

# Complex Analysis of Askaryan Radiation: Towards UHE- $\nu$ energy Reconstruction via the Hilbert Envelope of Observed Signals

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## I. INTRODUCTION

The introduction.

## II. UNITS, DEFINITIONS, AND CONVENTIONS

The units.

## III. COLLECTION OF MAIN RESULTS

Here is a list of the basic results and ideas for this paper.

- Let the signal model  $s(t)$  be

$$s(t) = -E_0 t e^{-\frac{1}{2}(t/\sigma_t)^2} \quad (1)$$

This is the off-cone field equation from [1]. The parameter  $\sigma_t$  is the pulse width, and it depends two quantities: the longitudinal length of the UHE- $\nu$ -induced cascade, and the angle at which the cascade is observed relative to the Cherenkov angle. The parameter  $E_0$  is the amplitude normalization, and it depends on two parameters:  $\sigma_t$ , and  $\omega_0$ , the cutoff frequency from the cascade form factor.

- Let  $\hat{s}(t)$  represent the Hilbert transform of  $s(t)$ . The *analytic signal* of  $s(t)$  is

$$s_a(t) = s(t) + j\hat{s}(t) \quad (2)$$

The magnitude of the analytic signal,  $|s_a(t)|$ , is the *envelope* of the signal. The Hilbert transform  $\hat{s}(t)$  is equivalent to the convolution of  $s(t)$  and the tempered distribution  $h(t) = 1/(\pi t)$ .

- Let  $S(f)$  be the Fourier transform of  $s(t)$ . The Fourier transform of the analytic signal is

$$\mathcal{F}\{s_a(t)\}_f = S_a(f) = S(f)(1 + \text{sgn } f) \quad (3)$$

The sign function,  $\text{sgn}$  gives  $-1$  if  $f < 0$ ,  $0$  if  $f = 0$ , and  $1$  if  $f > 0$ .

- Taking the inverse Fourier transform of Eq. 3, the analytic signal may be written in terms of  $S(f)$ :

$$s_a(t) = 2 \int_0^\infty S(f) e^{2\pi j f t} df \quad (4)$$

- The Fourier transform of Eq. 1 is

$$S(f) = E_0 \sigma_t^3 (2\pi)^{3/2} j f e^{-2\pi^2 f^2 \sigma_t^2} \quad (5)$$

- Using the gaussian spectral width  $\sigma_f$  from [2], and the gaussian width of  $s(t)$  from [1], it was shown in [1] that the uncertainty principle holds for off-cone signals:

$$\sigma_t \sigma_f \geq \frac{1}{2\pi} \quad (6)$$

The equality is reached in the limit the far-field parameter limits to zero:  $\eta \rightarrow 0$ . This makes the signal spectrum

$$S(f) = E_0 \sigma_t^3 (2\pi)^{3/2} j f e^{-\frac{1}{2}(f/\sigma_f)^2} \quad (7)$$

Inserting  $S(f)$  into Eq. 4,  $s_a(t)$  is

$$s_a(t) = \frac{E_0 \sigma_t^3 (2\pi)^{3/2}}{\pi} \frac{d}{dt} \int_0^\infty e^{-\frac{1}{2}(f/\sigma_f)^2} e^{2\pi j f t} df \quad (8)$$

- Let  $k^2/4 = \frac{1}{2}(f/\sigma_f)^2$ , and  $x = t/(\sqrt{2}\sigma_t)$ . Equation 8 can be broken into real and imaginary parts:

$$s_a(t) = \frac{E_0 \sigma_t}{\sqrt{2\pi}} \frac{dI}{dx} \quad (9)$$

$$\Re\{I\} = \int_0^\infty e^{-k^2/4} \cos(kx) dk \quad (10)$$

$$\Im\{I\} = \int_0^\infty e^{-k^2/4} \sin(kx) dk \quad (11)$$

The real part of  $I$  is even, so it can be extended to  $(-\infty, \infty)$  if it is multiplied by  $1/2$ . The result is

$$\Re\{I\} = \sqrt{\pi} e^{-x^2} \quad (12)$$

The imaginary part of  $I$  is proportional to *Dawson's integral*,  $D(x)$ :

$$\Im\{I\} = 2D(x) \quad (13)$$

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- The overall analytic signal,  $s_a(t)$ , is

$$s_a(t) = -E_0 \left( t e^{-\frac{1}{2}(t/\sigma_t)^2} - \frac{2j\sigma_t}{\sqrt{2\pi}} \frac{dD(x)}{dx} \right) \quad (14)$$

The signal envelope is  $|s_a(t)|$ . It is important to note that, though  $D(x)$  is not evaluated analytically, a high-precision algorithm for computing  $D(x)$  was given in [3]. Note that  $s_a(0) \neq 0$ , since  $dD(x)/dx = 1 - 2xD(x)$ .

- Signal data in detectors designed to observe Askaryan pulses is equivalent to the convolution of the signal and detector response functions. Signal models are convolved with measured detector responses to create *signal templates*. Signal templates are cross-correlated with observed data to identify UHE- $\nu$  signals. The oscillations of signal templates and observed data can introduce various uncertainties when cross-correlated. This problem intensifies when the signal-to-noise ratio between Askaryan pulse data and thermal noise decreases. To reduce these uncertainties, the Hilbert envelope of observed signals is used in cross-correlations instead of the original signals. We seek an analytic equation for the Hilbert envelope of the data. That is, we seek the envelope of the convolution of the analytic signal model with a typical detector response. The RLC damped oscillator is a standard circuit model for the RF dipole antennas used in RNO-G and the proposed IceCube Gen2 [4–6].
- There are two paths to calculating the final result. The first option involves three steps. First, the detector response,  $r(t)$  is convolved with  $s(t)$ . Second, the analytic signal of the result is found. Third, the magnitude of the analytic signal is computed, which can be compared to envelopes of observed signals. The second option involves computing the envelope of the convolution of  $r(t)$  with  $s(t)$  directly from  $s_a(t)$  and  $r_a(t)$ .
- Let  $x(t)*y(t)$  represent the convolution of two functions  $x(t)$  and  $y(t)$ . Let the envelope of the convolution be  $\mathcal{E}_{x*y}(t)$ .  $\mathcal{E}_{x*y}(t)$ ,  $x_a(t)$ , and  $y_a(t)$  are related by

$$\mathcal{E}_{x*y}(t) = \frac{1}{2} |x_a(t) * y_a(t)| \quad (15)$$

We should include the proof here (I have to track it down in my notebooks).

- The analytic signal of the Askaryan pulse is given by Eq. 14. The RLC damped oscillator response and corresponding analytic signal are

$$r(t) = R_0 e^{-\gamma t} \cos(2\pi f_0 t) \quad (16)$$

$$r_a(t) = R_0 e^{-\gamma t} e^{2\pi j f_0 t} \quad (17)$$

The parameter  $\gamma$  is the *decay constant*, and the parameter  $f_0$  is the resonance frequency. Note that the envelope of  $r(t)$ ,  $|r_a(t)|$ , is simply  $R_0 \exp(-\gamma t)$ , as it should be. The proof of Eq. 17 is as follows:

$$r(t) = R_0 e^{-\gamma t} \cos(2\pi f_0 t) \quad (18)$$

$$R(f) = \frac{R_0}{2\pi j} \left( \frac{f - \frac{j\gamma}{2\pi}}{(f - z_+)(f - z_-)} \right) \quad (19)$$

$$z_+ = f_0 + \frac{j\gamma}{2\pi} \quad (20)$$

$$z_- = -f_0 + \frac{j\gamma}{2\pi} \quad (21)$$

$$\hat{r}(t) = \mathcal{F}^{-1} \{ -j \operatorname{sgn}(f) R(f) \} \quad (22)$$

$$\hat{r}(t) = \frac{R_0}{2j} (e^{2\pi j f_0 t} - e^{-2\pi j f_0 t}) e^{-\gamma t} \quad (23)$$

$$\hat{r}(t) = R_0 \sin(2\pi f_0 t) e^{-\gamma t} \quad (24)$$

$$r_a(t) = R_0 (\cos(2\pi f_0 t) + j \sin(2\pi f_0 t)) e^{-\gamma t} \quad (25)$$

$$r_a(t) = R_0 e^{2\pi j f_0 t} e^{-\gamma t} \quad (26)$$

In evaluating the inverse Fourier transform in Eq. 22, the poles at  $z_+$  and  $z_-$  must be enclosed in separate contour integrals.

- The goal is now to apply Eq. 15 by convolving  $s_a(t)$  with  $r_a(t)$ . The real part of  $s_a(t)$  can be convolved with  $r_a(t)$  in the time domain. As in Eq. 14, real part of  $s_a(t)$  is  $s(t)$ . Since  $r_a(t < 0) = 0$ , the convolution is written

$$r_a(t) * \Re\{s_a(t)\} = \int_0^\infty r(\tau) s(t - \tau) d\tau \quad (27)$$

Let  $x = t/(\sqrt{2}\sigma_t)$ ,  $y = \tau/(\sqrt{2}\sigma_t)$ , and  $z_0 = \sqrt{2}\sigma_t(\gamma - 2\pi j f_0)$ . The convolution may be organized as follows

$$r_a(t) * \Re\{s_a(t)\} = R_0 s(t) I + R_0 E_0 \sigma_t^2 e^{-\frac{1}{2}(t/\sigma_t)^2} \frac{dI}{dt} \quad (28)$$

with

$$I(x, z_0) = \sqrt{2}\sigma_t \int_0^\infty e^{-y^2 + (2x - z_0)y} dy \quad (29)$$

Let  $b = x - z_0/2$ . Completing the square in the exponent produces

$$I(x, z_0) = \sqrt{2}\sigma_t e^{b^2} \int_0^\infty e^{-(y-b)^2} dy \quad (30)$$

The integral may be cast as a complementary error function by substituting  $u = y - b$  and letting  $b = jz$ :

$$I(x, z_0) = \sqrt{\frac{\pi}{2}} \sigma_t e^{-z^2} \operatorname{erfc}(-jz) \quad (31)$$

Thus, the integral is proportional to the *Faddeeva function*,  $w(z)$ :

$$I(x, z_0) = \sqrt{\frac{\pi}{2}} \sigma_t w(z) \quad (32)$$

Substituting Eq. 32 into Eq. 28, and simplifying, produces

$$r_a(t) * \Re\{s_a(t)\} = \sqrt{\frac{\pi}{2}} R_0 \sigma_t s(t) w(z) - j \frac{\sqrt{\pi}}{2} R_0 E_0 \sigma_t^2 e^{-\frac{1}{2}(t/\sigma_t)^2} \frac{dw(z)}{dz} \quad (33)$$

- Turning to the convolution of  $r_a(t)$  with  $\Im(s_a)$ , note that the imaginary part of Eq. 14 contains  $dD(x)/dx$ , the derivative of the Dawson function. To complete the convolution, the inverse Fourier transform of the product of the Fourier transforms of  $r_a(t)$  and  $\Im\{s_a(t)\}$  is calculated. To complete this step, the Fourier transform of  $dD(x)/dx$  is needed:

$$x = t/(\sqrt{2}\sigma_t) \quad (34)$$

$$f' = \sqrt{2}\sigma_t f \quad (35)$$

$$\mathcal{F}\left\{\frac{dD(x)}{dx}\right\} = \int_{-\infty}^{\infty} \frac{dD(x)}{dx} e^{-2\pi j f t} dt \quad (36)$$

$$\mathcal{F}\left\{\frac{dD(x)}{dx}\right\} = \sqrt{2}\sigma_t (2\pi j f') \int_{-\infty}^{\infty} D(x) e^{-2\pi j f' x} dx \quad (37)$$

$$\mathcal{F}\left\{\frac{dD(x)}{dx}\right\} = 2\pi^2 \sigma_t^2 f e^{-\frac{1}{2}(f/\sigma_f)^2} \text{sgn}(f) \quad (38)$$

Let  $z_0 = (\gamma - 2\pi j f_0)$ , so that the Fourier transform of  $r_a$  is

$$\mathcal{F}\{r_a(t)\} = R_0 \int_{-\infty}^{\infty} e^{-zt} e^{-2\pi j f t} dt \quad (39)$$

$$\mathcal{F}\{r_a(t)\} = \frac{R_0}{z_0 + 2\pi j f} \quad (40)$$

Let  $z_1 = f_0 + j\gamma/(2\pi)$ . Multiplying the Fourier transforms, and simplifying, yields

$$\mathcal{F}\{r_a(t) * \Im\{s_a(t)\}\} = -j\sqrt{2\pi} R_0 E_0 \sigma^3 \frac{f \text{sgn}(f) e^{-\frac{1}{2}(f/\sigma_f)^2}}{f - z_1} \quad (41)$$

The inverse Fourier transform now gives  $r_a(t) * \Im\{s_a(t)\}$ :

$$r_a(t) * \Im\{s_a(t)\} = -j\sqrt{2\pi} R_0 E_0 \sigma_t^3 \int_{-\infty}^{\infty} \frac{f \text{sgn}(f) e^{-\frac{1}{2}(f/\sigma_f)^2} e^{2\pi j f t}}{f - z_1} df \quad (42)$$

Let  $y = f/(\sqrt{2}\sigma_f)$ ,  $x = t/(\sqrt{2}\sigma_t)$ ,  $\text{sgn}(f) = \text{sgn}(y)$ , and  $z_1 \rightarrow f_0/(\sqrt{2}\sigma_f) + j\gamma/(2\pi\sqrt{2}\sigma_f)$ . The inverse Fourier transform can be organized as:

$$r_a(t) * \Im\{s_a(t)\} = -\frac{R_0 E_0 \sigma_t^3}{\sqrt{2\pi}} \frac{d}{dt} \int_{-\infty}^{\infty} \frac{\text{sgn}(y) e^{-y^2 + 2jxy}}{y - z_1} dy \quad (43)$$

Completing the square in the exponent, factoring a unit of  $j$  from the exponent, and changing the derivative from  $d/dt$  to  $d/dx$  gives

$$r_a(t) * \Im\{s_a(t)\} = -\frac{R_0 E_0 \sigma_t^2}{2\sqrt{\pi}} \frac{d}{dx} e^{-x^2} \int_{-\infty}^{\infty} \frac{\text{sgn}(y) e^{(x+jy)^2}}{y - z_1} dy \quad (44)$$

Let  $z = x + jy$ , with  $y = \Im(z)$ . In this substitution,  $x$  is treated as a constant, so  $dz = jy$ . Let  $z_2 = x + jz_1$ , while  $z_1 = f_0/(\sqrt{2}\sigma_f) + j\gamma/(2\pi\sqrt{2}\sigma_f)$ . The result is

$$r_a(t) * \Im\{s_a(t)\} = -\frac{R_0 E_0 \sigma_t^2}{2\sqrt{\pi}} \frac{d}{dx} e^{-x^2} \int_{x-j\infty}^{x+j\infty} \frac{\text{sgn}(\Im(z)) e^{z^2}}{y - z_2} dz \quad (45)$$

The integral may be finished by extending  $z$  to the complex plane and using contour integration. The location of the pole is in the first quadrant, unless  $t = (\gamma\sigma_t)\sigma_t$ . In that case,  $\Re(z_2) = 0$ , and the pole is located on the imaginary axis. For typical values of  $\sigma_t$  and  $\gamma$ ,  $t \approx 0.1$  ns. Since  $r(t) = 0$  if  $t < 0$ , the location of  $z_2$  is usually in the first quadrant. Let the line integrals of the rectangular contour be:  $I_1$ ,  $(x, -jR) \rightarrow (x, jR)$ ,  $I_2$ ,  $(x, jR) \rightarrow (0, jR)$ ,  $I_3$ ,  $(0, jR) \rightarrow (0, -jR)$ , and  $I_4$ ,  $(0, -jR) \rightarrow (x, -jR)$ .  $I_2$  and  $I_4$  are both proportional to  $\exp(-R^2)$ , so  $I_2 \rightarrow 0$  and  $I_4 \rightarrow 0$  when  $R \rightarrow \infty$ .  $I_1$  is the integral in Eq. 45. Let  $w = -y$ , so that  $I_3$  is

$$I_3 = -\int_0^R \frac{e^{-y^2}}{y - z_1} dy + \int_0^R \frac{e^{-w^2}}{w + z_1} dw \quad (46)$$

Integrals of the type in  $I_3$  are known as *Goodwin-Staton* functions, and are related to the other error functions. The Goodwin-Staton function,  $G(z_1)$ , is odd, so  $-G(-z_1) = G(z_1)$ . The result for  $I_3$  is

$$I_3 = -G(-z_1) + G(z_1) = 2G(z_1) \quad (47)$$

The residue of the enclosed contour  $I_C = I_1 + I_2 + I_3 + I_4$  is  $\exp(z_2^2)$ . With  $I_2$ ,  $I_3$ , and  $I_4$  completed, the solution for  $I_1$  is

$$I_1 = 2\pi j e^{z_2^2} - 2G(z_1) \quad (48)$$

Inserting this result into Eq. 45, and simplifying, gives

$$r_a(t) * \Im\{s_a(t)\} = -\frac{R_0 E_0 \sigma_t^2}{2\sqrt{\pi}} \frac{d}{dx} () \quad (49)$$

#### IV. CONCLUSION

The conclusion.

## Appendix A: Details

The details.

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