

Orthogonality Relations

We now explain the basic reason why the remarkable Fourier coefficient formulas work. We begin by repeating them from the last note:

$$\begin{aligned} a_0 &= \frac{1}{L} \int_{-L}^L f(t) dt, \\ a_n &= \frac{1}{L} \int_{-L}^L f(t) \cos\left(n \frac{\pi}{L} t\right) dt, \\ b_n &= \frac{1}{L} \int_{-L}^L f(t) \sin\left(n \frac{\pi}{L} t\right) dt. \end{aligned} \quad (1)$$

The key fact is the following collection of integral formulas for sines and cosines, which go by the name of **orthogonality relations**:

$$\frac{1}{L} \int_{-L}^L \cos\left(n \frac{\pi}{L} t\right) \cos\left(m \frac{\pi}{L} t\right) dt = \begin{cases} 1 & n = m \neq 0 \\ 0 & n \neq m \\ 2 & n = m = 0 \end{cases}$$

$$\frac{1}{L} \int_{-L}^L \cos\left(n \frac{\pi}{L} t\right) \sin\left(m \frac{\pi}{L} t\right) dt = 0$$

$$\frac{1}{L} \int_{-L}^L \sin\left(n \frac{\pi}{L} t\right) \sin\left(m \frac{\pi}{L} t\right) dt = \begin{cases} 1 & n = m \neq 0 \\ 0 & n \neq m \end{cases}$$

Proof of the orthogonality relations: This is just a straightforward calculation using the periodicity of sine and cosine and either (or both) of these two methods:

Method 1: use $\cos at = \frac{e^{iat} + e^{-iat}}{2}$, and $\sin at = \frac{e^{iat} - e^{-iat}}{2i}$.

Method 2: use the trig identity $\cos(\alpha) \cos(\beta) = \frac{1}{2}(\cos(\alpha + \beta) + \cos(\alpha - \beta))$, and the similar trig identities for $\cos(\alpha) \sin(\beta)$ and $\sin(\alpha) \sin(\beta)$.

Using the orthogonality relations to prove the Fourier coefficient formula

Suppose we know that a periodic function $f(t)$ has a Fourier series expansion

$$f(t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos\left(n \frac{\pi}{L} t\right) + b_n \sin\left(n \frac{\pi}{L} t\right) \quad (2)$$

How can we find the values of the coefficients? Let's choose one coefficient, say a_2 , and compute it; you will easily how to generalize this to any other coefficient. The claim is that the right-hand side of the Fourier coefficient formula (1), namely the integral

$$\frac{1}{L} \int_{-L}^L f(t) \cos\left(2 \frac{\pi}{L} t\right) dt.$$

is in fact the coefficient a_2 in the series (2). We can replace $f(t)$ in this integral by the series in (2) and multiply through by $\cos\left(2\frac{\pi}{L}t\right)$, to get

$$\frac{1}{L} \int_{-L}^L \frac{a_0}{2} \cos\left(2\frac{\pi}{L}t\right) + \sum_{n=1}^{\infty} \left(a_n \cos\left(n\frac{\pi}{L}t\right) \cos\left(2\frac{\pi}{L}t\right) + b_n \sin\left(n\frac{\pi}{L}t\right) \cos\left(2\frac{\pi}{L}t\right) \right) dt$$

Now the orthogonality relations tell us that almost every term in this sum will integrate to 0. In fact, the only non-zero term is the $n = 2$ cosine term

$$\frac{1}{L} \int_{-L}^L a_2 \cos\left(2\frac{\pi}{L}t\right) \cos\left(2\frac{\pi}{L}t\right) dt$$

and the orthogonality relations for the case $n = m = 2$ show this integral is equal to a_2 as claimed.

Why the denominator of 2 in $\frac{a_0}{2}$?

Answer: it is in fact just a convention, but the one which allows us to have the same Fourier coefficient formula for a_n when $n = 0$ and $n \geq 1$. (Notice that in the $n = m$ case for cosine, there is a factor of 2 only for $n = m = 0$.)

Interpretation of the constant term $\frac{a_0}{2}$.

We can also interpret the constant term $\frac{a_0}{2}$ in the Fourier series of $f(t)$ as the average of the function $f(t)$ over one full period: $\frac{a_0}{2} = \frac{1}{2L} \int_{-L}^L f(t) dt$.

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