

解:

阱内的归一化波函数 (归一化) 为:

$$\psi_n(x) = \sin[k_n(x+a)]$$

$$\text{其中 } k_n \cdot 2a = n \cdot \pi$$

$$\text{即 } k_n = \frac{n\pi}{2a} \quad n = 1, 2, \dots$$

归一化因子:

$$\int_{-a}^a |\psi_n(x)|^2 dx = \int_{-a}^a \sin^2[k_n(x+a)] dx$$

$$= \int_{-a}^a \left\{ \frac{1}{2} - \frac{1}{2} \cos[2k_n(x+a)] \right\} dx$$

$$= a$$

\Rightarrow 归一化后的波函数为:

$$\psi_n(x) = \frac{1}{\sqrt{a}} \sin[k_n(x+a)]$$

波粒能量为:

$$\hat{H}' = \begin{cases} V_0 & -a < x < 0, \quad 0 < t < T \\ 0 & \text{otherwise} \end{cases}$$

故 $H'_{21} = \langle \psi_2 | \hat{H}' | \psi_1 \rangle$

$$= \int_{-\infty}^{+\infty} \psi_2^*(x) \hat{H}' \psi_1(x) dx$$

$$= \frac{V_0}{a} \int_{-a}^0 \sin\left[\frac{\pi}{a}(x+a)\right] \sin\left[\frac{\pi}{2a}(x+a)\right] dx$$

$$= \frac{2V_0}{a} \int_{-a}^0 \sin^2\left[\frac{\pi}{2a}(x+a)\right] \cos\left[\frac{\pi}{2a}(x+a)\right] dx$$

$$= \frac{2a}{\pi} \frac{2V_0}{a} \int_{-a}^0 \sin^2\left[\frac{\pi}{2a}(x+a)\right] d\sin\left[\frac{\pi}{2a}(x+a)\right]$$

$$= \frac{4V_0}{\pi} \int_0^1 x^2 dx = \frac{4V_0}{3\pi}$$

故当 $t \leq T$ 时, 从 $|\psi_1\rangle$ 到 $|\psi_2\rangle$ 的几率为

$$W_{1 \rightarrow 2}(t) = \frac{1}{\hbar^2} |H'_{21}|^2 \frac{\sin^2(\omega_{12}t/2)}{(\omega_{12}/2)^2}$$

$$\text{其中 } \omega_{12} = \frac{E_1 - E_2}{\hbar} = \frac{\hbar^2(k_1^2 - k_2^2)}{2m\hbar}$$

$$= -\frac{\hbar}{2m} \frac{3\pi^2}{4a^2}$$

故 T 时刻,

$$W_{1 \rightarrow 2}(T) = \frac{1}{\hbar^2} \left(\frac{4V_0}{3\pi} \right)^2 \frac{\sin^2 \left(\frac{\hbar}{4m} \frac{3\pi^2}{4a^2} T \right)}{\left(\frac{\hbar}{4m} \frac{3\pi^2}{4a^2} \right)^2}$$

$$= \left[\frac{4V_0}{\hbar} \frac{4m}{3\pi} \frac{4a^2}{\hbar 3\pi^2} \sin \left(\frac{\hbar}{4m} \frac{3\pi^2}{4a^2} T \right) \right]^2$$

$$= \left[\frac{64V_0 m a^2}{9 \hbar^2 \pi^3} \sin \left(\frac{3\hbar \pi^2}{16 m a^2} T \right) \right]^2$$

T 时刻以后, 微扰为 0, $|\psi_1\rangle, |\psi_2\rangle$ 又

成为系统的本征态 (定态波函数), 故

处于 $|\psi_2\rangle$ 的几率不再变化, 为 $W_{1 \rightarrow 2}(T)$