

PH 201

OPTICS & LASERS

Lecture_Lasers_4

Conditions for Producing a Laser

Population Inversion, Gain, & Gain saturation

Absorption & Gain on a Homogeneously Broadened Radiative Transition

- ❖ Consider a beam of light having intensity per unit frequency $I(\nu)$ at frequency ν & frequency width $\Delta\nu$, passing through a medium of thickness L & cross-sectional area A .

$$I = I(\nu)\Delta\nu$$

I_0 = Intensity I just before it enters the medium

Within medium we assume there are atoms occupying at least two specific energy levels (u & l). Radiative transitions from u to l can occur over a narrow range of frequencies within emission linewidth (center frequency ν_0).

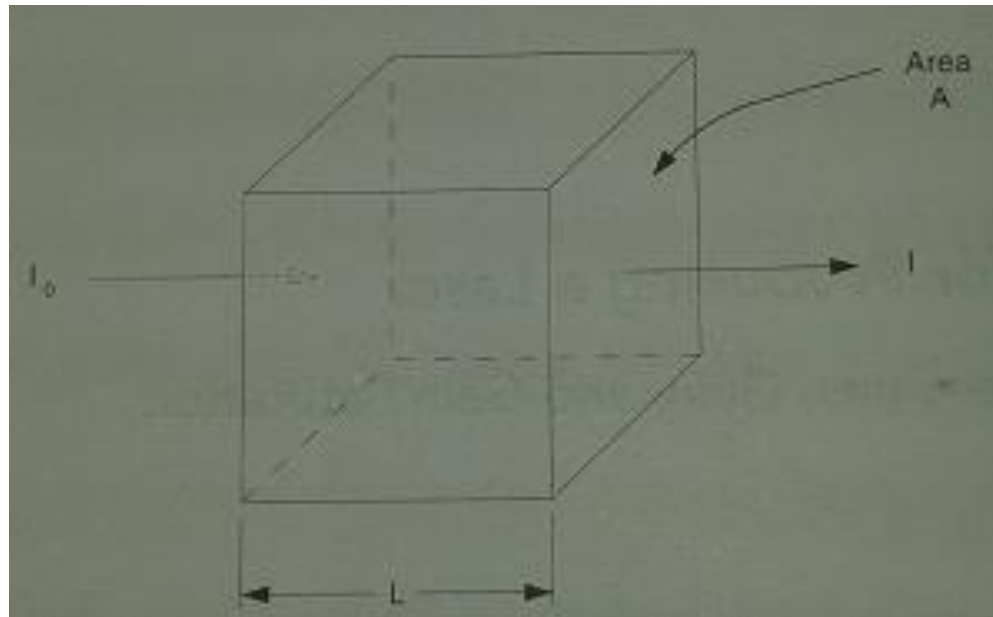
We also assume that frequency of incident beam falls somewhere within that linewidth.

$$\Delta E_{ul} = E_u - E_l = h\nu_{ul}$$

Determine the effect of medium on beam as beam passes through it.

How I is changed after it has passed through medium?

Assume that dominant broadening of transition under consideration is homogeneous, with a Lorentzian lineshape.



Light of intensity I_0 incident upon an absorptive material of length L & area A .

N_u = Total no. of atoms per unit volume in upper level

N_l = Total no. of atoms per unit volume in lower level

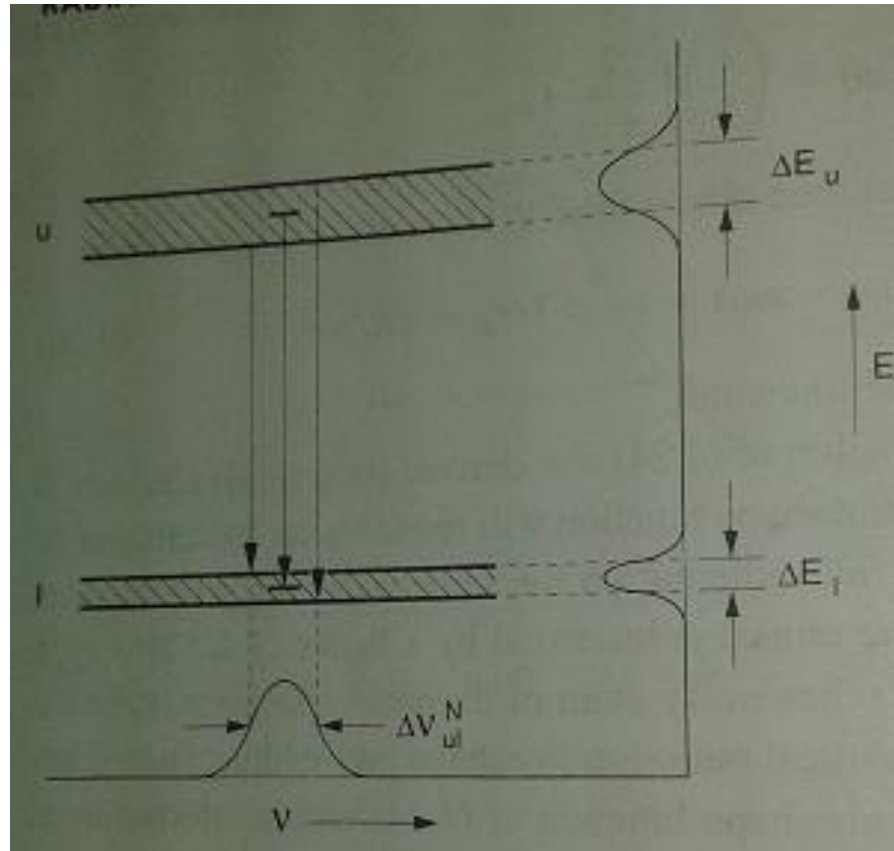
A_{ul} = Radiative transition probability for transitions occurring between energy levels u & l .

$$A_{ul}(\nu) = \frac{\chi_{ul}^T / 4\pi^2}{(\nu - \nu_0)^2 + (\chi_{ul}^T / 4\pi)^2} A_{ul}$$

χ_{ul}^T = Total decay rate of upper & lower levels (minimum emission linewidth)

$$\chi_{ul}^T = \chi_u + \chi_l$$

$$\chi_{ul}^T = 2\pi\Delta\nu_{ul}^N = \sum_i A_{ui} + \sum_j A_{lj}$$



Quantum mechanical description of natural linewidth of emission resulting from a radiating transition between two levels (u & l).

Stimulated emissions are proportional to energy density $u(\nu)$ of beam at frequency ν , as well as Einstein B coefficients.

No. of stimulated transitions (upward or downward) occurring between two levels per unit volume per unit time,

$$\begin{aligned} N_l B_{lu}(\nu) \Delta\nu \cdot u(\nu) &= N_l B_{lu}(\nu) I(\nu) \Delta\nu \eta / C \\ &= N_l B_{lu}(\nu) I \eta / C \end{aligned} \quad \text{(upward)}$$

&

$$\begin{aligned} N_u B_{ul}(\nu) \Delta\nu \cdot u(\nu) &= N_u B_{ul}(\nu) I(\nu) \Delta\nu \eta / C \\ &= N_u B_{ul}(\nu) I \eta / C \end{aligned} \quad \text{(downward)}$$

Using relation

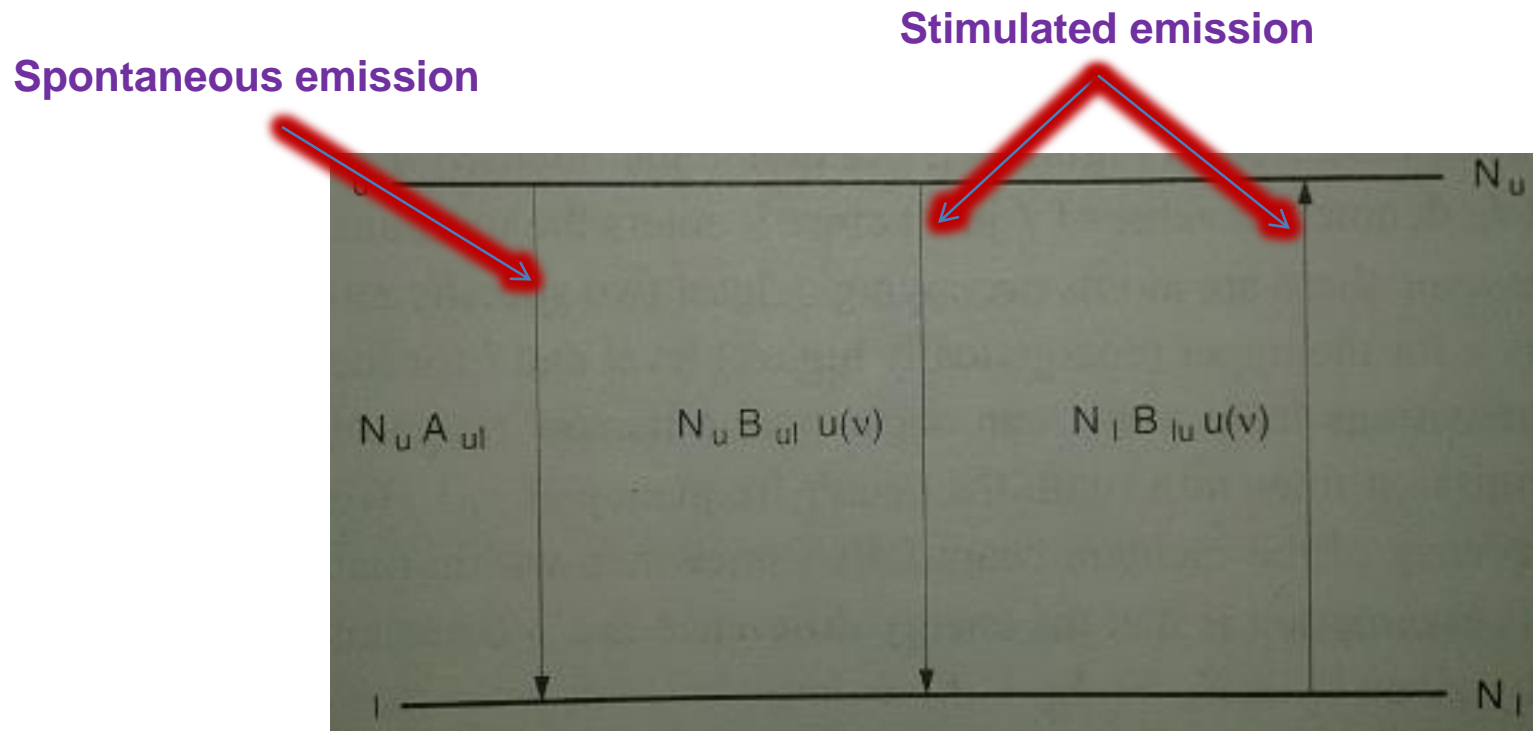
$$u(\nu) \Delta\nu = I(\nu) \Delta\nu \eta / C = I \eta / C$$

It is also assumed that B coefficients per unit frequency $B_{ul}(\nu)$ & $B_{lu}(\nu)$ have same frequency dependence as $A_{ul}(\nu)$.

Consider three possible radiative interactions between two levels u & l .

Spontaneous emission downward from u to l at a spontaneous emission rate = A_{ul}

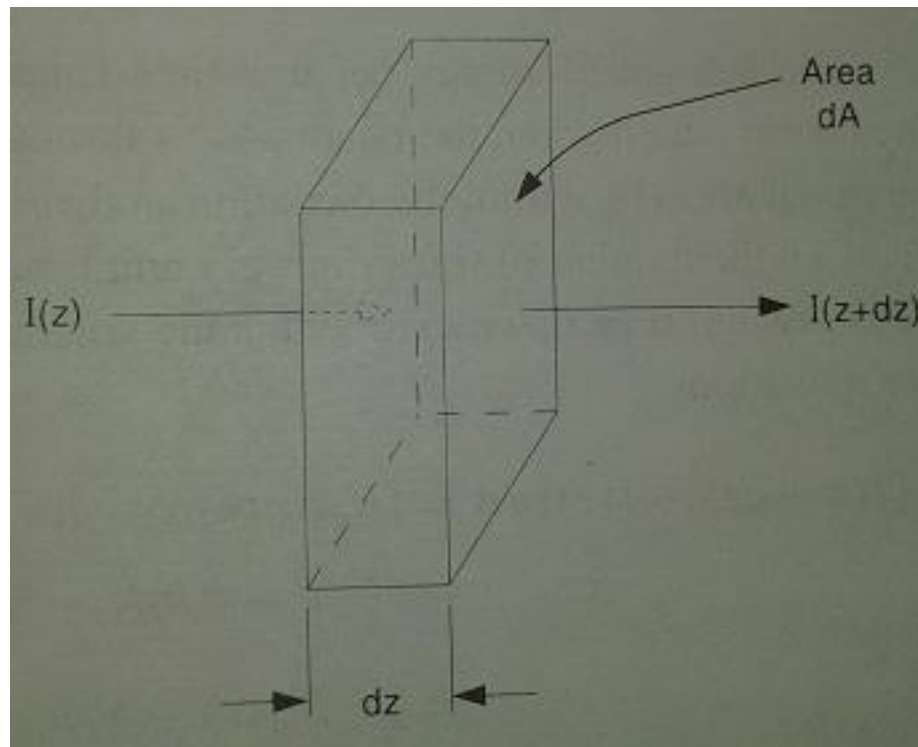
No. of transitions downward per unit volume per unit time at frequency ν for transition from u to l = $N_u A_{ul}$



Energy levels of a two-level system, depicting the population flux transferred between two levels via spontaneous emission, stimulated emission, & absorption.

Consider a small length dz of medium. In this case, we will estimate energy due to stimulated emissions.

Spontaneous emission contribution is neglected here because such emission is radiated in all directions into a solid angle 4π & thus contributes very little in direction of incident beam.



Light incident upon an incremental length dz of absorbing material.

Amount of energy per unit time added when beam passes through a region of length dz & cross-sectional area dA within medium can therefore be expressed as difference between no. of transitions per unit time upward & no. of transitions per unit time downward within volume, multiplied by photon energy per transition.

$$\begin{aligned}
 [I(z + dz) - I(z)]dA &= [N_u B_{ul}(\nu) \Delta\nu \cdot u(\nu) - N_l B_{lu}(\nu) \Delta\nu \cdot u(\nu)] h \nu dA dz \\
 &= [N_u B_{ul}(\nu) I \eta / c - N_l B_{lu}(\nu) I \eta / c] h \nu dA dz \\
 &= [N_u B_{ul}(\nu) - N_l B_{lu}(\nu)] \frac{h \nu I \eta}{c} dA dz
 \end{aligned}$$

Energy is added to or subtracted from beam in discrete amounts $h\nu$ as a result of the two terms within bracket.

Factor involving N_u leads to stimulated emission. Consequently, another photon is added to beam and N_u is decreased by 1 every time beam stimulates an electron from level u to l .

Hence, population within medium is transferred from level u to l . Factor containing N_l involves absorption of photons from beam with population moving from level l to u .

$$[I(z + dz) - I(z)]dA = [N_u B_{ul}(\nu) - N_l B_{lu}(\nu)] \frac{h\nu I \eta}{c} dA dz$$

Using following relation, $dI = I(z + dz) - I(z)$

Dividing through by dA & dz ,

$$\frac{dI}{dz} = [N_u B_{ul}(\nu) - N_l B_{lu}(\nu)] \frac{h\nu \eta}{c} I$$

It is a differential Eqn of form, $\frac{dI}{dz} = CI$

Solution of such Eqn, $I = I_0 e^{Cz}$

$$I = I_0 e^{g^H(\nu)z}$$

$$g^H(\nu) = [N_u B_{ul}(\nu) - N_l B_{lu}(\nu)] \frac{h\nu \eta}{c}$$

H – homogeneous broadening

Gain coefficient (dimension: 1/length)

Gain Coefficient & Stimulated Emission Cross-section for Homogeneous Broadening

$$g^H(\nu) = [N_u B_{ul}(\nu) - N_l B_{lu}(\nu)] \frac{h\nu\eta}{c}$$

Using following relationship of Einstein A & B coefficients,

$$\frac{A_{ul}(\nu)}{B_{ul}(\nu)} = \frac{8\pi h \eta^3 \nu^3}{c^3} \equiv \frac{A_{ul}}{B_{ul}}$$

$$\frac{g_l B_{lu}(\nu)}{g_u B_{ul}(\nu)} = 1 \quad \text{or} \quad g_l B_{lu}(\nu) = g_u B_{ul}(\nu)$$

Gain coefficient can be expressed as,

$$g^H(\nu) = \left[N_u - \frac{g_u}{g_l} N_l \right] \frac{c^2}{8\pi\eta^2\nu^2} A_{ul}(\nu)$$

$$A_{ul}(\nu) = \frac{\chi_{ul}^T / 4\pi^2}{(\nu - \nu_0)^2 + (\chi_{ul}^T / 4\pi)^2} A_{ul}$$

Gain coefficient,

$$g^H(\nu) = \left[N_u - \frac{g_u}{g_l} N_l \right] \frac{c^2}{8\pi\eta^2\nu^2} \left[\frac{\chi_{ul}^T / 4\pi^2}{(\nu - \nu_0)^2 + (\chi_{ul}^T / 4\pi)^2} \right] A_{ul}$$

Let us define,

$$\Delta N_{ul} = \left[N_u - \frac{g_u}{g_l} N_l \right]$$

$$\sigma_{ul}^H(\nu) = \frac{c^2}{8\pi\eta^2\nu^2} A_{ul}(\nu) = \frac{c^2}{8\pi\eta^2\nu^2} \left[\frac{\chi_{ul}^T / 4\pi^2}{(\nu - \nu_0)^2 + (\chi_{ul}^T / 4\pi)^2} \right] A_{ul}$$

This term is defined as stimulated emission cross-section & it has dimension of length² or area.

$$g^H(\nu) = \sigma_{ul}^H(\nu) \Delta N_{ul}$$

Since gain coefficient is highest at center of emission line frequency ($\nu = \nu_0$).

$$g^H(\nu = \nu_0) \equiv g^H(\nu_0) \equiv g_0^H$$

$$g_0^H = \left[N_u - \frac{g_u}{g_l} N_l \right] \frac{c^2}{2\pi\eta^2\nu_0^2\chi_{ul}} A_{ul}$$

Stimulated emission cross-section at center frequency ($\nu = \nu_0$)

$$\sigma_{ul}^H(\nu_0) = \frac{c^2 A_{ul}}{2\pi\eta^2\nu_0^2\chi_{ul}} = \frac{\lambda_{ul}^2 A_{ul}}{4\pi^2\eta^2\Delta\nu_{ul}^H}$$

$$\chi_{ul} = 2\pi\Delta\nu_{ul}^H$$

Relationship for intensity at a specific distance z into medium for homogeneous broadening:

$$I = I_0 e^{g^H(\nu)z} = I_0 e^{\sigma_{ul}^H(\nu)[N_u - (g_u/g_l)N_l]z}$$

$$I = I_0 e^{\sigma_{ul}^H(\nu)\Delta N_{ul}z}$$

Since $g^H(\nu)$, $\sigma_{ul}^H(\nu)$, & z are positive quantities, if ΔN_{ul} is positive

**Then beam intensity will increase exponentially with distance z & so provide amplification;
conversely, if ΔN_{ul} is negative then intensity of beam will decrease leading to absorption.**



How amplification occurs in a medium with energy levels u & l & population densities N_u & N_l in those levels?

Population inversion

$$I = I_0 e^{\sigma_{ul} [N_u - (g_u / g_l) N_l] z}$$

I_0 = Intensity of beam when it enters into medium

I = Intensity at some distance z into medium

If exponent is positive, intensity of beam will increase – amplification will occur.

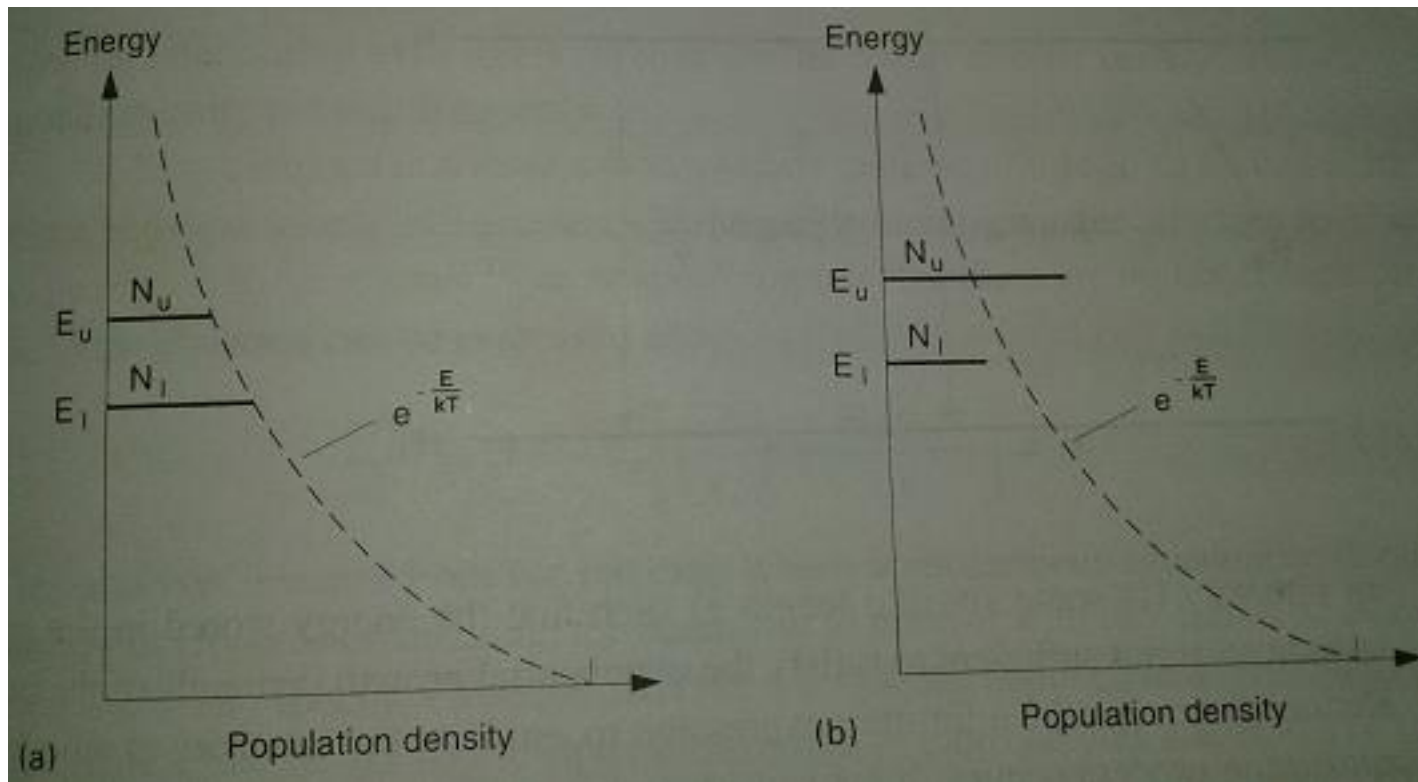
If exponent is negative, intensity of beam will decrease – absorption will occur.

Since values of σ_{ul} & z are always positive, the amplification will occur only if

$$N_u > \frac{g_u}{g_l} N_l$$

$$\frac{g_l N_u}{g_u N_l} > 1$$

The case of upper level being more populated than lower level is referred to as a population inversion. This relationship is not normal under conditions of thermal equilibrium.



Population distribution between levels u & l for conditions of (a) thermal equilibrium & (b) a population inversion.

A population inversion is a *necessary* condition for amplification or laser action to occur, but it is not a *sufficient* condition.

- ❖ Population inversion (necessary condition)
- ❖ Saturation intensity (sufficient condition)