Study guide: Approximation of functions

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Finite elements can handle complex geometry, adaptive meshes, higher-order approximations and has a firm theory

- Can with ease solve PDEs in domains with complex geometry
- Can with ease create varying spatial resolution to get accuracy where it is needed
- Can with ease provide higher-order approximations
- Has a rigorous mathematical analysis framework



Solving PDEs by the finite element method

Stationary PDEs:

- Transform the PDE problem to a variational form
- Openine function approximation over finite elements
- Use a computational machinery to derive linear systems
- Solve linear systems

Time-dependent PDEs:

- Finite elements in space
- \bullet Finite difference (or ODE solver) in time

We start with function approximation, then we treat PDEs

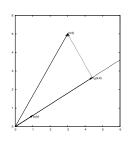
Learning strategy

- Start with approximation of functions, not PDEs
- Introduce finite element approximations
- See later how this machinery is applied to PDEs

Reason:

The finite element method has many concepts and a jungle of details. This learning strategy minimizes the mixing of ideas, concepts, and technical details.

Find a vector in some space that approximates a given vector in a space of higher dimension



The approximation is a linear combination of prescribed basis functions

General idea of finding an approximation u(x) to some given f(x):

$$u(x) = \sum_{i=0}^{N} c_i \psi_i(x)$$

where

- \bullet $\psi_i(x)$ are prescribed functions
- ullet $c_{i},\;i=0,\ldots,N$ are unknown coefficients to be determined

We have three methods to determine the coefficients

We shall address three approaches:

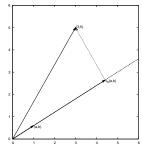
- The least squares method
- The projection (or Galerkin) method
- The interpolation (or collocation) method

Underlying motivation for our notation

Our mathematical framework for doing this is phrased in a way such that it becomes easy to understand and use the FEniCS software package for finite element computing.

Approximation of planar vectors; problem

Given a vector $\mathbf{f} = (3, 5)$, find an approximation to \mathbf{f} directed along a given line.



Approximation of planar vectors; vector space terminology

$$V=\operatorname{\mathsf{span}}\left\{oldsymbol{\psi}_0
ight\}$$

- ullet ψ_0 is a basis vector in the space V
- Seek $\boldsymbol{u} = c_0 \psi_0 \in V$
- Determine c_0 such that u is the "best" approximation to f
- Visually, "best" is obvious

Define

- the error e = f u
- the (Eucledian) scalar product of two vectors: (u, v)
- the norm of $e: ||e|| = \sqrt{(e, e)}$

The least squares method; principle

- Idea: find c_0 such that ||e|| is minimized
- Mathematical convenience: minimize $E = ||e||^2$

$$\frac{\partial E}{\partial c_0} = 0$$

The least squares method; calculations

$$E(c_0) = (\mathbf{e}, \mathbf{e}) = (\mathbf{f} - \mathbf{u}, (\mathbf{f} - \mathbf{u})) = (\mathbf{f} - c_0 \psi_0, \mathbf{f} - c_0 \psi_0)$$

= $(\mathbf{f}, \mathbf{f}) - 2c_0(\mathbf{f}, \psi_0) + c_0^2(\psi_0, \psi_0)$

$$\frac{\partial E}{\partial c_0} = -2(\mathbf{f}, \psi_0) + 2c_0(\psi_0, \psi_0) = 0 \tag{1}$$

$$c_0 = \frac{(\mathbf{f}, \psi_0)}{(\psi_0, \psi_0)} = \frac{3a + 5b}{a^2 + b^2}$$

Observation to be used later: the vanishing derivative (1) can be alternatively written as

$$(e, \psi_0) = 0$$

The projection (or Galerkin) method

- Last slide: min E is equivalent with $(e, \psi_0) = 0$
- $(e, \psi_0) = 0$ implies $(e, \mathbf{v}) = 0$ for any $\mathbf{v} \in V$
- ullet That is: instead of using the least-squares principle, we can require that ullet is orthogonal to $any\ ullet \in V$ (visually clear, but can easily be computed too)
- Precise math: find c_0 such that $(e, \mathbf{v}) = 0$, $\forall \mathbf{v} \in V$
- ullet Equivalent (see notes): find c_0 such that $(m{e}, \psi_0) = 0$
- ullet Insert $oldsymbol{e} = oldsymbol{f} c_0 \psi_0$ and solve for c_0
- ullet Same equation for c_0 and hence same solution as in the least squares method

Approximation of general vectors

Given a vector \boldsymbol{f} , find an approximation $\boldsymbol{u} \in V$:

$$V = \operatorname{span} \{\psi_0, \dots, \psi_N\}$$

We have a set of linearly independent basis vectors $\psi_0,\dots,\psi_N.$ Any $\pmb{u}\in V$ can then be written as

$$\boldsymbol{u} = \sum_{j=0}^{N} c_j \psi_j$$

The least squares method

Idea: find c_0, \ldots, c_N such that $E = ||e||^2$ is minimized, e = f - u.

$$E(c_0,\ldots,c_N) = (\boldsymbol{e},\boldsymbol{e}) = (\boldsymbol{f} - \sum_j c_j \psi_j, \boldsymbol{f} - \sum_j c_j \psi_j)$$
$$= (\boldsymbol{f},\boldsymbol{f}) - 2\sum_{j=0}^N c_j (\boldsymbol{f},\psi_j) + \sum_{p=0}^N \sum_{q=0}^N c_p c_q (\psi_p,\psi_q)$$

$$\frac{\partial E}{\partial c_i} = 0, \quad i = 0, \dots, N$$

After some work we end up with a linear system

$$\sum_{j=0}^{N} A_{i,j} c_j = b_i, \quad i = 0, \dots, N$$
 (2)

$$A_{i,j} = (\psi_i, \psi_j) \tag{3}$$

The projection (or Galerkin) method

Can be shown that minimizing ||e|| implies that e is orthogonal to all $\mathbf{v} \in \mathcal{V}$:

$$(\boldsymbol{e}, \boldsymbol{v}) = 0, \quad \forall \boldsymbol{v} \in V$$

which implies that e most be orthogonal to each basis vector:

$$(e, \psi_i) = 0, \quad i = 0, ..., N$$

This orthogonality condition is the principle of the projection (or Galerkin) method. Leads to the same linear system as in the least squares method.

Approximation of a function in a function space

Let V be a function space spanned by a set of basis functions ψ_0,\ldots,ψ_N .

$$V = \operatorname{span} \{\psi_0, \dots, \psi_N\}$$

Find $u \in V$ as a linear combination of the basis functions:

$$u = \sum_{j \in \mathcal{T}_s} c_j \psi_j, \quad \mathcal{I}_s = \{0, 1, \dots, N\}$$

The least squares method can be extended from vectors to functions

As in the vector case, minimize the (square) norm of the error, E, with respect to the coefficients c_i , $j \in \mathcal{I}_s$:

$$E = (e, e) = (f - u, f - u) = (f(x) - \sum_{j \in \mathcal{I}_s} c_j \psi_j(x), f(x) - \sum_{j \in \mathcal{I}_s} c_j \psi_j(x))$$

$$\frac{\partial E}{\partial c_i} = 0, \quad i = \in \mathcal{I}_s$$

But what is the scalar product when ψ_i is a function?

$$(f,g) = \int_{\Omega} f(x)g(x) \, dx$$

(natural extension from Eucledian product $(m{u},m{v})=\sum_{j}u_{j}v_{j})$

The least squares method; details

$$E(c_0, ..., c_N) = (e, e) = (f - u, f - u)$$

$$= (f, f) - 2 \sum_{j \in \mathcal{I}_s} c_j(f, \psi_i) + \sum_{p \in \mathcal{I}_s} \sum_{q \in \mathcal{I}_s} c_p c_q(\psi_p, \psi_q)$$

$$\frac{\partial E}{\partial c_i} = 0$$
, $i = \in \mathcal{I}_s$

The computations are identical to the vector case, and consequently we get a linear system

$$\sum_{i \in \mathcal{I}_s}^N A_{i,j} c_j = b_i, \ i \in \mathcal{I}_s, \quad A_{i,j} = (\psi_i, \psi_j), \ b_i = (f, \psi_i)$$

The projection (or Galerkin) method

As before, minimizing (e, e) is equivalent to

$$(e, \psi_i) = 0, \quad i \in \mathcal{I}_s$$

which is equivalent to

$$(e, v) = 0, \forall v \in V$$

which is the projection (or Galerkin) method

The algebra is the same as in the multi-dimensional vector case, and we get the same linear system as arose from the least squares method.

Example: fit a parabola by a straight line; problem

Problem

Approximate a parabola
$$f(x) = 10(x-1)^2 - 1$$
 by a straight line.

$$V = \text{span} \{1, x\}$$

That is,
$$\psi_0(x)=1$$
, $\psi_1(x)=x$, and $N=1$. We seek

$$u = c_0\psi_0(x) + c_1\psi_1(x) = c_0 + c_1x$$

Example: fit a parabola by a straight line; solution

$$A_{0,0} = (\psi_0, \psi_0) = \int_1^2 1 \cdot 1 \, dx = 1$$

$$A_{0,1} = (\psi_0, \psi_1) = \int_1^2 1 \cdot x \, dx = 3/2$$

$$A_{1,0} = A_{0,1} = 3/2$$

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$$A_{1,1} = (\psi_1, \psi_1) = \int_1^2 x \cdot x \, dx = 7/2$$

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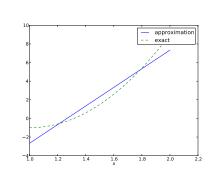
$$b_1 = (f, \psi_0) = \int_1^2 (10(x-1)^2 - 1) \cdot 1 \, dx = 7/3$$

$$b_2 = (f, \psi_1) = \int_1^2 (10(x-1)^2 - 1) \cdot x \, dx = 13/3$$

Solution of 2x2 linear system:

$$c_0 = -38/3$$
, $c_1 = 10$, $u(x) = 10x - \frac{38}{3}$

Example: fit a parabola by a straight line; plot



Ideas for implementing the least squares method via symbolic computations

Consider symbolic computation of the linear system, where

- f(x) is given as a sympy expression f (involving the symbol x),
- psi is a list of {ψ_i}_{i∈I_s}
- ullet Omega is a 2-tuple/list holding the domain Ω

Carry out the integrations, solve the linear system, and return $u(x) = \sum_{i} c_{i} \psi_{j}(x)$

Basic symbolic (SymPy) code for least squares

```
import sympy as sym
def least_squares(f, psi, Omega):
   N = len(psi) - 1
A = sym.zeros((N+1, N+1))
b = sym.zeros((N+1, 1))
x = sym.Symbol('x')
   (x, Omega[0], Omega[1]))
A[j,i] = A[i,j]
b[i,0] = sym.integrate(psi[i]*f, (x, Omega[0], Omega[1]))
c = A.LUsolve(b)
    for i in range(len(psi)):
    u += c[i, 0] *psi[i]
    return u, c
```

Observe: symmetric coefficient matrix so we can halve the integrations.

Improved code if symbolic integration fails

- If sym.integrate fails, it returns an sym.Integral object.
 We can test on this object and fall back on numerical integration.
- We can include a boolean argument symbolic to explicitly choose between symbolic and numerical computing.

```
def least_squares(f, psi, Omega, symbolic=True):
    for i in range(M+1):
        for j in range(i, M+1):
            integrand = psi[i]*psi[j]
        if symbolic:
            I = sym integrate(integrand, (x, Omega[0], Omega[i]))
        if not symbolic or isinstance(I, sym.Integral):
            # Could not integrate symbolically,
            # foll back on numerical integration
            integrand = sym lambdify(tx], integrand)
            I = sym mymath quad(integrand, [Omega[0], Omega[i]])
        A[i,j] = A[j,i] = I

integrand = psi[i]*f
    if symbolic:
        I = sym integrate(integrand, (x, Omega[0], Omega[i]))
    if not symbolic or isinstance(I, sym.Integral):
        integrand = sym.lambdify([x], integrand)
```

Plotting of the solution

Compare f and u visually:

```
def comparison_plot(f, u, Omega, filename='tmp.pdf'):
    x = sym.8ymbol('x')
    # Turn f and u to ordinary Python functions
    f = sym.lambdify([x], f, modules="numpy")
    u = sym.lambdify([x], u, modules="numpy")
    resolution = 401 # no of points in plot
    xcoor = linspace(Omega[0], Omega[1], resolution)
    exact = f(xcoor)
    approx = u(xcoor)
    plot(xcoor, approx)
    hold('on')
    plot(xcoor, exact)
    legend(['approximation', 'exact'])
    savefig(filename)
```

All code in module approx1D.py

The approximation is exact if $f \in V$

- What if we add $\psi_2 = x^2$ to the space V?
- That is, approximating a parabola by any parabola?
- (Hopefully we get the exact parabola!)

```
>>> from approx1D import *
>>> x = sym.Symbol('x')
>>> f = 10(*x-1)**2-1
>>> u, c = least_squares(f=f, psi=[1, x, x**2], Omega=[1, 2])
>>> print u
10*x**2 - 20*x + 9
>>> print sym expand(f)
10*x**2 - 20*x + 9
```

The general result: perfect approximation if $f \in V$

- What if we use $\psi_i(x) = x^i$ for i = 0, ..., N = 40?
- The output from least_squares is $c_i = 0$ for i > 2

General result

If $f \in V$, least squares and projection/Galerkin give u = f.

Proof of why $f \in V$ gives exact u

If $f \in V$, $f = \sum_{j \in \mathcal{I}_s} d_j \psi_j$, for some $\{d_i\}_{i \in \mathcal{I}_s}$. Then

$$b_i = (f, \psi_i) = \sum_{j \in \mathcal{I}_s} d_j(\psi_j, \psi_i) = \sum_{j \in \mathcal{I}_s} d_j A_{i,j}$$

The linear system $\sum_i A_{i,j} c_j = b_i, \ i \in \mathcal{I}_s$, is then

$$\sum_{j \in \mathcal{I}_s} c_j A_{i,j} = \sum_{j \in \mathcal{I}_s} d_j A_{i,j}, \quad i \in \mathcal{I}_s$$

which implies that $c_i = d_i$ for $i \in \mathcal{I}_s$ and u is identical to f.

Finite-precision in numerical computations; question

The previous computations were symbolic. What if we solve the linear system numerically with standard arrays?

That is, f is parabola, but we approximate with

$$u(x) = c_0 + c_1 x + c_2 x^2 + c_3 x^3 + \dots + c_N x^N$$

We expect $c_2 = c_3 = \cdots = c_N = 0$ since $f \in V$ implies u = f.

Will we get this result with finite precision computer arithmetic?

Finite-precision in numerical computations; results

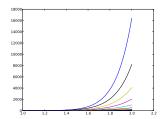
numpy 64	numpy 32	sympy	exact
8.98	5.57	9.62	9
-19.93	-7.65	-23.39	-20
9.96	-4.50	17.74	10
-0.26	4.13	-9.19	0
0.72	2.99	5.25	0
-0.93	-1.21	0.18	0
0.73	-0.41	-2.48	0
-0.36	-0.013	1.81	0
0.11	0.08	-0.66	0
-0.02	0.04	0.12	0
0.002	-0.02	-0.001	0

- Column 2: matrix and lu_solve from sympy.mpmath.fp
- Column 3: numpy matrix with 4-byte floats
- Column 4: numpy matrix with 8-byte floats

The ill-conditioning is due to almost linearly dependent basis functions for large ${\it N}$

- Significant round-off errors in the numerical computations (!)
- But if we plot the approximations they look good (!)

Source or problem: the \mathbf{x}^i functions become almost linearly dependent as i grows:



Ill-conditioning: general conclusions

- Almost linearly dependent basis functions give almost singular matrices
- Such matrices are said to be *ill conditioned*, and Gaussian elimination is severely affected by round-off errors
- ullet The basis $1,x,x^2,x^3,x^4,\ldots$ is a bad basis
- \bullet Polynomials are fine as basis, but the more orthogonal they are, $(\psi_i,\psi_j)\approx 0$, the better

Fourier series approximation; problem and code

Let's approximate f by a typical Fourier series expansion

$$u(x) = \sum_{i} a_{i} \sin i\pi x = \sum_{j=0}^{N} c_{j} \sin((j+1)\pi x)$$

which means that

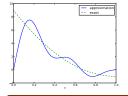
$$V = \text{span} \{ \sin \pi x, \sin 2\pi x, \dots, \sin (N+1)\pi x \}$$

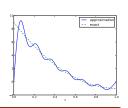
Computations using the least_squares function:

N = 3 from sympy import sin, pi psi = [sin(pi*(i+1)*x) for i in range(N+1)] f = 10*(x-1)**2 - 1 Omega = [0, 1] u, c = least_squares(f, psi, Omega) comparison_plot(f, u, Omega)

Fourier series approximation; plot

Left: N=3, right: N=11:





Problem:

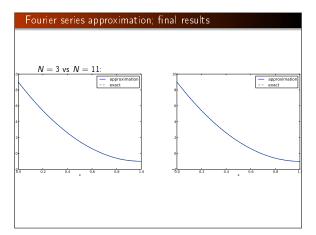
All $\psi_i(0) = 0$ and hence $u(0) = 0 \neq f(0) = 9$. Similar problem at x = 1. The boundary values of u are always wrong!

Fourier series approximation; improvements

- ullet Considerably improvement with ${\it N}=11$, but still undesired discrepancy at ${\it x}=0$ and ${\it x}=1$
- Possible remedy: add a term that leads to correct boundary
 values.

$$u(x) = f(0)(1-x) + xf(1) + \sum_{j \in \mathcal{I}_s} c_j \psi_j(x)$$

The extra terms ensure u(0) = f(0) and u(1) = f(1) and is a strikingly good help to get a good approximation!



Orthogonal basis functions

This choice of sine functions as basis functions is popular because

- the basis functions are orthogonal: $(\psi_i, \psi_i) = 0$
- \bullet implying that $A_{i,i}$ is a diagonal matrix
- implying that we can solve for $c_i = 2 \int_0^1 f(x) \sin((i+1)\pi x) dx$
- ullet and what we get is the standard Fourier sine series of f

In general, for an orthogonal basis, $A_{i,j}$ is diagonal and we can easily solve for c_i :

$$c_i = \frac{b_i}{A_{i,i}} = \frac{(f, \psi_i)}{(\psi_i, \psi_i)}$$

Function for the least squares method with orthogonal basis functions

```
def least_squares_orth(f, psi, Omega):
    N = len(psi) - 1
    A = (0)*(N+1)
    b = (0)*(N+1)
    x = sym.Symbol('x')
    for i in range(N+1):
        A[i] = sym integrate(psi[i]**2, (x, Omega[0], Omega[1]))
        b[i] = sym.integrate(psi[i]*f, (x, Omega[0], Omega[1]))
        c = [b[i]/A[i] for i in range(len(b))]
    u = 0
    for i in range(len(psi)):
        u += c[i]*psi[i]
    return u, c
```

Function for the least squares method with orthogonal basis functions; symbolic and numerical integration

Extensions:

- We can choose between symbolic or numerical integration (symbolic argument).
- If symbolic, we fall back on numerical integration after failure (sym.Integral is returned from sym.integrate).

```
for i in range(N+1):
    # Diagonal matrix term
    A[i] = sym.integrate(psi[i]**2, (x, Omega[0], Omega[1]))
    # Right-hand side term
    integrand = psi[i]**
    if symbolic:
        I = sym.integrate(integrand, (x, Omega[0], Omega[1]))
    if not symbolic or isinstance(I, sym.Integral):
        print 'numerical integration of', integrand
        integrand = sym.lambdify([x], integrand)
        I = sym.mpmath.quad(integrand, [Omega[0], Omega[1]))
b[i] = I
```

The collocation or interpolation method; ideas and math

Here is another idea for approximating f(x) by $u(x) = \sum_i c_i \psi_i$:

- Force $u(x_i) = f(x_i)$ at some selected *collocation* points $\{x_i\}_{i \in \mathcal{T}_-}$
- Then u is said to interpolate f
- The method is known as interpolation or collocation

$$u(x_i) = \sum_{j \in \mathcal{I}_s} c_j \psi_j(x_i) = f(x_i) \quad i \in \mathcal{I}_s, N$$

This is a linear system with no need for integration:

$$\sum_{j\in\mathcal{I}_s}A_{i,j}c_j=b_i,\quad i\in\mathcal{I}_s \tag{5}$$

$$A_{i,j} = \psi_i(x_i) \tag{6}$$

$$b_i = f(x_i) \tag{7}$$

No symmetric matrix: $\psi_i(x_i) \neq \psi_i(x_i)$ in general

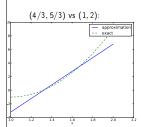
The collocation or interpolation method; implementation

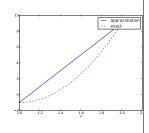
points holds the interpolation/collocation points

```
def interpolation(f, psi, points):
    N = len(psi) - 1
    A = sym.zeros((N+i, N+i))
    b = sym.zeros((N+i, N+i))
    b = sym.zeros((N+i, 1))
    x = sym.zeros((N+i, 1))
    x = sym.symbol('va')
    # Turn psi and f into Python functions
    psi = [sym.lambdify([x], psi[i]) for i in range(N+i)]
    f = sym.lambdify([x], fsi]
    for j in range(N+i):
        A[i,j] = psi[j](coints[i])
        b[i,0] = f(points[i])
    c = A.LUsolve(b)
    u = 0
    for i in range(len(psi)):
        u += c[i,0]*psi[i](x)
    return u
```

The collocation or interpolation method; approximating a parabola by linear functions

- Potential difficulty: how to choose x_i?
- The results are sensitive to the points!





Lagrange polynomials; motivation and ideas

Motivation:

- The interpolation/collocation method avoids integration
- ullet With a diagonal matrix $A_{i,j}=\psi_j(\mathbf{x}_i)$ we can solve the linear system by hand

The Lagrange interpolating polynomials ψ_j have the property that

$$\psi_i(\mathbf{x}_j) = \delta_{ij}, \quad \delta_{ij} = \begin{cases} 1, & i = j \\ 0, & i \neq j \end{cases}$$

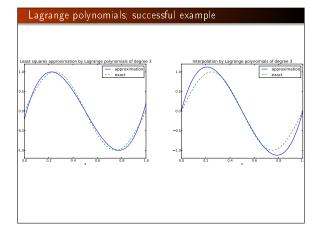
Hence, $c_i = f(x_i)$ and

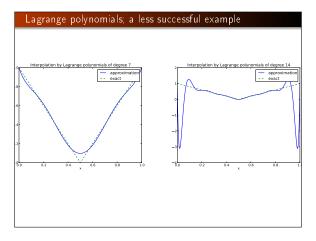
$$u(x) = \sum_{j \in \mathcal{I}_s} f(x_i) \psi_i(x)$$

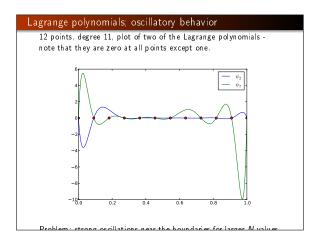
 Lagrange polynomials and interpolation/collocation look convenient

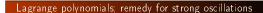
Lagrange polynomials; formula and code

$$\begin{split} \psi_i(\mathbf{x}) &= \prod_{j=0, j \neq i}^N \frac{\mathbf{x} - \mathbf{x}_j}{\mathbf{x}_i - \mathbf{x}_j} = \frac{\mathbf{x} - \mathbf{x}_0}{\mathbf{x}_i - \mathbf{x}_0} \cdots \frac{\mathbf{x} - \mathbf{x}_{i-1}}{\mathbf{x}_i - \mathbf{x}_{i-1}} \frac{\mathbf{x} - \mathbf{x}_{i+1}}{\mathbf{x}_i - \mathbf{x}_{i+1}} \cdots \frac{\mathbf{x} - \mathbf{x}_N}{\mathbf{x}_i - \mathbf{x}_N} \\ \text{def Lagrange_polynomial(x, i, points):} \\ & \quad \text{p = 1 in range(len(points)):} \\ & \quad \text{if k = 1 in permode } \\ & \quad \text{p = e (x - points[k])/(points[i] - points[k])} \end{split}$$





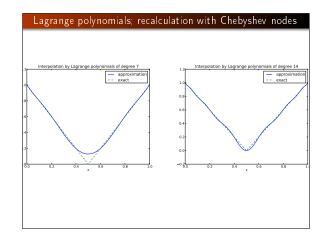


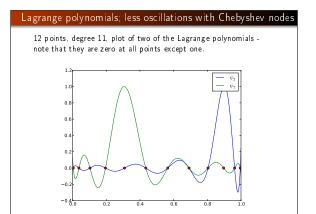


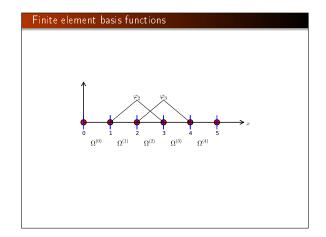
The oscillations can be reduced by a more clever choice of interpolation points, called the *Chebyshev nodes*:

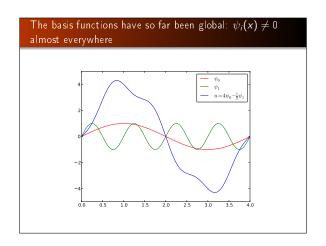
$$x_i = \frac{1}{2}(a+b) + \frac{1}{2}(b-a)\cos\left(\frac{2i+1}{2(N+1)}pi\right), \quad i = 0..., N$$

on an interval [a, b].



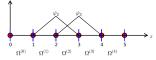


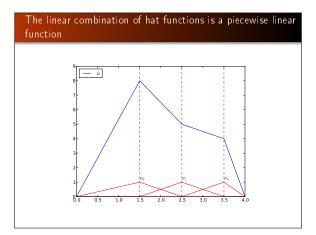




In the finite element method we use basis functions with local support

- Local support: $\psi_i(x) \neq 0$ for x in a small subdomain of Ω
- Typically hat-shaped
- ullet u(x) based on these ψ_i is a piecewise polynomial defined over many (small) subdomains
- ullet We introduce $arphi_i$ as the name of these finite element hat functions (and for now choose $\psi_i=arphi_i$)





Elements and nodes

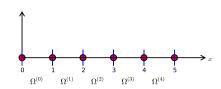
Split Ω into N_e non-overlapping subdomains called *elements*:

$$\Omega = \Omega^{(0)} \cup \dots \cup \Omega^{(N_e)}$$

On each element, introduce N_n points called *nodes*: x_0, \ldots, x_{N_n-1}

- The finite element basis functions are named $\varphi_i(x)$
- $ullet arphi_i = 1$ at node i and 0 at all other nodes
- ullet $arphi_i$ is a Lagrange polynomial on each element
- ullet For nodes at the boundary between two elements, $arphi_i$ is made up of a Lagrange polynomial over each element

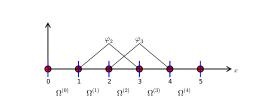
Example on elements with two nodes (P1 elements)



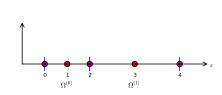
Data structure: nodes holds coordinates or nodes, elements holds the node numbers in each element

nodes = [0, 1.2, 2.4, 3.6, 4.8, 5] elements = [[0, 1], [1, 2], [2, 3], [3, 4], [4, 5]]

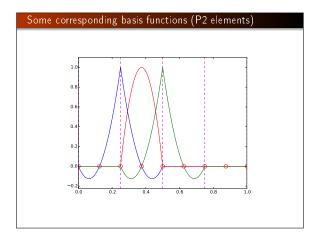
Illustration of two basis functions on the mesh

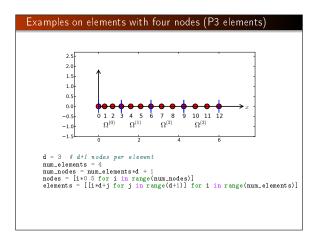


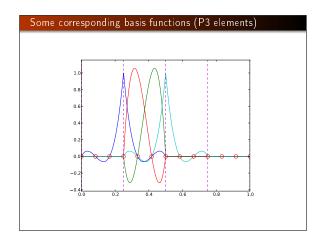
Example on elements with three nodes (P2 elements)

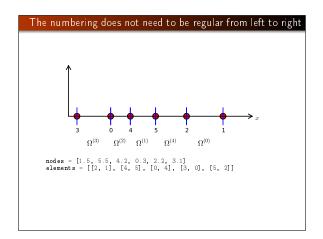


nodes = [0, 0.125, 0.25, 0.375, 0.5, 0.625, 0.75, 0.875, 1.0] elements = [[0, 1, 2], [2, 3, 4], [4, 5, 6], [6, 7, 8]]









Interpretation of the coefficients c_i

Important property: c_i is the value of u at node i, x_i :

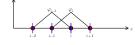
$$u(x_i) = \sum_{j \in \mathcal{I}_s} c_j \varphi_j(x_i) = c_i \varphi_i(x_i) = c_i$$

because $\varphi_i(x_i) = 0$ if $i \neq j$ and $\varphi_i(x_i) = 1$

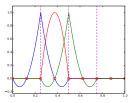
Properties of the basis functions

- $\varphi_i(x) \neq 0$ only on those elements that contain global node i
- ullet $\varphi_i(x)\varphi_j(x)
 eq 0$ if and only if i and j are global node numbers in the same element

Since $A_{i,j}=\int \varphi_i \varphi_j \,\mathrm{d} x$, most of the elements in the coefficient matrix will be zero

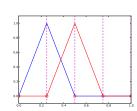


How to construct quadratic φ_i (P2 elements)



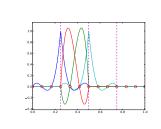
- Associate Lagrange polynomials with the nodes in an element
- When the polynomial is 1 on the element boundary, combine it with the polynomial in the neighboring element that is also 1 at the same point

Example on linear
$$\varphi_i$$
 (P1 elements)

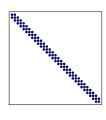


$$\varphi_i(x) = \begin{cases} 0, & x < x_{i-1} \\ (x - x_{i-1})/h & x_{i-1} \le x < x_i \\ 1 - (x - x_i)/h, & x_i \le x < x_{i+1} \\ 0, & x \ge x_{i+1} \end{cases}$$

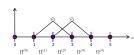
Example on cubic φ_i (P3 elements)



Calculating the linear system for c_i



Computing a specific matrix entry (1)

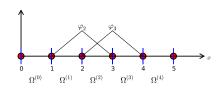


 $A_{2,3}=\int_{\Omega} arphi_2 arphi_3 dx\colon \, arphi_2 arphi_3
eq 0$ only over element 2. There,

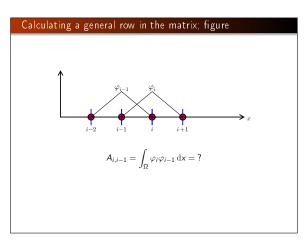
$$\varphi_3(x) = (x - x_2)/h, \quad \varphi_2(x) = 1 - (x - x_2)/h$$

$$A_{2,3} = \int_{\Omega} \varphi_2 \varphi_3 \, \mathrm{d}x = \int_{x_2}^{x_3} \left(1 - \frac{x - x_2}{h} \right) \frac{x - x_2}{h} \, \mathrm{d}x = \frac{h}{6}$$

Computing a specific matrix entry (2)



$$A_{2,2} = \int_{x_1}^{x_2} \left(\frac{x - x_1}{h} \right)^2 dx + \int_{x_2}^{x_3} \left(1 - \frac{x - x_2}{h} \right)^2 dx = \frac{2h}{3}$$



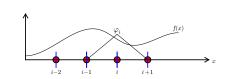
Calculating a general row in the matrix; details

$$A_{i,i-1} = \int_{\Omega} \varphi_i \varphi_{i-1} \, \mathrm{d}x$$

$$= \underbrace{\int_{x_{i-2}}^{x_{i-1}} \varphi_i \varphi_{i-1} \, \mathrm{d}x}_{\varphi_i = 0} + \underbrace{\int_{x_{i-1}}^{x_i} \varphi_i \varphi_{i-1} \, \mathrm{d}x}_{\varphi_i = 0} + \underbrace{\int_{x_i}^{x_{i+1}} \varphi_i \varphi_{i-1} \, \mathrm{d}x}_{\varphi_{i-1} = 0} + \underbrace{\int_{x_i}^{x_{i+1}} \varphi_i \varphi_{i-1} \, \mathrm{d}x}_{\varphi_{i-1} = 0} + \underbrace{\int_{x_i}^{x_{i+1}} \varphi_i \varphi_{i-1} \, \mathrm{d}x}_{\varphi_{i-1} = 0} + \underbrace{\int_{x_i}^{x_{i+1}} \varphi_i \varphi_{i-1} \, \mathrm{d}x}_{\varphi_i = 0} + \underbrace{\int_{x_i}^{x_{i+1}} \varphi_i \varphi_i \, \mathrm{d}x}_{\varphi_i = 0} + \underbrace{\int_{x_i}^{x_{i+1}} \varphi_i \varphi_i \, \mathrm{d}x}_{\varphi_i = 0} + \underbrace{\int_{x_i}^{x_{i+1}} \varphi_i \, \mathrm{d}x}_{$$

- ullet $A_{i,i+1}=A_{i,i-1}$ due to symmetry
- $A_{i,i} = 2h/3$ (same calculation as for $A_{2,2}$)
- ullet $A_{0,0}=A_{N,N}=h/3$ (only one element)

Calculation of the right-hand side



$$b_{i} = \int_{\Omega} \varphi_{i}(x) f(x) dx = \int_{x_{i-1}}^{x_{i}} \frac{x - x_{i-1}}{h} f(x) dx + \int_{x_{i}}^{x_{i+1}} \left(1 - \frac{x - x_{i}}{h}\right) f(x)$$

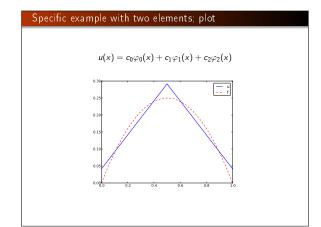
Need a specific f(x) to do more...

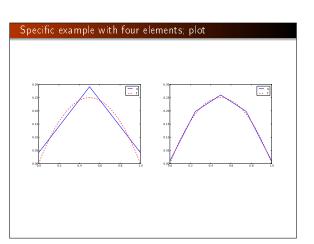
Specific example with two elements; linear system and solution

- f(x) = x(1-x) on $\Omega = [0,1]$
- ullet Two equal-sized elements [0,0.5] and [0.5,1]

$$A = \frac{h}{6} \begin{pmatrix} 2 & 1 & 0 \\ 1 & 4 & 1 \\ 0 & 1 & 2 \end{pmatrix}, \quad b = \frac{h^2}{12} \begin{pmatrix} 2 - 3h \\ 12 - 14h \\ 10 - 17h \end{pmatrix}$$

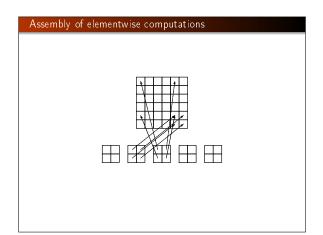
$$c_0 = \frac{h^2}{6}$$
, $c_1 = h - \frac{5}{6}h^2$, $c_2 = 2h - \frac{23}{6}h^2$





Specific example: what about P2 elements?

Recall: if $f \in V$, u becomes exact. When f is a parabola, any choice of P2 elements (1 or many) will give u = f exactly. The same is true for P3, P4, ... elements since all of them can represent a 2nd-degree polynomial exactly.



Split the integrals into elementwise integrals

$$A_{i,j} = \int_{\Omega} \varphi_i \varphi_j dx = \sum_{e} \int_{\Omega(e)} \varphi_i \varphi_j dx, \quad A_{i,j}^{(e)} = \int_{\Omega(e)} \varphi_i \varphi_j dx$$

Important observations:

- $A_{i,j}^{(e)} \neq 0$ if and only if i and j are nodes in element e (otherwise no overlap between the basis functions)
- \bullet All the nonzero elements in $A_{i,j}^{(e)}$ are collected in an element matrix
- \bullet The element matrix has contributions from the φ_i functions associated with the nodes in element
- ullet It is convenient to introduce a *local numbering* of the nodes in an element: $0,1,\ldots,d$



The element matrix and local vs global node numbers

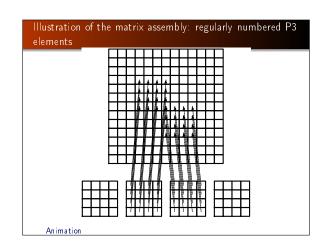
$$\tilde{A}^{(e)} = \{\tilde{A}^{(e)}_{r,s}\}, \quad \tilde{A}^{(e)}_{r,s} = \int_{\Omega^{(e)}} \varphi_{q(e,r)} \varphi_{q(e,s)} dx, \quad r,s \in I_d = \{0,\ldots,d\}$$

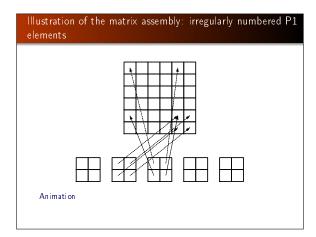
Now

- ullet r,s run over $local\ node\ numbers$ in an element: $0,1,\ldots,d$
- ullet i,j run over global node numbers $i,j\in\mathcal{I}_s=\{0,1,\ldots,N\}$
- i = q(e, r): mapping of local node number r in element e to the global node number i (math equivalent to i=elements[e][r])
- ullet Add $ilde{A}_{r,s}^{(e)}$ into the global $A_{i,j}$ (assembly)

$$A_{q(e,r),q(e,s)} := A_{q(e,r),q(e,s)} + \tilde{A}_{r,s}^{(e)}, \quad r,s \in I_d$$

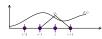
Illustration of the matrix assembly: regularly numbered P1 elements Animation





Assembly of the right-hand side

$$b_i = \int_{\Omega} f(x)\varphi_i(x)dx = \sum_{e} \int_{\Omega^{(e)}} f(x)\varphi_i(x)dx, \quad b_i^{(e)} = \int_{\Omega^{(e)}} f(x)\varphi_i(x)dx$$



Important observations:

- $ullet b_i^{(e)}
 eq 0$ if and only if global node i is a node in element e(otherwise $\varphi_i = 0$)
- ullet The d+1 nonzero $b_i^{(e)}$ can be collected in an element vector $\tilde{b}_{r}^{(e)} = \{\tilde{b}_{r}^{(e)}\}, r \in I_{d}$

Assembly:

$$b_{a(e,r)} := b_{a(e,r)} + \tilde{b}_r^{(e)}, \quad r, s \in I_d$$

Mapping to a reference element

Instead of computing

$$\tilde{A}_{r,s}^{(e)} = \int_{\Omega^{(e)}} \varphi_{q(e,r)}(x) \varphi_{q(e,s)}(x) dx = \int_{x_L}^{x_R} \varphi_{q(e,r)}(x) \varphi_{q(e,s)}(x) dx$$

we now map $[x_L,x_R]$ to a standardized reference element domain [-1,1] with local coordinate X

We use affine mapping: linear stretch of $X \in [-1, 1]$ to $x \in [x_L, x_R]$

$$x = \frac{1}{2}(x_L + x_R) + \frac{1}{2}(x_R - x_L)$$

$$x=\frac{1}{2}(x_L+x_R)+\frac{1}{2}(x_R-x_L)X$$
 or rewritten as
$$x=x_m+\frac{1}{2}hX, \qquad x_m=(x_L+x_R)/2, \quad h=x_R-x_L$$

Integral transformation

Reference element integration: just change integration variable from x to X. Introduce local basis function

$$\tilde{\varphi}_r(X) = \varphi_{q(e,r)}(x(X))$$

$$\tilde{A}_{r,s}^{(e)} = \int_{\Omega^{(e)}} \varphi_{q(e,r)}(x) \varphi_{q(e,s)}(x) dx = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{r}(X) \tilde{\varphi}_{s}(X) \underbrace{\frac{dx}{dX}}_{\det J = h/2} dX = \int_{-1}^{1} \tilde{\varphi}_{s}(X) \underbrace{$$

$$\tilde{b}_r^{(e)} = \int_{\Omega^{(e)}} f(x) \varphi_{q(e,r)}(x) dx = \int_{-1}^1 f(x(X)) \tilde{\varphi}_r(X) \det J dX$$

Advantages of the reference element

- Always the same domain for integration: [-1, 1]
- We only need formulas for $\tilde{\varphi}_r(X)$ over one element (no piecewise polynomial definition)
- \bullet $\tilde{\varphi}_r(X)$ is the same for all elements: no dependence on element length and location, which is "factored out" in the mapping and $\det J$

Standardized basis functions for P1 elements

$$\tilde{\varphi}_0(X) = \frac{1}{2}(1 - X)$$
 (8)

$$\tilde{\varphi}_1(X) = \frac{1}{2}(1+X) \tag{9}$$

Note: simple polynomial expressions (no need to consider piecewisely defined functions)

Standardized basis functions for P2 elements

$$\tilde{\varphi}_0(X) = \frac{1}{2}(X - 1)X$$
 (10)

$$\tilde{\varphi}_1(X) = 1 - X^2 \tag{11}$$

$$\tilde{\varphi}_2(X) = \frac{1}{2}(X+1)X$$
 (12)

Easy to generalize to arbitrary order!

How to find the polynomial expressions?

Three alternatives:

- lack lack lack Map the global basis function $arphi_i(x)$ over an element to X coordinates
- ② Compute $\tilde{\varphi}_r(X)$ from scratch using
 - a given polynomial order d
 - $\tilde{\varphi}_r(X) = 1 \text{ at local node } 1$
 - $\tilde{\varphi}_r(X)=1$ at all other local nodes
- Use formulas for Lagrange interpolating polynomials on the element

Integration over a reference element; element matrix

P1 elements and f(x) = x(1-x).

$$\begin{split} \tilde{A}_{0,0}^{(e)} &= \int_{-1}^{1} \tilde{\varphi}_{0}(X) \tilde{\varphi}_{0}(X) \frac{h}{2} dX \\ &= \int_{-1}^{1} \frac{1}{2} (1 - X) \frac{1}{2} (1 - X) \frac{h}{2} dX = \frac{h}{8} \int_{-1}^{1} (1 - X)^{2} dX = \frac{h}{3} \end{split}$$
(13)

$$\tilde{A}_{1,0}^{(e)} = \int_{-1}^{1} \tilde{\varphi}_{1}(X) \tilde{\varphi}_{0}(X) \frac{h}{2} dX
= \int_{-1}^{1} \frac{1}{2} (1+X) \frac{1}{2} (1-X) \frac{h}{2} dX = \frac{h}{8} \int_{-1}^{1} (1-X^{2}) dX = \frac{h}{6}$$
(14)

$$\tilde{A}_{0,1}^{(e)} = \tilde{A}_{1,0}^{(e)} \tag{15}$$

$$\tilde{A}_{1,1}^{(e)} = \int_{-1}^{1} \tilde{\varphi}_1(X) \tilde{\varphi}_1(X) \frac{h}{2} dX$$

Integration over a reference element; element vector

$$\tilde{b}_{0}^{(e)} = \int_{-1}^{1} f(x(X)) \tilde{\varphi}_{0}(X) \frac{h}{2} dX
= \int_{-1}^{1} (x_{m} + \frac{1}{2} hX) (1 - (x_{m} + \frac{1}{2} hX)) \frac{1}{2} (1 - X) \frac{h}{2} dX
= -\frac{1}{24} h^{3} + \frac{1}{6} h^{2} x_{m} - \frac{1}{12} h^{2} - \frac{1}{2} h x_{m}^{2} + \frac{1}{2} h x_{m}$$
(17)
$$\tilde{b}_{1}^{(e)} = \int_{-1}^{1} f(x(X)) \tilde{\varphi}_{1}(X) \frac{h}{2} dX
= \int_{-1}^{1} (x_{m} + \frac{1}{2} hX) (1 - (x_{m} + \frac{1}{2} hX)) \frac{1}{2} (1 + X) \frac{h}{2} dX
= -\frac{1}{24} h^{3} - \frac{1}{6} h^{2} x_{m} + \frac{1}{12} h^{2} - \frac{1}{2} h x_{m}^{2} + \frac{1}{2} h x_{m}$$
(18)

 x_m : element midpoint.

Tedious calculations! Let's use symbolic software

```
>>> import sympy as sym
>>> x, x_m, h, X = sym.symbols('x x_m h X')
>>> sym.integrate(h/8*(1-X)**2, (X, -1, 1))
h/3
>>> sym.integrate(h/8*(1+X)*(1-X), (X, -1, 1))
h/6
>>> x = x_m + h/2*X
>>> b_0 = sym.integrate(h/4*x*(1-x)*(1-X), (X, -1, 1))
>>> print b_0
-h**3/24 + h**2*x_m/6 - h**2/12 - h*x_m**2/2 + h*x_m/2
```

Can print out in LATEX too (convenient for copying into reports):

```
>>> print sym.latex(b_0, mode='plain')
- \frac(1){24} \ h^{3} + \frac{1}{6} \ h^{2} x_{m} - \frac{1}{12} \ h^{2} - \half \ h x_{m}^{2} + \half \ h x_{m} - \frac{1}{6} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \ h^{2} \ h^{2} \ half \ h x_{m} - \frac{1}{6} \ h^{2} \
```

Implementation

- Functions can operate in symbolic or numeric mode

def phi_r(r, X, d): if isinstance(X, sym.Symbol): h = sym.Rational(1, d) # node spacing nodes = [2*i*h - 1 for i in range(d+1)] Coming functions appear in fe_approx1D.py # assume I is numeric: use floats for nodes nodes = np.linspace(-1, 1, d+1) • The code documents all steps in finite element calculations! return Lagrange_polynomial(X, r, nodes) def Lagrange_polynomial(x, i, points): f = 1 for k in range(len(points)): p *= (x - points[k])/(points[i] - points[k]) return p def basis(d=1): """Return the complete basis.""" X = sym.Symbol('X') phi = [phi_r(r, X, d) for r in range(d+1)]

Compute the element matrix

```
def element_matrix(phi, Omega_e, symbolic=True):
    n = len(phi)
    A_e = sym.zeros((n, n))
      X = sym.Symbol('X')
if symbolic:
h = sym.Symbol('h')
      else:

h = 0mega_e[1] - 0mega_e[0]

det J = h/2 \# dx/dX
       for r in range(n):
             for s in range(r, n):

A_e[r,s] = sym.integrate(phi[r]*phi[s]*detJ, (X, -1, 1))

A_e[s,r] = A_e[r,s]
       return A e
```

Example on symbolic vs numeric element matrix

Compute finite element basis functions in the reference

Let $\tilde{\varphi}_r(X)$ be a Lagrange polynomial of degree d:

import sympy as sym import numpy as np

```
>>> from fe_approx1D import *
>>> phi = basis(d=1)
>>> phi - Usaisut-/

>>> phi | 1/2 - X/2, 1/2 + X/2|

>>> element_matrix(phi, Omega_e=[0.1, 0.2], symbolic=True)

[h/3, h/3]
```

Compute the element vector

```
def element_vector(f, phi, Omega_e, symbolic=True):
    n = len(phi)
          b_e = sym.zeros((n, 1))
         b_e = sym.zeros(\(\mathbf{n}\), \(\text{if}\)
# Make f a function of I
X = sym.Symbol('\(\mathbf{X}\'\))
if symbolic:
    h = sym.Symbol('\(\mathbf{h}\'\))
          else:
         else:  h = 0 \text{ mega\_e[i]} - 0 \text{ mega\_e[o]}   x = (0 \text{ mega\_e[o]} + 0 \text{ mega\_e[i]})/2 + h/2*X \text{ # mapping }   f = f. \text{ subs}('x', x) \text{ # substitute mapping formula for } x \text{ detJ} = h/2 \text{ # } ds/dI 
          for r in range(n):
b.e[r] = sym.integrate(f*phi[r]*detJ, (X, -1, 1))
return b_e
Note f.subs('x', x): replace x by x(X) such that f contains X
```

Fallback on numerical integration if symbolic integration of $\int f \tilde{\varphi}_r dx$ fails

- Element matrix: only polynomials and sympy always succeeds
- Element vector: $\int f \tilde{\varphi} dx$ can fail (sympy then returns an Integral object instead of a number)

```
def element_vector(f, phi, Omega_e, symbolic=True):
             I = sym.integrate(f*phi[r]*detJ, (X, -1, 1)) # try...
            if isinstance(I, sym Integral):
   h = Omega_e[1] - Omega_e[0] # Ensure h is numerical
           n = Umega_e1:1 - Umega_e1:0! # Insure h is ni
detJ = h/2
integrand = sym lambdify([X], f *phi[r] *detJ)
I = sym mpmath.quad(integrand, [-1, 1])
b_e[r] = I
```


Linear system solution

```
if symbolic:
    c = A.LUsolve(b)  # sympy arrays, symbolic Gaussian elim.
else:
    c = np.linalg.solve(A, b) # numpy arrays, numerical solve
```

Note: the symbolic computation of A, b and A.LUsolve(b) can be very tedious.

Example on computing symbolic approximations

```
>>> h, x = sym.symbols('h x')
>>> nodes = [0, h, 2*h]
>>> elements = [[0, 1], [1, 2]]
>>> pin = basis(a].
>>> f = x*(1.x)
>>> A

h/3, h/6, 0]

[h/6, 2*h/3, h/6]

[ 0, h/6, h/3]
>>> b

[ h**2/6 - h**3/12]

[ h**2/6 - 7*h**3/12]

[ 5*h*2/6 - 17*h*3/12]
>>> c = A.LUsolve(b)
>>> c

[12*(7*h**2/12 - 35*h**3/72)/(7*h)]

[ 7*(4*h**2/12 - 35*h**3/21)/(2*h)]
```

Example on computing numerical approximations

The structure of the coefficient matrix

Note: do this by hand to understand what is going on!

General result: the coefficient matrix is sparse

- Sparse = most of the entries are zeros
- Below: P1 elements

Exemplifying the sparsity for P2 elements

$$A = \frac{h}{30} \begin{pmatrix} 4 & 2 & -1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 2 & 16 & 2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ -1 & 2 & 8 & 2 & -1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 2 & 16 & 2 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 2 & 8 & 2 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 2 & 16 & 2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 & 2 & 8 & 2 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 2 & 16 & 2 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 2 & 4 & 4 \end{pmatrix}$$

Matrix sparsity pattern for regular/random numbering of P1 elements

- Left: number nodes and elements from left to right
- Right: number nodes and elements arbitrarily





Matrix sparsity pattern for regular/random numbering of P3 elements

- Left: number nodes and elements from left to right
- Right: number nodes and elements arbitrarily





Sparse matrix storage and solution

The minimum storage requirements for the coefficient matrix $A_{i,j}$:

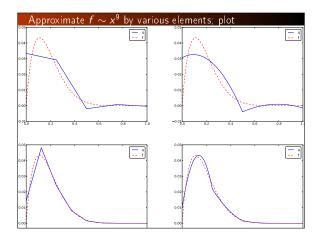
- P1 elements: only 3 nonzero entries per row
- P2 elements: only 5 nonzero entries per row
- P3 elements: only 7 nonzero entries per row
- It is important to utilize sparse storage and sparse solvers
- In Python: scipy.sparse package

Approximate $f \sim x^9$ by various elements; code

Compute a mesh with N_e elements, basis functions of degree d, and approximate a given symbolic expression f(x) by a finite element expansion $u(x) = \sum_i c_j \varphi_j(x)$:

```
import sympy as sym
from fe_approxID import approximate
x = sym.Symbol('x')'

approximate(f=x*(1-x)**8, symbolic=False, d=1, N_e=4)
approximate(f=x*(1-x)**8, symbolic=False, d=2, N_e=2)
approximate(f=x*(1-x)**8, symbolic=False, d=1, N_e=8)
approximate(f=x*(1-x)**8, symbolic=False, d=2, N_e=8)
```



Comparison of finite element and finite difference approximation

- Finite difference approximation u_i of a function f(x): simply choose $u_i = f(x_i)$
- This is the same as $u \approx \sum_i c_i \varphi_i$ + interpolation (see next slide)
- $u \approx \sum_i c_i \varphi_i + \mathsf{Galerkin/projection}$ or least squares method: must derive and solve a linear system
- ullet What is really the difference in the approximation u?

Interpolation/collocation with finite elements

Let $\{x_i\}_{i\in\mathcal{I}_s}$ be the nodes in the mesh. Collocation/interpolation means

$$u(x_i) = f(x_i), \quad i \in \mathcal{I}_s,$$

which translates to

$$\sum_{i\in\mathcal{I}_n}c_j\varphi_j(x_i)=f(x_i),$$

but $\varphi_j(\mathbf{x}_i)=0$ if $i\neq j$ so the sum collapses to one term $c_i\varphi_i(\mathbf{x}_i)=c_i$, and we have the result

$$c_i = f(x_i)$$

Same result as the standard finite difference approach, but finite elements define u also between the x_i points

Galerkin/project and least squares vs collocation/interpolation or finite differences

- Scope: work with P1 elements
- Use projection/Galerkin or least squares (equivalent)
- Interpret the resulting linear system as finite difference equations

The P1 finite element machinery results in a linear system where equation no i is

$$\frac{h}{6}(u_{i-1}+4u_i+u_{i+1})=(f,\varphi_i)$$

Note:

- ullet We have used u_i for c_i to make notation similar to finite differences
- The finite difference counterpart is just $u_i = f_i$

Expressing the left-hand side in finite difference operator notation

Rewrite the left-hand side of finite element equation no i:

$$h(u_i + \frac{1}{6}(u_{i-1} - 2u_i + u_{i+1})) = [h(u + \frac{h^2}{6}D_xD_xu)]_i$$

This is the standard finite difference approximation of

$$h(u+\frac{h^2}{6}u'')$$

Treating the right-hand side; Trapezoidal rule

$$(f,\varphi_i) = \int_{x_{i-1}}^{x_i} f(x) \frac{1}{h} (x - x_{i-1}) dx + \int_{x_i}^{x_{i+1}} f(x) \frac{1}{h} (1 - (x - x_i)) dx$$

Cannot do much unless we specialize \boldsymbol{f} or use $\boldsymbol{numerical\ integration}.$

Trapezoidal rule using the nodes:

$$(f,\varphi_i) = \int_{\Omega} f\varphi_i dx \approx h \frac{1}{2} (f(x_0)\varphi_i(x_0) + f(x_N)\varphi_i(x_N)) + h \sum_{j=1}^{N-1} f(x_j)\varphi_i(x_j)$$

 $arphi_i(\mathsf{x}_j) = \delta_{ij}$, so this formula collapses to one term:

$$(f, \varphi_i) \approx hf(x_i), \quad i = 1, \ldots, N-1.$$

Same result as in collocation (interpolation) and the finite difference method!

Treating the right-hand side; Simpson's rule

$$\int_{\Omega} g(x)dx \approx \frac{h}{6} \left(g(x_0) + 2 \sum_{j=1}^{N-1} g(x_j) + 4 \sum_{j=0}^{N-1} g(x_{j+\frac{1}{2}}) + f(x_{2N}) \right),$$

Our case: $g=f\varphi_i$. The sums collapse because $\varphi_i=0$ at most of the points.

$$(f,\varphi_i) \approx \frac{h}{3} (f_{i-\frac{1}{2}} + f_i + f_{i+\frac{1}{2}})$$

Conclusions:

- While the finite difference method just samples f at x_i , the finite element method applies an average (smoothing) of f around x_i
- ullet On the left-hand side we have a term $\sim hu''$, and u'' also contribute to smoothing
- There is some inherent smoothing in the finite element method

Finite element approximation vs finite differences

With Trap ezoidal integration of (f, φ_i) , the finite element method essentially solve

$$u + \frac{h^2}{6}u'' = f$$
, $u'(0) = u'(L) = 0$,

by the finite difference method

$$[u+\frac{h^2}{6}D_{x}D_{x}u=f]_i$$

With Simpson integration of (f, φ_i) we essentially solve

$$[u+\frac{h^2}{6}D_xD_xu=\bar{f}]_i,$$

where

$$\bar{f}_i = \frac{1}{3}(f_{i-1/2} + f_i + f_{i+1/2})$$

Note: as $h \to 0$, $hu'' \to 0$ and $\bar{f}_i \to f_i$.

Making finite elements behave as finite differences

- Can we adjust the finite element method so that we do not get the extra hu" smoothing term and averaging of f?
- This allows finite elements to inherit (desired) properties of finite differences

Result:

- Compute all integrals by the Trapezoidal method and P1 elements
- Specifically, the coefficient matrix becomes diagonal ("lumped") - no linear system (!)
- Loss of accuracy? The Trapezoidal rule has error $\mathcal{O}(h^2)$, the same as the approximation error in P1 elements

Limitations of the nodes and element concepts

So far,

- Nodes: points for defining φ_i and computing u values
- Elements: subdomain (containing a few nodes)
- This is a common notion of nodes and elements

One problem:

- Our algorithms need nodes at the element boundaries
- This is often not desirable, so we need to throw the nodes and elements arrays away and find a more generalized element concept

The generalized element concept has cells, vertices, nodes, and degrees of freedom

We introduce cell for the subdomain that we up to now called element

- A cell has vertices (interval end points)
- Nodes are, almost as before, points where we want to compute unknown functions
- ullet Degrees of freedom is what the c_j represent (usually function values at nodes)

The concept of a finite element

- a reference cell in a local reference coordinate system
- $oldsymbol{0}$ a set of basis functions $ilde{arphi}_r$ defined on the cell
- lack lack a set of degrees of freedom (e.g., function values) that uniquely determine the basis functions such that $\tilde{arphi}_r=1$ for degree of freedom number r and $\tilde{arphi}_r=0$ for all other degrees of freedom
- a mapping between local and global degree of freedom numbers (dof map)
- a geometric *mapping* of the reference cell onto to cell in the physical domain: $[-1,1] \Rightarrow [x_L,x_R]$

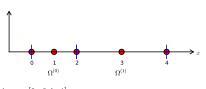
Basic data structures: vertices, cells, dof_map

- Cell vertex coordinates: vertices (equals nodes for P1 elements)
- Element vertices: cells[e][r] holds global vertex number of local vertex no r in element e (same as elements for P1 elements)
- dof_map[e,r] maps local dof r in element e to global dof number (same as elements for Pd elements)

The assembly process now applies dof_map:

A[dof_map[e][r], dof_map[e][s]] += A_e[r,s] b[dof_map[e][r]] += b_e[r]

Example: data structures for P2 elements



vertices = [0, 0.4, 1]
cells = [[0, 1], [1, 2]]
dof_map = [[0, 1, 2], [2, 3, 4]]

Example: P0 elements

Example: Same mesh, but u is piecewise constant in each cell (P0 element). Same vertices and cells, but

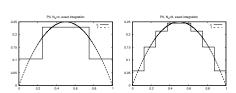
May think of one node in the middle of each element.

Note

We will hereafter work with cells, vertices, and dof_map.

A program with the fundamental algorithmic steps

Approximating a parabola by P0 elements



The approximate function automates the steps in the previous

```
from fe_approx1D_numint import *
x=sym.Symbol('x")
for N_e in 4, 8:
approximate(x*(1-x), d=0, N_e=N_e, Omega=[0,1])
```

Computing the error of the approximation; principles

$$L^2$$
 error: $||e||_{L^2}=\left(\int_\Omega e^2 dx
ight)^{1/2}$

Accurate approximation of the integral:

- Sample u(x) at many points in each element (call u_glob, returns x and u)
- Use the Trapezoidal rule based on the samples
- It is important to integrate u accurately over the elements
- (In a finite difference method we would just sample the mesh point values)

Computing the error of the approximation; details

Note

We need a version of the Trapezoidal rule valid for non-uniformly spaced points:

$$\int_{\Omega} g(x) dx \approx \sum_{j=0}^{n-1} \frac{1}{2} (g(x_j) + g(x_{j+1})) (x_{j+1} - x_j)$$

How does the error depend on h and d?

Theory and experiments show that the least squares or projection/Galerkin method in combination with Pd elements of equal length h has an error

$$||e||_{L^2} = Ch^{d+1}$$

where C depends on f, but not on h or d

Cubic Hermite polynomials; definition

• Can we construct $\varphi_i(x)$ with continuous derivatives? Yes!

Consider a reference cell [-1,1]. We introduce two nodes, X=-1and X = 1. The degrees of freedom are

- 0: value of function at X = -1
- 1: value of first derivative at X = -1
- ullet 2: value of function at X=1
- ullet 3: value of first derivative at X=1

Derivatives as unknowns ensure the same $\varphi'_i(x)$ value at nodes and thereby continuous derivatives.

Cubic Hermite polynomials; derivation

4 constraints on $\tilde{\varphi}_r$ (1 for dof r. 0 for all others):

•
$$\tilde{\varphi}_0(X_{(0)}) = 1$$
, $\tilde{\varphi}_0(X_{(1)}) = 0$, $\tilde{\varphi}'_0(X_{(0)}) = 0$, $\tilde{\varphi}'_0(X_{(1)}) = 0$

•
$$\tilde{\varphi}'_1(X_{(0)}) = 1$$
, $\tilde{\varphi}'_1(X_{(1)}) = 0$, $\tilde{\varphi}_1(X_{(0)}) = 0$, $\tilde{\varphi}_1(X_{(1)}) = 0$

$$ilde{arphi}_2(X_{(1)})=1$$
, $ilde{arphi}_2(X_{(0)})=0$, $ilde{arphi}_2'(X_{(0)})=0$, $ilde{arphi}_2'(X_{(1)})=0$

•
$$\tilde{\varphi}_3'(X_{(1)}) = 1$$
, $\tilde{\varphi}_3'(X_{(0)}) = 0$, $\tilde{\varphi}_3(X_{(0)}) = 0$, $\tilde{\varphi}_3(X_{(1)}) = 0$

This gives 4 linear, coupled equations for each $\tilde{\varphi}_r$ to determine the 4 coefficients in the cubic polynomial

Cubic Hermite polynomials; result

$$\tilde{\varphi}_0(X) = 1 - \frac{3}{4}(X+1)^2 + \frac{1}{4}(X+1)^3 \tag{19}$$

$$\tilde{\varphi}_1(X) = -(X+1)(1-\frac{1}{2}(X+1))^2 \tag{20}$$

$$\tilde{\varphi}_1(X) = -(X+1)(1-\frac{1}{2}(X+1))^2$$

$$\tilde{\varphi}_2(X) = \frac{3}{4}(X+1)^2 - \frac{1}{2}(X+1)^3$$
(20)

$$\tilde{\varphi}_3(X) = -\frac{1}{2}(X+1)(\frac{1}{2}(X+1)^2 - (X+1)) \tag{22}$$

(23)

Numerical integration

- ullet $\int_{\Omega}farphi_{i}dx$ must in general be computed by numerical
- Numerical integration is often used for the matrix too

Common form of a numerical integration rule

$$\int_{-1}^{1} g(X) dX \approx \sum_{j=0}^{M} w_{j} g(\bar{X}_{j}),$$

where

- \bullet \bar{X}_i are integration points
- wi are integration weights

Different rules correspond to different choices of points and weights

The Midpoint rule

Simplest possibility: the Midpoint rule,

$$\int_{-1}^{1} g(X)dX \approx 2g(0), \quad \bar{X}_{0} = 0, \ w_{0} = 2,$$

Exact for linear integrands

Newton-Cotes rules apply the nodes

- Idea: use a fixed, uniformly distributed set of points in [-1, 1]
- The points often coincides with nodes
- ullet Very useful for making $arphi_iarphi_j=0$ and get diagonal ("mass") matrices ("lumping")

The Trapezoidal rule:

$$\int_{-1}^{1} g(X)dX \approx g(-1) + g(1), \quad \bar{X}_{0} = -1, \ \bar{X}_{1} = 1, \ w_{0} = w_{1} = 1,$$

Simpson's rule:

$$\int_{-1}^{1} g(X)dX \approx \frac{1}{3} \left(g(-1) + 4g(0) + g(1) \right),$$

where

 $\bar{\mathbf{v}}_{-}$ 1 $\bar{\mathbf{v}}_{-}$ 0 $\bar{\mathbf{v}}_{-}$ 1 ... 4

Gauss-Legendre rules apply optimized points

- Optimize the location of points to get higher accuracy
- Gauss-Legendre rules (quadrature) adjust points and weights to integrate polynomials exactly

$$M=1: \quad \bar{X}_0=-\frac{1}{\sqrt{3}}, \ \bar{X}_1=\frac{1}{\sqrt{3}}, \ w_0=w_1=1$$
 (24)

$$M=2:$$
 $\bar{X}_0=-\sqrt{\frac{3}{5}}, \ \bar{X}_0=0, \ \bar{X}_2=\sqrt{\frac{3}{5}}, \ w_0=w_2=\frac{5}{9}, \ w_1=\frac{8}{9}$ (25)

- M=1: integrates 3rd degree polynomials exactly
- M=2: integrates 5th degree polynomials exactly
- \bullet In general, M-point rule integrates a polynomial of degree 2M+1 exactly.

See numint.py for a large collection of Gauss-Legendre rules.

Approximation of functions in 2D

Extensibility of 1D ideas.

All the concepts and algorithms developed for approximation of 1D functions f(x) can readily be extended to 2D functions f(x,y) and 3D functions f(x,y,z). Key formulas stay the same.

Quick overview of the 2D case

Inner product in 2D:

$$(f,g) = \int_{\Omega} f(x,y)g(x,y)dxdy$$

Least squares and project/Galerkin lead to a linear system

$$\sum_{j \in \mathcal{I}_s} A_{i,j} c_j = b_i, \quad i \in \mathcal{I}_s$$
 $A_{i,j} = (\psi_i, \psi_j)$

 $b_i = (f, \psi_i)$ Challenge: How to construct 2D basis functions $\psi_i(x,y)$?

2D basis functions as tensor products of 1D functions

Use a 1D basis for x variation and a similar for y variation:

$$V_{x} = \text{span}\{\hat{\psi}_{0}(x), \dots, \hat{\psi}_{N_{x}}(x)\}$$
 (26)

$$V_{\nu} = \operatorname{span}\{\hat{\psi}_0(y), \dots, \hat{\psi}_{N_{\nu}}(y)\} \tag{27}$$

The 2D vector space can be defined as a *tensor product* $V = V_x \otimes V_y$ with basis functions

$$\psi_{p,q}(x,y) = \hat{\psi}_p(x)\hat{\psi}_q(y) \quad p \in \mathcal{I}_x, q \in \mathcal{I}_y$$
.

Tensor products

Given two vectors $a=(a_0,\ldots,a_M)$ and $b=(b_0,\ldots,b_N)$ their outer tensor product, also called the dyadic product, is $p=a\otimes b$, defined through

$$p_{i,j} = a_i b_j, \quad i = 0, \dots, M, \ j = 0, \dots, N.$$

Note: p has two indices (as a matrix or two-dimensional array)

Example: 2D basis as tensor product of 1D spaces,

$$\psi_{p,q}(x,y) = \hat{\psi}_p(x)\hat{\psi}_q(y), \quad p \in \mathcal{I}_x, q \in \mathcal{I}_y$$

Double or single index?

The 2D basis can employ a double index and double sum:

$$u = \sum_{p \in \mathcal{I}_x} \sum_{q \in \mathcal{I}_y} c_{p,q} \psi_{p,q}(x,y)$$

Or just a single index:

$$u = \sum_{j \in \mathcal{I}_s} c_j \psi_j(x, y)$$

with an index mapping (p,q) o i:

$$\psi_i(x, y) = \hat{\psi}_p(x)\hat{\psi}_q(y), \quad i = p(N_y + 1) + q \text{ or } i = q(N_x + 1) + p$$

Example on 2D (bilinear) basis functions; formulas

In 1D we use the basis

$$\{1,x\}$$

2D tensor product (all combinations):

$$\psi_{0.0} = 1$$
, $\psi_{1.0} = x$, $\psi_{0.1} = y$, $\psi_{1.1} = xy$

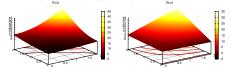
or with a single index:

$$\psi_0 = 1$$
, $\psi_1 = x$, $\psi_2 = y$, $\psi_3 = xy$

See notes for details of a hand-calculation.

Example on 2D (bilinear) basis functions; plot

Quadratic $f(x,y) = (1+x^2)(1+2y^2)$ (left), bilinear u (right):



Implementation; principal changes to the 1D code

Very small modification of approx1D.py:

- Omega = [[O, L_x], [O, L_y]]
- Symbolic integration in 2D
- Construction of 2D (tensor product) basis functions

Implementation; 2D integration

Implementation; 2D basis functions

Tensor product of 1D "Taylor-style" polynomials x^i :

```
def taylor(x, y, Nx, Ny):
    return [x**i*y**j for i in range(Nx+1) for j in range(Ny+1)]
```

Tensor product of 1D sine functions $\sin((i+1)\pi x)$:

```
def sines(x, y, Nx, Ny):
    return [sym.sin(sym.pi*(i+1)*x)*sym.sin(sym.pi*(j+1)*y)
    for i in range(Nx+1) for j in range(Ny+1)]
```

Complete code in approx2D.py

Implementation; application

$$\begin{split} f(x,y) &= (1+x^2)(1+2y^2) \\ >>> & \text{from approx2D import *} \\ >>> & f = (1+x**2)*(1+2*y***2) \\ >>> & \text{psi} = & \text{taylor}(x, y, 1, 1) \\ >>> & \text{omega} = & [10, 2], [0, 2]] \\ >>> & u, c = & \text{least_squares}(f, psi, 0mega) \\ >> & \text{print u} \\ & \text{8x*y} - & 2*x/3 + 4*y/3 - 1/9 \\ >>> & \text{print sym expand}(f) \\ & \text{2*x**2*y**2} + & \text{x**2} + & 2*y**2 + 1 \end{split}$$

Implementation; trying a perfect expansion

Add higher powers to the basis such that $f \in V$:

```
>>> psi = taylor(x, y, 2, 2)
>>> u, c = least_squares(f, psi, Omega)
>>> print u
>>> print u-f
```

Expected: u=f when $f\in V$

Generalization to 3D

Key idea:

$$V = V_x \otimes V_y \otimes V_z$$

Repeated outer tensor product of multiple vectors

$$\begin{aligned} a^{(q)} &= (a_0^{(q)}, \dots, a_{N_q}^{(q)}), \quad q = 0, \dots, m \\ p &= a^{(0)} \otimes \dots \otimes a^{(m)} \\ p_{i_0, i_1, \dots, i_m} &= a_{i_1}^{(0)} a_{i_1}^{(1)} \dots a_{i_m}^{(m)} \end{aligned}$$

$$\psi_{p,q,r}(x,y,z) = \hat{\psi}_p(x)\hat{\psi}_q(y)\hat{\psi}_r(z)$$

$$u(x,y,z) = \sum_{p \in \mathcal{I}_x} \sum_{q \in \mathcal{I}_y} \sum_{r \in \mathcal{I}_z} c_{p,q,r} \psi_{p,q,r}(x,y,z)$$

Finite elements in 2D and 3D

The two great advantages of the finite element method:

- Can handle complex-shaped domains in 2D and 3D
- Can easily provide higher-order polynomials in the approximation

Finite elements in 1D: mostly for learning, insight, debugging

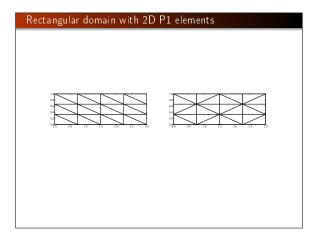
Examples on cell types

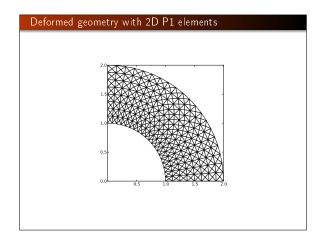
2D:

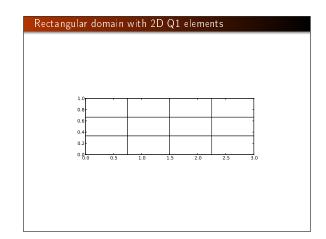
- triangles
- quadrilaterals

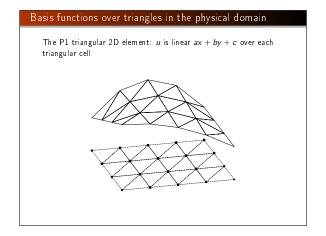
3D:

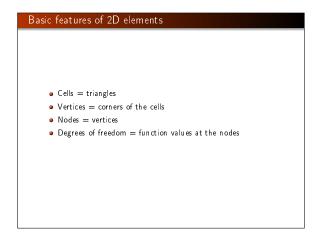
- tetrah edra
- h ex a h e d r a

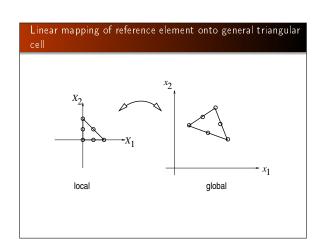


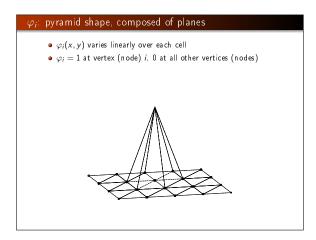


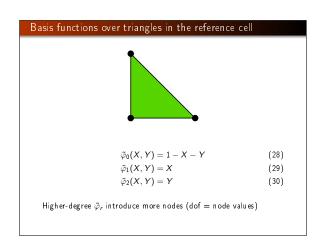


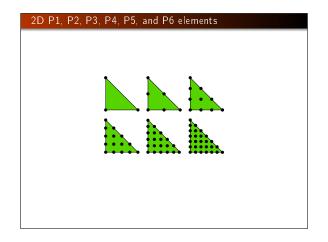


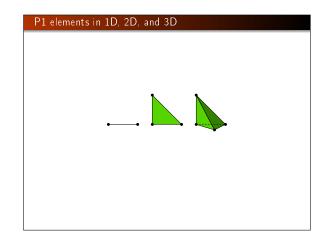


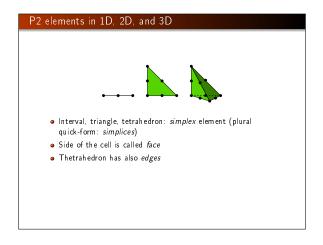












Affine mapping of the reference cell; formula

Mapping of local $\mathbf{X}=(X,Y)$ coordinates in the reference cell to global, physical $\mathbf{x}=(x,y)$ coordinates:

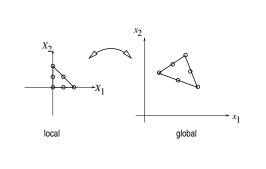
$$\mathbf{x} = \sum_{r} \tilde{\varphi}_{r}^{(1)}(\mathbf{X}) \mathbf{x}_{q(\mathbf{e},r)} \tag{31}$$

where

- r runs over the local vertex numbers in the cell
- x_i are the (x, y) coordinates of vertex i
- $\tilde{\varphi}_r^{(1)}$ are P1 basis functions

This mapping preserves the straight/planar faces and edges.

Affine mapping of the reference cell; figure

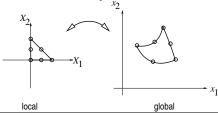


Isoparametric mapping of the reference cell

Idea: Use the basis functions of the element (not only the P1 functions) to map the element

$$\mathbf{x} = \sum_{r} \tilde{\varphi}_{r}(\mathbf{X}) \mathbf{x}_{q(e,r)}$$

Advantage: higher-order polynomial basis functions now map the reference cell to a *curved* triangle or tetrahedron.



Computing integrals

Integrals must be transformed from $\Omega^{(e)}$ (physical cell) to $\tilde{\Omega}^r$ (reference cell):

$$\int_{\Omega^{(e)}} \varphi_i(\mathbf{x}) \varphi_j(\mathbf{x}) \, d\mathbf{x} = \int_{\tilde{\Omega}^r} \tilde{\varphi}_i(\mathbf{X}) \tilde{\varphi}_j(\mathbf{X}) \det J \, d\mathbf{X}$$

$$\int_{\Omega^{(e)}} \varphi_i(\mathbf{x}) f(\mathbf{x}) \, d\mathbf{x} = \int_{\tilde{\Omega}^r} \tilde{\varphi}_i(\mathbf{X}) f(\mathbf{x}(\mathbf{X})) \det J \, d\mathbf{X}$$
(32)

where dx = dxdy or dx = dxdydz and $\det J$ is the determinant of the Jacobian of the mapping x(X).

$$J = \left[\begin{array}{cc} \frac{\partial x}{\partial X} & \frac{\partial x}{\partial Y} \\ \frac{\partial y}{\partial X} & \frac{\partial y}{\partial Y} \end{array} \right], \quad \det J = \frac{\partial x}{\partial X} \frac{\partial y}{\partial Y} - \frac{\partial x}{\partial Y} \frac{\partial y}{\partial X}$$

Affine mapping (31): det $J = 2\Delta$, $\Delta = \text{cell volume}$

Remark on going from 1D to 2D/3D

Finite elements in 2D and 3D builds on the same *ideas* and *concepts* as in 1D, but there is simply much more to compute because the specific mathematical formulas in 2D and 3D are more complicated and the book keeping with dof maps also gets more complicated. The manual work is tedious, lengthy, and error-prone so automation by the computer is a must.