

Nonlinear Control Theory

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Introduction



- 1 **Nonlinear models and nonlinear phenomena**
- 2 Examples



Nonlinear models and nonlinear phenomena

Consider dynamical systems that are modeled by a finite number of coupled first-order ordinary differential equations:

$$\begin{aligned}\dot{x}_1 &= f_1(t, x_1, \dots, x_n, u_1, \dots, u_p), \\ \dot{x}_2 &= f_2(t, x_1, \dots, x_n, u_1, \dots, u_p), \\ &\vdots \\ \dot{x}_n &= f_n(t, x_1, \dots, x_n, u_1, \dots, u_p),\end{aligned}$$

where

- x_1, \dots, x_n : state variables;
- $\dot{x}_1, \dots, \dot{x}_n$: derivatives w.r.t. time t ;
- u_1, \dots, u_p : specified input variables.



Vector notations:

$$x = \begin{bmatrix} x_1 \\ x_2 \\ \vdots \\ x_n \end{bmatrix}, \quad u = \begin{bmatrix} u_1 \\ u_2 \\ \vdots \\ u_p \end{bmatrix}, \quad f(t, x, u) = \begin{bmatrix} f_1(t, x, u) \\ f_2(t, x, u) \\ \vdots \\ f_n(t, x, u) \end{bmatrix}.$$

$\Downarrow \quad \Downarrow \quad \Downarrow$

- One n -dimensional first-order vector differential equation (state equation):

$$\dot{x} = f(t, x, u).$$

- Output equation (for particular interest):

$$y = h(t, x, u).$$



- Unforced state equation:

$$\dot{x} = f(t, x).$$

- * Working with an unforced state equation does not necessarily mean that the input to the system is zero.
- * Possibly $u = \gamma(x)$ or $u = \gamma(t, x)$.

- Autonomous system (time-invariant system):

$$\dot{x} = f(x).$$

- * If the system is not autonomous, then it is called nonautonomous or time varying.



Definition (Equilibrium point)

A point $x = x^*$ in the state space is said to be an equilibrium point of $\dot{x} = f(t, x)$, if it has the property that whenever the state of the system starts at x^* , it will remain at x^* for all future time.

- For the autonomous system $\dot{x} = f(x)$, the equilibrium points are the real roots of the equation

$$f(x) = 0.$$

- * An equilibrium point could be isolated; that is, there are no other equilibrium points in its vicinity, or there could be a continuum of equilibrium points.



For linear systems, the state model take the special form:

$$\begin{aligned}\dot{x} &= A(t)x + B(t)u, \\ y &= C(t)x + D(t)u.\end{aligned}$$

- Linear systems: superposition principle.
- For nonlinear systems, the superposition principle does not hold any longer.
- Linearization around some operation points.
- However, linearization alone will not be sufficient.
 - * Linearization predicts only “local” behavior; it cannot predict the “global” behavior throughout the state space.
 - * There are “essentially nonlinear phenomena” that cannot be described by linear models



Essentially nonlinear phenomena:

- **Finite escape time:** A nonlinear system's state can go to infinity in finite time.
- **Multiple isolated equilibria:** A nonlinear system can have more than one isolated equilibrium point. The state may converge to one of several operating points, depending on the initial state.
- **Limit cycles:** There are nonlinear systems that can go into an oscillation of fixed amplitude and frequency, irrespective of the initial state.
- **Subharmonic, harmonic, or almost-periodic oscillations:** A nonlinear system under periodic excitation can oscillate with frequencies of submultiples or multiples of the input frequency.
- **Chaos:** The steady-state behavior of a nonlinear system is possibly not equilibrium, periodic or almost-periodic oscillation, but exhibits randomness, despite the deterministic nature of the system.
- **Multiple modes of behavior:** More than one limit cycle; harmonic, subharmonic, or more complicated steady-state behavior; discontinuous jump, etc.



- ① Nonlinear models and nonlinear phenomena
- ② **Examples**



Examples – Pendulum

Parameters and variables:

- l denotes the length of the rod, and m denotes the mass of the bob. Assume the rod is rigid and has zero mass.
- θ denotes the angle subtended by the rod and the vertical axis through the pivot point.
- g is the acceleration due to gravity.
- There is a frictional force resisting the motion, which is proportional to the speed of the bob with a coefficient of friction k

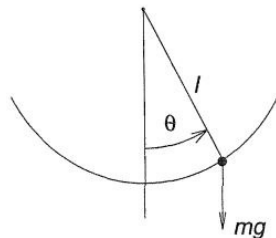


Figure: Pendulum



Using Newton's second law of motion:

$$ml\ddot{\theta} = -mg \sin \theta - k\dot{\theta}.$$

Take the state variables: $x_1 = \theta$ and $x_2 = \dot{\theta}$:

$$\dot{x}_1 = x_2,$$

$$\dot{x}_2 = -\frac{g}{l} \sin x_1 - \frac{k}{m} x_2.$$

Set $\dot{x}_1 = \dot{x}_2 = 0$ to find the equilibrium points:

$$0 = x_2, \quad 0 = -\frac{g}{l} \sin x_1 - \frac{k}{m} x_2 \quad \Rightarrow \quad (x_1, x_2)_{eq} = (n\pi, 0), \quad n = 0, \pm 1, \pm 2, \dots$$



If the frictional resistance is neglected ($k = 0$), the resulting system

$$\begin{aligned}\dot{x}_1 &= x_2, \\ \dot{x}_2 &= -\frac{g}{l} \sin x_1\end{aligned}$$

is conservative.

- * If the pendulum is given an initial push, it will keep oscillating forever with a nondissipative energy exchange between kinetic and potential energies.

If we can apply a torque T as the control input to the pendulum, then

$$\begin{aligned}\dot{x}_1 &= x_2, \\ \dot{x}_2 &= -\frac{g}{l} \sin x_1 - \frac{k}{m} x_2 + \frac{1}{ml^2} T.\end{aligned}$$



Example – Tunnel-Diode Circuit

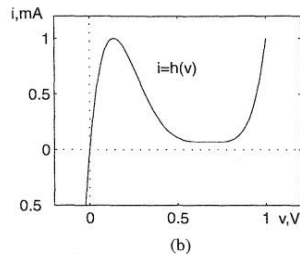
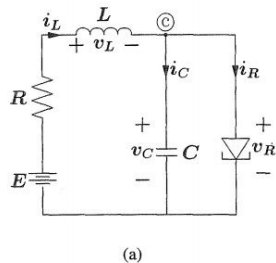


Figure: (a) Tunnel-diode circuit; (b) Tunnel-diode $v_R - i_R$ characteristic.



- $i_R = h(v_R)$, $i_C = C \frac{dv_C}{dt}$, $v_L = L \frac{di_L}{dt}$
- States: $x_1 = v_C$ and $x_2 = i_L$; Input: $u = E$
- Kirchhoff's current law: $i_C + i_R - i_L = 0 \Rightarrow i_C = -h(x_1) + x_2$
- Kirchhoff's voltage law: $v_C - E + Ri_L + v_L = 0 \Rightarrow v_L = -x_1 - Rx_2 + u$
- State model:

$$\begin{aligned}\dot{x}_1 &= \frac{1}{C}[-h(x_1) + x_2], \\ \dot{x}_2 &= \frac{1}{L}[-x_1 - Rx_2 + u]\end{aligned}$$



The equilibrium points are determined by setting $\dot{x}_1 = \dot{x}_2 = 0$:

$$0 = \frac{1}{C}[-h(x_1) + x_2],$$

$$0 = \frac{1}{L}[-x_1 - Rx_2 + u]$$

$$\Rightarrow h(x_1) = \frac{E}{R} - \frac{1}{R}x_1.$$

Three isolated equilibrium points!!

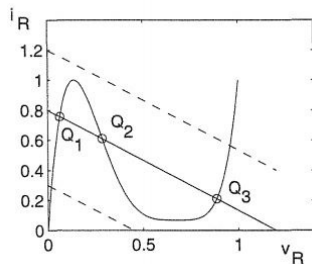


Figure: Equilibrium points of the tunnel diode circuit



Example – Mass-Spring System

- Newton's law of motion:

$$m\ddot{y} + F_f + F_{sp} = F$$

where F_f is a resistive force due to friction, and F_{sp} is the restoring force of the spring.

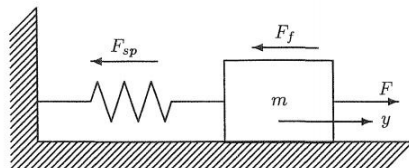


Figure: Mass-spring mechanical system



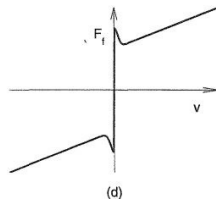
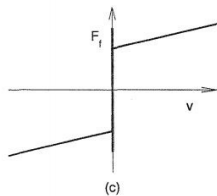
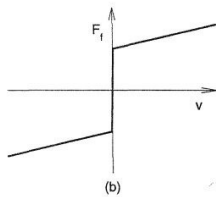
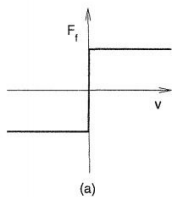
The restoring force of the spring:

- $F_{sp} = g(y)$, and $g(0) = 0$.
- For small displacement: $g(y) = ky$ with constant k .
- Large displacement, softening spring: $g(y) = k(1 - a^2y^2)y$ with $|ay| < 1$.
- Large displacement, hardening spring: $g(y) = k(1 + a^2y^2)y$.

The resistive force:

- Static friction $F_s = \pm\mu_s mg$ with $0 < \mu_s < 1$ to keep the mass at rest.
- Coulomb friction $F_c = -\mu_k mg$ for $v < 0$, and $F_c = \mu_k mg$ for $v > 0$.
- Viscous friction $F_v = h(v)$ with $h(0) = 0$ which is possibly nonlinear.





- (a) Coulomb friction
- (b) Coulomb plus linear viscous friction
- (c) Static, Coulomb, and linear viscous friction
- (d) Static, Coulomb, and linear viscous friction – Stribeck effect



- Hardening spring + linear viscous friction + a periodic external force $F = A \cos \omega t$

$$m\ddot{y} + c\dot{y} + ky + ka^2y^3 = A \cos \omega t$$

Duffing's equation, which is a typical example of periodic excitation.

- Linear spring + static friction + Coulomb friction + linear viscous friction

$$m\ddot{y} + c\dot{y} + ky + \eta(y, \dot{y}) = 0$$

where

$$\eta(y, \dot{y}) = \begin{cases} \mu_k mg \text{sign}(\dot{y}), & \text{for } |\dot{y}| > 0, \\ -ky, & \text{for } \dot{y} = 0 \text{ and } |y| \leq \mu_s mg/k, \\ -\mu_s mg \text{sign}(\dot{y}), & \text{for } \dot{y} = 0 \text{ and } |y| > \mu_s mg/k. \end{cases}$$



- Let $x_1 = y$ and $x_2 = \dot{y}$,

$$\dot{x}_1 = x_2$$

$$\dot{x}_2 = -\frac{k}{m}x_1 - \frac{c}{m}x_2 - \frac{1}{m}\eta(x_1, x_2).$$

- * Equilibrium set

$$\dot{x}_1 = 0, \dot{x}_2 = 0 \quad \Rightarrow \quad x_2 = 0, \quad -\mu_s mg/k \leq x_1 \leq \mu_s mg/k.$$

- * Discontinuity

$$\text{When } x_2 > 0, \eta = \mu_k mg; \quad \text{when } x_2 < 0, \eta = -\mu_k mg.$$

This is an example of *piecewise linear analysis*, where a system is represented by linear models in various regions of the state space, certain coefficients changing from region to region.

