

# Electromagnetic Theory: PHYS330

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## Summary

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## Week 4 Summary

1. Atoms, polarizations, and dipole moments
2.  $\vec{P}$ , dipole per unit volume, and bound charges
3.  $\vec{D}$ , the electric displacement
4. Linear dielectrics

## **Dipole Moment, and The Electric Field of a Dipole**

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# Dipole Moment, and The Electric Field of a Dipole

From the multipole expansion, the monopole and dipole terms are (asynchronous video content):

$$V(r, \theta)_{n=0} = \frac{1}{4\pi\epsilon_0} \frac{\int_{all} \rho(\vec{r}') d\tau'}{r} \quad (1)$$

$$V(r, \theta)_{n=1} = \frac{1}{4\pi\epsilon_0} \frac{\hat{r} \cdot \int_{all} \vec{r}' \rho(\vec{r}') d\tau'}{r^2} \quad (2)$$

$$\vec{p} = \hat{r} \cdot \int_{all} \vec{r}' \rho(\vec{r}') d\tau' \quad (3)$$

# Dipole Moment, and The Electric Field of a Dipole

Valid for any charge distribution, useful for getting far-field E-fields:

$$\vec{p} = \hat{r} \cdot \int_{all} \vec{r} \rho(\vec{r}) d\tau' \quad (4)$$

1. Discrete charge distribution example: asynch. video content
2. Continuous: “similar” to calculating the moment of inertia of a mass distribution (only it’s linear vector, not  $r^2$ ).

## Dipole Moment, and The Electric Field of a Dipole

For a dipole, with no other charges ( $\vec{p} = q\vec{d}$ ):

$$V_{n=1}(r, \theta) = \frac{\hat{r} \cdot \vec{p}}{4\pi\epsilon_0 r^2} \quad (5)$$

1. Using Eq. 5, calculate the electric field of a dipole in spherical coordinates.
2. Are there any regions for which the E-field is zero?
3. What happens when you set  $\partial E / \partial \theta = 0$ ?
4. Recall that the energy density of a field is  $u_E = \frac{1}{2}\epsilon_0 \vec{E} \cdot \vec{E}$ . For a given radius  $R$ , is the energy density higher at  $\theta = 0$  or  $\theta = \pi/2$ ?

*Breakout rooms.*

# Atoms, polarizations, and dipole moments

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## Atoms, polarizations, and dipole moments

Suppose an external field  $\vec{E}$  induces a dipole moment  $\vec{p}$  in an atomic charge distribution:

$$\boxed{\vec{p} = \alpha \vec{E}} \quad (6)$$

This statement is empirical, but it's true for “ordinary” field strengths: field isn't strong enough to ionize the atom.

## Atoms, polarizations, and dipole moments

**What is the electric field a distance  $d$  from the center of a uniformly charged sphere of radius  $a$ ? [Hint: use Gauss' law, and assume  $\rho$  is constant in spherical coordinates].**

# Atoms, polarizations, and dipole moments

Result:

$$E = \frac{1}{4\pi\epsilon_0} \frac{qd}{a^3} \quad (7)$$

But then assume that  $p = qd$ , so

$$\alpha = 4\pi\epsilon_0 a^3 \quad (8)$$

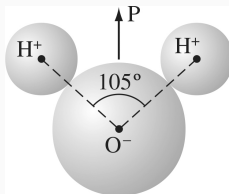
H	He	Li	Be	C	Ne	Na	Ar	K	Cs
0.667	0.205	24.3	5.60	1.67	0.396	24.1	1.64	43.4	59.4

**Figure 1:** Do you understand the units of this table? The numbers are quoted as  $\alpha/4\pi\epsilon_0$ , in units of  $10^{-30} \text{ m}^3$ . What would they be in nanometers cubed?

The trouble is that we cannot easily measure the volume of an atom (realm of quantum mechanics).

# Atoms, polarizations, and dipole moments

Molecules can also have a *permanent* dipole moment: polar molecules.



**Figure 2:** How would you calculate the dipole moment here?

Show that the torque on such a molecule in an external field is  $\vec{\tau} = \vec{p} \times \vec{E}$  (Professor example).

## Atoms, polarizations, and dipole moments

If you have (approximately) aligned polar molecules with an external field  $\vec{E} = E_0 \hat{x}$ , and then *reverse* the direction of the field, in what direction is the torque? Assume the dipole moments are in the xy-plane.

- A:  $-\hat{z}$
- B:  $\hat{y}$
- C:  $\hat{z}$
- D: The torque is zero

# Atoms, polarizations, and dipole moments

In summary, there are two reasons there could be dipole moments within a material:

1. The atoms are *stretched* and you get an  $\alpha = 4\pi\epsilon_0 a^3$
2. The atoms or molecules are *rotated* and you get a dipole moment  $\vec{p}$  per atom/molecule.

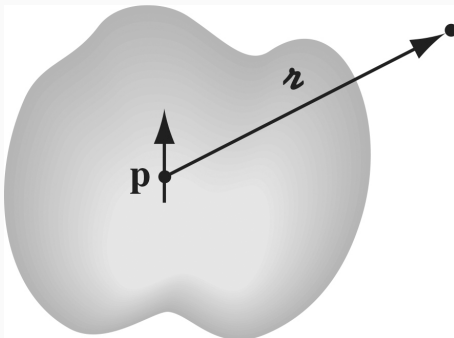
Macroscopically, it is easier to demonstrate the polar molecule effect: [https://youtu.be/riMrg\\_k0\\_\\_w](https://youtu.be/riMrg_k0__w)

$\vec{P}$ , dipole per unit volume, and  
bound charges

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## $\vec{P}$ , dipole per unit volume, and bound charges

We need to understand the field of a polarized material. Suppose we introduce the *dipole moment per unit volume*,  $\vec{P}$ .



**Figure 3:** The definition of the dipole moment per unit volume, and geometry.



## $\vec{P}$ , dipole per unit volume, and bound charges

This implies that

$$\vec{p} = \vec{P} d\tau' \quad (9)$$

$$V(r) = \frac{1}{4\pi\epsilon_0} \frac{\vec{p} \cdot \hat{r}}{r^2} \quad (10)$$

$$V(r) = \frac{1}{4\pi\epsilon_0} \int_V \frac{\vec{P}(\vec{r}') \cdot \hat{r}}{r^2} d\tau' \quad (11)$$

$$V(r) = \frac{1}{4\pi\epsilon_0} \int_V \vec{P}(\vec{r}') \cdot \nabla \left( \frac{1}{r} \right) d\tau' \quad (12)$$

$$\nabla \cdot (f\vec{A}) = f(\nabla \cdot \vec{A}) + \vec{A} \cdot \nabla(f) \quad (13)$$

$$V(r) = \frac{1}{4\pi\epsilon_0} \left\{ \int_V \nabla' \cdot \left( \frac{\vec{P}}{r} \right) d\tau' - \int_V \frac{1}{r} (\nabla' \cdot \vec{P}) d\tau' \right\} \quad (14)$$

## $\vec{P}$ , dipole per unit volume, and bound charges

$$V(r) = \frac{1}{4\pi\epsilon_0} \left\{ \int_V \nabla' \cdot \left( \frac{\vec{P}}{r} \right) d\tau' - \int_V \frac{1}{r} (\nabla' \cdot \vec{P}) d\tau' \right\} \quad (15)$$

$$V(r) = \frac{1}{4\pi\epsilon_0} \left\{ \oint_S \left( \frac{\vec{P}}{r} \right) \cdot d\vec{a}' - \int_V \frac{1}{r} (\nabla' \cdot \vec{P}) d\tau' \right\} \quad (16)$$

$$d\vec{a}' = da' \hat{n} \quad (17)$$

$$\sigma_b = \vec{P} \cdot \hat{n} \quad (18)$$

$$\rho_b = -\nabla \cdot \vec{P} \quad (19)$$

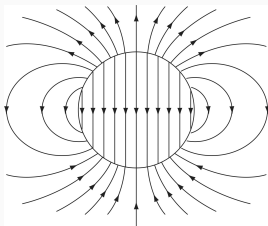
$$V(r) = \frac{1}{4\pi\epsilon_0} \left\{ \oint_S \frac{\sigma_b}{r} da' - \int_V \frac{\rho_b}{r} d\tau' \right\} \quad (20)$$

The appearance of *bound charge*.

## $\vec{P}$ , dipole per unit volume, and bound charges

Suppose we have a sphere with uniform polarization in the z-direction (and it is constant). The  $\rho_b$  is zero because

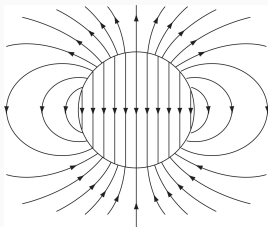
- A: There is no bound charge inside a sphere.
- B: The divergence of a constant is zero.
- C: By symmetry.
- D: Otherwise the integral over  $\rho_b$  would diverge.



**Figure 4:** The uniformly polarized sphere.  $\sigma_b = \vec{P} \cdot \hat{n} = P \cos \theta$ .

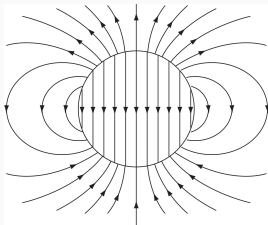
## $\vec{P}$ , dipole per unit volume, and bound charges

Think for a moment: in your own words, why do the field lines point in *opposite* directions just inside and just outside the surface of the sphere?

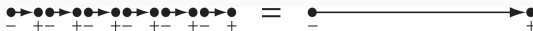


**Figure 5:** The uniformly polarized sphere.  $\sigma_b = \vec{P} \cdot \hat{n} = P \cos \theta$ .

## $\vec{P}$ , dipole per unit volume, and bound charges



**Figure 6:** The uniformly polarized sphere.  $\sigma_b = \vec{P} \cdot \hat{n} = P \cos \theta$ .



**Figure 7:** Whenever you think of bound charge density versus surface charge density, think of this picture.

## $\vec{P}$ , dipole per unit volume, and bound charges

**Conceptual question:** Given Eq. 20, what is the potential due a disk of surface bound charge  $\sigma_b$  at a point slightly above the surface? *[Hint: if it helps, think of  $\sigma_b = \vec{P}_0 \cdot \hat{z}$ , where  $P_0$  is a constant.]*

$\vec{D}$ , the electric displacement

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## Conclusion

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