

Rocket Engines

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Rocket Propulsion:

Thrust

Conservation of Momentum

Impulse & Momentum

Combustion & Exhaust Velocity

Specific Impulse

Rocket Engines

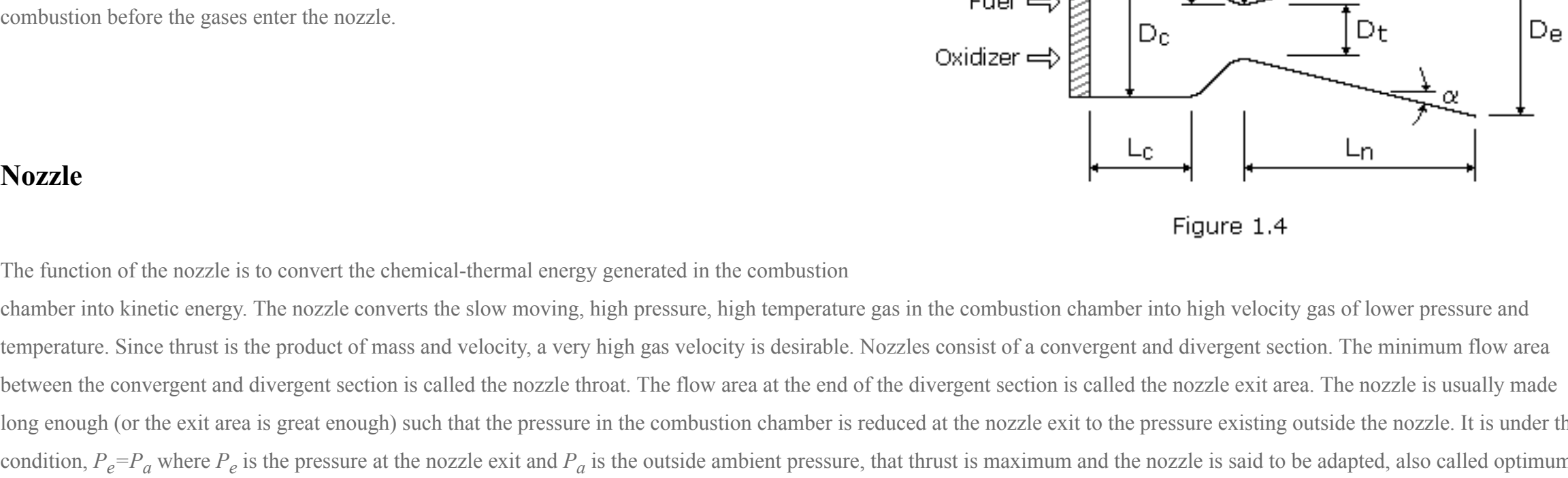
Power Cycles

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Staging



A typical rocket engine consists of the *nozzle*, the *combustion chamber*, and the *injector*, as shown in Figure 1.4. The *combustion chamber* is where the burning of propellants takes place in high pressure. The chamber must be strong enough to contain the high pressure generated by, and the high temperature resulting from, the combustion process. Because of the high temperature and heat transfer, the chamber and nozzle are usually cooled. The chamber must also be of sufficient length to ensure complete combustion before the gases enter the nozzle.

Nozzle

The function of the *nozzle* is to convert the chemical-thermal energy generated in the *combustion chamber* into kinetic energy. The *nozzle* converts the slow moving, high pressure, high temperature gas in the *combustion chamber* into high velocity gas of lower pressure and temperature. Since thrust is the product of mass and velocity, a very high gas velocity is desirable. Nozzles consist of a convergent and divergent section. The minimum flow area between the convergent and divergent section is called the *nozzle throat*. The flow area at the end of the divergent section is called the *nozzle exit area*. The *nozzle* is usually made long enough (or the exit area is great enough) such that the pressure in the *combustion chamber* is reduced at the *nozzle exit* to the pressure existing outside the *nozzle*. It is under this condition, $P_c > P_a$, where P_c is the pressure at the *nozzle exit* and P_a is the outside ambient pressure, that thrust is maximum and the *nozzle* is said to be adapted, also called optimum or correct expansion. When P_c is greater than P_a , the *nozzle* is under-extended. When the opposite is true, it is over-extended.

We see therefore, a *nozzle* is designed for the altitude at which it has to operate. At the Earth's surface, at the atmospheric pressure of sea level (0.1 MPa or 14.7 psi), the discharge of the exhaust gases is limited by the separation of the jet from the *nozzle wall*. In the cosmic vacuum, this physical limitation does not exist. Therefore, there have to be two different types of engines and nozzles, those which propel the first stage of the launch vehicle through the atmosphere, and those which propel subsequent stages or control the orientation of the spacecraft in the vacuum of space.

The *nozzle throat area*, A_t , can be found if the total propellant flow rate is known and the propellants and operating conditions have been selected. Assuming perfect gas law theory, we have

$$(1.26) \quad A_t = \frac{\dot{q}}{P_t} \sqrt{\frac{R^* T_t}{M k}}$$

where \dot{q} is the propellant mass flow rate, P_t is the gas pressure at the *nozzle throat*, T_t is the gas temperature at the *nozzle throat*, R^* is the universal gas constant, and k is the specific heat ratio. P_t and T_t are given by

$$(1.27) \quad P_t = P_c \left(1 + \frac{k-1}{2} \right)^{\frac{k+1}{k-1}}$$

$$(1.28) \quad T_t = \left(1 + \frac{k-1}{2} \right) T_c$$

where P_c is the *combustion chamber pressure* and T_c is the *combustion chamber flame temperature*.

The hot gases must be expanded in the diverging section of the *nozzle* to obtain maximum thrust. The pressure of these gases will decrease as energy is used to accelerate the gas. We must find that area of the *nozzle* where the gas pressure is equal to the outside atmospheric pressure. This area will then be the *nozzle exit area*.

Mach number N_a is the ratio of the gas velocity to the local speed of sound. The Mach number at the *nozzle exit* is given by the perfect gas expansion expression

$$(1.29) \quad N_a^2 = \left(\frac{2}{k-1} \right) \left[\left(\frac{P_c}{P_a} \right)^{\frac{k-1}{k}} - 1 \right]$$

where P_a is the pressure of the ambient atmosphere.

The *nozzle exit area*, A_e , corresponding to the exit Mach number is given by

$$(1.30) \quad A_e = \left(\frac{A_t}{N_a} \right) \left[1 + \frac{k-1}{2} \right] N_a^2 \left(\frac{k+1}{2(k-1)} \right)^{\frac{k+1}{k-1}}$$

The section ratio, or expansion ratio, is defined as the area of the exit A_e divided by the area of the throat A_t .

For launch vehicles (particularly first stages) where the ambient pressure varies during the burn period, trajectory computations are performed to determine the optimum exit pressure. However, an additional constraint is the maximum allowable diameter for the *nozzle exit cone*, which in some cases is the limiting constraint. This is especially true on stages other than the first, where the *nozzle diameter* may not be larger than the outer diameter of the stage below. For space engines, where the ambient pressure is zero, thrust always increases as *nozzle expansion ratio* increases. On these engines, the *nozzle expansion ratio* is generally increased until the additional weight of the longer *nozzle* costs more performance than the extra thrust it generates.

(For additional information, please see Supplement #1: [Optimizing Expansion for Maximum Thrust](#).)

Since the flow velocity of the gases in the converging section of the rocket *nozzle* is relatively low, any smooth and well-rounded convergent *nozzle* section will have very low energy losses. By contrast, the contour of the diverging *nozzle* section is very important to performance, because of the very high flow velocities involved. The selection of an optimum *nozzle shape* for a given expansion ratio is generally influenced by the following design considerations and goals: (1) uniform, parallel, axial gas flow at the *nozzle exit* for maximum momentum vector, (2) minimum separation and turbulence losses within the *nozzle*, (3) shortest possible *nozzle length* for minimum space envelope, weight, wall friction losses, and cooling requirements, and (4) ease of manufacturing.

Conical nozzle: In early rocket engine applications, the conical nozzle, which proved satisfactory in most respects, was used almost exclusively. A conical *nozzle* allows ease of manufacture and flexibility in converting an existing design to higher or lower expansion ratio without major redesign.

The configuration of a typical conical nozzle is shown in Figure 1.4. The *nozzle throat* section has the contour of a circular arc with radius R , ranging from 0.25 to 0.75 times the throat diameter, D_t . The half-angle of the *nozzle* convergent cone section, θ_c , can range from 20 to 45 degrees. The divergent cone half-angle, θ_e , varies from approximately 12 to 18 degrees. The conical *nozzle* with a 15-degree divergent half-angle has become almost a standard because it is a good compromise on the basis of weight, length, and performance.

Since certain performance losses occur in a conical nozzle as a result of the nonaxial component of the exhaust gas velocity, a correction factor, λ_c , is applied in the calculation of the exit-gas momentum. This factor (thrust efficiency) is the ratio between the exit-gas momentum of the conical *nozzle* and that of an ideal *nozzle* with uniform, parallel, axial gas-flow. The value of λ_c can be expressed by the following equation:

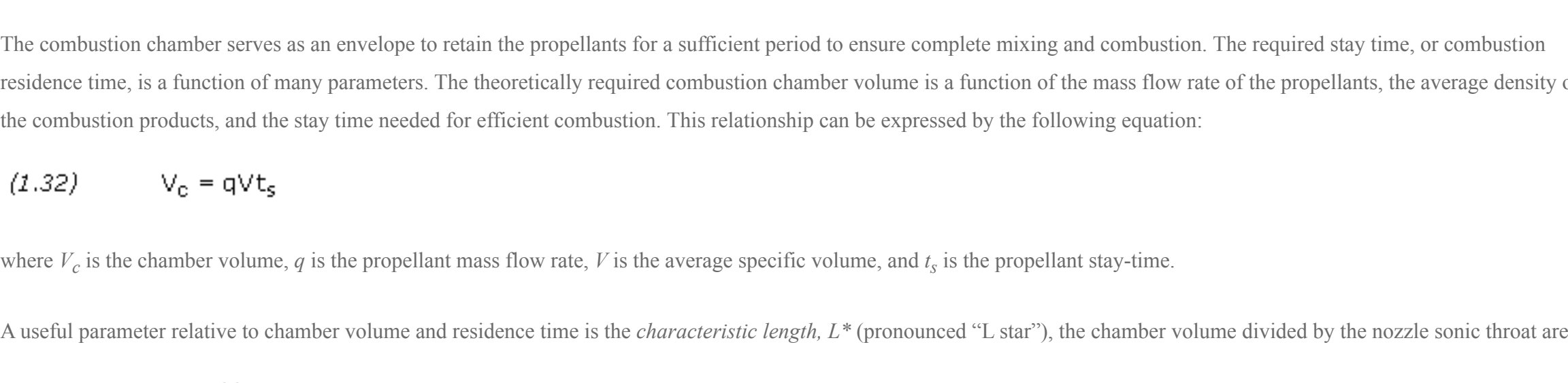
$$(1.31) \quad \lambda_c = \frac{1 + 0.05 \theta_e}{2}$$

Bell nozzle: To gain higher performance and shorter length, engineers developed the bell-shaped *nozzle*. It employs a fast-expansion (radial-flow) section in the initial divergent region, which leads to a uniform, axially directed flow at the *nozzle exit*. The wall contour is changed gradually enough to prevent oblique shocks.

An equivalent 15-degree half-angle conical nozzle is commonly used as a standard to specify bell nozzles. For instance, the length of an 80% bell *nozzle* (distance between throat and exit plane) is 80% of that of a 15-degree half-angle conical *nozzle* having the same throat area, radius below the throat, and area expansion ratio. Bell *nozzle* lengths beyond approximately 80% do not significantly contribute to performance, especially when weight penalties are considered. However, bell *nozzle* lengths up to 100% can be optimum for applications stressing very high performance.

One convenient way of designing a near optimum thrust bell *nozzle* contour uses the parabolic approximation procedures suggested by G.V.B. Rao. The design configuration of a parabolic approximation bell *nozzle* is shown in Figure 1.5. The *nozzle* contour immediately upstream of the throat T is a circular arc with a radius of $0.5 R_t$. The divergent section *nozzle* contour is made up of a circular entrance section with a radius of $0.82 R_t$ from the throat T to the point N and parabola from there to the exit E .

Design of a specific nozzle requires the following data: throat diameter D_t , axial length of the *nozzle* from throat to exit plane L_n (or the desired fractional length, L_f), based on a 15-degree conical nozzle), expansion ratio ϵ , initial wall angle of the parabola θ_p , and *nozzle exit* wall angle θ_e . The wall angles θ_p and θ_e are shown in Figure 1.6 for a function of the expansion ratio. Optimum *nozzle* contours can be approximated very accurately by selecting the proper inputs. Although no allowance is made for different propellant combinations, experience has shown only small effect of the specific heat ratio upon the contour.



Combustion Chamber

The *combustion chamber* serves as an envelope to retain the propellants for a sufficient period to ensure complete mixing and combustion. The required stay time, or *combustion residence time*, is a function of many parameters. The theoretically required *combustion chamber volume* is a function of the mass flow rate of the propellants, the average density of the *combustion products*, and the stay time needed for efficient combustion. This relationship can be expressed by the following equation:

$$(1.32) \quad V_c = \dot{q} V_t \tau_s$$

where V_c is the chamber volume, \dot{q} is the propellant mass flow rate, V is the average specific volume, and τ_s is the propellant stay-time.

A useful parameter relative to chamber volume and residence time is the *characteristic length*, L^* (pronounced "L, star"), the chamber volume divided by the *nozzle* sonic throat area:

$$(1.33) \quad L^* = \frac{V_c}{A_t}$$

The L^* concept is much easier to visualize than the more elusive "combustion residence time", expressed in small fractions of a second. Since the value of A_t is in nearly direct proportion to the product of \dot{q} and V , L^* is essentially a function of τ_s .

The customary method of establishing the L^* of a new thrust chamber design largely relies on past experience with similar propellants and engine size. Under a given set of operating conditions, such as type of propellant, mixture ratio, chamber pressure, injector design, and chamber geometry, the value of the minimum required L^* can only be evaluated by actual firings of experimental thrust chambers. Typical L^* values for various propellants are shown in the table below. With throat area and minimum required L^* established, the chamber volume can be calculated by equation (1.33).

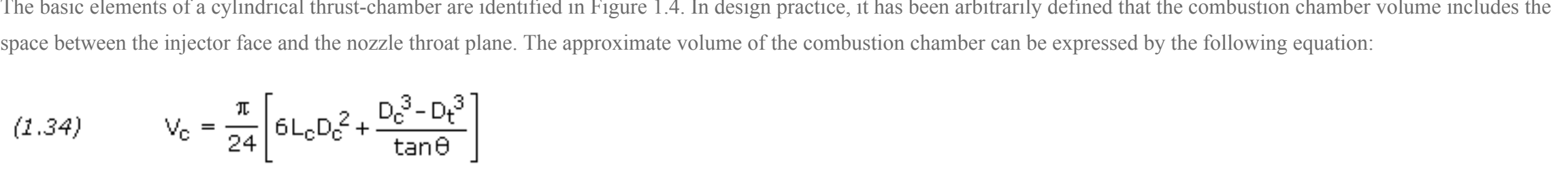
| Propellant Combination | L^* , cm |
|---|------------|
| Nitric acid/hydrazine-base fuel | 76-89 |
| Nitrogen tetroxide/hydrazine-base fuel | 76-89 |
| Hydrogen peroxide/RP-1 (including catalyst bed) | 152-178 |
| Liquid oxygen/RP-1 | 102-127 |
| Liquid oxygen/ammonia | 76-102 |
| Liquid oxygen/liquid hydrogen (GH ₂ injection) | 56-71 |
| Liquid oxygen/liquid hydrogen (LH ₂ injection) | 76-102 |
| Liquid fluorine/liquid hydrogen (GH ₂ injection) | 56-66 |
| Liquid fluorine/liquid hydrogen (LH ₂ injection) | 64-76 |
| Liquid fluorine/hydrazine | 61-71 |
| Chlorine trifluoride/hydrazine-base fuel | 51-89 |

Table 1: Chamber Characteristic Lengths, L^*

Three geometrical shapes have been used in *combustion chamber* design – spherical, near-spherical, and cylindrical – with the cylindrical chamber being employed most frequently in the United States. Compared to a cylindrical chamber of the same volume, a spherical or near-spherical chamber offers the advantage of less cooling surface and weight; however, the spherical chamber is more difficult to manufacture and has provided poorer performance in other respects.

The total *combustion process*, from injection of the reactants until completion of the chemical reactions and conversion of the products into hot gases, requires finite amounts of time and volume, as expressed by the characteristic length L^* . The value of this factor is significantly greater than the linear length between injector face and throat plane. The *contraction ratio* is defined as the major cross-sectional area of the chamber divided by the throat area. Typically, large engines are constructed with a low contraction ratio and a comparatively long length, and smaller chambers employ a large contraction ratio with a shorter length, while still providing sufficient L^* for adequate vaporization and combustion dwell-time.

As a good place to start, the process of sizing a new *combustion chamber* examines the dimensions of previously successful designs in the same size class and plotting such data in a rational manner. The throat size of a new engine can be generated with a fair degree of confidence, so it makes sense to plot the data from historical sources in relation to throat diameter. Figure 1.7 plots chamber length as a function of throat diameter (with approximating equation). It is important that the output of any modeling program not be slavishly applied, but be considered a logical starting point for specific engine sizing.



The basic elements of a cylindrical thrust-chamber are identified in Figure 1.4. In design practice, it has been arbitrarily defined that the *combustion chamber volume* includes the space between the injector face and the *nozzle* throat plane. The approximate volume of the *combustion chamber* can be expressed by the following equation:

$$(1.34) \quad V_c = \frac{\pi}{24} \left[6 L_c D_t^2 + \frac{D_t^3 - D_t^3}{\tan \theta} \right]$$

Rearranging equation (1.34) we get the following, which can be solved for the chamber diameter via iteration:

$$(1.35) \quad D_c = \sqrt{\frac{D_t^3 + \frac{24}{\tan \theta} L_c V_c}{D_c + 6 \tan \theta L_c}}$$

Injector

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The injector, as the name implies, injects the propellants into the combustion chamber and the right conditions to yield an efficient, stable combustion process. Placed at the forward, or upper, end of the combustor, the injector also performs the structural task of closing off the top of the *combustion chamber* against the high pressure and temperature it contains. The injector has been compared to the *carburetor* of an automobile engine, since it provides the fuel and oxidizer at the proper rates and in the correct proportions; this may be an appropriate comparison. However, the injector, heated directly over the high-pressure combustion, performs many other functions related to the combustion and cooling processes and is much more important to the function of the rocket engine than the carburetor is for an automobile engine.

No other component of a rocket engine has as great an impact upon engine performance as the injector. In various and different applications, well-designed injectors may have a fairly wide spread in combustion efficiency, and it is not uncommon for an injector with C^* efficiency as low as 92% to be considered acceptable. Small engines designed for special purposes, such as attitude control, may be optimized for response and light weight at the expense of combustion efficiency, and may be deemed very satisfactory even if efficiency falls below 90%. In general, however, recently well-designed injection systems have demonstrated C^* efficiencies so close to 100% of theoretical that the ability to measure this parameter is the limiting factor in its determination. High levels of combustion efficiency derive from uniform distribution of the desired mixture ratio and fine atomization of the liquid propellants. Local mixing within the injection-element spray pattern must take place at a virtually a microscopic level to ensure combustion efficiencies approaching 100%.

Combustion stability is also a very important requirement for a satisfactory injector design. Under certain conditions, shock and detonation waves are generated by local disturbances in the chamber, possibly caused by fluctuations in mixing or propellant flow. These may trigger pressure oscillations that are amplified and maintained by the combustion processes. Such high-amplitude waves – referred to as *combustion instability* – produce high levels of vibration and heat flux that can be very destructive. A major portion of the design and development effort therefore concerns stable combustion. High performance can become secondary if the injector is easily triggered into destructive instability, and many of the injector parameters that provide high performance appear to reduce the stability margins.

PROBLEM 1.7

A rocket engine uses the same propellant, mixture ratio, and combustion chamber pressure as that in problem 1.5. If the propellant flow rate is 500 kg/s, calculate the area of the exhaust *nozzle throat*.

SOLUTION:

Given: $P_c = 50 \times 0.101325 = 5.066 \text{ MPa}$
 $T_c = 3,470 \text{ °sup. °C}$ $\text{sup}^{\text{>}} \text{>} \text{K}$
 $M = 21.40$
 $k = 1.221$
 $\dot{q} = 500 \text{ kg/s}$

Equation (1.27),

$$P_t = P_c \left(1 + \frac{k-1}{2} \right)^{\frac{k+1}{k-1}}$$

$$P_t = 5.066 \times \left(1 + \frac{1.221-1}{2} \right)^{\frac{2.221}{1.221-1}}$$

$$P_t = 2,839 \text{ MPa} = 2,839 \times 10^6 \text{ N/m}^2$$

Equation (1.28),

$$T_t = T_c \left(1 + \frac{k-1}{2} \right)$$

$$T_t = 3,470 \left(1 + \frac{1.221-1}{2} \right)$$

$$T_t = 3,125 \text{ K}$$

Equation (1.26),

$$A_t = \left(\frac{\dot{q}}{P_t} \right) \sqrt{\frac{R^* T_t}{M \cdot k}}$$

$$A_t = \left(\frac{500}{2,839 \times 10^6} \right) \times \sqrt{\frac{8,314.5 \times 3,125}{21.40 \times 1.221}}$$

$$A_t = 0.1756 \text{ m}^2$$

PROBLEM 1.8

The rocket engine in problem 1.7 is optimized to operate at an elevation of 2000 meters. Calculate the area of the *nozzle exit* and the section ratio.

SOLUTION:

Given: $P_c = 5.066 \text{ MPa}$
 $A_t = 0.1756 \text{ m}^2$
 $k = 1.221$

From [Atmosphere Properties](#),

$P_a = 0.0795 \text{ MPa}$

Equation (1.29),

$$N_a^2 = \left(\frac{2}{k-1} \right) \times \left[\left(\frac{P_c}{P_a} \right)^{\frac{k-1}{k}} - 1 \right]$$

$$N_a^2 = \left(\frac{2}{1.221-1} \right) \times \left[\left(\frac{5.066}{0.0795} \right)^{\frac{1.221-1}{1.221}} - 1 \right]$$

$$N_a^2 = 10.15$$

$$N_a = (10.15)^{1/2} = 3.185$$

Equation (1.30),

$$A_e = \left(\frac{A_t}{N_a} \right) \left[1 + \frac{k-1}{2} \right] N_a^2 \left(\frac{k+1}{2(k-1)} \right)^{\frac{k+1}{k-1}}$$

$$A_e = \left(\frac{0.1756}{3.185} \right) \times \left[1 + \frac{1.221-1}{2} \right] \times 10.15 \left(\frac{1.221+1}{2} \right)^{\frac{1.221+1}{1.221-1}}$$

$$A_e = 1.426 \text{ m}^2$$

Section Ratio,

$$A_e / A_t = 1.426 / 0.1756 = 8.12$$

PROBLEM 1.9

For the rocket engine in problem 1.7, calculate the volume and dimensions of a possible combustion chamber. The convergent cone half-angle is 20 degrees.

SOLUTION:

Given: $A_t = 0.1756 \text{ m}^2 = 1,756 \text{ cm}^2$
 $D_t = 2 \times (1,756 / \pi)^{1/2} = 47.3 \text{ cm}$
 $\theta = 20^\circ$

From Table 1,

$L^* = 102\text{-}127 \text{ cm}$ for LOX/RP-1, RP's use 110 cm

Equation (1.33),

$$V_c = A_t \times L^*$$

$$V_c = 1,756 \times 110 = 193,160 \text{ cm}^3$$

From Figure 1.7,

$L_c = 66 \text{ cm}$ (second-order approximation)

Equation (1.35),

$$D_c = \sqrt{\frac{D_t^3 + 24 \pi \times \tan \theta \times V_c}{D_c + 6 \times \tan \theta \times L_c}}$$

$$D_c = \sqrt{\frac{47.3^3 + 24 \pi \times \tan(20^\circ) \times 193,160}{D_c + 6 \times \tan(20^\circ) \times 66}}$$

$$D_c = 56.6 \text{ cm (four iterations)}$$

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