# Solutions for Homework 4

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## 1 Problem 4.10

A sphere of radius R carries a polarization

$$\mathbf{P}(\mathbf{r}) = k\mathbf{r} \tag{1}$$

In Eq. 1, k is a constant and and  $\mathbf{r}$  is the vector from the center.

- (a) Calculate the bound charges  $\sigma_b$  and  $\rho_b$ .
- (b) Find the field inside and outside the sphere.
- (a) The surface bound charge is at radius R, so  $\sigma_b = \mathbf{P} \cdot \hat{\mathbf{n}} = kR$ . The volumetric bound charge is  $\rho_b = -\nabla \cdot \mathbf{P} = -3k$ . Note the total charge should add up to zero:  $(4\pi R^2)\sigma_b + (4/3)\pi R^3\rho_b = 4\pi R^3k 4\pi R^3k = 0$ . (b) The field of any constant volumetric charge density -3k should be calculable via Gauss' law. We find, after integrating over a Gaussian surface of radius r < R:

$$\oint \mathbf{E} \cdot d\mathbf{a} = \mathbf{E} \cdot \mathbf{A} = \frac{1}{\epsilon_0} \rho = \frac{-4k\pi r^3}{\epsilon_0} \tag{2}$$

$$\mathbf{E} = -\frac{3kr}{\epsilon_0}\hat{\mathbf{r}} = -\mathbf{P}/\epsilon_0 \tag{3}$$

(b) Note that, because the net charge is zero, the field outside the sphere is zero.

## 2 Problem 4.14

When you polarize a neutral dielectric, the charge moves a bit, but the total remains zero. This fact should be reflected in the bound charges  $\sigma_b$  and  $\rho_b$ . Prove from Eqs. 4.11 and 4.12 that the total bound charge vanishes.

First, let's integrate the total volumetric bound charge:

$$-q = \int_{\mathcal{V}} \rho_b d\tau = -\int_{\mathcal{V}} \nabla \cdot \mathbf{P} d\tau = -\oint \mathbf{P} \cdot \hat{\mathbf{n}} da = -\oint \sigma_b da$$
 (4)

Next, the total surface bound charge is

$$q = \oint_{S} \sigma_b da \tag{5}$$

Now we see that Q = -q + q = 0.

#### 3 Problem 4.15

 $A\ thick\ spherical\ shell\ (inner\ radius\ a,\ outer\ radius\ b)\ is\ made\ of\ dielectric\ material\ with\ a\ frozen-in\ polarization$ 

$$\mathbf{P}(\mathbf{r}) = \frac{k}{r}\hat{\mathbf{r}} \tag{6}$$

In Eq. 6, k is a constant and r is the distance from the center. There is no free charge in the problem. Find the electric field in all three regions by two different methods:

- (a) Locate all the bound charge, and use Gauss' Law to calculate the field it produces.
- (b) Use Eq. 4.23 to find **D**, and then get **E** from Eq. 4.21. [Notice the second method is faster, and it avoids any explicit reference to the bound charges].

For (a), locating all the bound charge:

 $\bullet$  r=a:

$$\sigma_a = \mathbf{P} \cdot (-\hat{\mathbf{r}}) = -\left(\frac{k}{a}\right)\hat{\mathbf{r}} \tag{7}$$

• a < r < b:

$$\rho = -\nabla \cdot \mathbf{P} = -\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{k}{r} \right) = -\frac{k}{r^2} \tag{8}$$

 $\bullet$  r=b:

$$\sigma_b = \frac{k}{b} \tag{9}$$

• What is the total charge for this neutral object?

$$Q = 4\pi(b-a) - 4\pi \int_{a}^{b} \rho d\tau = 4\pi(b-a) - 4\pi(b-a) = 0$$
 (10)

The neutral object has Q = 0. Using Gauss' Law,  $\mathbf{E} = 0$  for r < a,  $\mathbf{E} = -\mathbf{P}/\epsilon_0$ , and for r > b,  $\mathbf{E} = 0$  because total charge is 0.

For (b), we find that since there is no free charge,  $\mathbf{D} = 0$  and by the definition of  $\mathbf{D}$ ,  $\mathbf{E} = -\mathbf{P}/\epsilon_0$ . Thus, we achieve the same result by making reference to only the free charge.

## 4 Problem 4.18

The space between the plates of a parallel plate capacitor is filled with two slabs of linear dielectric material. Each slab has thickness a, so the total distance between the plates is 2a. Slab 1 has a dielectric constant of 2, and slab 2 has a dielectric constant of 1.5. The free charge density on the top plate is  $\sigma$  and on the bottom plate  $-\sigma$ .

- (a) Using Gauss' law for **D**, we find that  $\mathbf{D} = -\sigma \hat{\mathbf{z}}$  throughout the system.
- (b) Using  $\mathbf{D} = \epsilon \mathbf{E}$ , we find slab 1 has  $\mathbf{E} = -(\sigma/2\epsilon_0)\hat{\mathbf{z}}$  and slab 2 has  $\mathbf{E} = -(2\sigma/3\epsilon_0)\hat{\mathbf{z}}$ .
- (c) Using the definition of **D**, and the results from (b), we find  $\mathbf{P} = -(\sigma/2)\hat{\mathbf{z}}$  for slab 1. For slab 2 we find  $\mathbf{P} = -(\sigma/3)\hat{\mathbf{z}}$ .
- (d) We know that  $\Delta V = E \delta z = E_1 a + E_2 a = (E_1 + E_2) a = (7\sigma a)/(6\epsilon_0)$ .
- (e) There is no  $\rho_b$  because the polarization densities are constant. For the bound surface charges, there is  $-\sigma/2$  at the top of slab 1, and  $\sigma/2$  at the bottom of slab 1 from medium 1. There is  $-\sigma/3$  at the top of slab 2 from medium 2, and  $\sigma/3$  at the bottom of slab 2 from medium 2. At the interface, if we add surface charges, we find  $\sigma/6$ .
- (f) To find the same fields as above, we must combine total charge from free and bound charges at the top and bottom surfaces.

## 5 Problem 4.26

A spherical conductor, of radius a, carries a charge Q. It is surrounded by linear dielectric material of susceptibility  $\chi_e$ , out to radius b. Find the energy of this configuration.

The goal is to find  $\mathbf{E}$  and  $\mathbf{D}$ , so that we can apply

$$W = \frac{1}{2} \int \mathbf{D} \cdot \mathbf{E} \, d\tau \tag{11}$$

The results are:

- Displacement inside the conductor (r < a): 0
- Displacement outside the conductor (r > a):  $Q/(4\pi r^2)\hat{\mathbf{r}}$
- Field inside the conductor: 0
- Field outside the conductor but inside the dielectric (a < r < b):  $Q/(4\pi\epsilon r^2)\hat{\mathbf{r}}$ .
- Field outside the dielectric (r > b):  $Q/(4\pi\epsilon_0 r^2)\hat{\mathbf{r}}$ .

Inserting all these into into Eq. 11, and integrating:

$$W = \frac{Q^2}{8\pi\epsilon_0(1+\gamma_e)} \left(\frac{1}{a} + \frac{\chi_e}{b}\right) \tag{12}$$

## 6 Problem 4.35

A point charge q is in embedded at the center of a sphere of linear dielectric material (with susceptibility  $\chi_e$  and radius R). Find the electric field, the polarization, and the bound charge densities,  $\sigma_b$  and  $\rho_b$ . What is the total bound charge on the surface? Where is the compensating negative bound charge located?

- Use Gauss' Law with electric displacement:  $\mathbf{D} = q/(4\pi r^2)\hat{\mathbf{r}}$ .
- The electric field is then  $\mathbf{E} = \mathbf{D}/\epsilon$ , or

$$\mathbf{E} = \frac{q}{4\pi\epsilon_0 (1 + \chi_e)r^2} \hat{\mathbf{r}} \tag{13}$$

• For polarization, note that linear dielectrics have  $\mathbf{P} = \epsilon_0 \chi_e \mathbf{E}$ , so

$$\mathbf{P} = \left(\frac{\chi_e}{1 + \chi_e}\right) \mathbf{D} = \left(\frac{\chi_e}{1 + \chi_e}\right) \left(\frac{q}{4\pi r^2}\right) \hat{\mathbf{r}}$$
(14)

• The surface and volume bound charge densities are:

$$\sigma_b = \mathbf{P} \cdot \hat{\mathbf{r}} = \left(\frac{\chi_e}{1 + \chi_e}\right) D \tag{15}$$

$$\rho_b = -\nabla \cdot \mathbf{P} = -q \left( \frac{\chi_e}{1 + \chi_e} \right) \delta^3(\mathbf{r}) \tag{16}$$

• The total surface bound charge is  $Q = q\left(\frac{\chi_e}{1+\chi_e}\right)$ , but the negative compensating charge is at the origin:

$$Q = \int \rho_b d\tau = -q \left( \frac{\chi_e}{1 + \chi_e} \right) \tag{17}$$