

Spin Qubit Control using NV Centres in Diamond

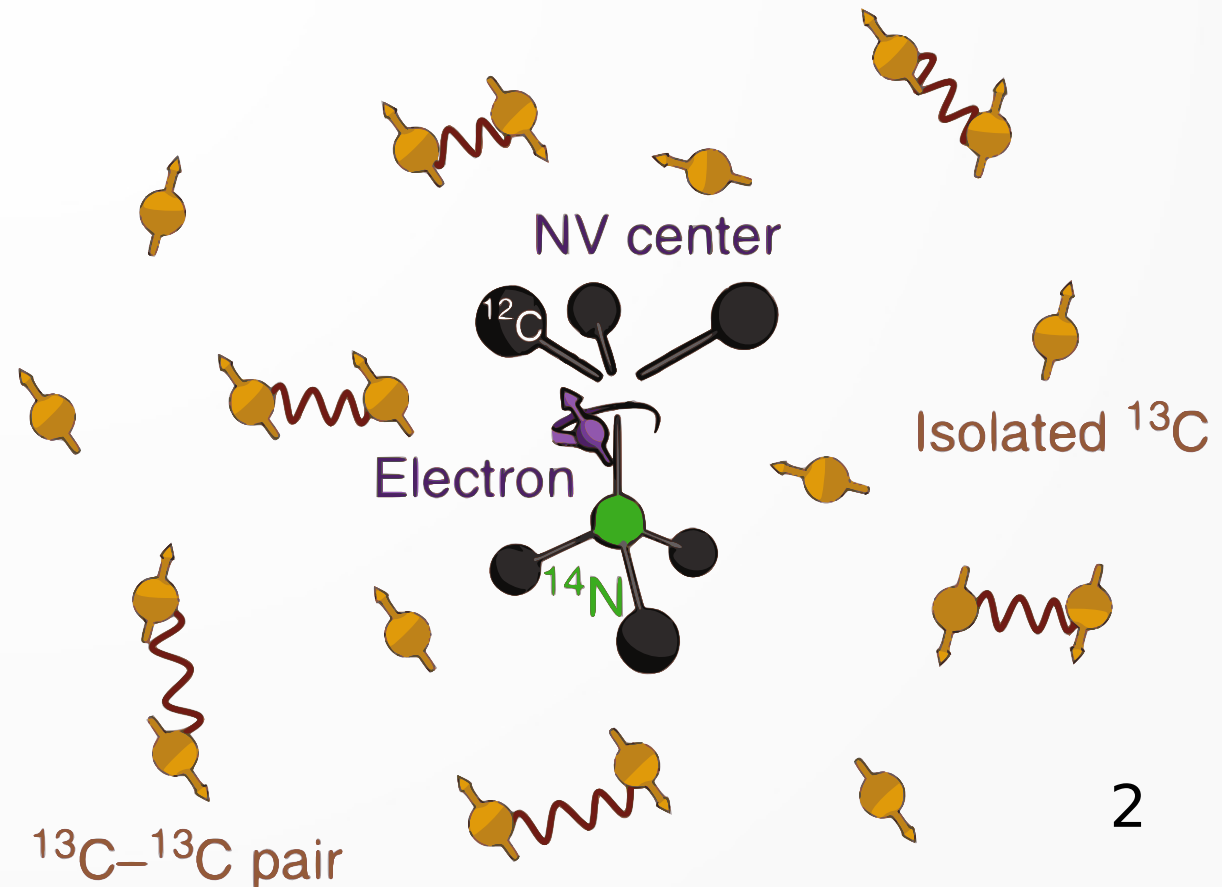
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Outline

- **NV Centres** – defect centres in diamond which, along with surrounding ^{13}C , constitute promising platforms for effective qubits
- **Model and simulation** – implemented an error model to simulate decoherence
- **Dynamical decoupling** – with perfect and imperfect π -rotations to extend decoherence time for spin qubits

Nitrogen-Vacancy Centres

- NV centres are spin-1 defects in the crystal structure of diamond, **effective 2-level system**
- ^{12}C bath protects spin of NV
- Surrounding ^{13}C atoms (spin-1/2) can be used as storage qubits
- **Main Challenge** – mitigating decoherence through dephasing noise due to magnetic fluctuations from surrounding ^{13}C spin bath¹



¹ van der Sar, T., Wang, Z., Blok, M. et al. Decoherence-protected quantum gates for a hybrid solid-state spin register. Nature 484, 82–86 (2012). <https://doi.org/10.1038/nature10900>

² Abobeih, M.H., Cramer, J., Bakker, M.A. et al. One-second coherence for a single electron spin coupled to a multi-qubit nuclear-spin environment. Nat Commun 9, 2552 (2018). <https://doi.org/10.1038/s41467-018-04916-z>

System Simulation

Evolution of state:

$$\frac{d\rho}{dt} = -\frac{i}{\hbar} [H, \rho]$$

ρ Density matrix.

H Hamiltonian.

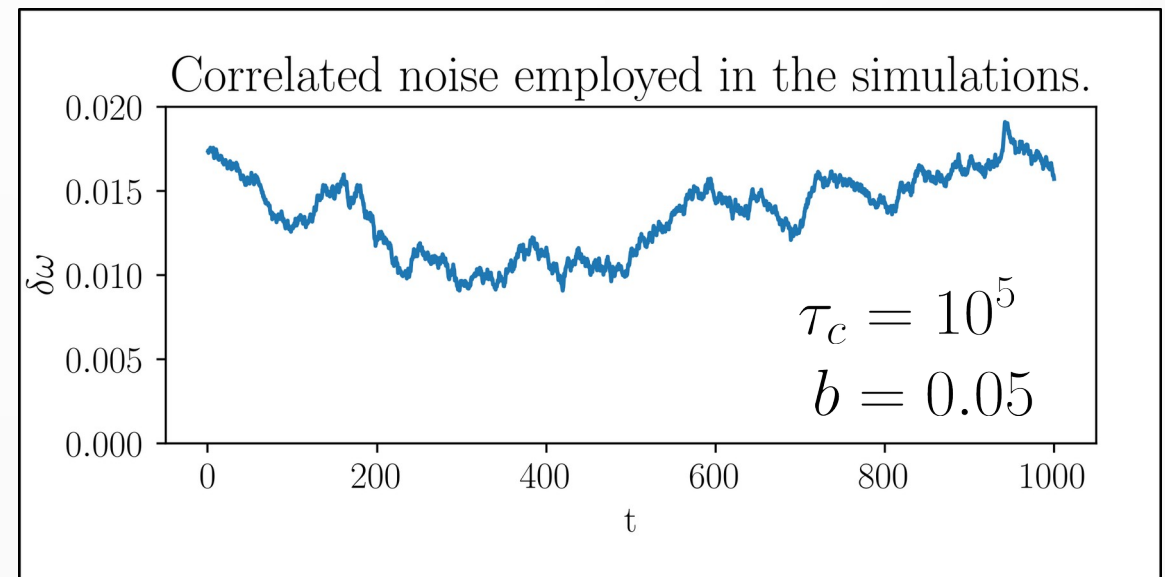
Numerical simulations performed with *Runge-Kutta 4*.

Error model:[‡]

$$H(t) = \delta\omega(t)\sigma_Z$$

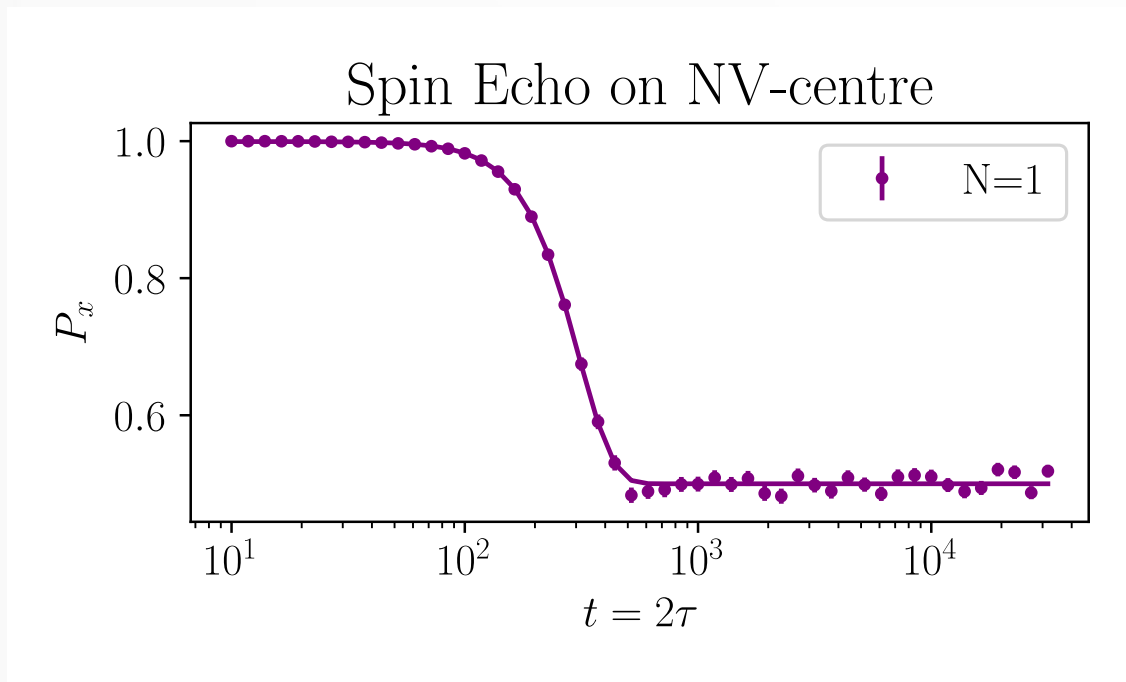
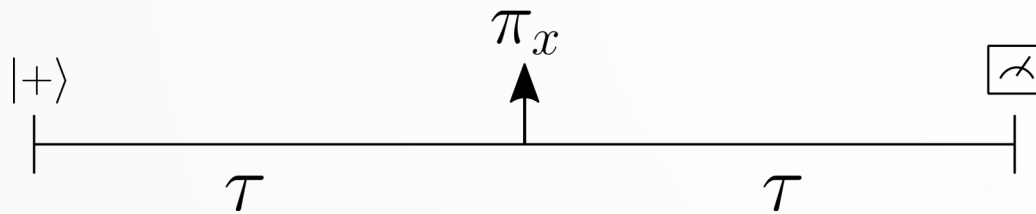
Auto-correlated normal noise
(average final values 960 times),

$$\langle \delta\omega(t)\delta\omega(0) \rangle = b^2 e^{-|t|/\tau_c}$$



[‡] Wang, Zhi-Hui et al. "Comparison of Dynamical Decoupling Protocols for a Nitrogen-Vacancy Center in Diamond." Physical Review B 85.15 (2012): <https://doi.org/10.1103/PhysRevB.85.155204>

Spin echo



Simulation parameters:

$$\langle \delta\omega(t)\delta\omega(0) \rangle = b^2 e^{-|t|/\tau_c}$$

$$\tau_c = 10^5 \quad b = 0.05$$

Fit function:

$$P_x = \frac{1 + e^{-(t/T_2)^n}}{2}$$

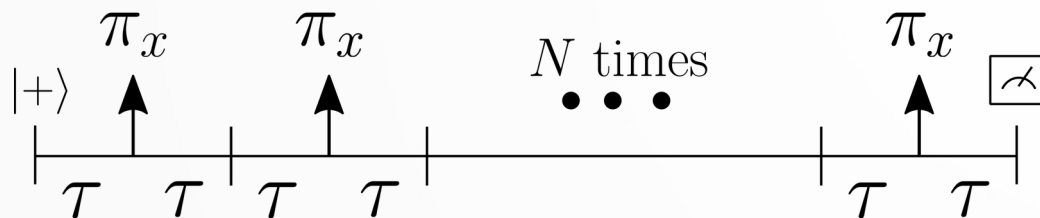
Theory:

$$\left\{ \begin{array}{l} T_2 = 782 \\ n = 3.0 \end{array} \right.$$

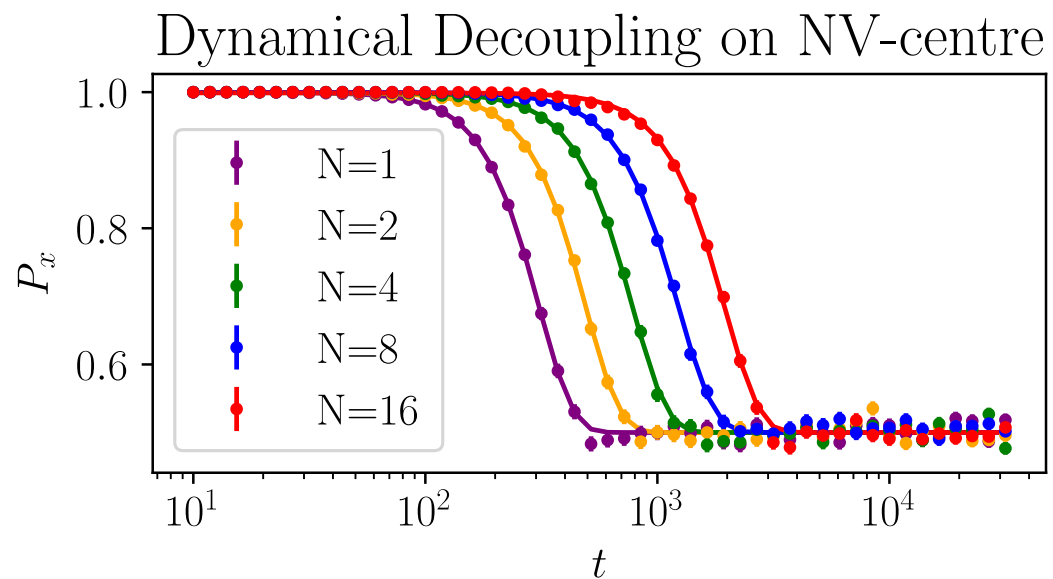
Simulation:

$$\left\{ \begin{array}{l} T_2 = 310(3) \\ n = 3.0(1) \end{array} \right.$$

Dynamical Decoupling



Same simulation parameters

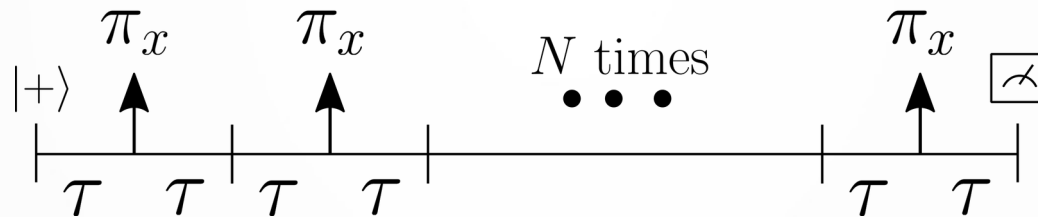


Same fit function: $P_x = \frac{1 + e^{-\left(t/T_{1/e}\right)^n}}{2}$

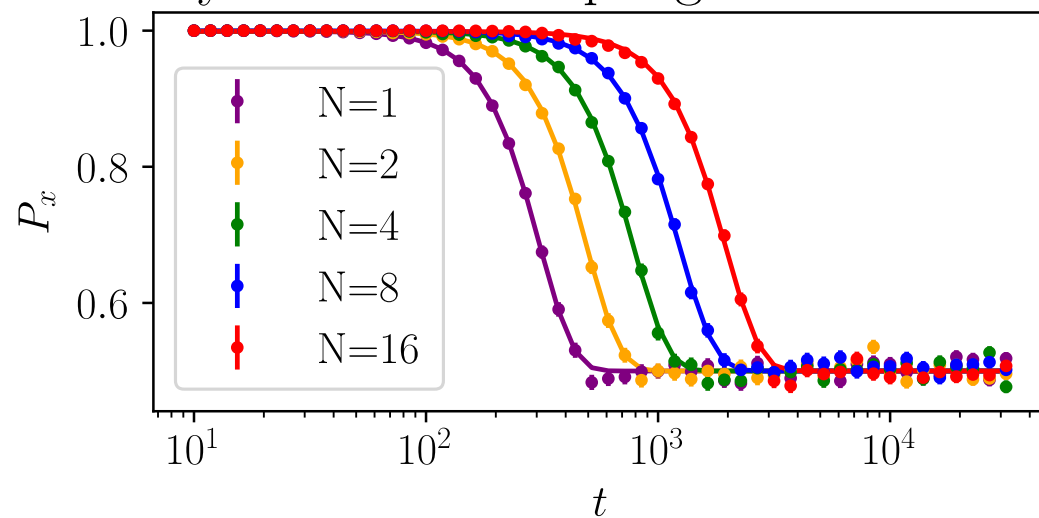
Theory: $n = 3.0$

Worst simulation: $n = 2.90(7)$

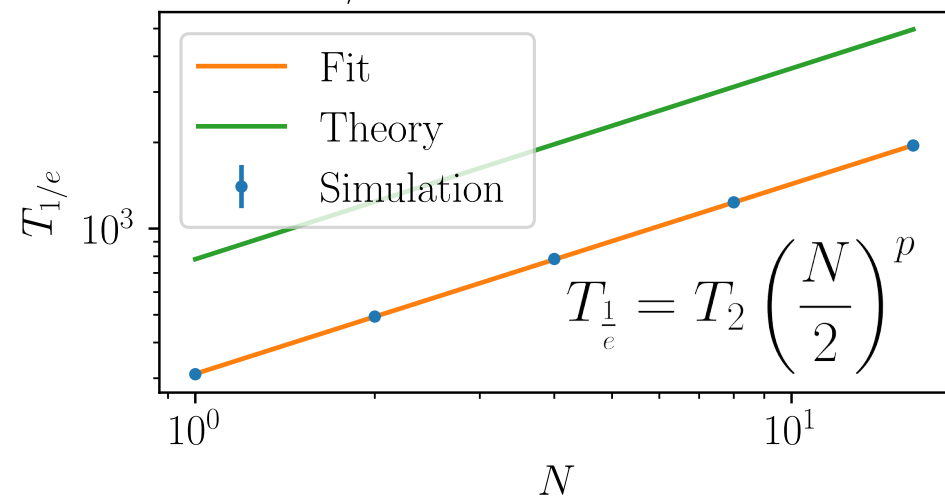
Dynamical Decoupling



Dynamical Decoupling on NV-centre



Scaling of $T_{1/e}$ with Dynamical Decoupling

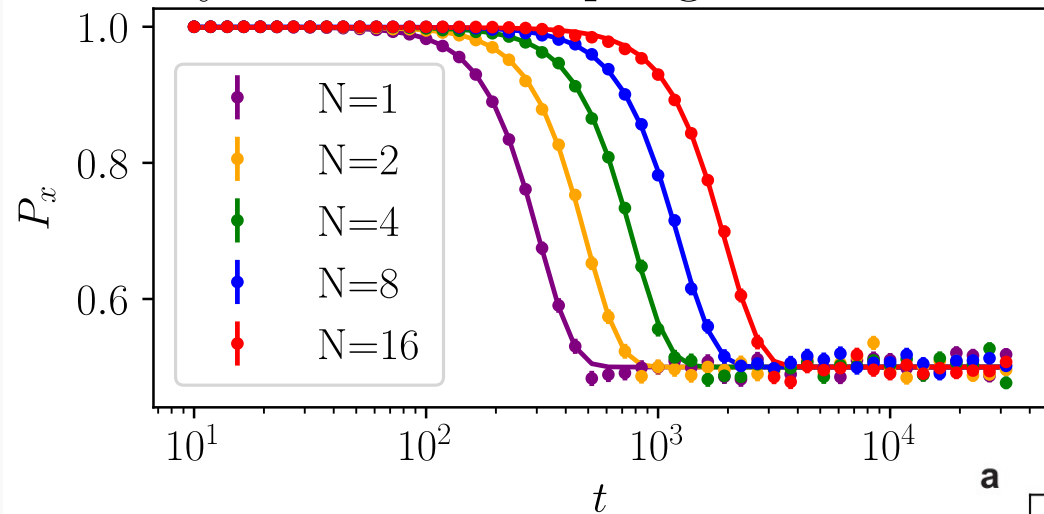


$$p_{\text{theory}} = 2/3$$

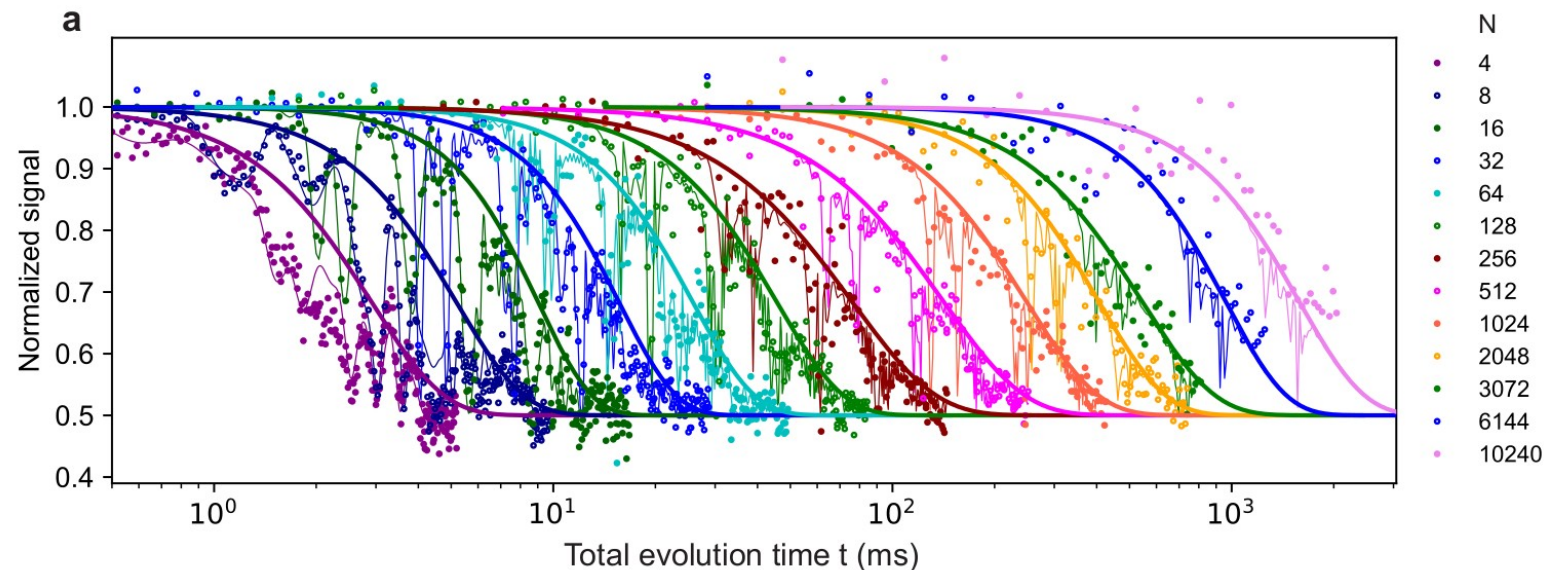
$$p_{\text{simulation}} = 0.663(1)$$

Dyn Decoupling

Dynamical Decoupling on NV-centre



- Same qualitative features between simulation and experiment (envelope)
- Values of n for the fit vary from 1.8 to 2.8



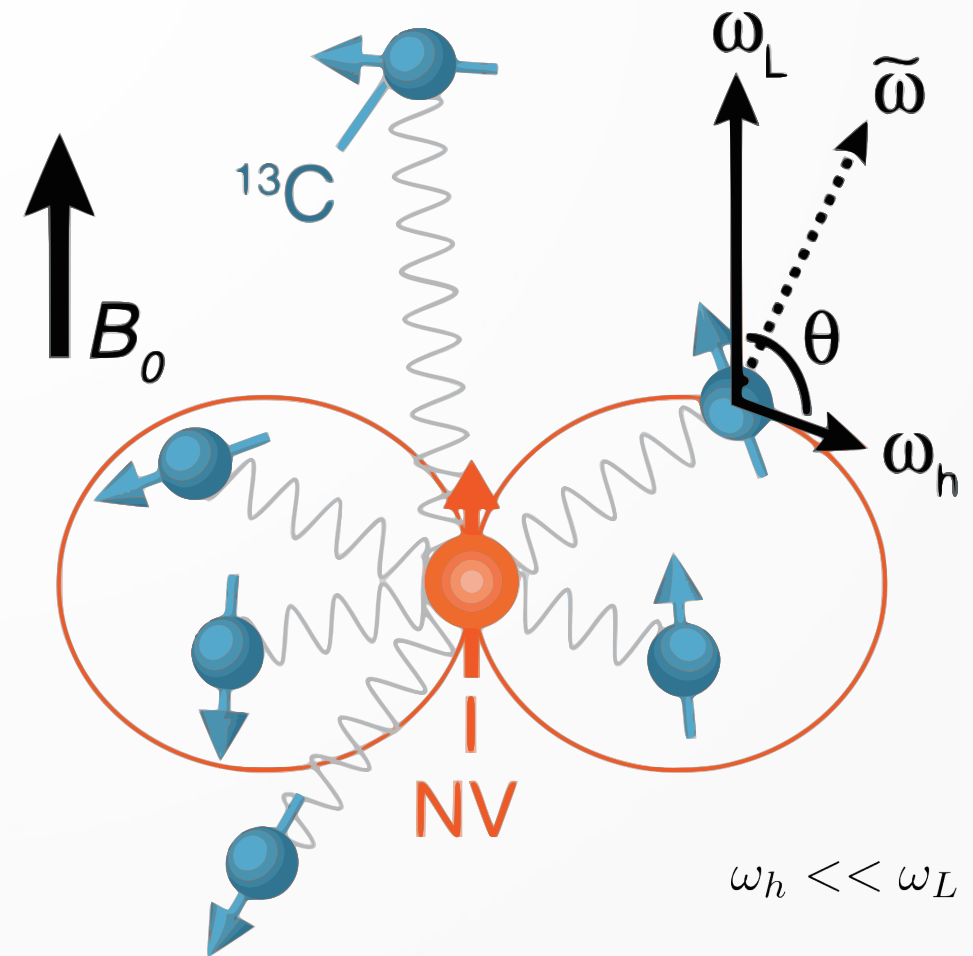
NV-Carbon Interaction

- Hamiltonian of the ^{13}C atoms is dependent on the NV centre spin state

$$H_0 = \omega_L \hat{S}_z$$

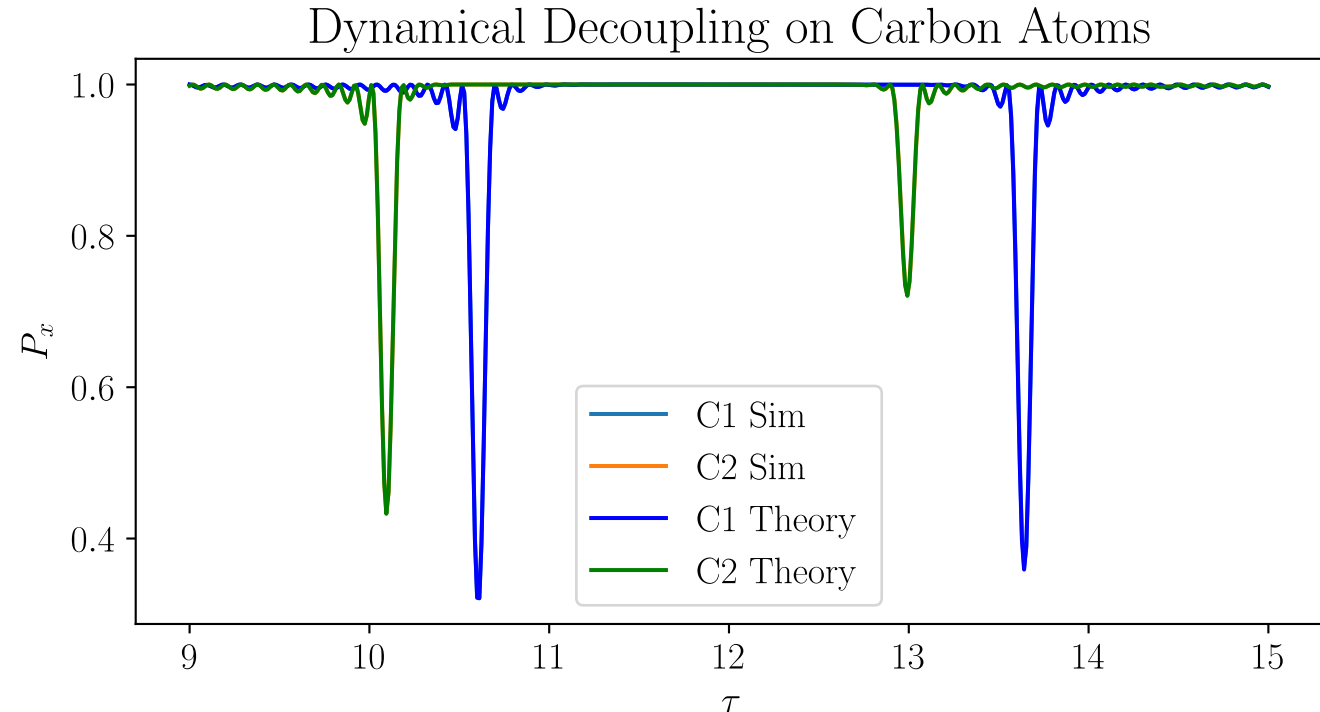
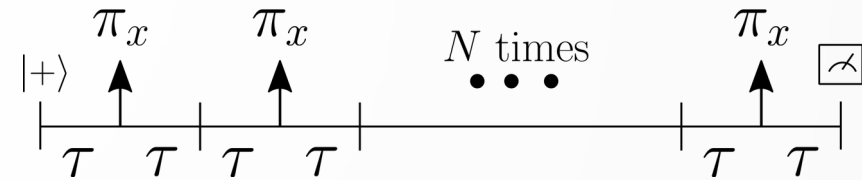
$$H_1 = (\omega_L + w_h \cos\theta) \hat{S}_z + w_h \sin\theta \hat{S}_x$$

- Raises possibility of **individual control** because each ^{13}C has its own resonance with the NV centre



Dynamical Decoupling and Entanglement Generation

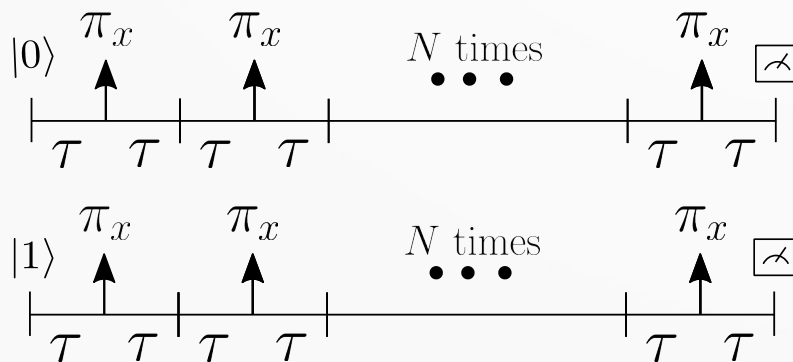
- Initialize NV in $|+\rangle$ and the ^{13}C in the mixed state ρ_m
- Due to conditional interaction between NV and Carbon only at special resonant frequencies, entanglement appears as small peaks
- Simulation coincides perfectly with the theoretical prediction*



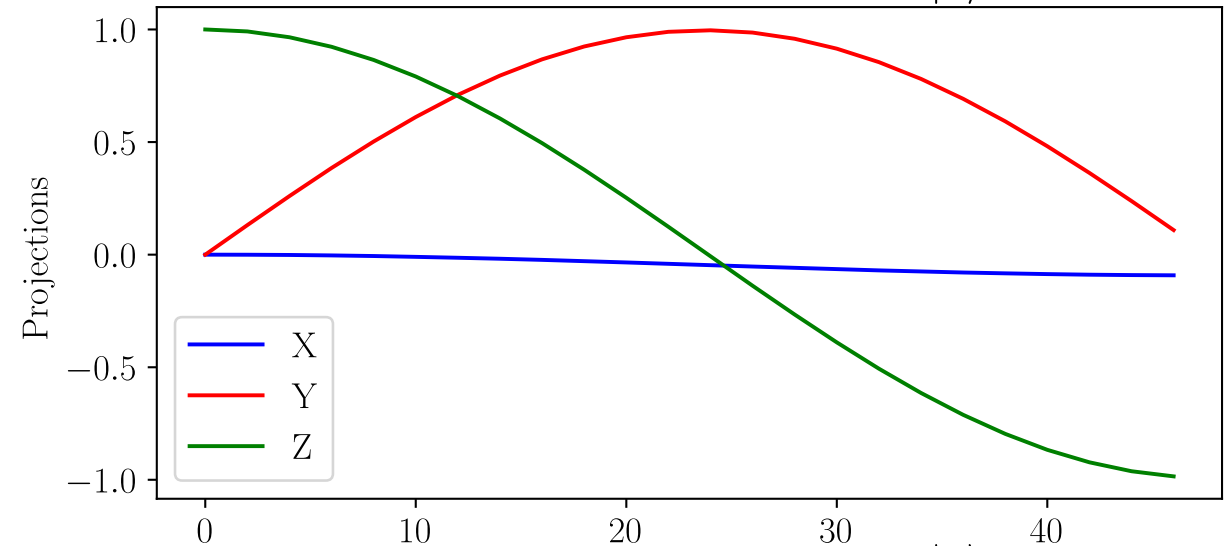
* Taminiau, T. H. et al. "Detection and Control of Individual Nuclear Spins Using a Weakly Coupled Electron Spin." Physical Review Letters 109.13 (2012) <https://doi.org/10.1103/PhysRevLett.109.137602>

Controlled Rotation of Carbons

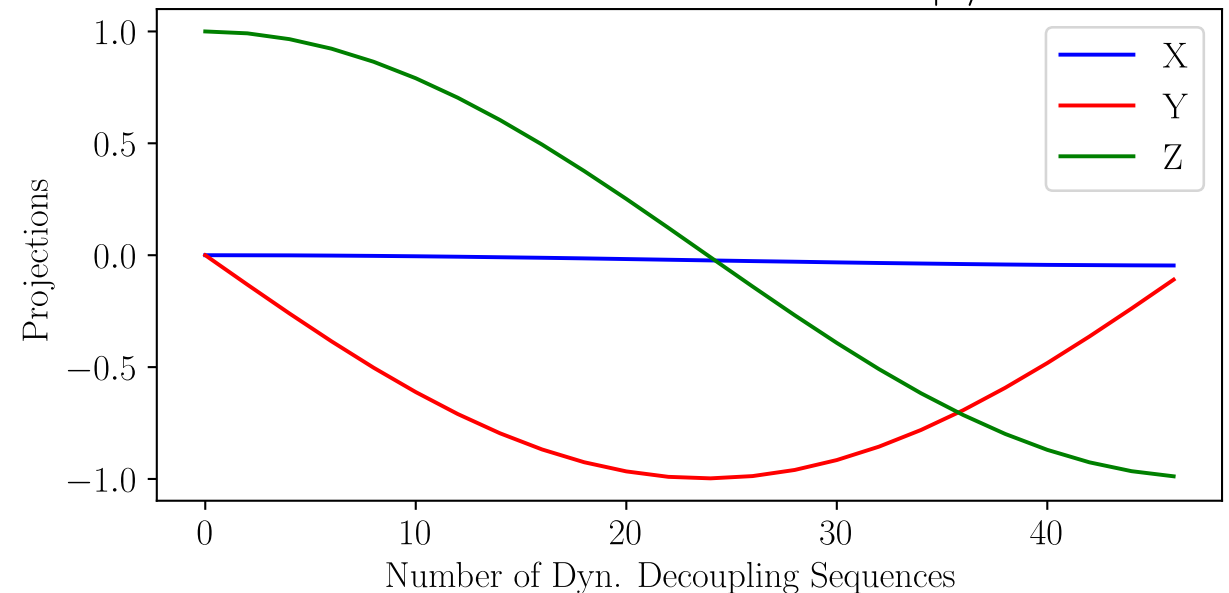
- By flipping the spin state of the NV centre at the resonant precession frequency for a ^{13}C , we can implement coherent rotations
- Direction of rotation is **conditional on initial spin** of the NV
- Even with just ~ 40 sequences, capable of performing π -rotation with over 99% fidelity



Rotation of Carbon with NV in $|0\rangle$ state

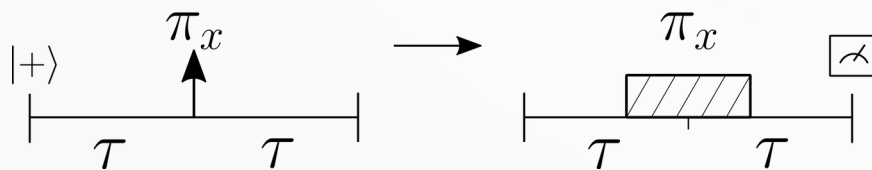


Rotation of Carbon with NV in $|1\rangle$ state

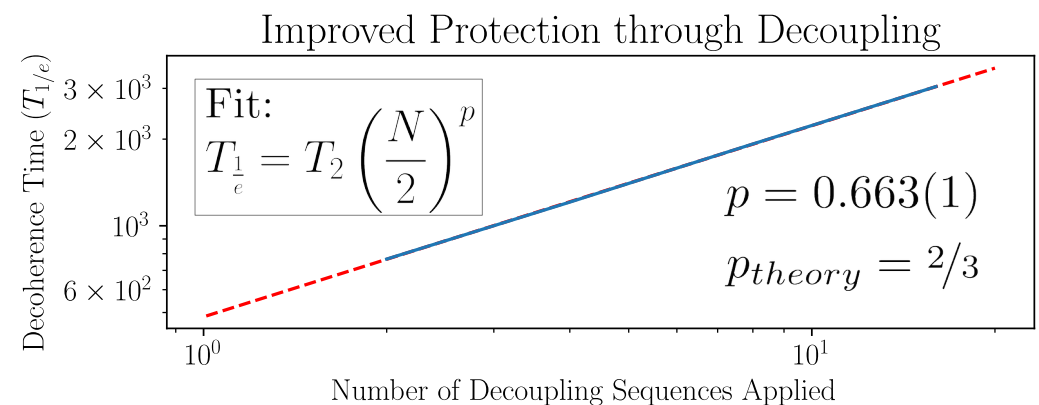
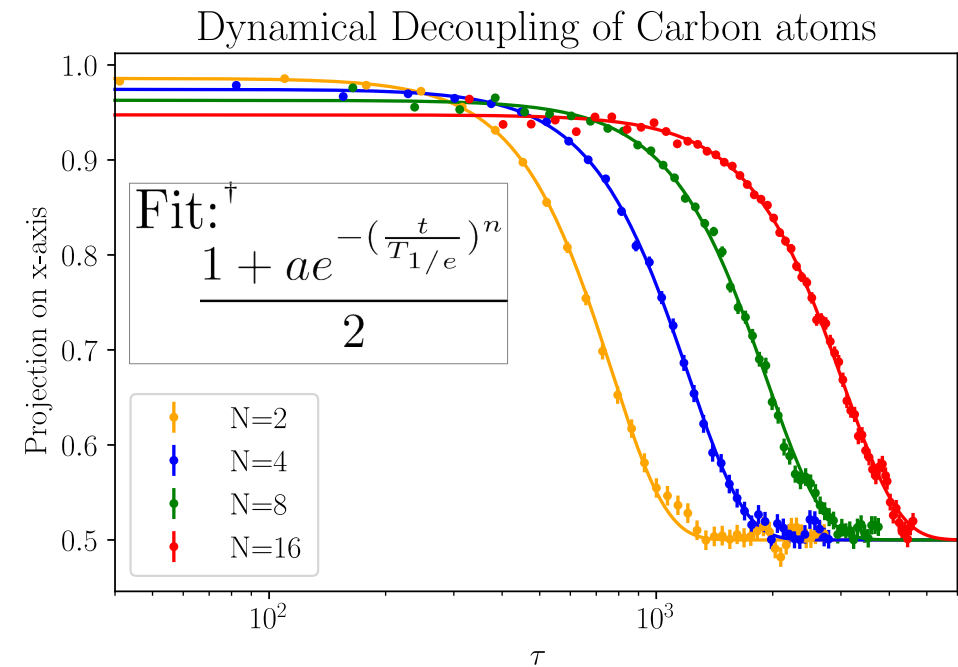


Dynamical Decoupling on Carbon

- To protect the ^{13}C from dephasing noise, we can implement the spin-echo protocol similar to that on the NV
- **Main Difference** - cannot implement π -rotations directly, must be mediated by NV centre, and thus they take much longer



- More rotations lead to better protection (higher fidelity for longer time)



[†] Abobeih, M.H., Cramer, J., Bakker, M.A. et al. One-second coherence for a single electron spin coupled to a multi-qubit nuclear-spin environment. Nat Commun 9, 2552 (2018). <https://doi.org/10.1038/s41467-018-04916-z>

Conclusion

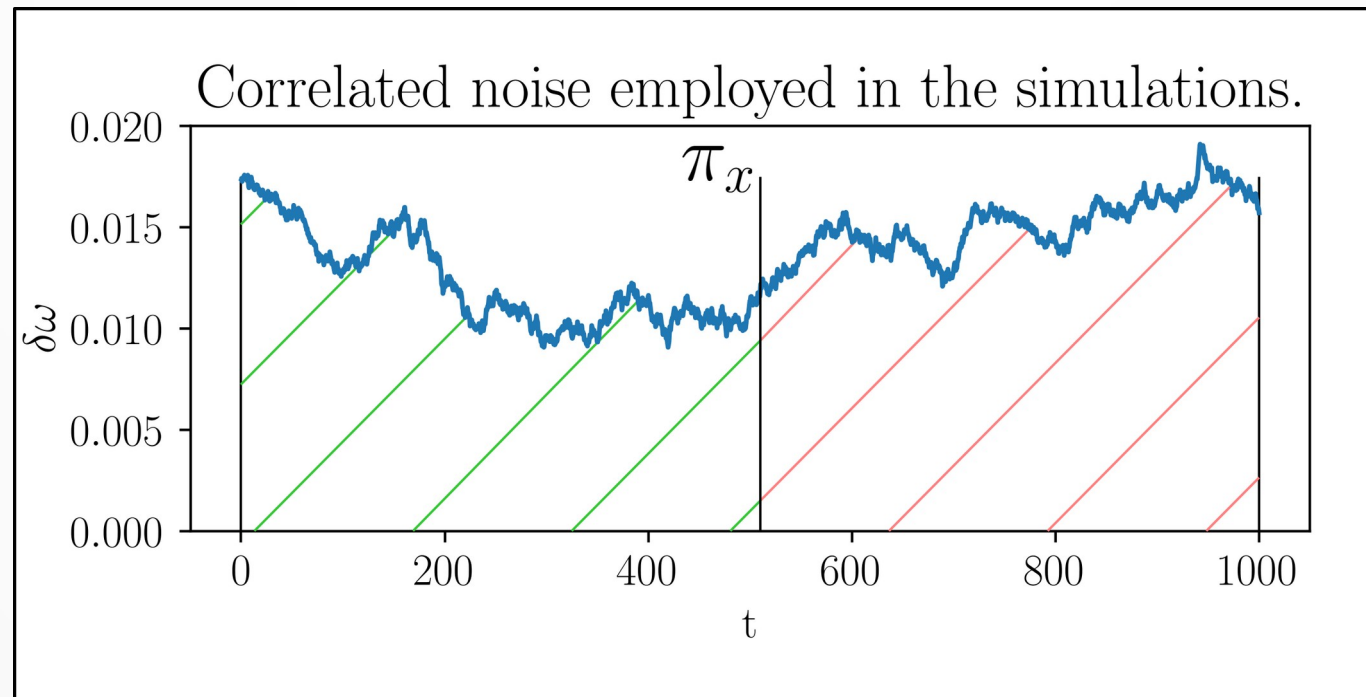
- Decoherence can be mitigated by to some extent using dynamical decoupling sequences, addressing a key challenge in the use of NV centres as nodes for quantum computing
- Our model qualitatively describes the system's evolution under decoupling sequences, obtaining the values for several parameters correctly
- Due to the way the simulation scales with time, simulating for high N values was not feasible. Performance could be improved by implementing the unitary for time evolution directly rather than solving the differential equation shown above
- Future Outlook – probing the effect of different types of decoupling sequences on the spin qubits, as done in recent papers[†]

[†] Wang, Zhi-Hui et al. “Comparison of Dynamical Decoupling Protocols for a Nitrogen-Vacancy Center in Diamond.” Physical Review B 85.15 (2012): <https://doi.org/10.1103/PhysRevB.85.155204>

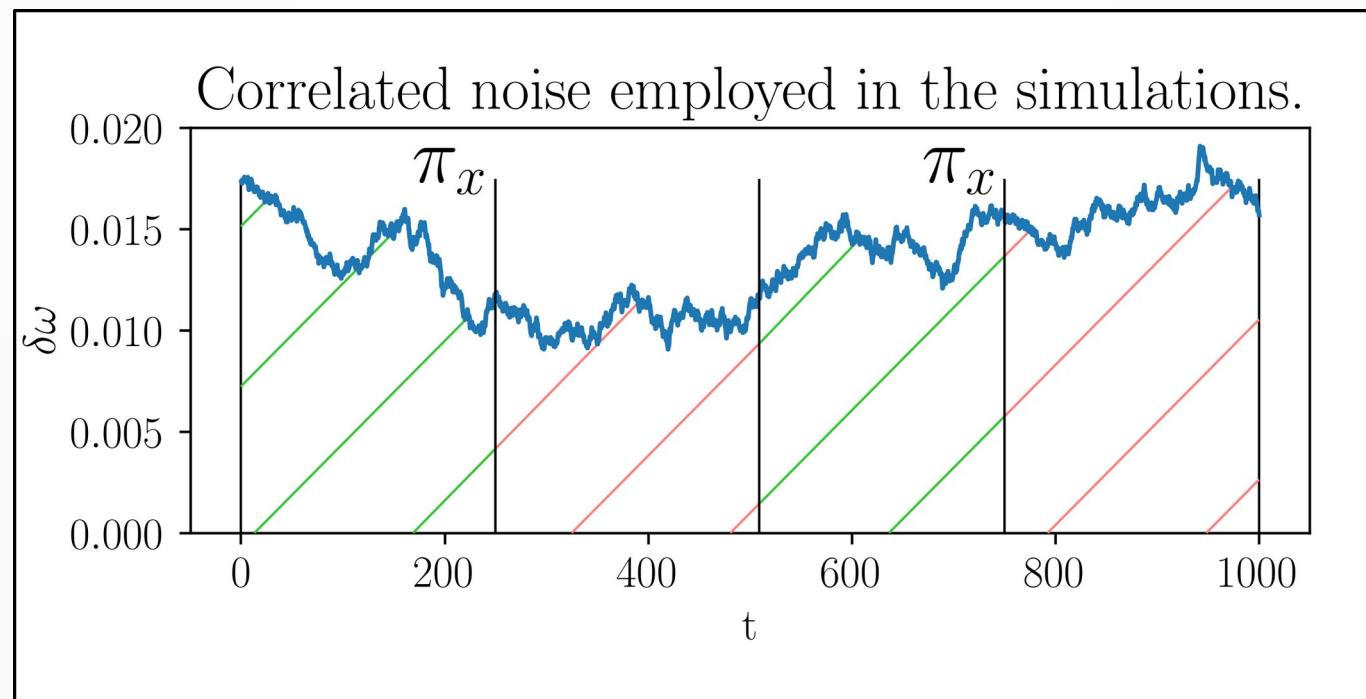
Thank you!

Questions?

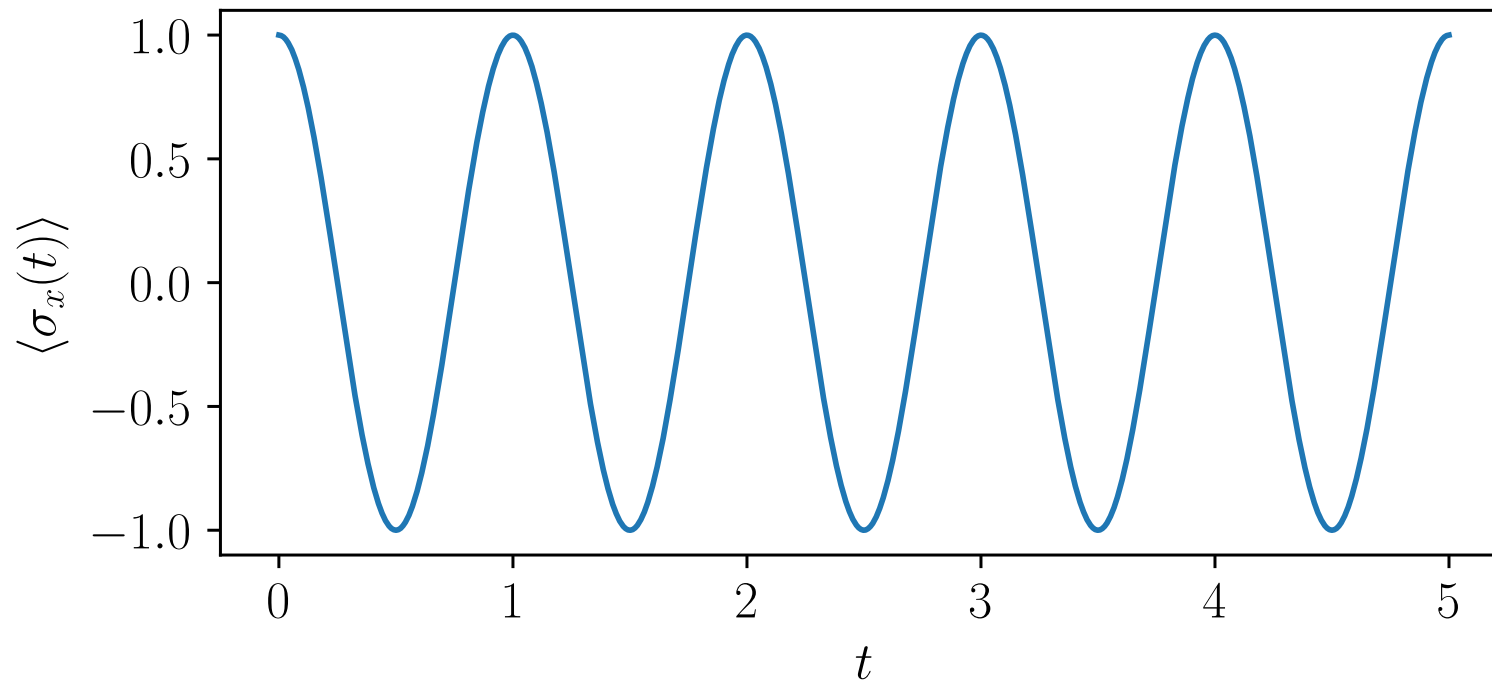
Appendix



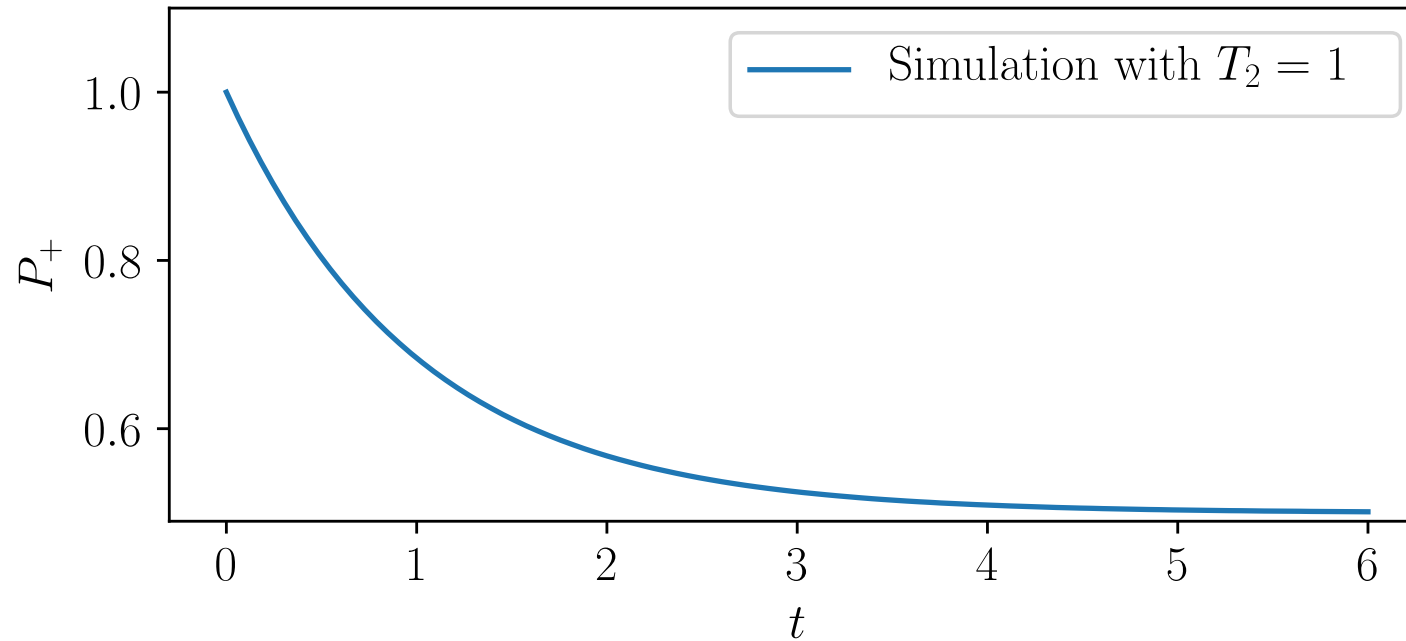
Appendix



Appendix



Appendix



$$\frac{d\rho}{dt} = -\frac{i}{\hbar} [H, \rho] + \sum_i \gamma_i \left(L_i \rho L_i^\dagger - \frac{1}{2} \{ L_i L_i^\dagger, \rho \} \right)$$