

Bifurcation of coherent structures in nonlocally coupled system

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Coherent structures

- ① spikes, fronts, wave trains... (Maybe some pictures?)

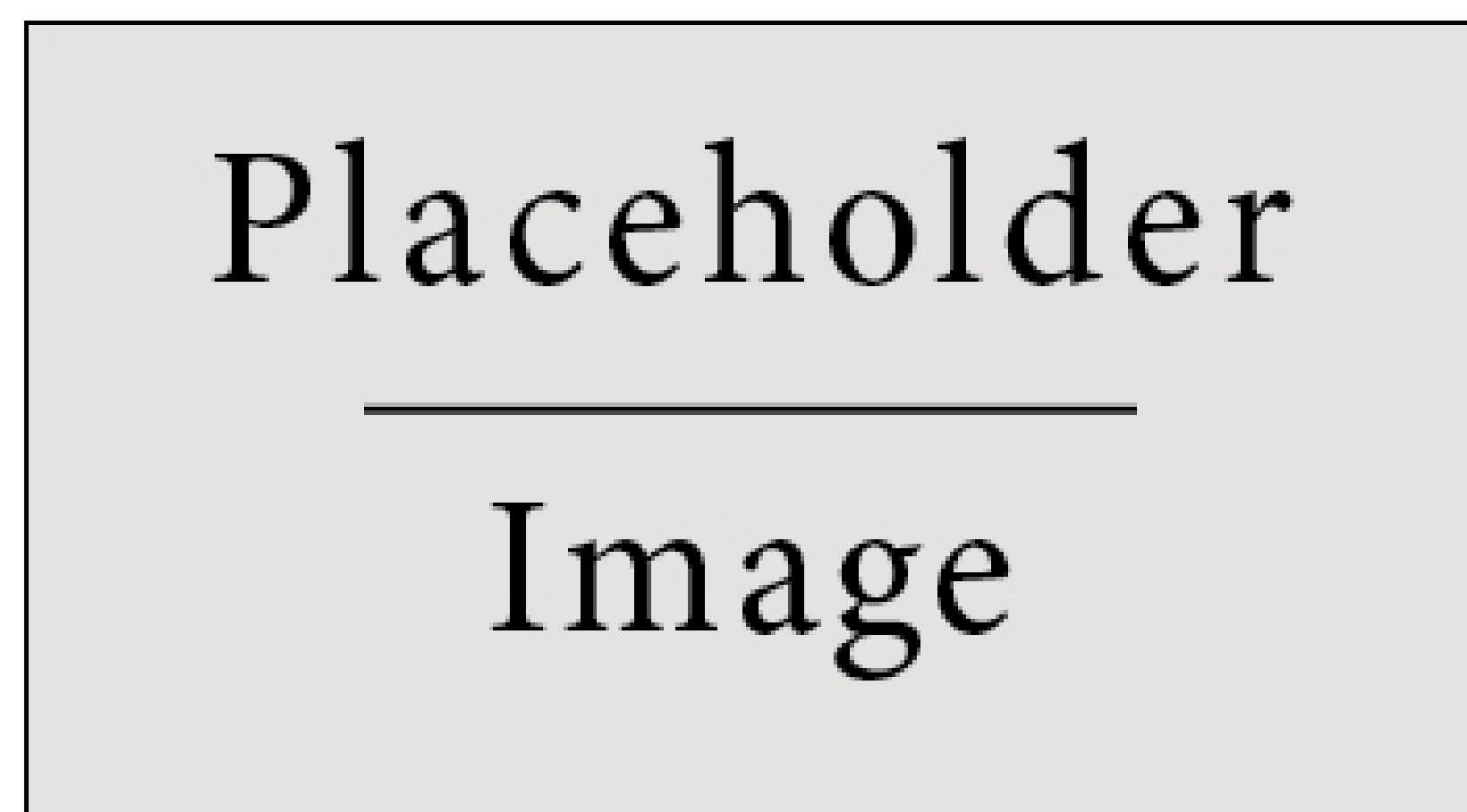


Figure 1: Figure caption

Reaction-Diffusion Models

Bifurcation to coherent structures in RD system

$$U_t = D\Delta_x U + F(U; \mu).$$

Assume

- $U = U(x) \in \mathbb{R}^k, x \in \mathbb{R}^n$;
- $\mu \in \mathbb{R}^p$ parameter, nonlinearity N satisfies $N(0; \mu) = 0$.

Radially symmetric patterns are studied in [1] using spatial dynamic techniques.

- ① rewrite the stationary equation as a nonautonomous system of ODE;
- ② carefully construct a center manifold;
- ③ study reduced dynamics on the center manifold, get complete classification of small bounded solutions.

Nonlocal Model Equations

Nonlocal diffusion are ubiquitous in modeling of natural phenomena [2]

- ① neural field models: $u_t = -K * S(u)$,
- ② water wave equations: $u_t = (Mu - u^2)_x$.

Can we find similar patterns in these models? will they have different properties due to nonlocality?

A Bifurcation Problem

We study the following system of equations for $U(x) \in \mathbb{R}^k$ with $x \in \mathbb{R}^n$

$$U + K * U - N(U; \mu) = 0. \quad (1)$$

Effective linear diffusive coupling (L):

- K is a matrix of convolution kernel with finite second moments $K(x), |x|^2 K(x) \in L^1(\mathbb{R}^n)$
- K is symmetric, $K(\gamma x) = K(x)$ for all $\gamma \in \Gamma \subset O(n)$. The fixed point set of Γ is $\{0\}$ only.
- the Fourier determinant $D(\xi) = \det(I + \bar{K}(\xi))$ has $D(0) = 0, D'(0) = 0, D''(0) \neq 0$.

Let e span the kernel of $I + \bar{K}(0)$, choose e^* span the cokernel.

A Bifurcation Problem-Continued

Transcritical bifurcation in kinetics (TC):

- We assume a transcritical bifurcation scenario:

$$\begin{aligned} N(0; \mu) &= 0, \\ \langle e^*, D_{u\mu} N(0; 0)e \rangle &\neq 0, \\ \langle e^*, D_{uu} N(0; 0)[e, e] \rangle &\neq 0 \end{aligned}$$

- N is smooth, so the superposition operator $U(\cdot) \mapsto N(U(\cdot))$ is smooth.

We then find a pseudo-differential operator L and invertible matrix P, Q so that

$$LP(I + K*)Q = \text{diag}(M, I_{k-1}),$$

with $M = (1 - \Delta)^{-1}\Delta$ and system decouples into two equations for a scalar function v_c and a \mathbb{R}^{k-1} valued function v_h .

Properties of the Spike

- ① The spikes constructed here are not necessarily exponentially localized, it depends on the localization of the convolution kernel K , for K algebraically localized we get algebraically localized spikes due to the corrector w .
- ② Typical examples of Γ can be all of $O(n)$ or subgroups generated by reflection across a hyperplane invariant under the action of Γ . Generalize the studies on radial symmetry in the local case.

Further Directions

- ① Stability: probably unstable, notice no Evans function techniques available due to nonlocality. It is possible to use the information on the spectrum of the ground state and a perturbative argument.
- ② Transition to large μ ?: For μ large, in some examples it can be shown many discontinuous solutions exist by implicit function theorem. How do things look in between large μ and small μ ?

References

- [1] Arnd Scheel. Radially symmetric patterns of reaction-diffusion systems. *Mem. Amer. Math. Soc.*, 165(786):viii+86, 2003.
- [2] Paul C. Bressloff. Spatiotemporal dynamics of continuum neural fields. *J. Phys. A*, 45(3):033001, 109, 2012.
- [3] Juncheng Wei and Matthias Winter. *Mathematical aspects of pattern formation in biological systems*, volume 189 of *Applied Mathematical Sciences*. Springer, London, 2014.

Main Result

Fix $n < 6$ and $\ell > n/2$. Assume Hypotheses (L) and (TC), a solution of the form

$$U(x; \mu) = -\beta^{-1}\alpha\mu[v_*(\sqrt{\alpha\mu}x) + w(x; \mu)]e + v_\perp(x; \mu)$$

where v_* is the unique positive ground state of $\Delta v - v + v^2 = 0$, $w \in H^\ell(\mathbb{R}^n)$ is a corrector which converges to 0 as $\mu \rightarrow 0$, lastly, v_\perp satisfy $\langle e, v_\perp \rangle = 0$ and $\|v_\perp\| = O(\mu^2)$.

Sketch of Proof

- Work with function space

$$H_\Gamma^\ell = \{u \in H^\ell, u(\gamma x) = u(x), \gamma \in \Gamma\};$$

- Set $\varepsilon = \sqrt{\alpha\mu}$, then **rescale** by $v_c = -\beta^{-1}\varepsilon^2\tilde{v}_c(\varepsilon x)$, $v_h = \varepsilon^2\tilde{v}_h(\varepsilon x)$, reduce equation

$$\varepsilon^{-2}m^\varepsilon\Delta\tilde{v}_c = \tilde{v}_c - \tilde{v}_c^2 + O(\varepsilon^2) \quad (2)$$

$$\tilde{v}_h = O(\varepsilon^2) \quad (3)$$

with m^ε the rescaled pseudo-differential operator, with symbol $(1 + \varepsilon^2|\xi|^2)^{-1}$;

- solve (3) to express \tilde{v}_h in terms of \tilde{v}_c by a fixed point argument. Get a reduced scalar bifurcation equation;
- substitute the ansatz $\tilde{v}_c = v_* + w$, get an equation in w .

Sketch of Proof-Continued

- Observe $(m^\varepsilon)^{-1} : H^\ell \rightarrow H^{\ell-2}$ is well-defined, with $\|(m^\varepsilon)^{-1} - 1\| = O(\varepsilon^2)$ as $\varepsilon \rightarrow 0$;
- **precondition** equation \tilde{v}_c by the operator $(m^\varepsilon)^{-1}$, simplify, get

$$0 = \Delta w - (w - 2v_*w - w^2) - ((m^\varepsilon)^{-1} - 1)(w - 2v_*w - w^2 + \Delta v_*) + O(\varepsilon^2) \quad (4)$$

denote the right hand of (4) by $F(w; \varepsilon)$, note $F(w; \varepsilon) \rightarrow 0$ in L^2 as $(w; \varepsilon) \rightarrow (0; 0)$, also, $D_w F$ is continuous near $(0, 0)$, with

$$D_w F(0; 0) = \Delta - 1 + 2v_*,$$

which is nondegenerate and invertible from H_Γ^ℓ to $H_\Gamma^{\ell-2}$ ([3]), these facts allow the set up of an Newton iteration scheme to continue $F(w; \varepsilon) = 0$ from H_Γ^ℓ to $H_\Gamma^{\ell-2}$ near $(0; 0)$.