

# 角动量算符

$$\left. \begin{aligned} x &= r \sin \theta \cos \varphi, & y &= r \sin \theta \sin \varphi, & z &= r \cos \theta; \\ r^2 &= x^2 + y^2 + z^2, & \cos \theta &= \frac{z}{r}, & \tan \varphi &= \frac{y}{x}. \end{aligned} \right\} \text{球坐标系}$$

$$\left. \begin{aligned} \hat{L}_x &= y\hat{p}_z - z\hat{p}_y = \frac{\hbar}{i} \left( y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right), \\ \hat{L}_y &= z\hat{p}_x - x\hat{p}_z = \frac{\hbar}{i} \left( z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right), \\ \hat{L}_z &= x\hat{p}_y - y\hat{p}_x = \frac{\hbar}{i} \left( x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right). \end{aligned} \right\}$$

$$\left. \begin{aligned} \hat{L}_x &= i\hbar \left( \sin \varphi \frac{\partial}{\partial \theta} + \cot \theta \cos \varphi \frac{\partial}{\partial \varphi} \right), \\ \hat{L}_y &= -i\hbar \left( \cos \varphi \frac{\partial}{\partial \theta} - \cot \theta \sin \varphi \frac{\partial}{\partial \varphi} \right), \\ \hat{L}_z &= -i\hbar \frac{\partial}{\partial \varphi}; \end{aligned} \right\}$$

$$\begin{aligned} \hat{L}^2 &= \hat{L}_x^2 + \hat{L}_y^2 + \hat{L}_z^2 = -\hbar^2 \left[ \left( y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right)^2 \right. \\ &\quad \left. + \left( z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right)^2 + \left( x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right)^2 \right]. \end{aligned}$$

$$\hat{L}^2 = -\hbar^2 \left[ \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \varphi^2} \right].$$

# 角动量算符

角动量算符的各分量互相不对易！

$L_x$   $L_y$   $L_z$ 不可同时观测

但 $\hat{L}^2$ 和 $\hat{L}_z$ 可以同时观测（注意 $z$ 的选取是任意的）

# 角动量算符

$\hat{L}^2$ 和 $\hat{L}_z$ 的共同本征值与本征函数:

$$\hat{L}^2 Y_{lm}(\theta, \varphi) = l(l+1)\hbar^2 Y_{lm}(\theta, \varphi)$$

$$\hat{L}_z Y_{lm}(\theta, \varphi) = m\hbar Y_{lm}(\theta, \varphi)$$

$$Y_{lm}(\theta, \varphi) = N e^{im\varphi} P_\ell^m(\cos\theta) \quad \text{球谐函数}$$

下面列出前面几个球谐函数:

$$Y_{0,0} = \frac{1}{\sqrt{4\pi}},$$

$$Y_{1,1} = -\sqrt{\frac{3}{8\pi}} \sin\theta e^{i\varphi} = -\sqrt{\frac{3}{8\pi}} \frac{x+iy}{r},$$

$$Y_{1,0} = \sqrt{\frac{3}{4\pi}} \cos\theta = \sqrt{\frac{3}{4\pi}} \frac{z}{r},$$

$$Y_{1,-1} = \sqrt{\frac{3}{8\pi}} \sin\theta e^{-i\varphi} = \sqrt{\frac{3}{8\pi}} \frac{x-iy}{r},$$

$$Y_{2,2} = \sqrt{\frac{15}{32\pi}} \sin^2\theta e^{2i\varphi} = \sqrt{\frac{15}{32\pi}} \left(\frac{x+iy}{r}\right)^2,$$

角动量的量子化 (离散化)

$$l = 0, 1, 2, \dots$$

$$m = -l, -(l-1), \dots, (l-1), l$$

$\hat{L}^2$ 算符与本征值 $l(l+1)\hbar^2$ 对应的本征态的简并度:  $2l+1$

$$Y_{2,1} = -\sqrt{\frac{15}{8\pi}} \sin\theta \cos\theta e^{i\varphi} = -\sqrt{\frac{15}{8\pi}} \frac{(x+iy)z}{r^2},$$

$$Y_{2,0} = \sqrt{\frac{5}{16\pi}} (3\cos^2\theta - 1) = \sqrt{\frac{5}{16\pi}} \frac{(2z^2 - x^2 - y^2)}{r^2},$$

$$Y_{2,-1} = \sqrt{\frac{15}{8\pi}} \sin\theta \cos\theta e^{-i\varphi} = \sqrt{\frac{15}{8\pi}} \frac{(x-iy)z}{r^2},$$

$$Y_{2,-2} = \sqrt{\frac{15}{32\pi}} \sin^2\theta e^{-2i\varphi} = \sqrt{\frac{15}{32\pi}} \left(\frac{x-iy}{r}\right)^2.$$

## 2.6 自旋与二能级体系

### 1. 电子自旋

自旋表征了电子的内部状态，是独立于位置的运动变量

$$[S_x, S_y] = i\hbar S_z, \quad [S_y, S_z] = i\hbar S_x, \quad [S_z, S_x] = i\hbar S_y, \quad [S_z, S^2] = 0$$

对于电子：  $S^2 = \frac{1}{2}(\frac{1}{2} + 1)\hbar^2$ ,  $S_z = \pm \frac{1}{2}\hbar$ ,  $\hat{S}^2$ 和 $S_z$ 的共同本征态是电子内部状态的一组基矢

如果只考虑自旋，电子状态是一个二维矢量空间（Hilbert空间）

## 2.6 自旋与二能级体系

### 2. 泡利矩阵

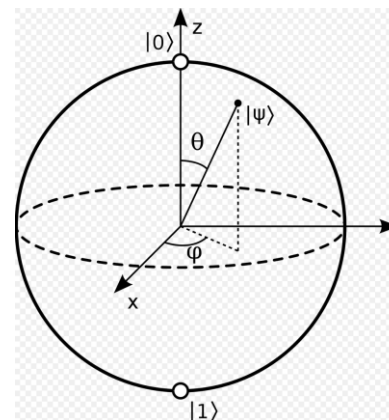
- ① 给定维度的厄米矩阵构成一个实矢量空间
- ②  $2 \times 2$ 厄米矩阵空间的基:  $I = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$ ,  $\sigma_x = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$ ,  $\sigma_y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$ ,  $\sigma_z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$
- ③  $\sigma_x^2 = \sigma_y^2 = \sigma_z^2 = -i\sigma_x\sigma_y\sigma_z = I$
- ④  $[\sigma_x, \sigma_y] = 2i\sigma_z$ ,  $[\sigma_y, \sigma_z] = 2i\sigma_x$ ,  $[\sigma_z, \sigma_x] = 2i\sigma_y$ ,  $\{\sigma_x, \sigma_y\} = \{\sigma_y, \sigma_z\} = \{\sigma_z, \sigma_x\} = 0$
- ⑤  $S_x = \frac{\hbar}{2}\sigma_x$ ,  $S_y = \frac{\hbar}{2}\sigma_y$ ,  $S_z = \frac{\hbar}{2}\sigma_z$
- ⑥ 验证:  $\sigma_x, \sigma_y, \sigma_z$ 的本征值是 $\pm 1$ , 对应于 $S_x, S_y, S_z$ 只能取 $\pm\hbar/2$ 两个本征值
- ⑦ 本征态:  $\psi_x^+ = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ 1 \end{bmatrix}$ ,  $\psi_x^- = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -1 \end{bmatrix}$ ,  $\psi_y^+ = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ i \end{bmatrix}$ ,  $\psi_y^- = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -i \end{bmatrix}$ ,  $\psi_z^+ = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$ ,  $\psi_z^- = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$
- ⑧  $\vec{\sigma}$ 矢量

## 2.6 自旋与二能级体系

### 3. Bloch球面

- ① 自旋算符 (矢量) :  $\mathbf{S} = \frac{\hbar}{2} \boldsymbol{\sigma} = \frac{\hbar}{2} (\sigma_x, \sigma_y, \sigma_z)$
- ② 任意方向矢量:  $\mathbf{n} = (n_x, n_y, n_z) = (\sin\theta\cos\varphi, \sin\theta\sin\varphi, \cos\theta)$
- ③ 自旋在该方向的投影:  $S_n = \mathbf{S} \cdot \mathbf{n}$

- ④ 本征值  $\pm\hbar/2$ , 其中  $\hbar/2$  对应的本征态为 
$$\begin{bmatrix} e^{-i\frac{\varphi}{2}} \cos \frac{\theta}{2} \\ e^{i\frac{\varphi}{2}} \sin \frac{\theta}{2} \end{bmatrix}$$



### 4. 拉莫进动

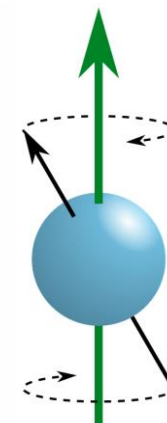
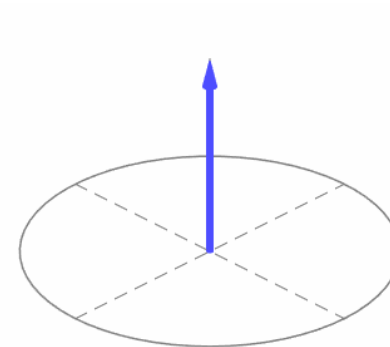
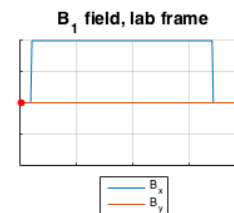
- ① 自旋与磁矩的关系:  $\boldsymbol{\mu} = -g_s \mu_B \frac{\mathbf{S}}{\hbar} = -g_s \mu_B \frac{\hbar}{2} \frac{\boldsymbol{\sigma}}{\hbar} = -\frac{1}{2} g_s \mu_B \boldsymbol{\sigma}$

- ② 外磁场下的哈密顿量:  

$$\hat{H} = -\boldsymbol{\mu} \cdot \mathbf{B} = \frac{1}{2} g_s \mu_B \mathbf{B} \cdot \boldsymbol{\sigma} = \frac{1}{2} g_s \mu_B B \sigma_z \equiv \frac{1}{2} \hbar \omega \sigma_z \text{ (磁场方向为z)}$$

- ③ 解薛定谔方程: 任意方向的自旋状态 
$$\begin{bmatrix} e^{-i\frac{\varphi+\omega t}{2}} \cos \frac{\theta}{2} \\ e^{i\frac{\varphi+\omega t}{2}} \sin \frac{\theta}{2} \end{bmatrix}, \quad \omega = \frac{g_s \mu_B}{\hbar} B$$

- ④ 经典理解; 与陀螺进动的类比



## 2.6 自旋与二能级体系

### 5. 自旋的历史

- ① 泡利：电子的能态数应该乘以2，与某种隐藏的旋转有关
- ② 1925 乌伦贝克、古兹密特：这种隐藏的旋转可以认为对应于电子的自转
- ③ 1928 狄拉克：狄拉克方程
  - a. 相对论质能关系：
$$E^2 = m^2 c^4 + p^2 c^2, \quad E = mc^2 \left(1 + \frac{p^2}{m^2 c^2}\right)^{1/2} \cong mc^2 \left(1 + \frac{p^2}{2m^2 c^2}\right) = mc^2 + \frac{p^2}{2m} \sim \frac{p^2}{2m}$$
  - b. 薛定谔方程： $\hat{H} = i\hbar \frac{\partial}{\partial t}, \quad \hat{p}_j = -i\hbar \frac{\partial}{\partial r_j}$  加上相对论质能关系
  - c. 强行要求一阶方程的形式关于时间和空间对称：方程的系数成为矩阵，态矢空间具有内部维度

## 2.6 自旋与二能级体系

### 6. 包含自旋的量子态

- ① 态矢空间维度加倍，算符对应的矩阵阶数也加倍

位置表象下：波函数： $\psi'(x, s_z) = \begin{bmatrix} \psi_1(x) \\ \psi_2(x) \end{bmatrix}$

- ② 自旋与位置没有耦合（自旋轨道耦合）时：哈密顿算符对应的矩阵是直积（克罗内克积）

$$\hat{H}' = \begin{bmatrix} \hat{H}'_{11} & \hat{H}'_{12} \\ \hat{H}'_{21} & \hat{H}'_{22} \end{bmatrix} = (aI + b\sigma_x + c\sigma_y + d\sigma_z) \otimes \hat{H}$$

从而波函数也可以表达成直积的形式  $\psi'(x, s_z) = \begin{bmatrix} s_1 \\ s_2 \end{bmatrix} \psi(x)$

- ③ 自旋轨道耦合简介
- ④ 存在自旋轨道耦合的情况：例如Rashba Hamiltonian













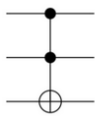
$$\hat{H}_R = \alpha(\boldsymbol{\sigma} \times \boldsymbol{p}) \cdot \hat{\boldsymbol{z}}$$



# 2.6 自旋与二能级体系

## 7. 量子计算简介

- ① 经典比特与逻辑
- ② n qubit ~ 2<sup>n</sup> 经典比特  
代价：结果是几率的
- ③ 量子逻辑门
- ④ 量子纠缠与测量
- ⑤ 消相干

Operator	Gate(s)		Matrix
Pauli-X (X)			$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$
Pauli-Y (Y)			$\begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$
Pauli-Z (Z)			$\begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$
Hadamard (H)			$\frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$
Phase (S, P)			$\begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}$
$\pi/8$ (T)			$\begin{bmatrix} 1 & 0 \\ 0 & e^{i\pi/4} \end{bmatrix}$
Controlled Not (CNOT, CX)			$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$
Controlled Z (CZ)	 		$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}$
SWAP	 		$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$
Toffoli (CCNOT, CCX, TOFF)			$\begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{bmatrix}$

Common quantum logic gates by name (including abbreviation), circuit form(s) and the corresponding unitary matrices.

## 2.6 自旋与二能级体系

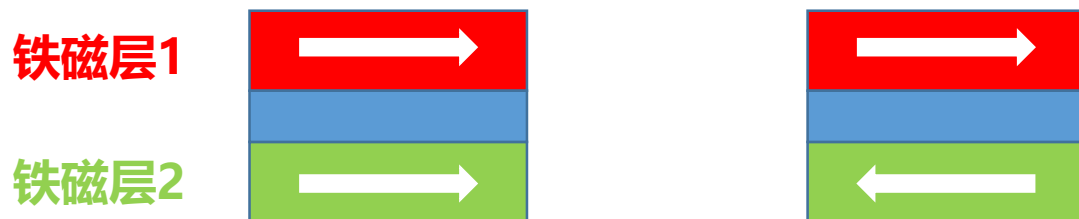
### 8. 自旋电子学简介

- ① 电子自旋导致宏观磁矩
- ② 宏观磁矩的交换相互作用导致磁性
- ③ 各向异性与剩磁——非易失特性

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- ④ 巨磁电阻：Fert & Grunberg 1988



### The Nobel Prize in Physics 2007



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**Albert Fert**  
Prize share: 1/2



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**Peter Grünberg**  
Prize share: 1/2

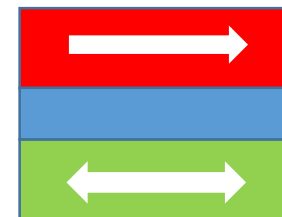
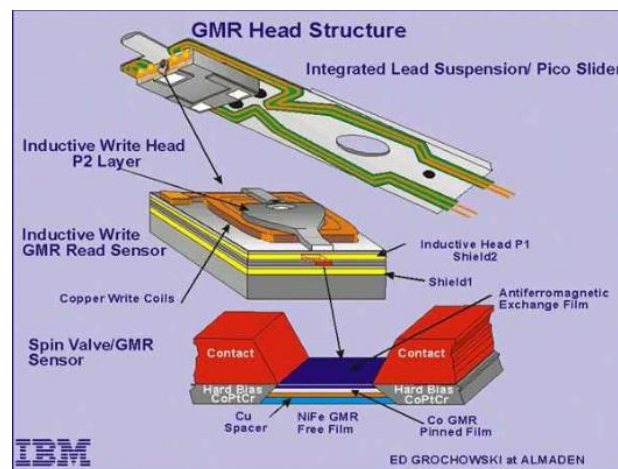
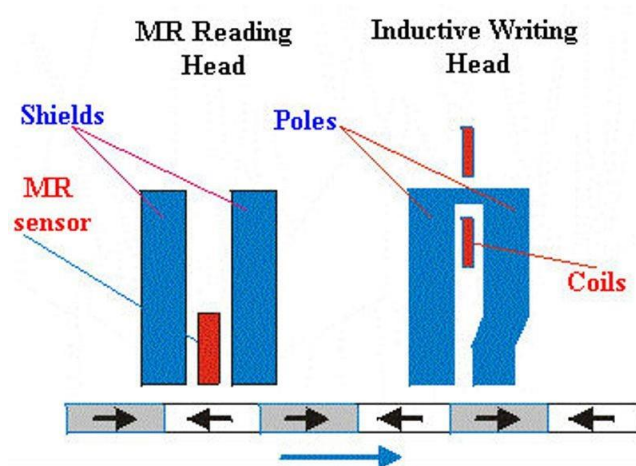
## 2.6 自旋与二能级体系

### 8. 自旋电子学简介



高密度机械硬盘

2019全球硬盘市值  
约600亿美元



## 2.6 自旋与二能级体系

### 8. 自旋电子学简介

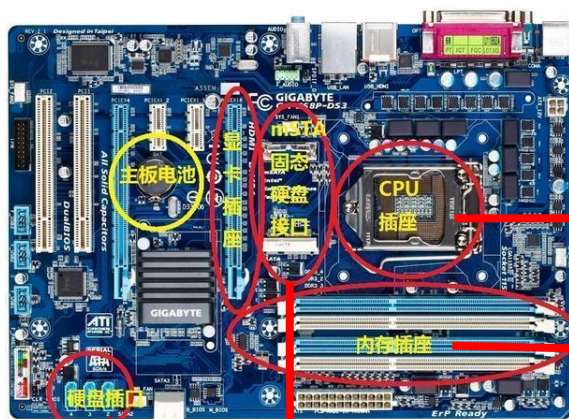
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- ④ 巨磁电阻：Fert & Grunberg 1988
- ⑤ STT & SOT
- ⑥ MRAM

## 2.6 自旋与二能级体系

### 8. 自旋电子学简介



黄页88网  
www.huangye88.com



SRAM 静态随机存储器

DRAM 动态随机存储器

SSD: 浮栅存储器

读写速度快

数据断电保持: 非易失

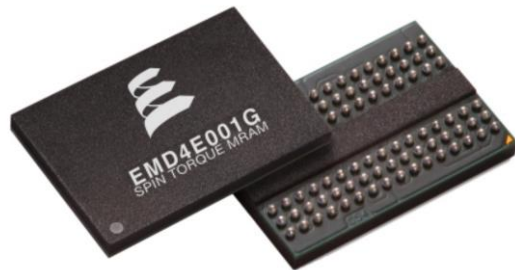
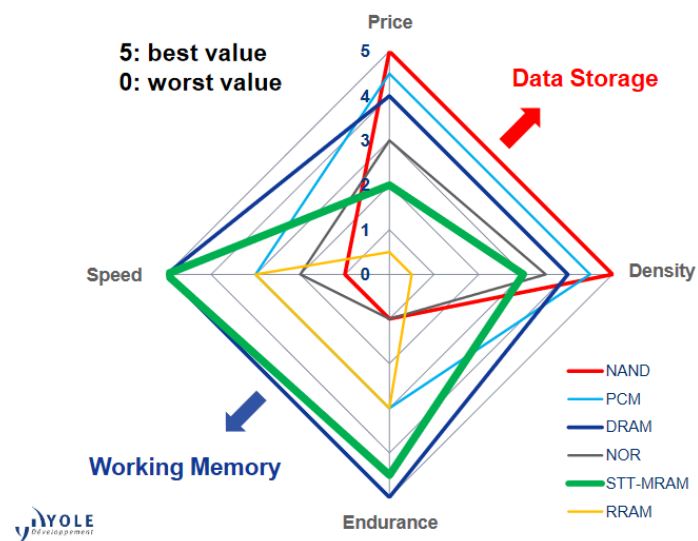
HDD: 机械硬盘

自旋电子器件

MRAM

## 2.6 自旋与二能级体系

### 8. 自旋电子学简介



MRAM IP and Design	 
Embedded MRAM manufacturers	Players in mass production or close to mass production 
Stand-alone MRAM manufacturers	<div>40nm, up to 128Mb  </div> <div>40nm, 28nm (256M, 1Gb)  </div> <div>Toggle Manufacturing  </div> <div>Expected: 28nm, 22nm  </div>

**STT-MRAM**  
预期到2029年全球市值达到200亿美元