Electromagnetc Theory: PHYS330

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Summary

Week 3 Summary

- 1. Laplace's Equation
 - One-dimension
 - Two-dimensions, three dimensions, uniqueness, boundaries
- 2. Separation of Variables: Boundary-value problems
 - · Cartesian coordinates
 - Spherical coordinates
- 3. Multipole Expansions
 - Far-fields
 - Monopole and dipole terms
 - · Electric Field of a Dipole

Laplace's Equation in one dimension:

$$\frac{d^2V}{dx^2} = 0 (1)$$

What is the solution?

$$V(x) = mx + b (2)$$

What is the magnitude of the E-field?

- · A: V(x)
- B: x
- C: b
- D: m

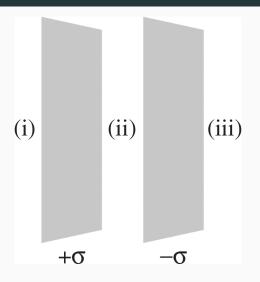


Figure 1: The setup of a parallel plate capacitor.

Suppose the negative side of the parallel plate capacitor is grounded, and the positive side is at a potential V_0 . Let the separation between the plates be x_0 . Further, let the positive plate occupy the yz plane, passing through the origin. Find the E-field magnitude and direction by solving Laplace's equation.

Show that the potential of a point charge at the origin satisfies Laplace's Equation for $r \neq 0$. Use the form of the Laplacian in spherical coordinates.

Let V(x) = mx + b. If $V(-a) = V_0$, and $V(a) = -V_0$, what are valid expressions for m and b?

- A: b = 0, and $m = -2V_0$
- B: b = a, and $m = V_0/a$
- C: b = 0, and $m = -V_0/a$
- D: $b = V_0$, and $m = -V_0/a$

Let V(x) = mx + b. If $V(-a) = V_0$, and $V(a) = -V_0$, what is the electric field?

- A: $\frac{V_0}{a}\hat{x}$
- B: $-\frac{V_0}{a}\hat{x}$
- C: $V_0\hat{x}$
- D: $-V_0\hat{x}$

Suppose a potential function $V(x,y) \propto (A \exp(-kx) + B \exp(kx))$. Which of the following is true, if $V \to 0$ as $x \to \infty$?

- A: A is 0
- B: B is 0
- · C: A and B are 0
- D: Neither A nor B is 0

Suppose a potential function $V(x,y) \propto (A\sin(kx) + B\cos(kx))$. Which of the following is true, if V = 0 as x = 0, and V = 0 as x = a?

- A: B is 0, and $k = n\pi$
- B: A is 0, and $k = n\pi/(2a)$
- · C: A and B are 0
- D: B is 0, and $k = n\pi/a$

Hyperbolic trigonometric functions:

•
$$\sinh(x) = \frac{1}{2} (e^x - e^{-x})$$

$$\cdot \cosh(x) = \frac{1}{2} \left(e^x + e^{-x} \right)$$

•
$$tanh(x) = \frac{\sinh(x)}{\cosh(x)}$$

Which of the following is zero?

- A: sinh(0)
 - B: cosh(0)
 - · C: tanh(0)
 - · D: None

Which of the following is one?

- A: sinh(0)
- B: cosh(0)
- C: tanh(0)
- · D: None

Hyperbolic trigonometric functions are solutions to which equation?

- A: $\frac{df}{dx} = k$
- B: $\frac{d^2f}{dx^2} = kx$
- C: $\frac{d^2f}{dx^2} = k^2f$
- D: $\frac{d^2f}{dx^2} = 0$

Fourier's Trick: Imagine a vector with *n* components:

$$\vec{\mathbf{v}} = \sum_{i=1}^{n} c_i \hat{\mathbf{x}}_i \tag{3}$$

In words, how do you solve for some c_m ?

- A: Divide by \hat{x}_i
- B: Take the dot product of both sides with \hat{x}_m
- \cdot C: Take the dot product \vec{v} and \vec{u} , and the sum the series
- D: Integrate both sides with respect to x

Fourier's Trick: Imagine a vector with *n* components:

$$\vec{V} = \sum_{i=1}^{n} c_i \hat{x}_i \tag{4}$$

In words, how do you solve for some c_m ? Note that:

$$\vec{\mathbf{v}} \cdot \hat{\mathbf{x}}_m = \sum_{i=1}^n c_n \hat{\mathbf{x}}_i \cdot \hat{\mathbf{x}}_m = c_m \tag{5}$$

Why? Because

$$\hat{\mathbf{x}}_i \cdot \hat{\mathbf{x}}_j = 0 \tag{6}$$
$$\hat{\mathbf{x}}_i \cdot \hat{\mathbf{x}}_i = 1 \tag{7}$$

$$\hat{\mathbf{x}}_i \cdot \hat{\mathbf{x}}_i = 1 \tag{7}$$

Fourier's Trick: Imagine a known function that happens to be equal to a sum:

$$f(x) = \sum_{i=1}^{\infty} c_i g_i(x)$$
 (8)

In words, how do you solve for some c_m ?

- A: Multiply both sides by $g_m(x)$
- B: Multiply both sides by $g_m(x)$ and integrate both sides with respect to x
- \cdot C: Sum the infinite series and solve for c_m with algebra
- D: Integrate both sides with respect to x

If it's true that a function can be written as an infinite series of functions with coefficients:

$$f(x) = \sum_{i=1}^{\infty} c_i g_i(x) \tag{9}$$

Then the functions $g_n(x)$ are said to be **complete**, or a complete basis (just like vectors are a sum of basis vectors. Examples of complete sets of functions:

- sines and cosines (Fourier series)
- exponentials with the right rates multiplying x
- Hyperbolic trigonometric functions (follows from exponentials)
- · Taylor series (polynomials with derivatives as coefficients)

The functions $f_n(x)$ are said to be orthogonal for $x \in [a, b]$ if

$$\int_{a}^{b} f_{n}(y) f_{m}(y) dy = \delta_{n,m}$$
 (10)

One example:

$$I_{n,m} = \int_{-L}^{L} \frac{\sin(n\pi x/L)}{\sqrt{L}} \frac{\sin(m\pi x/L)}{\sqrt{L}} dx$$
 (11)

What is the result of this integral? How would you approach solving this?

The Fourier series representation of a function f(x) is written:

$$S(x) = \frac{A_0}{2} + \sum_{i=1}^{\infty} (A_n \cos(nx) + B_n \sin(nx))$$
 (12)

with

$$A_n = \frac{1}{\pi} \int_0^{2\pi} f(x) \cos(nx) dx$$
 (13)

$$B_n = \frac{1}{\pi} \int_0^{2\pi} f(x) \sin(nx) dx$$
 (14)

Let's obtain the Fourier series coefficients A_n and B_n for a square-wave signal:

$$f(x) = 1, \quad 0 \le x \le \pi, \quad 0, \pi < x \le 2\pi$$
 (15)

(Observe on board). The result: $A_0 = 1.0$, all other $A_n = 0$, odd B_n values follow $2/(n\pi)$, even $B_n = 0$ as well.

Laplaces' Equation:

$$\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0 \tag{16}$$

Assume the solution follows

$$V(x, y, z) = X(x)Y(y)Z(z)$$
(17)

The Laplace equation then breaks into three separate ordinary differential equations. Application of boundary conditions to solve them (Asynchronous video content on Moodle).

Laplaces' Equation in spherical coordinates:

$$\frac{1}{r^2}\frac{\partial}{\partial r}\left(r^2\frac{\partial V}{\partial r}\right) + \frac{1}{r^2\sin\theta}\frac{\partial}{\partial\theta}\left(\sin\theta\frac{\partial V}{\partial\theta}\right) + \frac{1}{r^2\sin\theta}\left(\frac{\partial^2 V}{\partial\phi^2}\right) = 0 \qquad (18)$$

Assuming azimuthal symmetry means $V(r, \theta, \phi) = V(r, \theta)$ and $\partial V/\partial \phi = 0$. Thus, Eq. 18 reduces and admits general solutions:

$$\frac{1}{r^2}\frac{\partial}{\partial r}\left(r^2\frac{\partial V}{\partial r}\right) + \frac{1}{r^2\sin\theta}\frac{\partial}{\partial\theta}\left(\sin\theta\frac{\partial V}{\partial\theta}\right) = 0 \tag{19}$$

Let the general solutions be separable:

$$V(r,\theta) = R(r)\Theta(\theta) \tag{20}$$

The radial equation is

$$\frac{1}{R(r)}\frac{d}{dr}\left(r^2\frac{dR}{dr}\right) = l(l+1) \tag{21}$$

Exercise: show that the solution is

$$R(r) = Ar^{l} + Br^{-(l+1)}$$
 (22)

(The derivative operator distributes over addition, so the two solutions can be checked separately, or together).

What are the units of R(r)? What are the units of A and B?

The polar equation is

$$\frac{1}{\Theta(\theta)} \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right) = -l(l+1) \sin \theta \tag{23}$$

The solutions are complete, and orthogonal, and known as Legendre polynomials:

$$\Theta(\theta) = P_l(\cos \theta) \tag{24}$$

Defined by the Rodrigues formula:

$$P_l(x) = \frac{1}{2^l l!} \left(\frac{d}{dx}\right)^l (x^2 - 1)^l$$
 (25)

Exercise: show that

$$P_3(x) = (5x^3 - 3x)/2 (26)$$

What is the result of the following integrals?

$$I_1 = \int_{-1}^{1} P_1(x) P_2(x) dx \tag{27}$$

$$I_2 = \int_{-1}^{1} P_2(x) P_2(x) dx \tag{28}$$

The general solution is a sum of individual solutions:

$$V(r,\theta) = \sum_{l=0}^{\infty} \left(Ar^{l} + B/r^{l+1} \right) P_{l}(\cos \theta)$$
 (29)

The coefficients may be found via Fourier's Trick.

Example 3.9: Professor on Board

A specified charge density $\sigma_0(\theta)$ is glued over the surface of a spherical shell of radius R. Find the resulting potential inside and outside the sphere.

- 1. Inside the sphere, B=0 to avoid a singularity at the origin (center of sphere).
- 2. Outside the sphere, A = 0 to ensure $V \to 0$ as $r \to \infty$.
- 3. General boundary conditions at r = R: potential is continuous $(-\int \vec{E} \cdot d\vec{l} = 0)$.
- 4. Coefficients of same order *l* have a relationship.
- 5. E-field has a discontinuity at the boundary.
- 6. Fourier's trick to get the coefficients, after specifying $\sigma_0(\theta)$.

Imagine a physical dipole with q at $\hat{r}' = d/2\hat{z}$ and -q at $\hat{r}' = -d/2\hat{z}$. Show that (professor example)

$$V(r,\theta) = \frac{kqd}{r^2} P_1(\cos\theta)$$
 (30)

- 1. Far-field on script-r's
- 2. Subtract
- 3. Simplify
- 4. Note that $P_1(x) = x$.

Can't you break *any* charge distribution into a collection of monopoles, dipoles, quadrupoles, ... ? We can show in general that:

$$\boxed{\frac{1}{\imath} = \frac{1}{r} \sum_{n=0}^{\infty} \left(\frac{r'}{r}\right)^n P_n(\cos \theta)}$$
 (31)

Can't you break *any* charge distribution into a collection of monopoles, dipoles, quadrupoles, ... ?

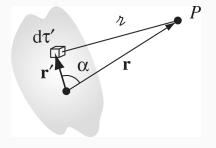


Figure 2: The general scheme for the multipole expansion.

Find the Law of Cosines from the definition of the separation vector:

Then we let

$$\epsilon = \left(\frac{r'}{r}\right)\left(\frac{r'}{r} - 2\cos\alpha\right) \tag{32}$$

Find the Taylor series of $f(\epsilon) = (1 + \epsilon)^{-1/2}$:

Remember that

$$\epsilon = \left(\frac{r'}{r}\right)\left(\frac{r'}{r} - 2\cos\alpha\right)$$

(33)

- 1. After computing the Taylor series, substitute $\epsilon = \left(\frac{r'}{r}\right)\left(\frac{r'}{r} 2\cos\alpha\right)$
- 2. Collecting like powers of $\left(\frac{r'}{r}\right)$ together will lead to

$$\boxed{\frac{1}{\imath} = \frac{1}{r} \sum_{n=0}^{\infty} \left(\frac{r'}{r}\right)^n P_n(\cos \theta)}$$
 (34)

Recall that the potential for any charge distribution is

$$V(\vec{r}) = \int \frac{k\rho(\vec{r'})}{\imath} d\tau' \tag{35}$$

Substitute the expansion for $1/\imath$, and reverse the order of summation and integration:

$$V(\vec{r}) = k \sum_{n=0}^{\infty} \frac{1}{r^{n+1}} \int (r')^n P_n(\cos \alpha) \rho(\vec{r'}) d\tau'$$
 (36)

The Monopole and Dipole Terms

The Monopole and Dipole Terms

Using the $\rho(\vec{r'})$ of a dipole oriented along the z-axis, reproduce Eq. 30 using the multipole expansion.

Hint: obtain the first few terms, but which ones vanish and why?

Conclusion

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