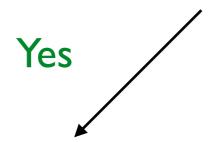
FREEZEING-IN CANNIBAL DARK SECTORS

Andrzej Hryczuk



based on:

E. Cervantes and AH 2407.12104 (to appear in JHEP)



At high temperature DM thermalises, has very high abundance



An efficient depletion is needed (freeze-out)

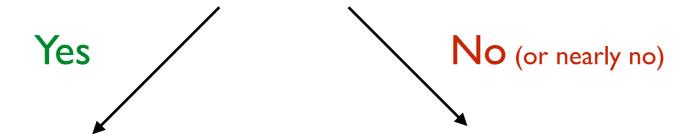


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Is initial population (e.g. from reheating) small?



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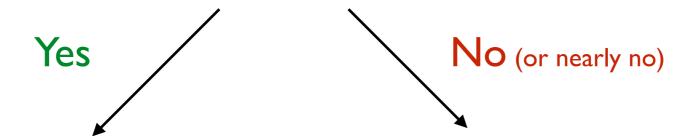
An efficient depletion is needed (freeze-out)

A slow production is needed (freeze-in)

Yes



"Initial condition" dependence



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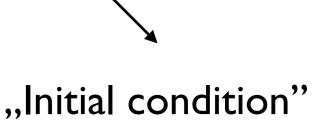
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Yes



dependence

No

AS OF THE STATE OF

"WIMP"



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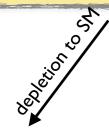
An efficient depletion is needed (freeze-out)

Yes

A slow production is needed (freeze-in)



"Initial condition" dependence



Olye State

"WIMP"

Secluded dark sector



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Yes

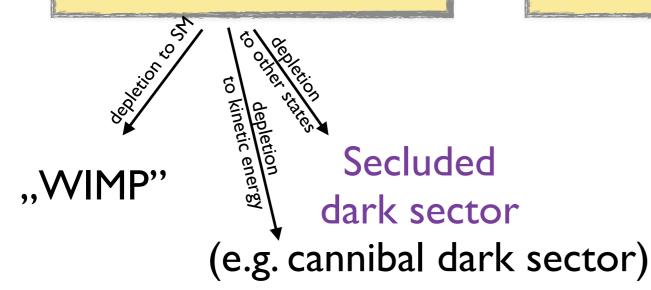
A slow production

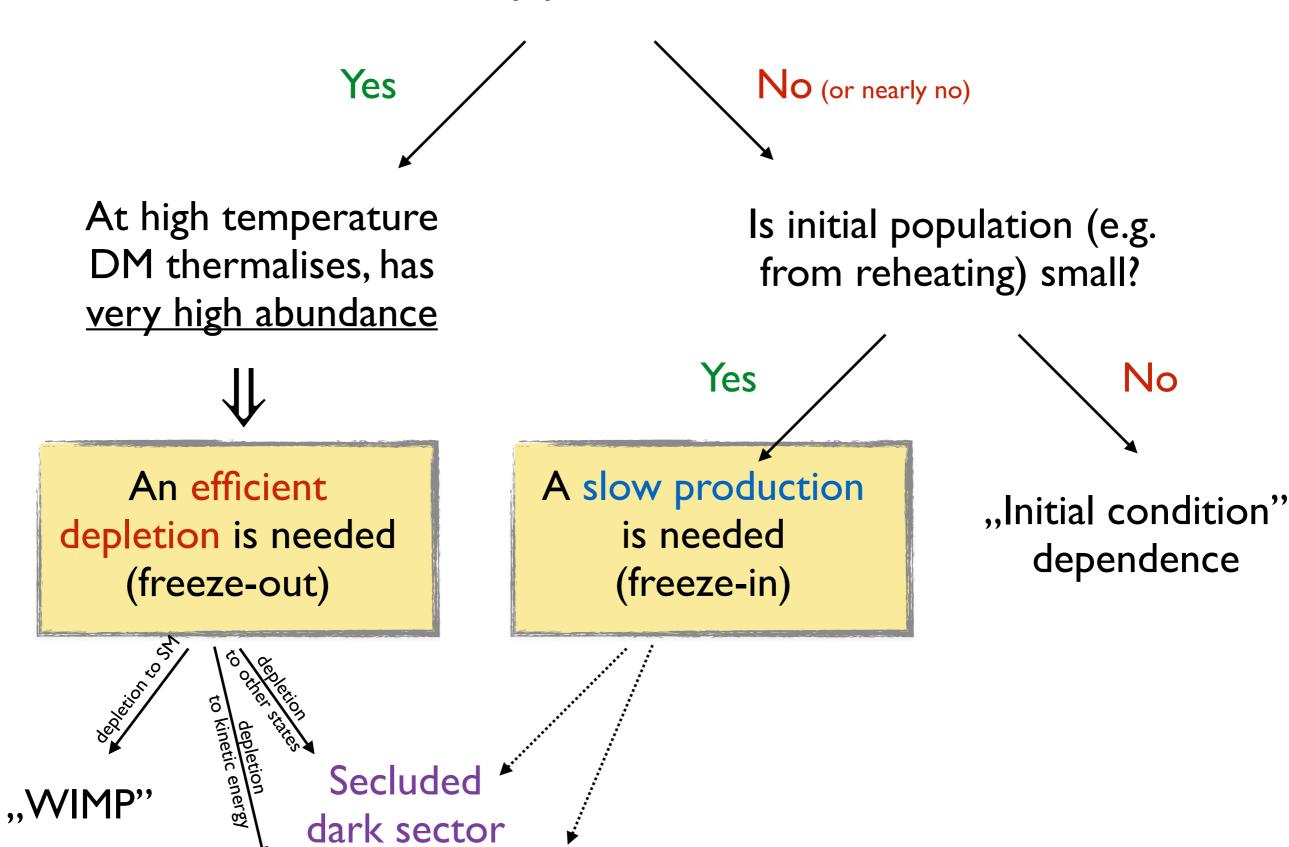
is needed

(freeze-in)

No

"Initial condition" dependence





(e.g. cannibal dark sector)

CANNIBAL DM

Explains depletion of DM solely through self number changing reactions!

SELF-INTERACTING DARK MATTER

ERIC D. CARLSON

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MARIE E. MACHACEK

Department of Physics, Northeastern University, Boston, MA 02115

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LAWRENCE J. HALL

Department of Physics, University of California; and Theoretical Physics Group, Physics Division, Lawrence Berkeley Laboratory, 1 Cyclotron Road, Berkeley, CA 94720

*Received 1992 March 17; accepted 1992 April 20

..., Hochberg et al. '14; ...

Simplest example, scalar ϕ with interactions (no coupling to SM!):

$$\frac{g}{3!}\phi^3 + \frac{\lambda}{4!}\phi^4 \quad \Longrightarrow \quad {}^{\phi} \qquad {}^{\phi}$$

To obtain correct relic abundance: $m_{\phi} \sim \mathcal{O}(10-100\,\mathrm{MeV})$

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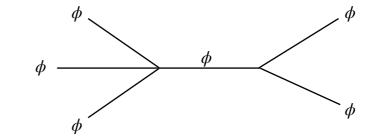
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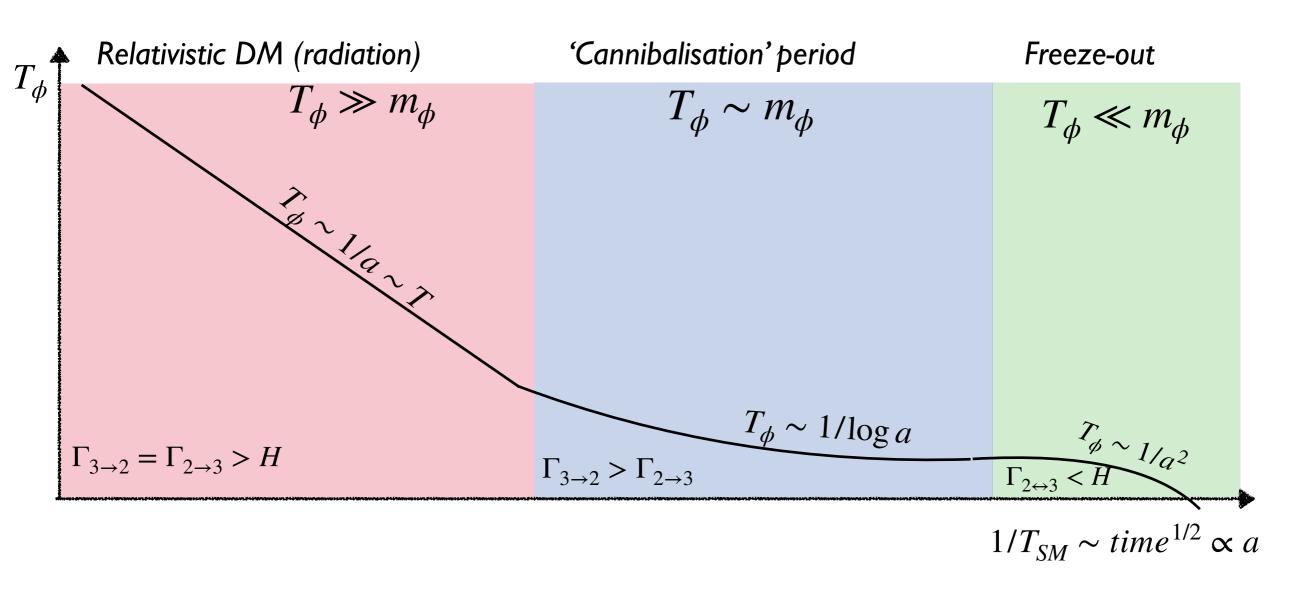
To obtain correct relic abundance: $m_{\phi} \sim \mathcal{O}(10-100\,\mathrm{MeV})$

cannibalisation: mass \rightarrow kinetic energy

expansion: kinetic energy redshifts \rightarrow depletion of E of the dark sector

CANNIBAL FREEZE-OUT

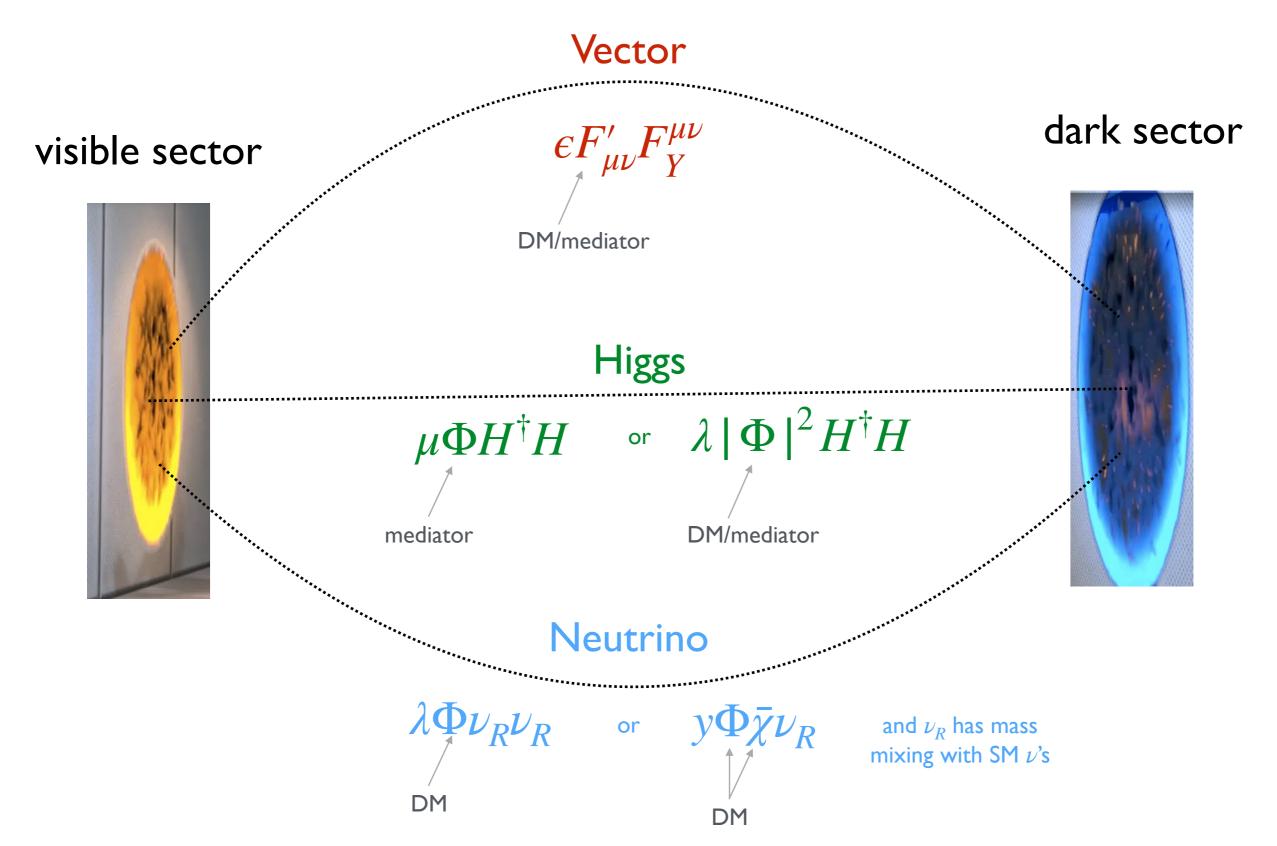
Self-heating of DM makes it warmer, erasing the formation of structures! Tracking its temperature evolution is essential!



- o Initially DM is relativistic $(T_{DM} \gg m_{DM})$;
- O During freeze-out the dark sector uses its rest mass as fuel to keep itself warm;
- o The system decouples and behaves as a non-relativistic gas.

PORTALS





EXAMPLE: SCALAR SINGLET DM

The simplest DM model:

$$\mathcal{L}_{SM} \supset -\frac{m^2}{2} \varphi^2 - \frac{\lambda}{4!} \varphi^4 - \frac{1}{2} \lambda_{h\varphi} \varphi^2 |H|^2$$

$$\langle \varphi \rangle = 0 \& \lambda_{h\varphi} \gtrsim 10^{-4}$$

"WIMP" (thermalizes, undergoes freeze-out)

$$\langle \varphi \rangle = 0 \& \lambda_{h\varphi} \lesssim 10^{-9}$$

FIMP (feebly-interacting, undergoes freeze-in)

$$\langle \varphi \rangle \neq 0$$

Cannibal DM: spontaneous \mathbb{Z}_2 breaking \rightarrow Higgs mixing terms

Hufnagel, Tytgat '22

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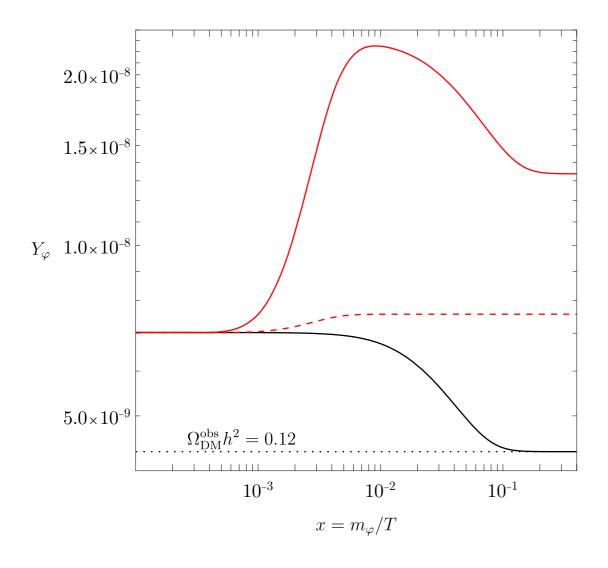
Evolution of number density $Y_{\varphi}:=n_{\varphi}/s$ and 'temperature' $x_{\varphi}:=m_{\varphi}/T_{\varphi}$ governed by set of Boltzmann equations:

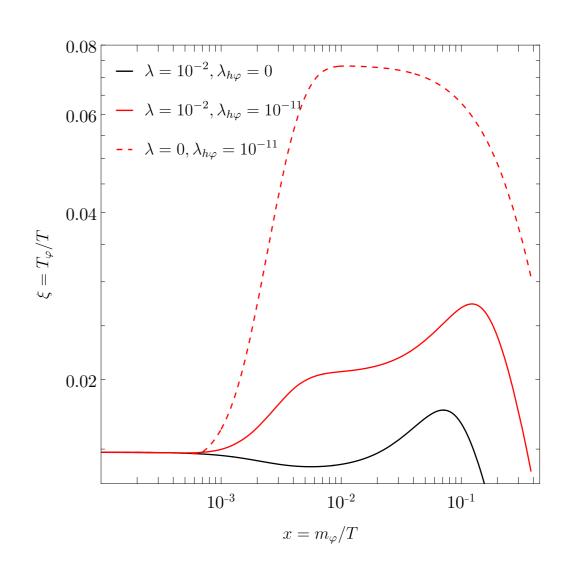
$$\frac{Y'_{\varphi}}{Y_{\varphi}} = \frac{1}{x\,\tilde{H}} \left(\langle C_{h\to\varphi\varphi} \rangle + \langle C_{hh\to\varphi\varphi} \rangle + \langle C_{3\leftrightarrow2} \rangle \right)$$
 Cannibal
$$-\frac{x'_{\varphi}}{x_{\varphi}} = \frac{1}{x\,\tilde{H}} \left(\langle C_{h\to\varphi\varphi} \rangle_2 + \langle C_{hh\to\varphi\varphi} \rangle_2 + \langle C_{\phi h\leftrightarrow\phi h} \rangle_2 + \langle C_{3\leftrightarrow2} \rangle_2 \right) - \frac{Y'_{\varphi}}{Y_{\varphi}} + \frac{H}{x\,\tilde{H}} \frac{\langle p^4/E^3 \rangle}{3T_{\varphi}} + \frac{2s'}{3s}$$

Freeze-in Freeze-in/out El. scattering

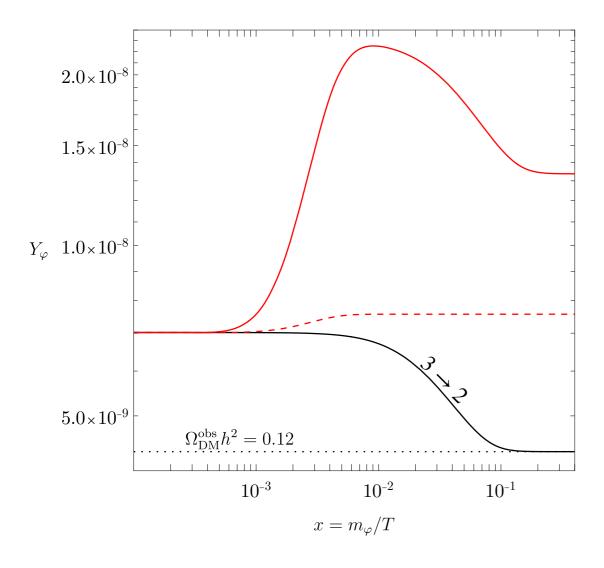
see also Ghosh et al. '22

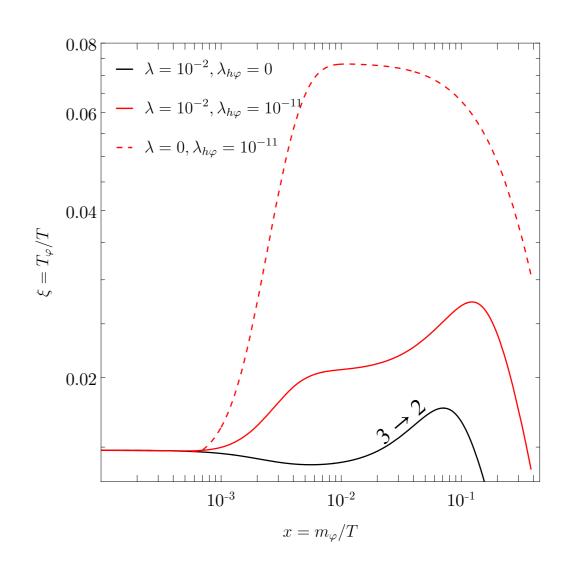
Consider an initial cold, low populated dark sector: $T_{DM}^i = 10^{-2} T_{SM}^i$





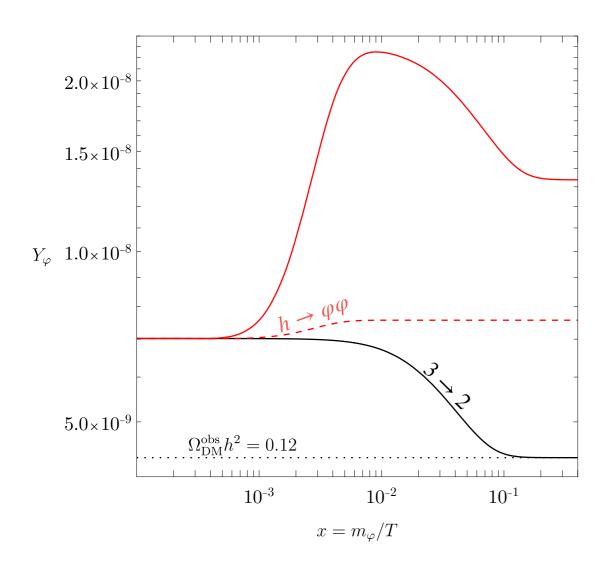
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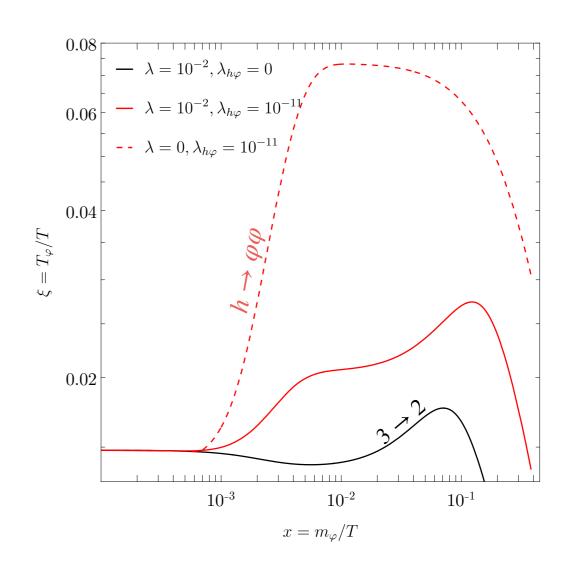




only cannibalization

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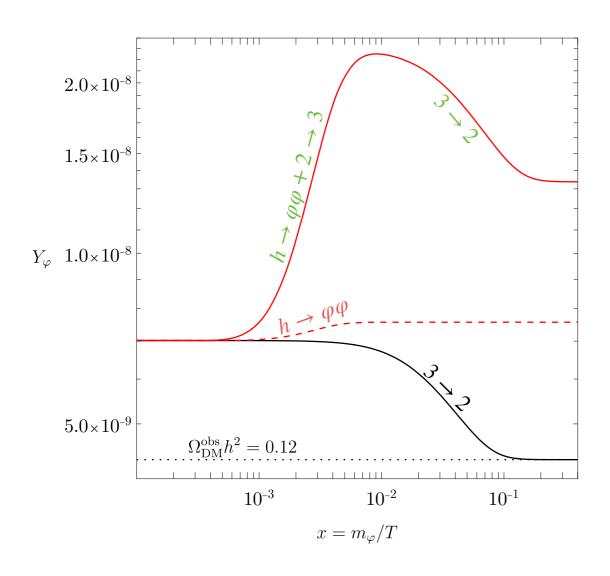


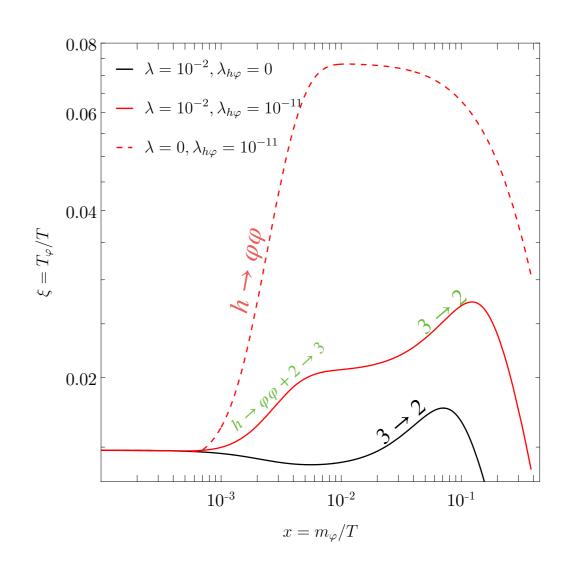


only cannibalization

only freeze-in

Consider an initial cold, low populated dark sector: $T_{DM}^i = 10^{-2} T_{SM}^i$





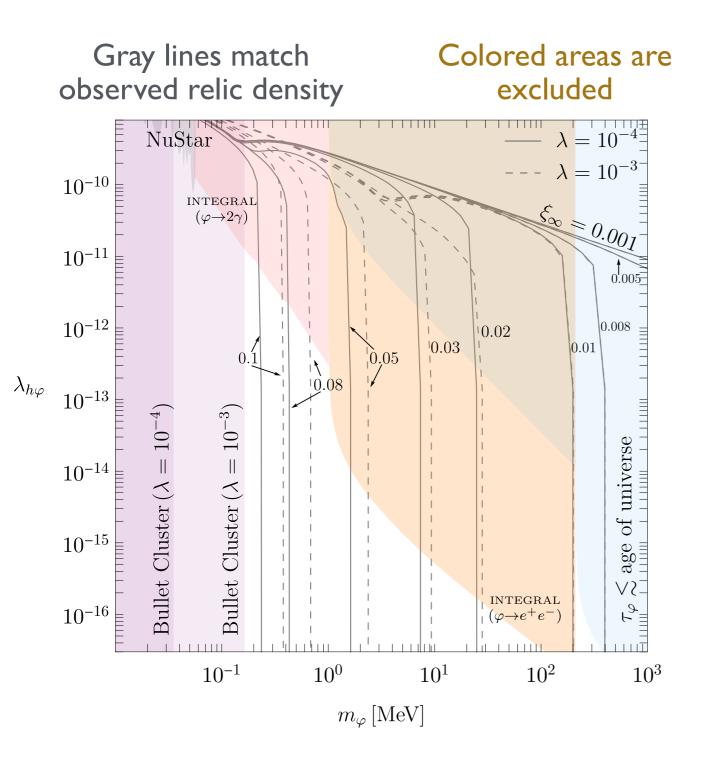
only cannibalization

only freeze-in

cannibal+freeze-in

THIS MODEL HAS AN ISSUE...

 \mathbb{Z}_2 breaking $\Rightarrow \varphi$ unstable with $\tau_{\varphi} \propto 1/\lambda_{h\varphi}^2$



For decaying DM only very long lifetimes are allowed $\Rightarrow \lambda_{h\omega}$ needs to be very small

The impact of our improved treatment negligible in the allowed region

... but still it is a viable model (in the asymmetric reheating scenario)

goes to the scenario of Hufnagel, Tytgat '22

Modified Model

Natural extension:

a complex (\mathbb{Z}_3 stabilised) DM candidate S with a real singlet mediator ϕ

$$\mathcal{L} \supset -\frac{1}{3!}g_s(S^3 + (S^*)^3) - \frac{\lambda_s}{4}|S|^4 - A_{\phi s}\phi|S|^2 - \frac{\lambda_{\phi s}}{2}\phi^2|S|^2 - B_{\phi h}\phi|H|^2$$
DM self interactions
thermalisation between ϕ
and S

Portal to the visible sector $B_{\phi h}\phi \left| H \right|^2$ induces mixing between ϕ and the Higgs post-EWPT:

$$\phi \rightarrow \phi + \theta h$$
 $h \rightarrow h - \theta \phi$
 ϕ couples to matter in a Higgs-like way

Such a model exhibits new dynamics + provides chance of detection

BOLTZMANN EQUATIONS

DM evolution

Freeze-in

DM-mediator interactions

DM self-interactions

$$\frac{Y'_S}{Y_S} = \frac{1}{x\,\tilde{H}} \left(\langle C_{h\to\phi SS^*} \rangle + \langle C_{h\to SS^*} \rangle + \langle C_{\phi\phi \leftrightarrow SS^*} \rangle + \langle C_{3\leftrightarrow 2} \rangle \right),$$

$$-\frac{x'_S}{x_S} = \frac{1}{x\,\tilde{H}} \left(\langle C_{h\to\phi SS^*} \rangle_2 + \langle C_{h\to SS^*} \rangle_2 + \langle C_{\phi S\leftrightarrow\phi S} \rangle_2 + \langle C_{3\leftrightarrow 2} \rangle_2 \right) - \frac{Y'_S}{Y_S} + \frac{H}{x\,\tilde{H}} \frac{\langle p^4/E^3 \rangle}{3T_S} + \frac{2s'}{3s}$$

Mediator evolution

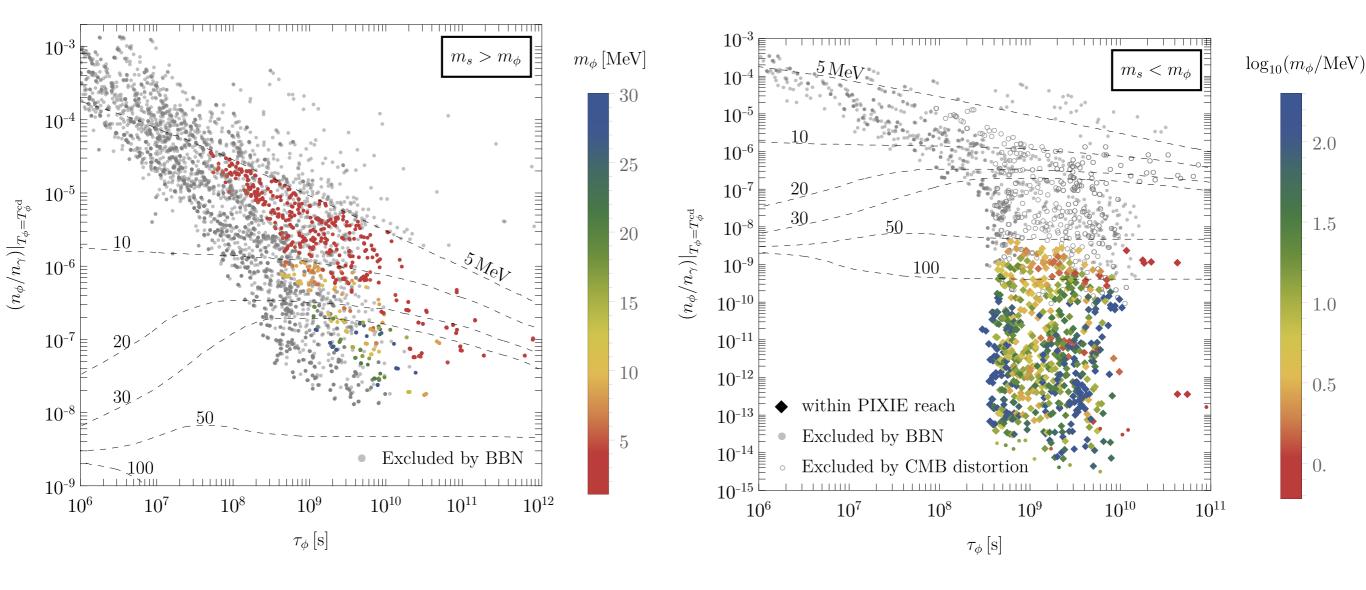
$$\frac{Y'_{\phi}}{Y_{\phi}} = \frac{1}{x \,\tilde{H}} \left(\langle C_{h \to \phi SS^*} \rangle + \langle C_{\text{sm sm} \to \text{sm} \phi} \rangle + \langle C_{\phi \phi \leftrightarrow SS^*} \rangle \right) ,$$

$$-\frac{x'_{\phi}}{x_{\phi}} = \frac{1}{x\,\tilde{H}}\left(\langle C_{h\to\phi SS^*}\rangle_2 + \langle C_{\text{sm sm}\to\text{sm}\,\phi}\rangle_2 + \langle C_{\phi S\leftrightarrow\phi S}\rangle_2\right) - \frac{Y'_{\phi}}{Y_{\phi}} + \frac{H}{x\,\tilde{H}}\frac{\langle p^4/E^3\rangle}{3T_{\phi}} + \frac{2s'}{3s}.$$

BBN & CMB CONSTRAINTS

Mediator is decaying to SM:

different phenomenology for two mass hierarchies between S (the DM) and ϕ

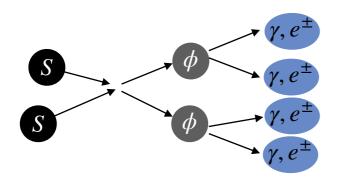


Resulting constraints are significant, but many viable points found in a scan & some with prospects of CMB distortion in reach of PIXIE

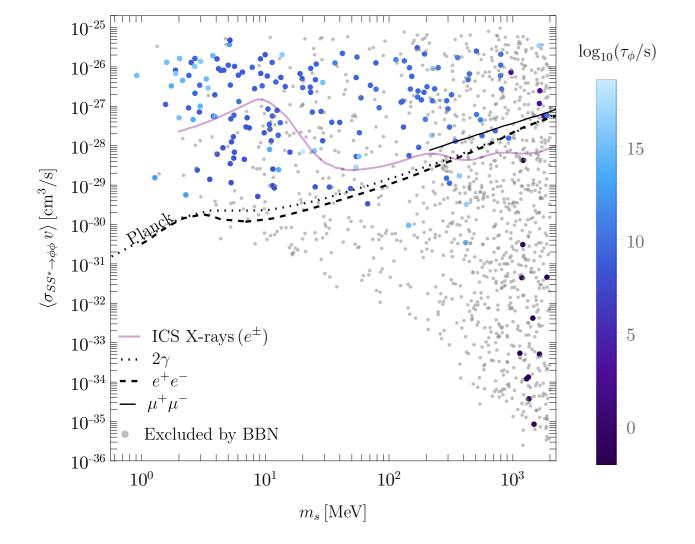
(The Primordial Inflation Explorer)

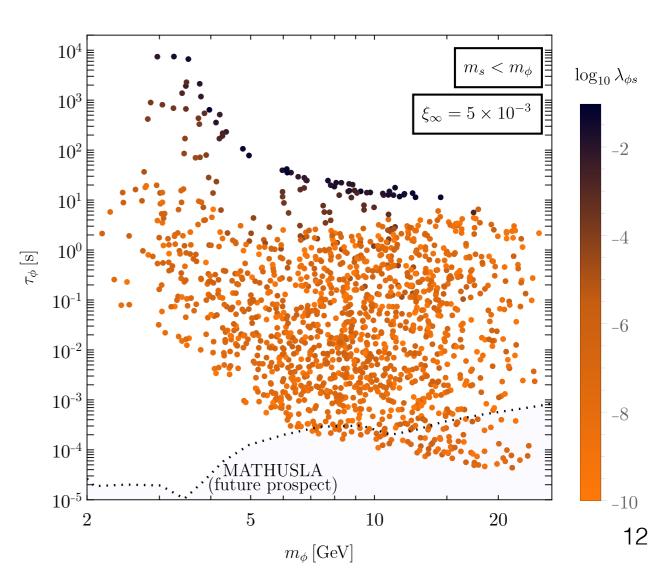
DM & MEDIATOR DETECTION

 $m_S > m_\phi$: indirect DM detection feasible



 $m_S < m_\phi$: weaker BBN limits allow for points that are potentially testable in LLP searches





CONCLUSIONS

- I. Frozen-in Cannibal Dark Matter is a viable scenario, naturally avoiding the large scale structure formation limits plaguing cannibalising DM.
- 2. It has unique evolution in the Early Universe & potentially detectable signals in indirect searches.
- 3. Temperature (and momentum distribution) can have a non-trivial impact in such scenarios and a need to be studied carefully.

