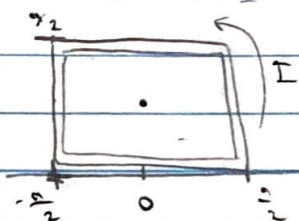


MW#S 4 S.4 S.7 S.11 S.12 S.16 S.19 S.21 S.22 S.23 S.27

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- 5.4) Magnetic field in some region is  $\vec{B} = k z \hat{x}$  where  $k$  is constant. Find the force on a square loop (side  $a$ ), lying in the  $yz$  plane centered @ origin, carries current  $I$  counterclockwise from  $x$  axis.



$$\begin{aligned} F_{\text{mag}} &= I(\vec{a} \times \vec{B}) \\ &= I(a \times B) \\ &= I a B \\ &= I a k z \hat{z} \end{aligned}$$

- 5.7) For a configuration of charges and currents confined within volume  $V$  show that  $\int_V \vec{J} d\tau = d\vec{p}/dt$  where  $\vec{p}$  is total dipole moment

$$\int_V \nabla \cdot (x \vec{J}) d\tau = \int_V (-x \frac{\partial J_x}{\partial x} + J_x) d\tau$$

$$\nabla \cdot (x \vec{J}) = x(\nabla \cdot \vec{J}) + \vec{J} \cdot (\nabla x) = (-x \frac{\partial J_x}{\partial x} + J_x)$$

$$\int_V \nabla \cdot (x \vec{J}) d\tau = \oint_S x \vec{J} \cdot d\vec{a} = 0 \quad \text{Gauss's law}$$

$$\int_V (-x \frac{\partial J_x}{\partial x} + J_x) d\tau = \oint_S x \vec{J} \cdot d\vec{a} = 0$$

$$\int_V (-x \frac{\partial J_x}{\partial x} + J_x) d\tau = 0$$

$$-\int_V x \frac{\partial J_x}{\partial x} d\tau + \int_V J_x d\tau = 0$$

$$\int_V J_x d\tau = \int_V x \frac{\partial J_x}{\partial x} d\tau$$

same process with  $\int_V \nabla \cdot (y \vec{J}) d\tau$  &  $\int_V \nabla \cdot (z \vec{J}) d\tau$

$$\int_V J_y d\tau = \int_V y \frac{\partial J_y}{\partial y} d\tau \quad \int_V J_z d\tau = \int_V z \frac{\partial J_z}{\partial z} d\tau$$

$$\int_V \vec{J} d\tau = \int_V \vec{r} \frac{\partial J}{\partial r} d\tau = \frac{d}{dt} \int_V \vec{r} \rho d\tau = \frac{d\vec{p}}{dt}$$

- 5.11) Magnetic field at point  $P$  on axis of solenoid w/  $n$  turns per unit length around cylinder tube w/ rad =  $a$  and current  $I$ . Field in infinite solenoid?



$$B(z) = \frac{\mu_0 I}{4a} \left( \frac{\cos \theta_1}{\sin^2 \theta_1} \right) 2\pi R = \frac{\mu_0 I}{2} \frac{R^2}{(R^2 + z^2)^{3/2}}$$

$$B(z) = \frac{\mu_0 I}{2} \frac{a^2}{(a^2 + z^2)^{3/2}} \quad I \rightarrow I n \quad z = a \cot(\theta)$$

$$\vec{B} = \int \frac{\mu_0 I n}{2} \frac{a^2}{(a^2 + z^2)^{3/2}} \left( \frac{-a}{\sin^2 \theta} \right) d\theta$$

$$dz = \frac{-a}{\sin^2 \theta} d\theta$$

$$\vec{B} = \frac{\mu_0 I n}{2} \int \frac{a^3 \sin^3 \theta}{a^3 \sin^2 \theta} d\theta$$

$$\vec{B} = \frac{\mu_0 I n}{2} \int_0^{\theta_2} \sin \theta d\theta$$

$$\vec{B} = \frac{\mu_0 I n}{2} [-(\cos \theta_2 - \cos \theta_1)]$$

$$\vec{B} = \frac{\mu_0 I n}{2} (\cos \theta_2 - \cos \theta_1)$$

for infinite solenoid  $\theta_2 = 0$ ,  $\theta_1 = 180^\circ \rightarrow \vec{B} = \frac{\mu_0 I n}{2} (1 - (-1))$

$$\vec{B} = \mu_0 I n \quad \text{for infinite solenoid}$$

S.12) Calculate magnetic field at center of uniformly charged spherical shell of radius  $R$  and total charge  $Q$  spinning at constant velocity  $\omega$

$$B(z) = \frac{\mu_0 I}{2} \frac{R^2}{(R^2 + z^2)^{3/2}}$$



$$R \rightarrow R \sin \theta$$

$$z \rightarrow R \cos \theta$$

$$I \rightarrow dI$$

$$\sigma = \frac{Q}{4\pi R^2}$$

$$v = \omega R = \omega R \sin \theta$$

$$K = \sigma v = \frac{Q}{4\pi R} \omega R \sin \theta = \frac{Q}{4\pi R} \omega \sin \theta$$

$$dI = K R d\theta = \frac{Q}{4\pi R} \omega \sin \theta R d\theta$$

$$dI = \frac{Q\omega}{4\pi} \sin \theta d\theta$$

$$\vec{B} = \frac{\mu_0}{2} \int dI \frac{(R \sin \theta)^2}{(R^2 \sin^2 \theta + R^2 \cos^2 \theta)^{3/2}}$$

$$= \frac{\mu_0}{2} \int \frac{Q\omega}{4\pi} \sin \theta d\theta \frac{R^2 \sin^2 \theta}{(R^2 \sin^2 \theta + R^2 \cos^2 \theta)^{3/2}}$$

$$= \frac{\mu_0}{2} \int \frac{Q\omega}{4\pi} \frac{R^2 \sin^3 \theta}{R^3} d\theta$$

$$= \frac{\mu_0 Q \omega}{8\pi R} \int \sin^3 \theta d\theta$$

$$= \frac{\mu_0 Q \omega}{8\pi R} \left[ -(\cos \theta - \frac{1}{3} \cos^3 \theta) \right]$$

$$= \frac{\mu_0 Q \omega}{8\pi R} \left[ \left( -1 + \frac{1}{3} \right) - \left( 1 - \frac{1}{3} \right) \right]$$

$$= \frac{\mu_0 Q \omega}{8\pi R} \left( -2 + \frac{1}{3} + \frac{1}{3} \right) = \frac{\mu_0 Q \omega}{8\pi R} \left( -\frac{4}{3} \right)$$

$$\vec{B} = \frac{\mu_0 Q \omega}{6\pi R} \hat{z}$$

S.16) Two long coaxial solenoids each of current  $I$  but in opposite directions. Inner solenoid (rad  $a$ ) has  $n_1$  turns per unit length, outer (rad  $b$ ) has  $n_2$ . Find  $\vec{B}$  in each of three regions (i) inside inner (ii) between (iii) outside



$$(iii) \text{ outside: } \vec{B} = 0$$

$$(i) \text{ inside: } \vec{B} = \mu_0 I (n_2 - n_1) \hat{z}$$

$$(ii) \text{ between: } \vec{B} = \mu_0 n_1 I \hat{z}$$

outside field points right, inside points left ( $-\hat{z}$ )

$$S.19) I = \int \vec{J} \cdot d\vec{a} \quad \text{What surface?}$$

$$\int \vec{J} \cdot d\vec{a} = \oint \nabla \cdot \vec{J} \quad \text{Gauss}$$

$\nabla \cdot \vec{J} = 0$  so the surface does not matter



5.21) Is Ampere's law consistent with the general rule Eq. 1.46 that divergence of curl is always zero? Show that Ampere's law cannot be valid outside magnetostatics. Is there any "trick" in the other 3 Maxwell equations?

$$1.46: \nabla \cdot (\nabla \times \vec{v}) = 0$$

$$\text{Ampere's Law: } \nabla \times \vec{B} = \mu_0 \vec{J}$$

$$\text{continuity: } \nabla \cdot \vec{J} = -\frac{\partial \rho}{\partial t}$$

$$\nabla \cdot (\nabla \times \vec{B}) = \mu_0 (\nabla \cdot \vec{J})$$

$$\nabla \cdot (\nabla \times \vec{B}) = -\mu_0 \frac{\partial \rho}{\partial t}$$

It is not consistent with divergence of curl always being 0, only true when  $\rho$  is constant which only happens in magnetostatics.

There are no other like tricks in the other Maxwell equations

Ex. 5.12) Find the vector potential of an infinite solenoid with  $n$  turns per unit length, radius  $R$ , and current  $I$ .

$$\vec{A} = \frac{\mu_0}{4\pi} \int \frac{\vec{J}(\vec{r}')}{r} dV'$$

$s$  is radius  $s$  inside solenoid

$$\oint \vec{A} \cdot d\vec{l} = \int \vec{B} \cdot d\vec{a}$$

$$\vec{A}(2\pi s) = B\pi s^2$$

$$2\pi s \vec{A} = \mu_0 n I \pi s^2 \quad \vec{B} = \mu_0 n I \text{ inside solenoid}$$

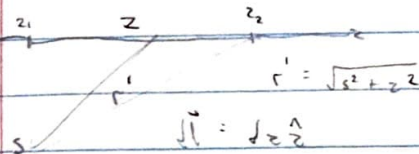
$$\boxed{\vec{A} = \frac{\mu_0 n I s}{2} \text{ for } s \leq R}$$

$$2\pi s \vec{A} = \mu_0 n I (\pi R^2)$$

$$\boxed{\vec{A} = \frac{\mu_0 n I R^2}{2s} \text{ for } s \geq R}$$

5.22) Magnetic vector potential of finite segment of straight wire of current  $I$ .

$$\vec{A} = \frac{\mu_0 I}{4\pi} \int \frac{1}{r} d\vec{l} \quad \vec{A} = \frac{\mu_0 I}{4\pi} \int \frac{\vec{r}}{r^2} d\vec{l}$$



$$\vec{A} = \frac{\mu_0 I}{4\pi} \int_{z_1}^{z_2} \frac{1}{\sqrt{s^2 + z^2}} dz \hat{z}$$

$$\vec{A} = \frac{\mu_0 I}{4\pi} \ln \left( \frac{z + \sqrt{s^2 + z^2}}{z_1 + \sqrt{s^2 + z_1^2}} \right) \hat{z}$$

$$\vec{A} = \frac{\mu_0 I}{4\pi} \ln \left[ \frac{z_2 + \sqrt{s^2 + z_2^2}}{z_1 + \sqrt{s^2 + z_1^2}} \right] \hat{z}$$

$$\vec{B} = \nabla \times \vec{A} = -\frac{\partial A_z}{\partial s} \hat{\phi}$$

$$= -\frac{\partial}{\partial s} \left[ \frac{\mu_0 I}{4\pi} \ln \left[ \frac{z_2 + \sqrt{s^2 + z_2^2}}{z_1 + \sqrt{s^2 + z_1^2}} \right] \right] \hat{\phi}$$

$$= -\frac{\mu_0 I}{4\pi} \left[ \frac{s}{\sqrt{s^2 + z_2^2} (\sqrt{s^2 + z_2^2} + z_2)} - \frac{s}{\sqrt{s^2 + z_1^2} (\sqrt{s^2 + z_1^2} + z_1)} \right] \hat{\phi}$$

where  $\frac{s}{\sqrt{s^2 + z^2} (\sqrt{s^2 + z^2} + z)}$   
 $= \sin \theta_z$

$$\text{eqn 5.22} \quad \vec{B} = \frac{\mu_0 I}{4\pi s} [\sin \theta_2 - \sin \theta_1] \hat{\phi}$$

S.22) Find vector potential above & below the plane surface current

$$\vec{K} = K \hat{x} \text{ over } xy \text{ plane}$$

$$\vec{B} = \pm \frac{\mu_0 K}{2} \hat{y} \quad \vec{A} \text{ only depends on } z \text{ so } \vec{A} = A(z) \hat{x}$$

$$\vec{B} = \nabla \times \vec{A} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ A(z) & 0 & 0 \end{vmatrix}$$

$$= \frac{\mu_0 K}{2} \hat{y} = -\frac{\partial}{\partial z} A \hat{y}$$

$$\frac{\partial}{\partial z} A = -\frac{\mu_0 K}{2}$$

$$A(z) \hat{x} = -\frac{\mu_0 K}{2} z \hat{x}$$

$$A = -\frac{\mu_0 K}{2} |z| \hat{x}$$