

Description of a first order multi-plane model. The different planes are supposed to be independant and interact with the incoming plane wave by assuming that this latter was not modified by the upstream planes.

Notations:

- $A \times B$ : matrix multiplication
- $A.B$ : element wise product
- $\frac{A}{B}$ : element wise division
- $e^A$ : element wise exponentiation
- $A^*$ : hermitian transposition
- $\bar{A}$ : complex conjugate
- $\mathcal{R}(A)$ : element wisereal part of a complex matrix
- $P_{z_p}$ :  $nb_{plane}$  propagators on distances equal to  $z_p$
- $\theta$ : variable with  $nb_{plane}$  complex planes  $\theta_p$  describing the complex optical length of the different planes
- $pad$ : padding operators
- $U_{BP}$ : background wavefront

## 1 Coherent interferences

Direct model:

$$\begin{aligned} U_{tot}(\theta) &= U_{BP} + \sum_{p=1}^{nb_{plane}} e^{-ikz_p} \left( pad^{-1} \left( P_{z_p} \otimes pad \left( e^{\frac{2i\pi}{\lambda} \theta_p} - 1 \right) \right) \right) \\ &= U_{BP} + \sum_{p=1}^{nb_{plane}} O_p \times \left( e^{\frac{2i\pi}{\lambda} \theta_p} - 1 \right) \end{aligned}$$

with  $O_p = e^{-ikz_p} pad^{-1} \circ P_{z_p} \otimes pad$  a linear operator.

Application of the hermitian transposition of the Jacobian at a given point  $\theta$  for a given plane  $p$ :

$$J_p^*(\theta) = -\frac{2i\pi}{\lambda} e^{-\frac{2i\pi}{\lambda} \theta_p} \cdot O_p^*$$

## 2 Incoherent interferences

Note on the interference of two incoherent waves  $U_1$  and  $U_2$ :

$$\begin{aligned} |U_{tot}|^2 &= (U_1 + U_2) \overline{(U_1 + U_2)} \\ &= |U_1|^2 + |U_2|^2 + 2\mathcal{R}(\overline{U_1}U_2) \end{aligned}$$

$2\mathcal{R}(\overline{U_1}U_2)$  is the interference term. For coherent waves, its temporal average is not null. But it disappears with temporal integration for incoherent waves.

Let's now assume  $nb_{plane}$  incoherent waves  $U_p$  but coherent with a background wave  $U_{BP}$ . Let's determine the resulting intensity:

$$\begin{aligned} |U_{tot}|^2 &= \left( U_{BP} + \sum_{p=1}^{nb_{plane}} U_p \right) \overline{\left( U_{BP} + \sum_{p=1}^{nb_{plane}} U_p \right)} \\ &= |U_{BP}|^2 + \sum_{p=1}^{nb_{plane}} (|U_p|^2 + 2\mathcal{R}(\overline{U_{BP}}U_p)) + \sum_{p \neq p'} \mathcal{R}(\overline{U_p}U_{p'}) \end{aligned}$$

According to the previous remark,  $\forall (p, p')$ , it comes  $\mathcal{R}(\overline{U_p}U_{p'}) = 0$ . Then :

$$\begin{aligned} |U_{tot}|^2 &= (1 - nb_{plane}) |U_{BP}|^2 + \sum_{p=1}^{nb_{plane}} (|U_p|^2 + 2\mathcal{R}(\overline{U_{BP}}U_p) + |U_{BP}|^2) \\ &= (1 - nb_{plane}) |U_{BP}|^2 + \sum_{p=1}^{nb_{plane}} |U_p + U_{BP}|^2 \end{aligned}$$

The direct model for incoherent waves is then:

$$\begin{aligned} |U_{tot}(\theta)|^2 &= (1 - nb_{plane}) |U_{BP}|^2 + \\ &\quad \dots \\ &\quad \sum_{p=1}^{nb_{plane}} \left| U_{BP} + e^{-ikz_p} \left( pad^{-1} \left( P_{z_p} \otimes pad \left( e^{\frac{2i\pi}{\lambda} \theta_p} - 1 \right) \right) \right) \right|^2 \\ &= (1 - nb_{plane}) |U_{BP}|^2 + \sum_{p=1}^{nb_{plane}} \left| U_{BP} + O_p \times \left( e^{\frac{2i\pi}{\lambda} \theta_p} - 1 \right) \right|^2 \end{aligned}$$

with  $O_p = e^{-ikz_p} pad^{-1} \circ P_{z_p} \otimes pad$  a linear operator.

Application of the hermitian transposition of the Jacobian at a given point  $\theta$  for a given plane  $p$  computed in  $y$ :

$$J_p^*(\theta) = -\frac{2i\pi}{\lambda} \frac{\overline{U_{BP} + O_p \times \left(e^{\frac{2i\pi}{\lambda}\theta_p} - 1\right)}}{\left|U_{BP} + O_p \times \left(e^{\frac{2i\pi}{\lambda}\theta_p} - 1\right)\right|} \cdot e^{-\frac{2i\pi}{\lambda}\theta_p} \cdot O_p^*$$

Important note: this approach can correctly deal with the spatial coherence length which can be different for each plane.

### 3 Cost function

It is possible to take into account a global spatial coherence length by convolving the simulated intensity  $|U_{tot}|^2$  by a filter  $F_{coh}$  corresponding to a diaphragm of the source seen by the sample and the sensor. This is a strong approximation as it is supposed to be equal for all the planes because it is applied after their interference. In the following,  $F_{coh}$  being a linear operator, it is seen as its matrix shape.

Intensity cost function for a given  $\theta$  and intensity data  $I_d$ :

$$D(\theta) = \left\| F_{coh} \times |U_{tot}(\theta)|^2 - I_d \right\|_W^2$$

Getting inspired by the appendix D of my PhD, let's now find the first order development of  $D(\theta + \delta\theta)$  according to  $\delta\theta$ . Assuming that  $\delta\theta$  is real by decomposing the complex values on their real and imaginary part, it comes:

$$\begin{aligned}
D(\theta + \delta\theta) &= \|F_{coh} \times |U_{tot}(\theta + \delta\theta)|^2 - I_d\|_W^2 \\
&= \|F_{coh} \times |U_{tot}(\theta) + J(\theta) \times \delta\theta|^2 - I_d\|_W^2 \\
&= \|F_{coh} \times (|U_{tot}(\theta)|^2 + |J(\theta) \times \delta\theta|^2 + \dots \\
&\quad 2\mathcal{R}(\overline{U_{tot}(\theta)} \cdot (J(\theta) \times \delta\theta))) - I_d\|_W^2 \\
&= \|F_{coh} \times |U_{tot}(\theta)|^2 - I_d\|_W^2 + \dots \\
&\quad 2\langle F_{coh} \times |U_{tot}(\theta)|^2 - I_d, F_{coh} \times \dots \\
&\quad 2\mathcal{R}(\overline{U_{tot}(\theta)} \cdot (J(\theta) \times \delta\theta)) \rangle_{W, \mathbb{R}} + o(\|\delta\theta\|) \\
&= D(\theta) + 4\mathcal{R}\langle F_{coh} \times |U_{tot}(\theta)|^2 - I_d, F_{coh} \times \dots \\
&\quad \overline{U_{tot}(\theta)} \cdot (J(\theta) \times \delta\theta) \rangle_{W, \mathbb{C}} + o(\|\delta\theta\|) \\
&= D(\theta) + 4\mathcal{R}\langle W \cdot (F_{coh} \times |U_{tot}(\theta)|^2 - I_d), F_{coh} \times \dots \\
&\quad \overline{U_{tot}(\theta)} \cdot (J(\theta) \times \delta\theta) \rangle_{\mathbb{C}} + o(\|\delta\theta\|) \\
&= D(\theta) + 4\mathcal{R}\langle U_{tot}(\theta) \cdot F_{coh}^* \times W \cdot (F_{coh} \times |U_{tot}(\theta)|^2 - I_d), \dots \\
&\quad J(\theta) \times \delta\theta \rangle_{\mathbb{C}} + o(\|\delta\theta\|) \\
&= D(\theta) + 4\mathcal{R}\langle J^*(\theta) \times [U_{tot}(\theta) \cdot F_{coh}^* \times W \cdot \dots \\
&\quad (F_{coh} \times |U_{tot}(\theta)|^2 - I_d)] , \delta\theta \rangle_{\mathbb{C}} + o(\|\delta\theta\|) \\
&= D(\theta) + \langle 4\mathcal{R}[J^*(\theta) \times [U_{tot}(\theta) \cdot F_{coh}^* \times W \cdot \dots \\
&\quad (F_{coh} \times |U_{tot}(\theta)|^2 - I_d)]] , \delta\theta \rangle_{\mathbb{R}} + o(\|\delta\theta\|)
\end{aligned}$$

The gradient is then:

$$\nabla D(\theta) = 4\mathcal{R}[J^*(\theta) \times [U_{tot}(\theta) \cdot F_{coh}^* \times W \cdot (F_{coh} \times |U_{tot}(\theta)|^2 - I_d)]]$$

Note here that it assumes that  $\theta$  is decomposed on its real and imaginary part. Using the formalism presented in the chapter 3, section 3.3 for my PhD, it is possible to directly express the gradient in terms of complex values:

$$\nabla D(\theta) = 4J^*(\theta) \times [U_{tot}(\theta) \cdot F_{coh}^* \times W \cdot (F_{coh} \times |U_{tot}(\theta)|^2 - I_d)]$$