

# Benchmark tests on pressure boundary conditions

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## 1. Analytical solutions

In the following are reported the formula relatives to the analytical solutions of the stress and displacement fields around a thick-walled pipe and sphere in elastic conditions and a thick walled sphere in elasto-plastic conditions. The full derivation of the analytical solutions can be found in [1].

### 1.1. Thick-walled elastic cylinder

The stress field around a thick walled elastic cylinder in plain strain conditions is expressed by the set of equations

$$\sigma_{rr} = \frac{p_i R_i^2 - p_a R_a^2}{R_a^2 - R_i^2} - \frac{(p_i - p_a) R_a^2 R_i^2}{(R_a^2 - R_i^2) r^2}, \quad (1)$$

$$\sigma_{\theta\theta} = \frac{p_i R_i^2 - p_a R_a^2}{R_a^2 - R_i^2} + \frac{(p_i - p_a) R_a^2 R_i^2}{(R_a^2 - R_i^2) r^2}, \quad (2)$$

$$\sigma_{zz} = 2\nu \frac{p_i R_i^2 - p_a R_a^2}{R_a^2 - R_i^2}, \quad (3)$$

defining the radial  $\sigma_{rr}$ , circumferential  $\sigma_{\theta\theta}$  and longitudinal  $\sigma_{zz}$  stress components in cylindrical coordinates, where the  $z$  axis is along the cylinder directive. The thick walled cylinder is subjected to an internal and an external pressure, respectively  $p_i$  and  $p_a$ , and has internal and external radii of  $R_i$  and  $R_a$ . The radial coordinate is  $r$  and  $\nu$  is Poisson's ratio. The radial displacement  $u_r$  writes

$$u_r = \frac{R_i^2 p_i (1 + \nu) r}{E (R_i^2 - R_a^2)} \left\{ \left[ \frac{p_a}{p_i} \left( \frac{R_a}{R_i} \right)^2 - 1 \right] (1 - 2\nu) + \left( \frac{p_a}{p_i - 1} \right) \left( \frac{R_a}{r} \right)^2 \right\} \quad (4)$$

where  $E$  is Young's modulus.

### 1.2. Thick-walled elastic sphere

The stress field around a thick walled elastic sphere of internal and external radii  $R_i$  and  $R_a$ , respectively, and subjected to and internal and external pressure of  $p_i$  and  $p_a$ , respectively, is given by

$$\sigma_{rr} = \frac{p_a R_a^3 - p_i R_i^3}{R_i^2 - R_a^2} - \frac{(p_a - p_i) R_a^3 R_i^3}{(R_i^3 - R_a^3) r^3}, \quad (5)$$

$$\sigma_{\theta\theta} = \sigma_{\phi\phi} = \frac{p_a R_a^3 - p_i R_i^3}{R_i^2 - R_a^2} + \frac{(p_a - p_i) R_a^3 R_i^3}{2 (R_i^3 - R_a^3) r^3}, \quad (6)$$

where in spherical coordinates,  $\theta$  and  $\phi$  define the angular coordinates and  $r$  the radial one. The radial displacement field writes

$$u_r = \frac{R_i^3 p_i r}{E (R_i^3 - R_a^3)} \left\{ \left[ \frac{p_a}{p_i} \left( \frac{R_a}{R_i} \right)^3 - 1 \right] (1 - 2\nu) + \left( \frac{p_a}{p_i} - 1 \right) \frac{1 + \nu}{2} \left( \frac{R_a}{r} \right)^3 \right\}. \quad (7)$$

### 1.3. Thick walled plastic sphere

The solution obtained is based on perfect plasticity theory (no hardening or softening) and the Von Mises plastic yield surface  $F$ , which reads

$$F = \frac{1}{2} s_{ij} s_{ij} - \frac{1}{3} \sigma_F^2 = 0, \quad (8)$$

where  $s_{ij} = \sigma_{ij} - 1/3 \sigma_{kk} \delta_{ij}$  is the deviatoric stress tensor and  $\sigma_F$  is the yield stress. A new variable  $r_p$  defines the boundary between the elastic and the plastic region and the yield pressure reads

$$p_{pl} = 2\sigma_F \ln \frac{r_p}{R_i} + \frac{2}{3} \sigma_F \left[ 1 - \left( \frac{r_p}{R_a} \right)^3 \right]. \quad (9)$$

In the elastic region, the stress components are

$$\sigma_{rr} = \frac{2}{3} \sigma_F \left[ \left( \frac{r_p}{R_a} \right)^3 - \left( \frac{r_p}{r} \right)^3 \right], \quad (10)$$

$$\sigma_{\theta\theta} = \sigma_{\phi\phi} = \frac{2}{3} \sigma_F \left[ \left( \frac{r_p}{R_a} \right)^3 + \frac{1}{2} \left( \frac{r_p}{r} \right)^3 \right], \quad (11)$$

and the radial displacement is

$$u_r = \frac{2\sigma_F}{3E} \left( \frac{r_p}{R_a} \right)^3 r \left[ (1 - 2\nu) + \frac{1 + \nu}{2} \left( \frac{r_a}{r} \right)^3 \right]. \quad (12)$$

In the plastic region, stress components write

$$\sigma_{rr} = 2\sigma_F \ln \frac{r_p}{R_i} + \frac{2}{3}\sigma_F \left[ \left( \frac{r_p}{R_a} \right)^3 - 1 \right], \quad (13)$$

$$\sigma_{\theta\theta} = \sigma_{\phi\phi} = 2\sigma_F \ln \frac{r_p}{R_i} + \frac{2}{3}\sigma_F \left[ \left( \frac{r_p}{R_a} \right)^3 + \frac{1}{2} \right], \quad (14)$$

while the radial displacement is

$$u_r = \frac{2(1 - 2\nu)\sigma_F}{3E} r \left[ \frac{3(1 - \nu)}{2(1 - 2\nu)} \left( \frac{r_p}{r} \right)^3 - 1 + \left( \frac{r_p}{R_a} \right)^3 + 3 \ln \left( \frac{r}{r_p} \right) \right]. \quad (15)$$

## 2. Numerical analyses and comparison

### 2.1. Elastic cylinder

The solution was obtained with the values reported in Table 2.1. The comparison between analytical and numerical solution is reported in Figure 1 for plain strain conditions and in Figure 2 for axisymmetric conditions. The numerical solution matches well the analytical ones.

Table 1: Values of parameters for elastic cylinder problem.

Parameter	Value	Units
$R_i$	1	mm
$R_a$	2	mm
$p_i$	52.2	MPa
$p_a$	0.1	MPa
$\nu$	0.3	-
$E$	210	GPa

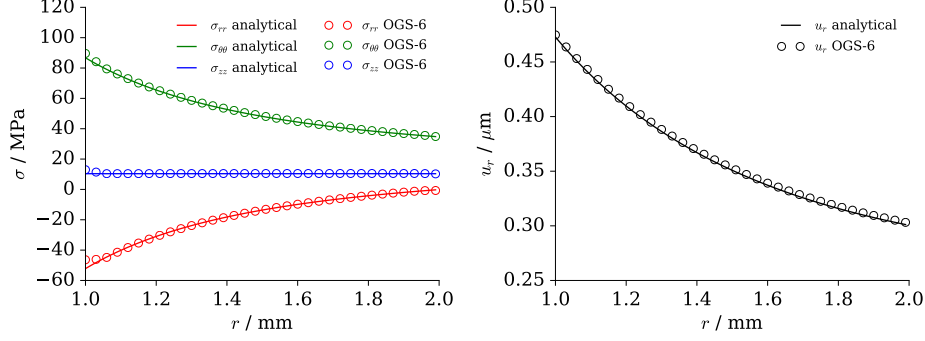


Figure 1: Plain strain elastic cylinder comparison between numerical and analytical results.

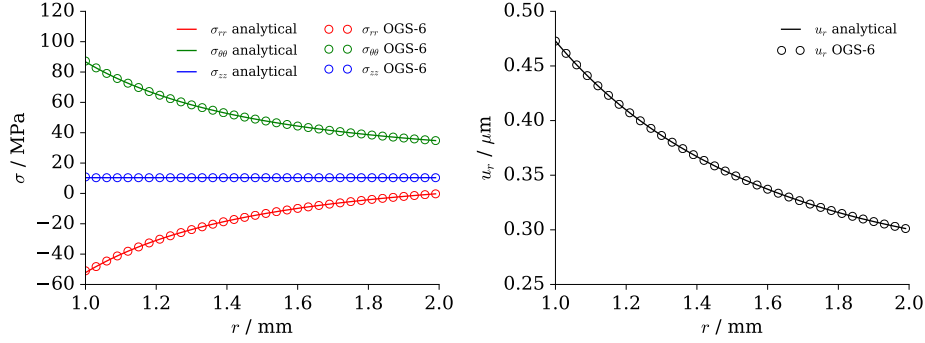


Figure 2: Axisymmetric elastic cylinder comparison between numerical and analytical results.

## 2.2. Elastic sphere

The comparison is carried out for the case of an elastic sphere of which properties are reported in Table 2.2. Two models are used for the numerical computations, a bi-dimensional axisymmetric one and a full tri-dimensional model. Results comparison for the axisymmetric model are reported in Figure 3 and for the tri-dimensional one in Figure 4. There is good agreement between numerical computations and the analytical solution. The slight discrepancy at the inner boundary is caused by nodal interpolation error.

Table 2: Values of parameters for elastic sphere problem.

Parameter	Value	Units
$R_i$	1	mm
$R_a$	2	mm
$p_i$	1	kPa
$p_a$	100	kPa
$\nu$	0.35	-
$E$	125	GPa

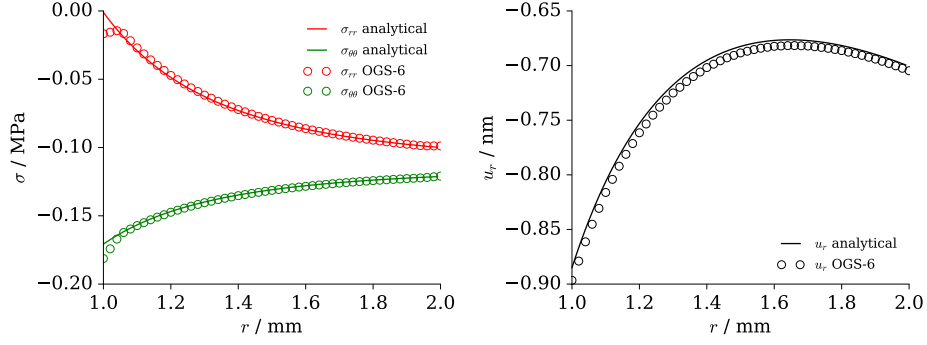


Figure 3: Axisymmetric elastic sphere comparison between numerical and analytical results.

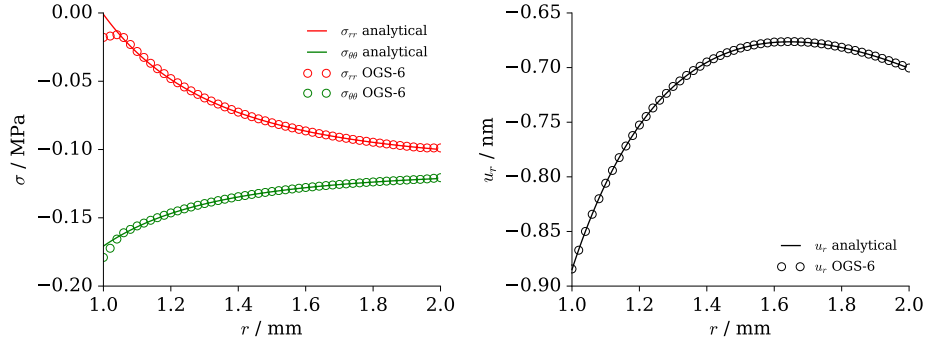


Figure 4: Tri-dimensional elastic sphere comparison between numerical and analytical results.

### 2.3. Plastic sphere

The final benchmark consists in simulating an elasto-plastic sphere subjected to an internal pressure. The material properties, geometry and boundary conditions are reported in Table 2.3. The plastic radius  $r_p$  was set to be at the middle point of the sphere wall thickness. Based on this value, through Equation 9 the plastic stress  $p_{pl}$  was computed. In the numerical analysis, the internal pressure was increased linearly to the value of the computed plastic pressure  $p_{pl}$ . Results comparison are shown in Figure 5 and once again, the numerical solution well compares to the analytical one, validating the pressure boundary conditions implementation in OGS-6.

Table 3: Values of parameters for plastic sphere problem.

Parameter	Value	Units
$R_i$	1	mm
$R_a$	2	mm
$p_i$	239.27	MPa
$p_a$	0	MPa
$\nu$	0.35	-
$E$	125	GPa
$\sigma_F$	200	MPa
$r_p$	1.5	mm
$p_{pl}$	239.27	MPa
$E$	125	GPa

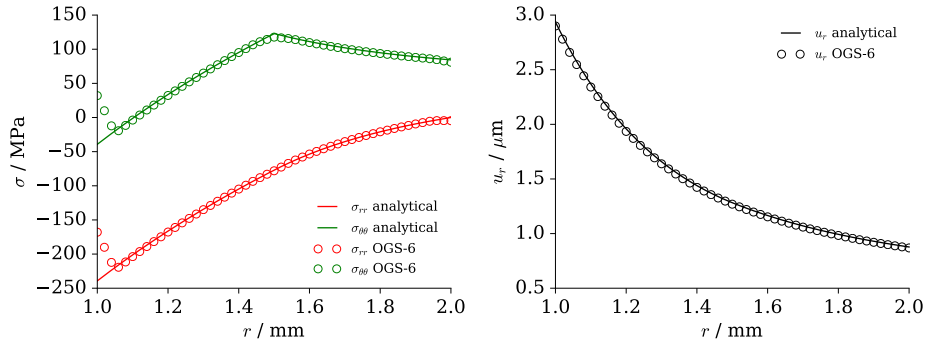


Figure 5: Axisymmetric plastic sphere comparison between numerical and analytical results.

### **3. Bibliography**

- [1] O. Kolditz, H. Shao, W. Wang, S. Bauer, Thermo-hydro-mechanical-chemical processes in fractured porous media: modelling and benchmarking, Springer, 2016.