

Nonlinear physics, dynamical systems and chaos theory



jean-christophe. loiseau@ensam. eu DynFluid, Arts et Métiers ParisTech, France





Flows on the line





An apparently simple system

▶ Let us consider the following first-order dynamical system

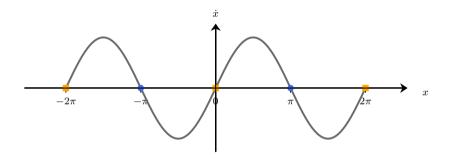
$$\dot{x} = \sin(x)$$
.

Its analytical solution is given by

$$t = \ln \left| \frac{\csc(x_0) + \cot(x_0)}{\csc(x) + \cot(x)} \right|.$$



Phase line



Phase line of the first-order dynamical system considered.





Fixed points

ightharpoonup Fixed points x^* are equilibrium solutions characterized by

$$f(x^*) = 0.$$

▶ In the present case, these are given by

$$x^* = n\pi$$
 for $n \in \mathbb{N}$.



Linear stability

▶ The dynamics of a perturbation $\eta(t) = x(t) - x^*$ is given by

$$\dot{\eta} = f(x^* + \eta).$$

lacktriangleq If η is small enough, $f(x^*+\eta)$ can be approximated by its first-order Taylor expansion around x^*

$$f(x^* + \eta) = f(x^*) + f'(x^*)\eta + \mathcal{O}(\eta^2).$$



Linear stability

▶ Given that $f(x^*) = 0$, the dynamics of η are governed by

$$\dot{\eta} = f'(x^*)\eta.$$

▶ Its analytical solution is given by

$$\eta(t) = \exp\left(f'(x^*)t\right)\eta_0.$$

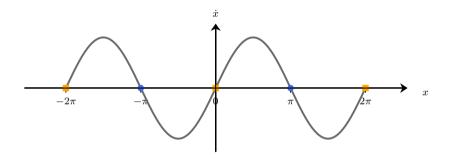


Linear stability

- ▶ The linear stability of a fixed point x^* is determined by the sign of $f'(x^*)$:
 - \hookrightarrow if $f'(x^*) > 0$, $\eta(t)$ growths exponentially fast. The fixed point is said to be linearly unstable.
 - \hookrightarrow if $f'(x^*) < 0$, $\eta(t)$ decays exponentially fast. The fixed point is said to be linearly stable.
 - \hookrightarrow if $f'(x^*) = 0$, one can not conclude and nonlinear analyses are required.
- Let us now re-analyze our original system and sketch the evolution of x(t).



Phase line

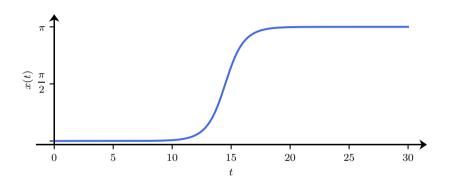


Phase line of the first-order dynamical system considered.





Evolution of x(t)



Evolution of x(t) for the initial condition $x_0 = 10^{-6}$.





Warning!

For a first-order system, the trajectories can only vary monotonically: either they end up on a stable fixed point, or they diverge to $\pm \infty$.



Oscillators, but not only...





Second-order systems are dynamical systems which can be described by

$$\dot{x} = f(x, y)$$

$$\dot{y} = g(x, y).$$

Having two degrees of freedom, they can exhibit dynamics much richer than simple first-order systems.



Our working example

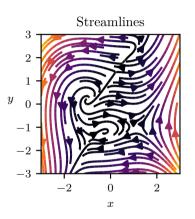
For the rest of this section, let us consider the following system

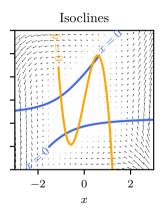
$$\dot{x} = x - y^2 + 1.28 + 1.4xy$$
$$\dot{y} = 0.2y - x + x^3.$$

▶ Note that this system is considered only for illustration purposes. To the best of my knowledge, it does not model any particular physics.



Phase plane and isoclines





Phase plane and isoclines (i.e. $\dot{x}=0$ and $\dot{y}=0$) of the system considered.





Phase plane and isoclines

- For a second-order system, phase plane and isoclines are fairly easy to plot and provide valuable insights into the dynamics of the system.
 - → Phase plane : general overview of the dynamics of the system.
 - Direct visualization of the stable and unstable fixed point of the system.

Question

How to compute these fixed points?





Interlude

How to compute fixed points?





How to compute fixed points?

Different techniques

- ► Fixed points are structuring the phase space of the dynamical system under scrutiny. Unfortunately, it may not be easy (nor possible) to compute them analytically.
- ► A number of different numerical techniques exist for that purpose. The following list is by no means exhaustive:
 - → Newton-Raphson method.
 - → Selective Frequency Damping,
 - → BoostConv.
 - \hookrightarrow ...



$$f: \mathbb{R} \to \mathbb{R}$$

Originally proposed by Isaac Newton (1645–1727) and Joseph Raphson (1648 – 1715) to solve

$$f(x) = 0.$$

Given an initial guess x_0 , the idea is to approximate f(x) by its first-order Taylor expansion around x_0 , i.e.

$$f(x) \simeq f(x_0) + f'(x_0)(x - x_0).$$

A better estimate x_1 of the root of f can then be obtained by solving

$$0 = f(x_0) + f'(x_0)(x_1 - x_0).$$



$$f: \mathbb{R} \to \mathbb{R}$$

 \triangleright After k iterations, the basic iteration scheme can be written as

$$x_{k+1} = x_k - \frac{f(x_k)}{f'(x_k)}.$$

▶ The iterative procedure stops when a user-defined criterion is fulfilled, usually

$$||f(x_k)|| \le \epsilon \text{ or } ||x_{k+1} - x_k|| \le \epsilon,$$

with $\epsilon \simeq 10^{-10}$.



 $f: \mathbb{R} \to \mathbb{R}$

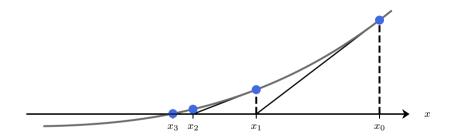


Illustration of the Newton-Raphson for $f(x) = x^3 - 2x - 5$ and $x_0 = 3.8$.





$$oldsymbol{f}: \mathbb{R}^n
ightarrow \mathbb{R}^n$$

- lackbox Generalization of the Newton-Raphson method to the case $f:\mathbb{R}^n o\mathbb{R}^n$ is quite straightforward.
- ightharpoonup Given an estimate x_k , the basic iteration reads

$$egin{aligned} oldsymbol{J}\deltaoldsymbol{x} &= -oldsymbol{f}(oldsymbol{x}_k) \ oldsymbol{x}_{k+1} &= oldsymbol{x}_k + \deltaoldsymbol{x}, \end{aligned}$$

where J is the Jacobian matrix of f evaluated at x_k .



Limitations

- ▶ Although efficient, Newton-Raphson method suffers from a number of limitations:
 - \hookrightarrow The fixed points computed may depend on the initial guess x_0 .
 - \hookrightarrow Evaluating f(x) might be computationally expensive.
 - \hookrightarrow At each iteration, the Jacobian matrix J needs to be evaluated and inverted ($\mathcal{O}(n^3)$ operations).
- ► A number of variants of the Newton-Raphson method exist as to address these different limitations. This however is algorithmic refinement beyond the scope of the present course.



Back to our example





Fixed point computation

- ▶ Visual inspection of the isoclines has revealed that the system exhibits six different fixed points.
- ▶ In the rest, we will use the Newton-Raphson method to compute these different fixed points.
- ▶ The state-dependent Jacobian matrix required for Newton method reads

$$J = \begin{bmatrix} 1 + 1.4y & -2y + 1.4x \\ -1 + 3x^2 & 0.2 \end{bmatrix}$$



Infinitesimal perturbations

The dynamics of an infinitesimal perturbation q evolving in the vicinity of a given fixed point x^* is given by

$$\dot{q} = Lq$$
,

where L is the Jacobian matrix of the system evaluated at $x=x^*$.

The analytical solution to this linear time-invariant dynamical system is given by

$$\mathbf{q}(t) = e^{\mathbf{L}t}\mathbf{q}_0,$$

where q_0 is the initial condition.



Interlude

Some elements of linear algebra





Matrix decompositions

Stewart (Comput. Sci. Engrg., 2:50-59, 2000) has listed the big six matrix decompositions:

- 1. Cholesky decomposition,
- 2. Pivoted LU decomposition,
- 3. QR decomposition,
- 4. Spectral decomposition, also known as Eigendecomposition,
- 5. Schur decomposition,
- 6. Singular Value decomposition.





Eigendecomposition of a matrix

- ▶ Eigendecomposition is the factorization of a matrix into its canonical form.
 - → The matrix is represented in terms of its eigenvalues and eigenvectors.
- \blacktriangleright Given a $n \times n$ (i.e. square) matrix L, eigendecomposition aims to rewrite it as

$$L = V\Lambda V^{-1},$$

where Λ is the diagonal matrix of eigenvalues and V is a $n \times n$ matrix whose ith column if the eigenvector v_i associated to $\Lambda_{ii} = \lambda_i$.



Matrix exponential

Warning, unless the matrix is diagonal,

$$\left(e^{\boldsymbol{L}}\right)_{ij} \neq e^{l_{ij}}$$



Matrix exponential

The infinitie series Taylor expansion of the classical exponential function $f(x) = e^x$ is given by

$$f(x) = 1 + x + \frac{x^2}{2!} + \frac{x^3}{3!} + \cdots$$

The same definition carries over for matrices.

$$e^{L} = I + L + \frac{L^{2}}{2!} + \frac{L^{3}}{3!} + \cdots,$$

where I is the identity matrix.



Matrix exponential

- ightharpoonup Evaluating e^L by brute force is quite impractical and extremely expansive from a computational point of view.
- \triangleright Using the eigendecomposition of L, the matrix exponential can be rewritten as

$$e^{\mathbf{L}} = \mathbf{V}e^{\mathbf{\Lambda}}\mathbf{V}^{-1}.$$

where e^{Λ} can easily be evaluated as it is a diagonal matrix.



Some elements of matrix algebra

Matrix exponential

$$e^{L} = I + L + \frac{L^{2}}{2!} + \frac{L^{3}}{3!} + \cdots,$$

$$= I + V\Lambda V^{-1} + \frac{V\Lambda V^{-1}V\Lambda V^{-1}}{2!} + \frac{V\Lambda V^{-1}V\Lambda V^{-1}V\Lambda V^{-1}}{3!} + \cdots,$$

$$= VIV^{-1} + V\Lambda V^{-1} + \frac{V\Lambda^{2}V^{-1}}{2!} + \frac{V\Lambda^{3}V^{-1}}{3!} + \cdots,$$

$$= V\left(I + \Lambda + \frac{\Lambda^{2}}{2!} + \frac{\Lambda^{3}}{3!} + \cdots\right)V^{-1},$$

$$= Ve^{\Lambda}V^{-1}.$$



Back to our example





Infinitesimal perturbations

The dynamics of an infinitesimal perturbation q evolving in the vicinity of a given fixed point x^* is given by

$$\dot{q} = Lq$$
,

where L is the Jacobian matrix of the system evaluated at $x=x^*$.

The analytical solution to this linear time-invariant dynamical system is given by

$$\mathbf{q}(t) = e^{\mathbf{L}t}\mathbf{q}_0,$$

where q_0 is the initial condition.



Linear instability

- Assuming the eigenvalues of L have been sorted by decreasing real parts, it is clear that if
 - $\Re(\lambda_1) > 0$, then $\lim_{t \to +\infty} e^{Lt} = +\infty$, i.e. there exist a perturbation v_1 that growths exponentially fast in time. The fixed point considered is said to be linearly unstable.
 - $\hookrightarrow \Re(\lambda_1) < 0$, then $\lim_{t \to +\infty} e^{Lt} = 0$, i.e. all perturbations $\underline{\text{decay}}$ exponentially fast in time. The fixed point considered is said to be linearly stable.
- ▶ The case $\Re(\lambda_1)$ is peculiar. The fixed point is said to be **marginally stable** and one needs to consider the actual nonlinear system in order to conclude.
 - → Weakly nonlinear analysis is beyond the scope of this lesson and will be addressed later on.

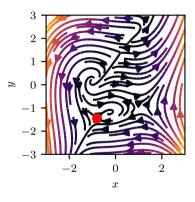


Different types of fixed points

- \triangleright Depending on the eigenspectrum of L, the perturbation may exhibits different kind of dynamics.
- One can thus classify the type of fixed points and the typical trajectory exhibited by the system in the vicinity of said fixed point solely based on the eigenvalues of the associated Jacobian matrix **L**.



Fixed point n°1



Phase plane of the nonlinear system considered.

► The first fixed point is given by

$$\boldsymbol{x}^* = \begin{bmatrix} -0.798 & -1.449 \end{bmatrix}^T$$

► The corresponding Jacobian matrix reads

$$\boldsymbol{J} = \begin{bmatrix} -1.03 & 1.78\\ 0.91 & 0.2 \end{bmatrix}$$

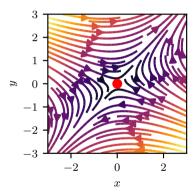


Fixed point n°1

 \triangleright J has two reals eigenvalues of opposite sign:

$$\lambda_1=0.99$$
 and $\lambda_2=-1.83$

- ► Such a fixed point is called a saddle.
 - \hookrightarrow Its first eigendirection v_1 is repulsive while its second, v_2 , is attractive.

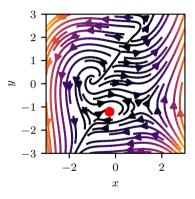


Trajectories in the vicinity of the fixed point x_1^* .





Fixed point n°2



Phase plane of the nonlinear system considered.

► The first fixed point is given by

$$\boldsymbol{x}^* = \begin{bmatrix} -0.26 & -1.2 \end{bmatrix}^T$$

► The corresponding Jacobian matrix reads

$$\boldsymbol{J} = \begin{bmatrix} -0.69 & 2.05 \\ -0.79 & 0.2 \end{bmatrix}$$

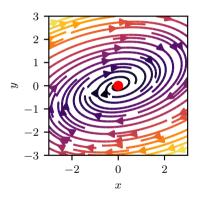


Fixed point n°2

ightharpoonup J has complex conjugate eigenvalues:

$$\lambda_{1.2} = -0.24 \pm 1.2i$$

- ► Such a fixed point is called a spiral.
 - \hookrightarrow In this case, $\Re(\lambda) < 0$, so it is a spiral sink.



Trajectories in the vicinity of the fixed point x_1^* .





Classifying the fixed points

 \blacktriangleright For a 2×2 matrix, its eigenvalues are solution of

$$\lambda^2 - \operatorname{tr}(\boldsymbol{L})\lambda + \det(\boldsymbol{L}) = 0,$$

where
$$tr(\mathbf{L}) = l_{11} + l_{22}$$
 and $det(\mathbf{L}) = l_{11}l_{22} - l_{21}l_{12}$.

ightharpoonup All the different types of fixed points can be placed onto a ${
m tr}(L)$ - ${
m det}(L)$ diagram.



Classifying the fixed points

