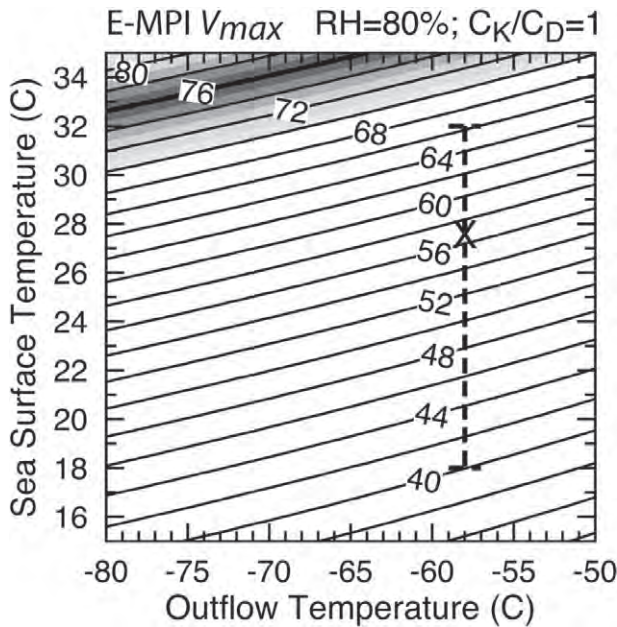


**FIG. 5.** SST derived from Advanced Very High Resolution Radiometer (AVHRR) satellite (averaged from 4 to 10 Sep, in color), Argos buoys, and NOAA P-3 downward-pointing radiometer (thin line, from 1800 UTC 13 Sep). Tracks of Hurricanes Fabian (heavy dashed, from 2 to 5 Sep) and Isabel (thin dashed, with thick solid line indicating analysis period from 1600 to 2300 UTC 13 Sep) are shown for reference.



**FIG. 6.** (a) GOES infrared satellite imagery at 2215 UTC (courtesy of NRL/Monterey) showing extent of hurricane outflow, and the dropwindsonde profile locations (red dots with UTC time labels) used in the outflow temperature calculation; "x" indicates location of additional outflow jet sample at 553-km radius from center. (b) High-altitude wind speed ( $\text{m s}^{-1}$ ) and  $\theta_e$  (K) from NOAA Gulfstream-IV dropwindsonde at 341-km radius from center at 2227 UTC. Tangential (red) and radial (blue) winds show anticyclonic outflow, with  $\theta_e$  in green. Black dashed line indicates linearly interpolated data.





**FIG. 7.** Theoretically predicted azimuthal mean tangential velocity  $V_{\max}$  at the boundary layer top assuming  $C_K/C_D = 1$  for varying outflow temperature and near-core SST with a constant  $RH = 80\%$ . “x” indicates the primary potential intensity estimate for the observed environment around Isabel (using a SST of  $27.5^\circ\text{C}$  near the eyewall region associated with the Fabian wake), which yields a  $56.6 \text{ m s}^{-1}$  mean  $V_{\max}$ . The dark solid curve represents the average storm-relative tangential wind speed at the top of the boundary layer derived from the dropwindsonde measurements, and the shading represents the  $6 \text{ m s}^{-1}$  standard deviation of this mean value. The vertical dashed line presents a range of mean intensities employing the same outflow temperature, SST, and RH used to obtain “x” so as to reflect uncertainties in the bulk exchange ratio and the ocean-cooling effect. The upper bound ( $\sim 66 \text{ m s}^{-1}$ ) entirely neglects the ocean cooling, includes dissipative heating, and assumes that  $C_K/C_D = 1$ . The lower bound ( $\sim 40 \text{ m s}^{-1}$ ) assumes  $C_K/C_D = 0.5$ , and assumes that the ocean-cooling effect is compensated by the dissipative heating effect. The discrepancy between these predicted values and the observed intensity suggests a significant limitation of E-MPI theory for this case.

theory. In “Sea-to-air exchange coefficients,” estimates of the impact of storm-induced cooling on the intensity of Isabel are presented to provide upper- and lower-bound intensity estimates.

**Outflow temperature.** Locations of soundings used to calculate the outflow temperature needed for the E-MPI calculation are shown in Fig. 6a, overlaid on infrared satellite imagery indicating the minimal extent of the upper-level outflow. Without the ability to track individual air parcels through eyewall

ascent and into the outflow, the outflow temperature calculation requires some modification from previous studies with numerical models (e.g., RE87). For completeness, the outflow temperature is calculated in the following three ways: i) as a  $d \ln(\theta_e)$ -weighted temperature, following the original definition in E86; ii) as an equilibrium-level temperature (the temperature at which a virtual parcel starting from an ambient surface state achieves the same environmental temperature after lifting by pseudoadiabatic ascent); and iii) as a radial wind-weighted temperature across the storm outflow (PM03). Table 1 presents a summary of these calculations.

The E-MPI theory (E86; RE87) uses the  $d \ln(\theta_e)$ -weighted integral definition

$$T_{\text{out}} \equiv \frac{1}{\ln \left[ \frac{\theta_{\text{ec}}}{\theta_{\text{ea}}} \right]} \int_{\theta_{\text{ea}}}^{\theta_{\text{ec}}} T d \ln(\theta_e), \quad (1)$$

where  $\theta_{\text{ea}}$  and  $\theta_{\text{ec}}$  denote the ambient subcloud layer  $\theta_e$  and that of a parcel in the eyewall updraft core ( $r \sim 45 \text{ km}$ , see Fig. 3b), respectively, consistent with the single-cycle Carnot model. Because  $\theta_e$  is approximately conserved during ascent in the eyewall and in the upper-tropospheric outflow,  $d \ln(\theta_e)$  is approximately zero until a parcel begins its descent due to radiational cooling at large ( $> 300 \text{ km}$ ) radii. Assuming the  $\theta_e$  profile of the sounding represents a family of parcels advected from the storm in the outflow, the integral in (1) can be performed on a single environmental sounding. The  $\theta_e$  at the outflow-layer bottom is assumed to be equal to  $\theta_{\text{ea}}$  (E86; RE87). As an example, the outflow jet at  $\sim 340 \text{ km}$  radius was sampled by a dropwindsonde (Fig. 6b) released by the NOAA Gulfstream-IV. Anticyclonic outflow extends from 11.5-km altitude up to  $\sim 14 \text{ km}$ , with peak anticyclonic tangential wind and radial wind velocities of  $13$  and  $17 \text{ m s}^{-1}$ , respectively. The  $\theta_e$  is assumed to increase monotonically with height from  $350$  to  $359 \text{ K}$  through the outflow layer. This assumption is supported by an additional measurement at  $\sim 550 \text{ km}$  radius along the same radial (the location indicated by an “x” in Fig. 5a), where the peak outflow wind speed was  $\sim 12 \text{ m s}^{-1}$  at  $13.5\text{-km}$  altitude and a similar thermodynamic structure existed (not shown). In contrast, most observations of the outflow around the storm have radial wind speeds  $< 10 \text{ m s}^{-1}$ , confirming that the hurricane outflow is a complex structure comprised of asymmetric outflow jets instead of the broad and homogeneous exhaust system predicted in axisymmetric numerical models (e.g., Holland 1987; Flatau and Stevens 1989; Vladimirov et al. 2001).

Note the  $\theta_e$  of 354 K in the center of the outflow jet is slightly less than that observed for the subcloud layer at the RMW (see Fig. 4b), and is approximately equal to the average  $\theta_e$  across the eyewall base (cf. Houze 1993). These limits from 350 to 359 K were used in the numerical integration of the seven environmental soundings<sup>3</sup> encircling the hurricane (Fig. 6a).

The temperature at the top of the boundary layer ( $T_B$ ), the outflow temperature ( $T_O$ ), and the RH from four representative environmental soundings are shown in Table 1. A range of outflow-layer depths and heights are observed in four quadrants of the hurricane, with an average of the seven  $d \ln(\theta_e)$ -weighted outflow temperatures of  $-58^\circ\text{C}$  (215.3 K). This value will be used as the primary system-scale outflow temperature.<sup>4</sup> The equilibrium-level and radial wind-weighted outflow temperatures have been shown previously to serve as reasonable proxies for the  $d \ln(\theta_e)$ -weighted outflow temperature, and are found here to be consistent with the outflow tem-

perature defined by (1). As an additional check, the axisymmetric thermal wind equation appropriate for a moist hurricane vortex was also used to assess the outflow temperature. This yielded consistent results, and details of this calculation are summarized in the appendix for interested readers.

**Relative humidity.** The average RH in the lowest 500 m in these seven environmental soundings is near 80%, consistent with the boundary layer mean structure (Fig. 4d). In the original and revised a priori E-MPI formulations (E86, E95), a simple closure is invoked for the near-surface RH between the environment and the region under the eyewall by assuming a constant RH between these regions. Clearly, the near-saturated conditions throughout the subcloud layer underneath the eyewall violate this assumption. However, E-MPI also assumes that the air parcels maintain constant temperatures as they travel into the storm core. Because the observed cooling tends

<sup>3</sup> Missing data in the  $\theta_e$  and wind fields were interpolated linearly between available observations in order to perform the integration. Though this smooths the radial wind in the PM03 outflow proxy calculation, the monotonic  $\ln(\theta_e)$  profile is not sensitive to this.

<sup>4</sup> The San Juan sounding appears to be an outlier. This could be due to the high tropopause height at low latitudes, the high equilibrium level because the rawinsonde was launched from land, and/or the different sensor type from the GPS dropwindsondes.

<b>TABLE 1. Observed and calculated environmental temperatures and RH for Hurricane Isabel on 13 Sep 2003. RH is the average in the lowest 500 m. Temperature at the boundary layer top (<math>T_B</math>), radial wind-weighted outflow temperature (<math>T_O</math> PM), equilibrium-level outflow temperature (<math>T_O</math> EL), <math>d \ln(\theta_e)</math>-weighted outflow temperature (<math>T_O</math> E-MPI), and <math>\theta_e</math> at the outflow-layer top and bottom are also shown.</b>										
Sounding	Radius from center (km)	RH (%)	$T_B$ ( $^\circ\text{C}$ )	$T_O$ PM ( $^\circ\text{C}$ )	$T_O$ EL ( $^\circ\text{C}$ )	$T_O$ E-MPI ( $^\circ\text{C}$ )	$\theta_e$ bottom (K)	Radial outflow layer		
								Altitude bottom (m)	$\theta_e$ top (K)	Altitude top (m)
Gulf Stream-IV 2227:14 UTC (northeast)	341	72.6	24.9	-54.2	-59	-52.9	349.7	11,140	358.8	14,060
Gulfstream-IV 2004:13 UTC (northwest)	525	82.5	24.5	-43.4	-44	-54.4	349.7	11,700	354.9	13,210
Gulfstream-IV 0033:36 UTC (southeast)	347	77.7	23.8	-55.7	-50	-56.5	349.8	12,090	359.2	14,359
TJSJ 2400:00 UTC (southwest)	584	84.2	24.3	-70.7	-78	-79.5	349.7	14,109	359.2	16,539
Average of seven soundings	653.8	80.7	24.1	-52.9	-56.7	-57.8	349.7	12,389	358.0	14,801

to offset the moistening, this yields an approximately similar effect in the moist entropy (e.g., Camp and Montgomery 2001). For the present calculations, a simple compromise is adopted (E95; PM03) by specifying a constant 80% RH.

**Sea-to-air exchange coefficients.** For the primary intensity estimate, the ratio of bulk sea-to-air enthalpy and angular momentum exchange coefficients is assumed to be unity ( $C_K/C_D = 1$ ). However, observations collected during the 2002–04 Coupled Boundary Layers Air–Sea Transfer (CBLAST) field program suggest a value of unity is an upper bound; current estimates of this ratio for a near-surface ( $\sim 10$  m) wind speed of approximately  $56 \text{ m s}^{-1}$  (see Figs. 4a,b) vary from approximately 0.5 to 0.7 (Black et al. 2006). Among the four parameters needed for the a priori MPI estimate (SST, outflow temperature, ambient RH, and  $C_K/C_D$ ), the last is thought to be the most uncertain. The complex air–sea interaction processes in extreme winds are not well understood, and further research is needed to elucidate these effects more clearly (e.g., Black 2004; Sullivan et al. 2004a,b; Black et al. 2006).

**MPI estimates: Azimuthal mean  $V_{\max}$  at boundary layer top.** The predicted  $V_{\max}$  from the a priori E-MPI theory for varying outflow temperatures and near-core SSTs with a constant RH = 80% and  $C_K/C_D = 1$  are shown in Fig. 7, with the “x” representing the primary intensity estimate. In these calculations,  $T_B$  is assumed to be  $\sim 4^\circ\text{C}$  lower than the SST, as suggested by the data in Table 1, and the thermodynamic efficiency is defined by

$$\varepsilon = (T_B - T_O)/T_B. \quad (2)$$

The average efficiency for Isabel on 13 September with an SST of  $27.5^\circ\text{C}$  is found to be 0.275. This original definition of the efficiency of the Carnot cycle by E86 provides a clear distinction between the SST and the air temperature at the top of the boundary layer (e.g., Houze 1993). The predictive formula used here for  $V_{\max}$  is otherwise identical to a revised formulation of the Carnot model (E95) and is the basis for the values shown in Fig. 7.

*A major conclusion of this study is that the observed azimuthal mean tangential wind speed at the boundary layer top ( $76 \text{ m s}^{-1}$ , with a  $6 \text{ m s}^{-1}$  standard deviation) is clearly stronger than the theoretically predicted  $56.6 \text{ m s}^{-1}$  E-MPI. Warm ocean eddies do not seem to be a plausible explanation for this discrepancy given that SST observations clearly depict the cool wake left*

earlier by Fabian (cf. E04). Satellite data (not shown) indicate that the SST was an additional 1–2 K lower in this region after the passage of Isabel, suggesting that the storm-induced ocean cooling (E04) and turbulent dissipative heating (BE98) tend in opposite ways.

The (turbulent) dissipative heating effect (BE98) relative to the original a priori MPI (E86; E95) is calculated according to  $\Delta V_{\max} = V_{\max} [(T_B/T_O)^{1/2} - 1]$ , which here represents a  $9.6 \text{ m s}^{-1}$  boost. Ocean-cooling effects depend primarily on the storm translation speed and the oceanic mixed layer depth. E04 predicted a  $\sim 13 \text{ m s}^{-1}$  intensity reduction for a storm moving at  $7 \text{ m s}^{-1}$  with a 30-m mixed layer depth (their section 2c and Fig. 3). Because Isabel crossed the wake left by Fabian (Fig. 6), Isabel likely encountered a deeper mixed layer, resulting in a somewhat reduced cooling and negative feedback effect. The estimated upper and lower bounds for the maximum intensity presented in Fig. 7 reflect uncertainties in both the bulk exchange ratio ( $C_K/C_D$ ) and the storm-induced ocean cooling. The upper-bound estimate ( $\sim 66 \text{ m s}^{-1}$ ) entirely neglects the ocean cooling caused by Isabel, includes dissipative heating, and assumes a ratio of enthalpy and angular momentum exchange coefficients of unity. The lower-bound estimate ( $\sim 40 \text{ m s}^{-1}$ ) assumes a  $C_K/C_D$  ratio of 0.5, and that the ocean cooling is compensated by the dissipative heating effect. The discrepancy between all of these predicted values and the observed intensity of  $76 \text{ m s}^{-1}$  suggests a significant limitation of current E-MPI theory for this case.

**CONCLUSIONS.** For two decades, E-MPI theory often has been presented as a rigorous upper bound on hurricane intensity that is not exceeded within observational error (E86, E88, E91, E95; RE87; BE98; E04). Whereas the single-cycle Carnot model provides a good first approximation for the maximum potential intensity, the current work presents evidence that the E-MPI theory does not include a critical aspect of the inner-core structure of mature hurricanes. The in situ observations in Hurricane Isabel suggest that the “turboboost” mechanism (PM03) may play an important role in providing more power to tropical cyclones and, thus, explain how  $V_{\max}$  may exceed this previous upper bound for storms that lack adverse environmental influences. For the case of Hurricane Isabel on 13 September this  $V_{\max}$  discrepancy between the a priori MPI theory and the observed intensity is consistent with recent high-resolution numerical model simulations, supporting the hypothesis that Isabel was superintense as defined therein. A similar vortex-scale structure, deduced from sorties flown on 12 September when the storm was even stronger



(peak mean tangential wind speed at the boundary layer top  $> 80 \text{ m s}^{-1}$ , not shown), suggests to us that the turboboost mechanism is the most likely candidate to explain the extended period of Isabel's intense state. More examples from this extraordinary set of observations are presented in Part II to shed insight into the exotic and complex nature of the eye–eyewall mixing phenomenon in Isabel. Whereas the observations in Part II were serendipitous, targeted observations in other intense tropical cyclones are encouraged, which, in conjunction with numerical, theoretical, and laboratory work, will lead to better understanding of the heat and momentum exchange between the eye and the eyewall and to an improved maximum intensity theory.

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## APPENDIX: THERMAL WIND BALANCE.

This appendix presents an examination of whether the deduced mean tangential wind and entropy fields presented in Fig. 3 are dynamically consistent and, given the result, a check on whether the outflow temperature computed in “Outflow temperature” is sufficiently cold to support the observed winds.

A thermal wind equation relating the maximum mean tangential wind speed and the radial gradient of moist entropy at the top of the boundary layer has been derived previously [E86, Eq. (13); E97, Eq. (13)], as

$$V_{\max}^2 \approx -r_m (T_B - T_O) \frac{dS_B}{dr}, \quad (\text{A1})$$

where  $S_B$  is the moist entropy at the boundary layer top,  $r$  is the radius from the vortex center,  $r_m$  is the RMW, and other symbols are as defined in the text. Equation (A1) is a diagnostic equation that depends on inter alia the validity of the hydrostatic and axisymmetric

cyclostrophic balance approximations (e.g., Holton 2004) in the eyewall of the storm. Cyclostrophic balance is a valid first approximation in the eyewall above the boundary layer where the Rossby number is large compared to unity and the flow is rotationally dominant (e.g., Willoughby 1990). The superintensity mechanism will, of course, increase the radial gradient of  $S_B$  (and thus the radial gradient of  $\theta_e$ ) and increase the resulting tangential wind speed. Therefore, (A1) should be valid for superintensity, as long as the winds are not appreciably supergradient. Because all quantities appearing in (A1) are calculable from the observations, the applicability of the thermal wind constraint can be assessed, and, conversely, it can be used to determine whether the outflow temperature computed in “Outflow temperature” is sufficiently cold to support the intensity observed. Inserting the  $\theta_e$  data, the radial moist entropy gradient near  $r \sim 42 \text{ km}$  (RMW) and  $z = 1 \text{ km}$  (boundary layer top) is approximately

$$\begin{aligned} \frac{dS_B}{dr} &= \frac{c_p}{\theta_e} \frac{d\theta_e}{dr} \\ &\approx \frac{1004 \text{ J}(\text{kgK})^{-1}}{357 \text{ K}} \left( \frac{-6 \text{ K}}{10 \text{ km}} \right) \\ &= -1.7 \times 10^{-3} \text{ m}(\text{s}^2 \text{ K})^{-1}. \end{aligned} \quad (\text{A2})$$

Substituting this gradient into (A1), and using the values of  $T_B$  and  $T_O$  computed previously, yields  $V_{\max} \approx 74 \text{ m s}^{-1}$ , which is close to the observed  $V_{\max} \approx 76 \text{ m s}^{-1}$  deduced from the dropwindsondes. If, instead, the observed  $V_{\max} \approx 76 \text{ m s}^{-1}$  is used and the outflow temperature for the above radial entropy gradient near the RMW is inferred, one obtains  $T_O \approx 216 \text{ K}$ , close to the outflow temperature calculated from the dropwindsonde and the Gulfstream-IV data as described in “Outflow temperature.”

It should be noted that this outflow temperature estimate is sensitive to the value chosen for the radial  $\theta_e$  gradient. Errors from humidity sensor biases (Bell and Montgomery 2006, manuscript submitted to *Mon. Wea. Rev.*) and spatial and temporal averaging limit the accuracy of the calculated radial entropy gradient. Deviations from axisymmetric thermal wind balance (A1) near the boundary layer top are also expected due to the mesovortices observed near the RMW. Given all of these potential errors, the consistency between the outflow temperature, the maximum tangential wind, and the radial entropy gradient during this analysis period is encouraging.

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