Previous Lecture

* Applications of Glauss haw

-) Planar symmetry (Infinite sheet)
- 11) Cylindrical symmetry (Infinitely long cylindrical rad)
- 111) spherical symmetry
 - D Planar Symmetry

Non conducting sheet
$$\Rightarrow$$
 $E = \frac{\sigma}{2\epsilon_0}$
 \Rightarrow conducting sheet \Rightarrow $E = \frac{\sigma}{\epsilon_0}$

by two parallel conducting | non conducting plates
at their central point.

11) Cylindrical Symmetry

Goal of study -: We have to find an expression for the electric field (E) at a distance (r) from the central axis of the rod.

Condition on charges and field

Assuming that charge distribution and the field have cylindrical symmetry

Explanation:

For this purpose, we have a section of long eylindrical plastic rod with a uniform charge density (1).

BNIF BT

E-field

To find the field at radius 'r', we consider a cylindrical gaussian surface' concentric with a rod. (see figure in Book)

>E (skims)

> Graussian

> plastic rod.

Flux

Since the charges are tive so for each patch element, the E-field must be radially outward.

Curved

Surface

we know that

Φ = EACOSO

For end caps (top bottom).

At both top and bottom ends E-field skims the surface (piercing) of the surface and makes an angle of 90° with area A. so, we have.

Otop = EACOS 900 = 0

Dottom = EA COS 90° = 0

For curved surface

E-field and A are parallel, so

Pourved = EA cos 0° = EA

where A = 2TTxh

 $\Phi = E(2\pi Yh)$

using Grauss Law

Where Venc = 1h -> (i

using (i & (ii in equ @) we get.

$$E = \frac{\lambda}{a\pi\epsilon_{o}\gamma}$$

 $E = \frac{1}{2\pi\epsilon_0 r}$ (magnitule of \vec{E})

Direction of E

The direction of E's radially outward for the charges and radially inward for - we charge.

111) Spherical symmetry

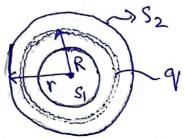
A shell of uniform charge attracts or replets a charged particle that is outside the shell as if all the shell's charge were concentrated at the center of the shell.

To understand spherical symmetry, first we a check E-field inside outsite the charged spherical shell by drawing crassian surfaces.

SI -> surface SI

S2 -> surface S2

Y -> radius from center to outer surface Sz.



R > ratios from center to total charge (Shell).

V > total charged enclosed by shell.

Glauss Law for St (Y > R)

For the assumption, r>R the amount of total charge is enclosed so we can write electric field value equal to Electric field hue to a point charge.

e.g -:

Day SZ \ E = K CV

Glauss Law for SI (r< R)

under, rer assumption, we can see, no charge enclosed by SI, so

E=0

E-field due to any arbitrary sphere enclosing charge

>> spherical Gaussian surface

Total charge enclosed = 9
radius from center to charge
distribution = R

radius from center to any point of craussian surface = T

B = suxpose charge density.

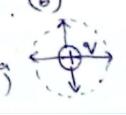
Craussian curface

A ssumption

> 8 can vary with central radius but not vary from surface to surface.

Graussian surface (a) (Y>R)

For this surface, entire charge of ties inside the surface for, YSR. we have $E = \frac{q \cdot k}{rz}$ (behaves as a point charge)



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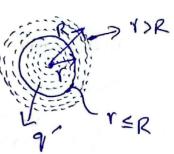
Craussian surface (b) (YER)

For this Craussian surface, total charge is divided into two regions, inside and outside the surface (b). So,

Q = total charge enclosed by surface r > R. Q' = charge enclosed by surface $r \le R$.

we have to find total E-field due to both surfaces.

For outside the surface MAN , though that is 1>R, there is no Grassian surface so outside charge does not set up field.



For inside the surface, rsp.

the charge enclosed is ay and the electric field due to this charge will be

E = Kay -> (1)

If total charge is uniform, then of enclosed is also uniform so we can write

$$V = \frac{4}{3}\pi R^3$$

$$V' = \frac{4}{3}\pi r^3$$
and

$$E = \frac{k \cdot \cancel{y}}{\cancel{y}^2} = \frac{k \cdot \cancel{y} \cdot \cancel{y}^3}{\cancel{x}^2}$$

$$\begin{bmatrix} E = k \frac{9}{R^3} \cdot r \end{bmatrix}$$