FYS4150 - Computational Physics Project 4

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Introduction

Method

Simple 2×2 lattice

We will first consider a 2×2 lattice and use that to test our analytical expression to the numerics. We assume that every spin has two directions, i.e. our states can be either be in spin up state or spin down state (shorthand notation as \uparrow or \downarrow respectively).

The energy of the Ising model, without an external magnetic field, is given by

$$E_i = -J \sum_{\langle kl \rangle}^N s_k s_l$$

Where J>0 is a coupling constant and N is the total number of spins. The symbol < kl> indicates that we only sum over the neighbours only. The values $s_k=\pm 1$ depends on which state it is in. We let $s_{\downarrow}=-1$ and $s_{\uparrow}=1$. We also have the magnetic moment is given as

$$M_i = \sum_{\langle k \rangle}^N s_k$$

Since we have a $2 \times 2 = 4$ lattice, and we have two spin directions, then the number of micro state (or configuration) is $2^4 = 16$. What this means is

Combinations of	(s_1, s_2, s_3, s_4)	$s_j = \{\uparrow, \downarrow\} = \{1, -1\}$	
$(\uparrow,\uparrow,\uparrow,\uparrow)$	$(\uparrow,\uparrow,\uparrow,\downarrow)$	$(\uparrow,\uparrow,\downarrow,\uparrow)$	$(\uparrow,\downarrow,\uparrow,\uparrow)$
$(\downarrow,\uparrow,\uparrow,\uparrow)$	$(\uparrow,\uparrow,\downarrow,\downarrow)$	$(\uparrow,\downarrow,\uparrow,\downarrow)$	$(\downarrow,\uparrow,\uparrow,\downarrow)$
$(\downarrow,\uparrow,\downarrow,\uparrow)$	$(\downarrow,\downarrow,\uparrow,\uparrow)$	$(\uparrow,\downarrow,\downarrow,\uparrow)$	$(\uparrow,\downarrow,\downarrow,\downarrow)$
$(\downarrow,\uparrow,\downarrow,\downarrow)$	$(\downarrow,\downarrow,\uparrow,\downarrow)$	$(\downarrow,\downarrow,\downarrow,\uparrow)$	$(\downarrow,\downarrow,\downarrow,\downarrow)$

Table 1: All the micro states possible.

that our we can have 16 different energies, as well as 16 different magnetic moment, for each respective micro state. Table 1 shows all the possible micro states.

Figure (ADD FIGURE OF GRID HERE) shows a 2×2 lattice. We see that the point s_1 has s_2 and s_3 as the closest neighbours. The energy term will then give the term $(s_1s_2 + s_2s_3)$ for the point s_1 . It does not include s_4 as it is not the closest neighbour to s_1 . We can then continue to add more terms for the three other points, but we need to be careful to not double count connections we already have. Doing this, the energy for each micro state i will be

$$E_i = -J \sum_{s_1 = +1} \sum_{s_2 = +1} \sum_{s_2 = +1} \sum_{s_4 = +1} (s_1 s_2 + s_1 s_3 + s_2 s_4 + s_3 s_4)$$
 (1)

Similarly for the magnetic moment we get when we sum over all micro states

$$M_i = \sum_{s_1 = \pm 1} \sum_{s_2 = \pm 1} \sum_{s_3 = \pm 1} \sum_{s_4 = \pm 1} (s_1 + s_2 + s_3 + s_4)$$
 (2)

Let us now determine both the energies and magnetic moments for all micro states. Using table 1, we can determine equation (1) and (2) to their respective micro state. Table 2 and 3 shows the energies and momenta (using the same combinations in table 1) respectively.

Let us now find an analytical expression for the partition function Z. It is defined as

$$Z = \sum_{i} e^{-\beta E_i}$$

It sums over all micro states i and $\beta = \frac{1}{k_b T}$, with k_b as the Boltzmann constant and T as the temperature. We already have all the energies given in table 2.

$E_i =$			
-4J	0	0	0
0	0	0	0
0	0	0	0
0	0	0	4J

Table 2: Energies for each respective micro state.

$M_i =$			
4	2	2	2
2	0	0	0
0	0	0	-2
-2	-2	-2	-4

Table 3: Magnetic moments for each respective micro state.

Using that, the partition function becomes

$$Z = e^{-4\beta J} + e^{4\beta J} + 14 \times e^0 = 2\cosh(\beta J) + 14$$

Implementation

Results

Conclusion

Reference