

PH 201

OPTICS & LASERS

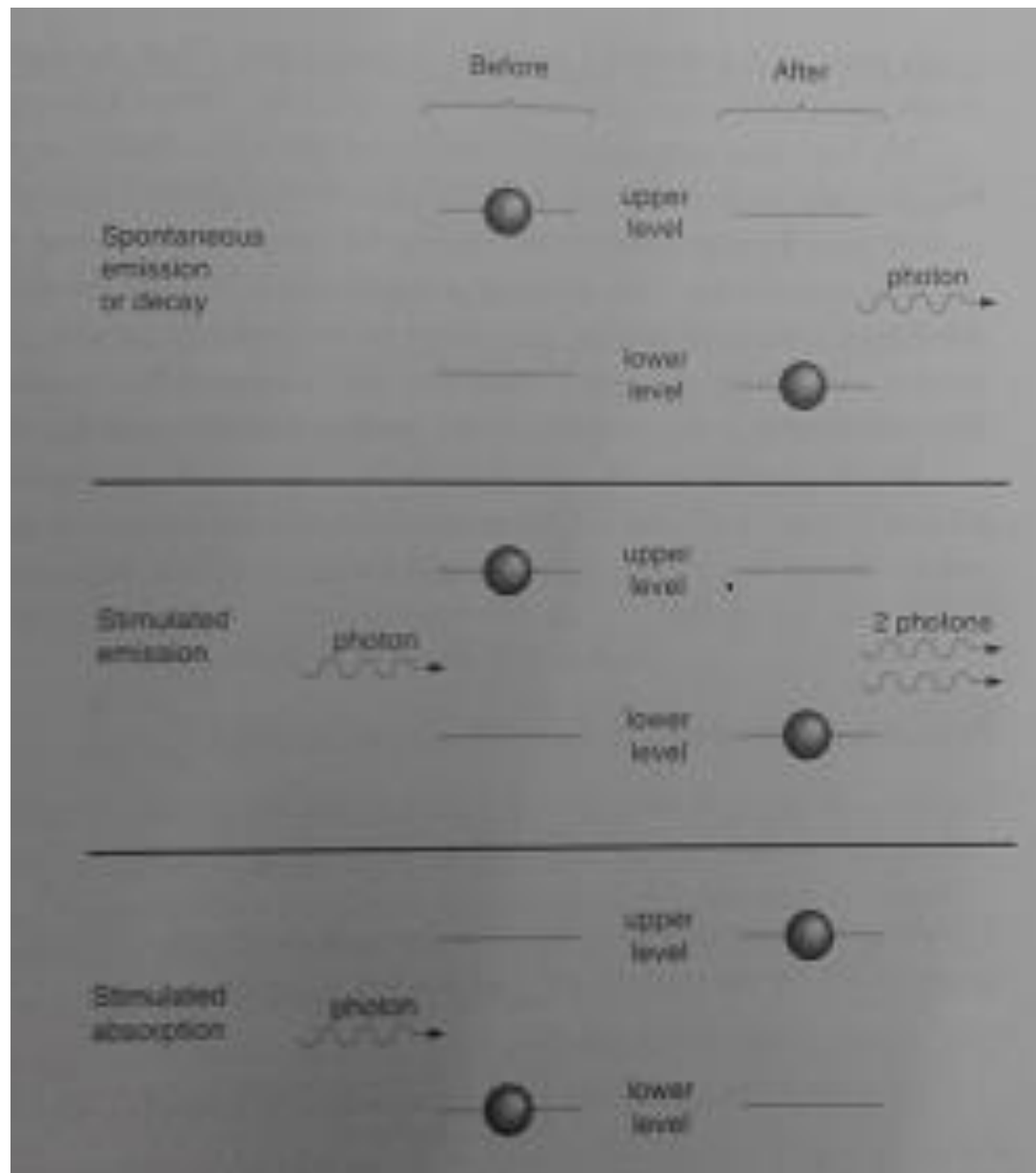
Lecture_Lasers_3

Absorption & Stimulated Emission Coefficients

Einstein suggested concept of stimulated emission in 1917.

- ❖ If a photon can stimulate an electron to move from a lower energy state l to higher energy state u by means of absorption, then a photon should also be able to stimulate an electron from same upper state u to lower state l .
- ❖ In case of absorption, photon disappears, with energy being transferred to absorbing species.
- ❖ In case of stimulated emission, the species would have to radiate an additional photon to conserve energy.
Such process must occur in order to keep population of two energy levels in thermal equilibrium.

Radiative processes producing interaction between two bound levels in a material: spontaneous emission, absorption, & stimulated emission.



Fundamental radiation processes associated with interaction of light with matter.

Consider a group of atoms having electrons occupying either energy levels u or l with population densities N_u & N_l (no. of atoms per unit volume).

Assuming atoms in thermal equilibrium with each other & must therefore be related by Boltzmann distribution,

$$\frac{N_u}{N_l} = \frac{g_u}{g_l} e^{-(E_u - E_l)/kT} = \frac{g_u}{g_l} e^{-\Delta E_{ul}/kT}$$

$$\Delta E_{ul} = E_u - E_l = h\nu_{ul}$$

Spontaneous transition probability, rate at which spontaneous transitions occur from level u to level l (no. per unit time) = A_{ul}

No. of spontaneous transitions from u to l per unit time per unit volume = $N_u A_{ul}$

Stimulation process would be proportional to photon energy density $u(\nu)$ at frequency ν_{ul}

Assuming proportionality constant for such stimulated transition B , then the upward flux – the no. of stimulated upward transitions per unit volume per unit time per unit frequency = $N_l B_{lu} u(\nu)$

Similarly downward flux = $N_u B_{ul} u(\nu)$

Constants A_{ul} , B_{ul} , & B_{lu} are referred to as Einstein A & B coefficients.

For populations N_u & N_l to be in radiative thermal equilibrium & for the principle of detailed balance to apply, downward radiative flux should equal the upward radiative flux between two levels.

$$N_u A_{ul} + N_u B_{ul} u(\nu) = N_l B_{lu} u(\nu)$$

$$u(\nu) = \frac{N_u A_{ul}}{N_l B_{lu} - N_u B_{ul}}$$

Dividing top & bottom terms in right hand side by N_u & then for ratio N_u/N_l

$$u(\nu) = \frac{A_{ul}}{B_{ul}} \left(\left[\frac{g_l B_{lu}}{g_u B_{ul}} \right] e^{h\nu_{ul}/kT} - 1 \right)^{-1}$$

$$u(\nu) = \frac{A_{ul}}{B_{ul}} \left(\left[\frac{g_l B_{lu}}{g_u B_{ul}} \right] e^{h\nu_{ul}/kT} - 1 \right)^{-1}$$

This Eq. has a familiar form with the following Eq.

Rayleigh-Jeans law: Energy density of radiation per unit volume

$$u(\nu) = \frac{8\pi h \eta^3 \nu^3}{c^3 (e^{h\nu/kT} - 1)} \quad \text{[Planck's law]}$$

Both Eqs. Concern radiation in thermal equilibrium, if true then they must be equivalent. Equivalence follows if

$$\frac{g_l B_{lu}}{g_u B_{ul}} = 1 \quad \text{or} \quad g_l B_{lu} = g_u B_{ul}$$

$$\frac{A_{ul}}{B_{ul}} = \frac{8\pi h \eta^3 \nu^3}{c^3}$$

Relationship between stimulated emission (B_{ul}) & absorption coefficients (B_{lu}).

$$u(\nu) = \frac{A_{ul}}{B_{ul}} \left(\left[\frac{g_l B_{lu}}{g_u B_{ul}} \right] e^{h\nu_{ul}/kT} - 1 \right)^{-1}$$

This Eq. has a familiar form with the following Eq.

Rayleigh-Jeans law: Energy density $u(\nu)$ of cavity radiation per unit volume within frequency ν to $\nu + d\nu$.

$$u(\nu) = \frac{8\pi\eta^3\nu^2}{c^3} kT$$

This result suggests that there is a continuous increase in energy density with frequency for a given temp T .

It suggests that energy density approaches infinity as frequency is increased.

It agrees with experiments for lower frequencies but does not predict experimentally observed maximum value for a given temp at higher frequencies.

Planck's law:

Planck explored the possibility of quantizing mode energy, postulating that an oscillator of frequency ν could have only discrete values $m h \nu$ of energy, where $m = 0, 1, 2, 3, \dots$

Planck referred to this unit of energy $h \nu$ as a quantum that could not be further divided.

$$u(\nu) = \frac{8\pi h \eta^3 \nu^3}{c^3 (e^{h\nu/kT} - 1)}$$

Both Eqs. Concern radiation in thermal equilibrium, if true then they must be equivalent. Equivalence follows if

$$\frac{g_l B_{lu}}{g_u B_{ul}} = 1 \quad \text{or} \quad g_l B_{lu} = g_u B_{ul}$$

$$\frac{A_{ul}}{B_{ul}} = \frac{8\pi h \eta^3 \nu^3}{c^3}$$

Relationship between stimulated emission (B_{ul}) & absorption coefficients (B_{lu}).

Relationship with spontaneous emission coefficient (A_{ul}).

$$B_{ul} = \frac{c^3}{8\pi h \eta^3 \nu^3} A_{ul}$$

Ratio of stimulated to spontaneous emission rates from level u .

$$\frac{B_{ul} u(\nu)}{A_{ul}} = \frac{1}{e^{h\nu_{ul}/kT} - 1}$$

Thus, stimulated emission plays a significant role only for temps in which kT is of, or greater than, the order of photon energy $h\nu_{ul}$.

Ratio is unity when

$$\frac{h\nu_{ul}}{kT} = \ln 2 = 0.693$$

For visible transitions, in green portion of spectrum (photons of order of 2.5 eV), such a relationship would be achieved for a temp of 33,500 K. Thus, in visible spectrum, dominance of stimulated emission over spontaneous emission normally happens only in stars, in high-temp & density laboratory plasmas such as laser produced plasma or in lasers.

In low-pressure plasmas the radiation can readily escape, so there is no opportunity for radiation density to build up to a value where stimulated decay rate is comparable to radiative decay rate.

In lasers, the ratio can be significantly greater than unity.

- ❖ A He-Ne laser operating at 632.8 nm has an output power 1.0 mW with 1.0 mm beam diameter. The beam passes through a mirror that has 99% reflectivity and 1% transmission at the laser wavelength. For this laser calculate the ratio

$$\frac{B_{ul}u(\nu)}{A_{ul}}$$

What is the effective blackbody temp of the laser beam as it emerges from the laser output mirror?

Assume beam diameter is 1.0 mm inside the laser cavity & that power is uniform over beam cross-section. Also, laser linewidth is approximately one tenth of Doppler width for transition.

Laser frequency

$$\nu = \frac{c}{\lambda} = \frac{3 \times 10^8 \text{ m/s}}{6.328 \times 10^{-7} \text{ m}} = 4.74 \times 10^{14} \text{ Hz}$$

$$\frac{A_{ul}}{B_{ul}} = \frac{8\pi\hbar\eta^3\nu^3}{c^3} = \frac{(8\pi)(6.63 \times 10^{-34} \text{ J-s})(1)^3(4.74 \times 10^{14} \text{ Hz})^3}{(3 \times 10^8 \text{ m/s})^3}$$

$$= 6.57 \times 10^{-14} \text{ J-s/m}^3$$

$$\frac{B_{ul}}{A_{ul}} = 1.52 \times 10^{13} \text{ m}^3 / \text{J-s}$$

Energy density $u(\nu)$ is related to intensity per unit frequency. Intensity $I(\nu)$ can be obtained by dividing laser beam power in cavity by beam cross-sectional area & frequency width of beam.

Power of beam within cavity traveling toward output mirror must be 100 mW & that reflected would be 99 mW (1 mW passes through mirror). Thus, total power in cavity is 199 mW.

Doppler width of He-Ne 632.8 nm transition is $1.5 \times 10^9 \text{ Hz}$.

$$u(\nu) = \frac{I(\nu)}{c} = \frac{[(199 \times 1.0 \text{ mW}) / (\pi \cdot (5 \times 10^{-4} \text{ m})^2)] / (0.1)(1.5 \times 10^9 \text{ Hz})}{3 \times 10^8 \text{ m/s}}$$

$$= 5.63 \times 10^{-12} \text{ J-s/m}^3$$

$$\frac{B_{ul}u(\nu)}{A_{ul}} = (1.52 \times 10^{13} \text{ m}^3 / \text{J} - \text{s})(5.63 \times 10^{-12} \text{ J} - \text{s} / \text{m}^3)$$

$$= 85.6$$

The stimulated emission rate is therefore almost 86 times the spontaneous emission rate or transitions from upper to lower laser level at 632.8 nm.

$$\frac{B_{ul}u(\nu)}{A_{ul}} = \frac{1}{e^{h\nu_{ul}/kT} - 1} = 85.6$$

$$e^{h\nu_{ul}/kT} - 1 = 1/85.6 = 0.0117$$

$$e^{h\nu_{ul}/kT} = \ln(1.0117) = 1.16 \times 10^{-2}$$

$$T = \frac{h\nu_{ul}}{(1.16 \times 10^{-2})k} = \frac{(6.63 \times 10^{-34} \text{ J} - \text{s})(4.74 \times 10^{14} \text{ Hz})}{(1.16 \times 10^{-2})(1.38 \times 10^{-23} \text{ J} / \text{K})}$$

$$= 1.96 \times 10^6 \text{ K} = 19,60,000 \text{ K}$$

Radiation intensity inside laser cavity has a value equivalent to that of a nearly 20,00,000 K blackbody – if we consider only radiation emitted from blackbody in frequency (or wavelength interval) over which laser operates.

Einstein A & B coefficients are associated with interaction of radiation with two specific energy levels, where radiation has exact frequency corresponding to energy separation between two levels.



How gain (amplification) & absorption of radiation can occur in a medium containing populations in those two levels?