

Assignment 3 AFT

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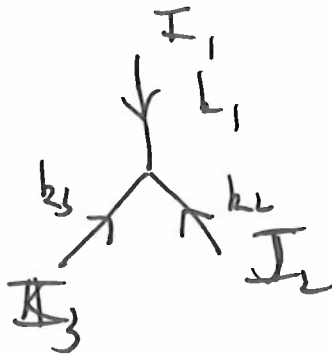
1. Each scalar has propagator

$$\langle 0 | T \phi_I(x) \phi_J(y) | 0 \rangle = \delta_{IJ} D(x, y) = D_{IJ}(x, y)$$

with:

$$D(x, y) = \int \frac{d^4 k}{(2\pi)^4} e^{-ik(x-y)} \frac{i}{k^2 - m^2 + i\epsilon}$$

The vertex



is

~~$$V_{I_1 I_2 I_3}(k_1, k_2, k_3) = i S_{int}$$~~

$$V_{I_1 I_2 I_3}(k_1, k_2, k_3) (2\pi)^4 \delta^4(k_1 + k_2 + k_3) = \frac{\delta^3}{\delta \tilde{\phi}_{I_1}(k_1) \delta \tilde{\phi}_{I_2}(k_2) \delta \tilde{\phi}_{I_3}(k_3)} i S_{int}$$



$$= -i g_{I_1 I_2 I_3} (2\pi)^4 \delta^4(k_1 + k_2 + k_3)$$

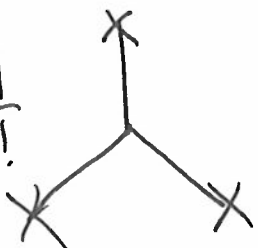
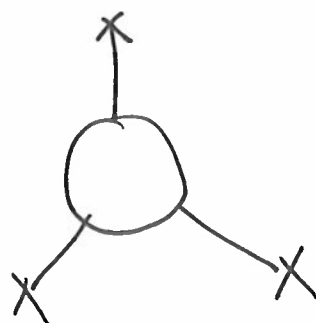
(2)

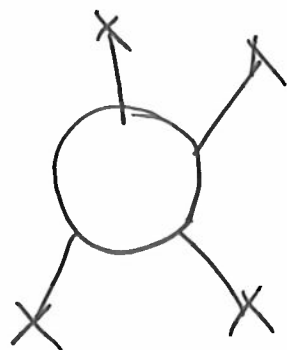
So $V_{I_1 I_2 I_3}(k_1, k_2, k_3) = -ig_{I_1 I_2 I_3}$

2. Sum of 1PI diagrams

$i\Gamma(\phi) =$ 

$+\frac{1}{2}$  $+$ $\frac{1}{2^2}$ 

$+\frac{1}{3!}$  $+$ $\frac{1}{3!}$ 

$+$ $\frac{1}{2^3}$  $+$ higher order

'Symmetry factors from Stueckelberg'.

In DeWitt notation

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$$* \text{---} \bigcirc = \phi_I (-ig_{IJK}) D_{JK}$$

In full this is $-i \int d^4x \sum_{IJK} \phi_I(x) g_{IJK} \delta_{JK} D(x,x)$

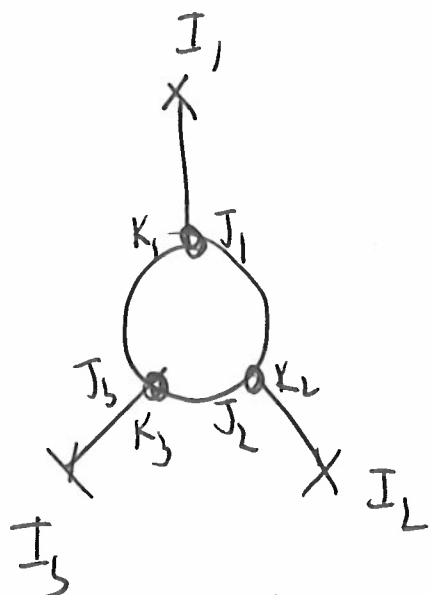
$$\frac{1}{2} * \text{---} * = \frac{1}{2} \phi_I (\overset{\text{from computational externals}}{D^{-1}})^{IJ} \phi_J$$

$$= i \int d^4x \left[\sum_I \frac{1}{2} (\partial \phi_I)^2 - \frac{1}{2} (m^2 - i\epsilon) \phi_I^2 \right]$$

$$\frac{1}{3!} * \text{---} * \text{---} * = \frac{-i}{3!} g_{IJK} \phi^I \phi^J \phi^K$$

$$= \frac{-i}{3!} \int d^4x \sum_{IJK} g_{IJK} \phi^I(x) \phi^J(x) \phi^K(x)$$

(4)

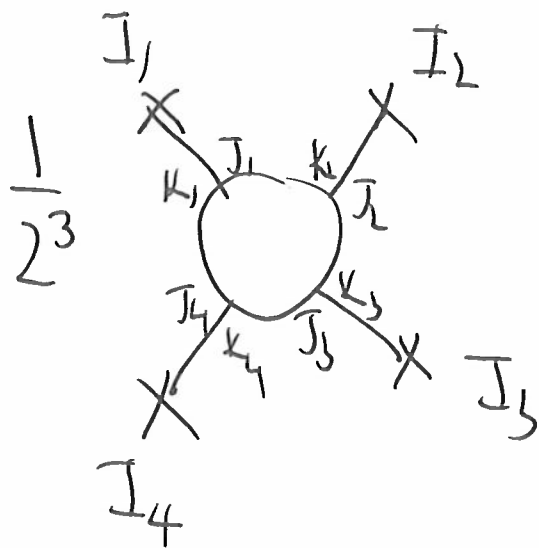
 $\frac{1}{3!}$ 

$$= \frac{1}{3!} \phi_{I_1} \phi_{I_2} \phi_{I_3} (-ig_{I_1 J_1 K_1}) (-ig_{I_2 J_2 K_2}) (-ig_{I_3 J_3 K_3})$$

$$\times D_{J_1 K_2} D_{J_2 K_3} D_{J_3 K_1}$$

$$= \frac{i}{3!} \int d^4 x_1 \int d^4 x_2 \int d^4 x_3 \sum_{I_1, I_2, I_3, J_1, J_2, K_1} \phi_{I_1}(x_1) \phi_{I_2}(x_2) \phi_{I_3}(x_3) g_{I_1 J_1 K_1} g_{I_2 J_2 K_2} g_{I_3 J_3 K_3} \\ \times D(x_1, x_2) D(x_2, x_3) D(x_3, x_1)$$

(5)



$$= \frac{1}{2^3} \phi_{I_1} \phi_{I_2} \phi_{I_3} \phi_{I_4} (-ig_{I_1 J_1 k_1}) (-ig_{I_2 J_2 k_2}) (-ig_{I_3 J_3 k_3}) (ig_{I_4 J_4 k_4})$$

$$\times D_{J_1 k_2} D_{J_2 k_3} D_{J_3 k_4} D_{J_4 k_1}$$

$$= \frac{1}{2^3} \int d^4 x_1 \int d^4 x_2 \int d^4 x_3 \int d^4 x_4 \sum_{I_1 \dots I_4, J_1 \dots J_4} \phi_{I_1}(x_1) \phi_{I_2}(x_2) \phi_{I_3}(x_3) \phi_{I_4}(x_4)$$

$$\times g_{I_1 J_1 J_4} g_{I_2 J_2 J_1} g_{I_3 J_3 J_2} g_{I_4 J_4 J_3}$$

$$\times D(x_1, x_2) D(x_2, x_3) D(x_3, x_4) D(x_4, x_1)$$

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3.

$$Z = \int d^4 \eta \left(1 + \frac{1}{2} \sum \eta_i A_{ij} \eta_j + \frac{1}{g} \sum_{ijkl} \eta_i \eta_j \eta_k \eta_l A_{ij} A_{kl} \right. \\ \left. + g \sum_{ijkl} \eta_i \eta_j \eta_k \eta_l B_{ijkl} + \text{higher order stuff vanishes} \right)$$

Since $\int d^4 \eta$ = 4 derivatives only the terms with 4 η 's matter.

Now $\int d^4 \eta \eta_i \eta_j \eta_k \eta_l = \epsilon_{ijkl}$ ↖ Levi-Civita!

$$\therefore \boxed{Z = \sum_{ijkl} \epsilon_{ijkl} \left(B_{ijkl} + \frac{1}{g} A_{ij} A_{kl} \right)}$$

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$$Z = \sum_{ijklabcd} \epsilon_{ijklabcd} \left[\begin{aligned} &g^2 B_{ijkl} B_{abcd} \\ &+ \frac{g}{8} B_{ijkl} A_{ab} A_{cd} \\ &+ \frac{1}{2^4 4!} A_{ij} A_{kl} A_{ab} A_{cd} \end{aligned} \right]$$

5.
$$Z = \int d\eta e^{\frac{1}{2} \sum \eta_i A_{ij} \eta_j} \left[\begin{aligned} &1 + \sum_{ijkl} g \eta_i \eta_j \eta_k \eta_l B_{ijkl} \\ &+ \sum_{ijklabcd} g^2 \eta_i \eta_j \eta_k \eta_l \eta_a \eta_b \eta_c \eta_d \\ &\quad \times (B_{ijkl} B_{abcd}) \end{aligned} \right]$$

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Wicks theorem

$$\overline{\eta_i \eta_j} = (A^{-1})_{ij}$$

For $g=0$ $Z = \int d\eta e^{\frac{1}{2} \sum \eta_i A_{ij} \eta_j}$

$$= \sqrt{\det(A)} \quad (\text{from lecture notes}).$$

$$\frac{Z}{\sqrt{\det(A)}} = 1 + g \sum_{ijkl} \langle \eta_i \eta_j \eta_k \eta_l \rangle B_{ijkl}$$

$$+ g^2 \sum_{ijklabcd} \langle \eta_i \eta_j \eta_k \eta_l \eta_a \eta_b \eta_c \eta_d \rangle B_{ijkl} B_{abcd}$$

$$= 1 + g \text{ (diagram: a single loop with a dot)} + g^2 \left[\begin{array}{c} \text{diagram: two separate loops, each with a dot} \\ + \\ \text{diagram: two loops connected by a line, each with a dot} \end{array} \right]$$

(10)

Now $\langle \eta_i \eta_j \eta_k \eta_l \rangle = (A^{-1})_{ij} (A^{-1})_{kl}$
 $- (A^{-1})_{ik} (A^{-1})_{jl}$

$$\langle \eta_i \eta_j \eta_k \eta_l \eta_a \eta_b \eta_c \eta_d \rangle$$

$$= (A^{-1})_{ij} (A^{-1})_{kl} (A^{-1})_{ab} (A^{-1})_{cd}$$

+ antisymmetrized permutations.

(11)

$$\int d^N \eta = \int d\eta_N \dots \int d\eta_1$$

$$\delta(\eta) = \eta_1 \dots \eta_N$$

$$\int d\eta_1 \delta(\eta) = \eta_2 \dots \eta_N$$

$$\int d\eta_2 \int d\eta_1 \delta(\eta) = \eta_3 \dots \eta_N$$

repeating $\int d^N \eta \delta(\eta) = 1$.

7. Consider

$$\int d^N \eta f(\eta) \delta(\eta)$$

Write $f(\eta)$ as a Taylor expansion. Since $\delta(\eta)$ has N etas and there are only N etas being integrated over, any more gives a repetition which must vanish.

$$\therefore \int d^N \eta f(\eta) \delta(\eta) = \int d^N \eta f(0) \delta(\eta) = f(0).$$

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8.

$$\delta(\eta) = C \int d^N J e^{-i\eta_N J_N} \dots e^{-i\eta_1 J_1}$$

$$= C \left(\int dJ_N e^{-i\eta_N J_N} \right) \dots \left(\int dJ_1 e^{-i\eta_1 J_1} \right)$$

Now $\int dJ_i e^{-i\eta_i J_i} = \int dJ_i (1 - i\eta_i J_i)$

$$= \int dJ_i (1 + iJ_i \eta_i)$$

$$= i\eta_i$$

$$\therefore \delta(\eta) = i^N C \eta_N \dots \eta_1 = i^N C \delta(\eta)$$

$$\therefore \boxed{C = \frac{1}{i^N}}$$