

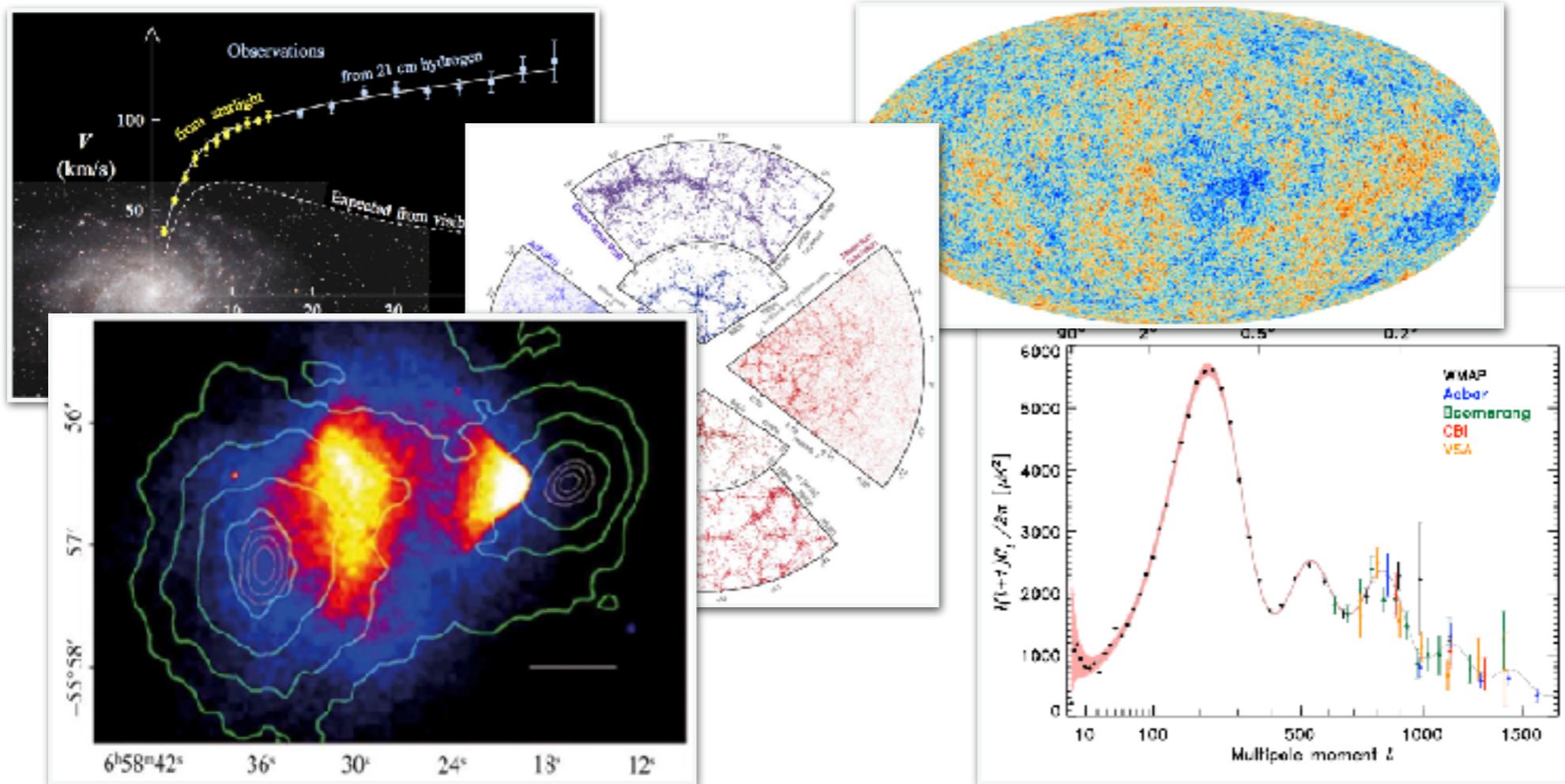
PARTICLE DARK MATTER AT A CROSSROADS

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University of Oslo



DARK MATTER

I don't think there is any need for convincing you that **DM exists...**



⇒ Evidence
on all scales!

... but perhaps I should argue why particle DM



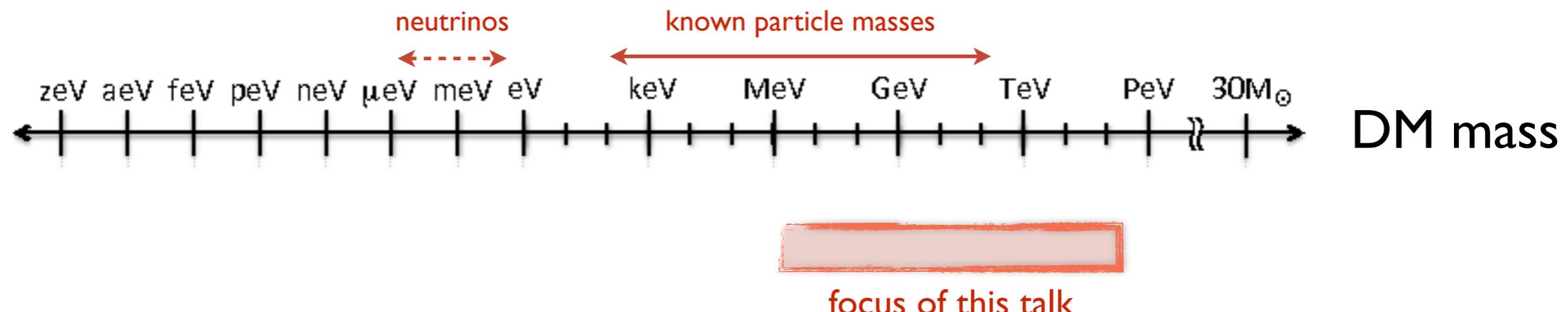
PARTICLE PHYSICIST'S PERSPECTIVE:

We know that the Standard Model (of particle physics) is not complete

its extension could *in principle* be extremely **minimal**... but it is far more likely that there are (many?) **new particles** we do not know yet

it is quite possible that some of them are stable and then they are **a dark matter**

if so it is very natural to expect that they are also **the** dark matter



OUTLINE

1. Vanilla WIMP and why is (was?) it so attractive?
 - hope for new physics around the weak scale
 - clear path for detection techniques
 - thermal production mechanism
2. Going heavier - is it all bad?
 - possible avenues
 - new challenges at the TeV scale
3. Going lighter - is it all contrived?
 - confronting experiments
 - new input from astrophysics & cosmology
4. A new hope?

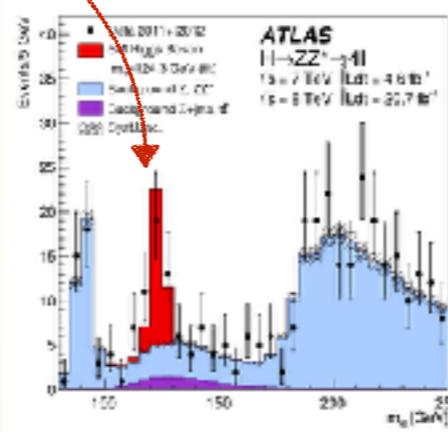
CHAPTER #1

VANILLA WIMP

NEW PHYSICS

(IS ALWAYS) AROUND THE CORNER

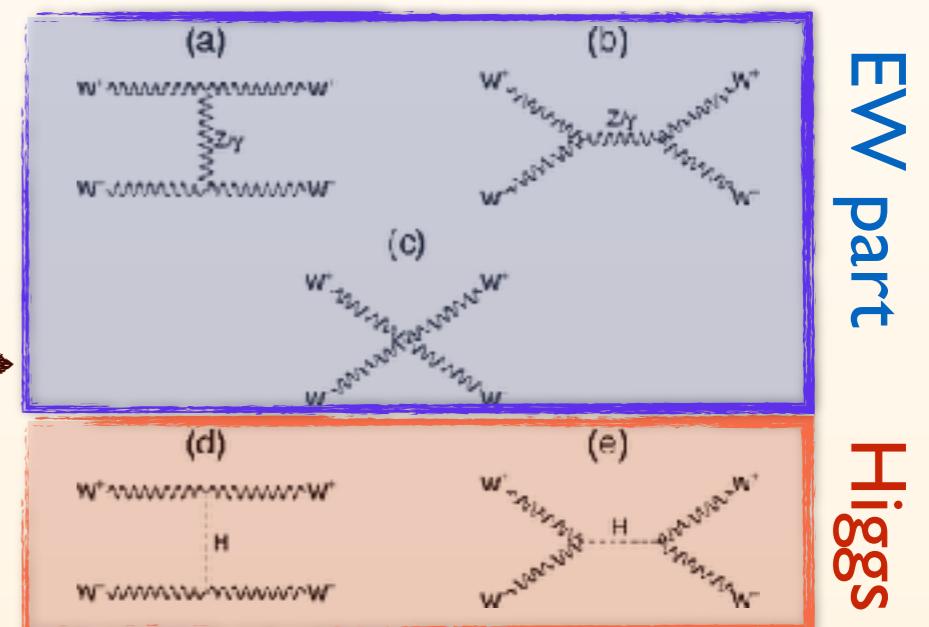
July 2012 - the Higgs boson



since then:



but then we knew sth is there: vide so-called
unitarization of the WW scattering cross section



Now, after the Higgs was found - **The Hierarchy Problem**

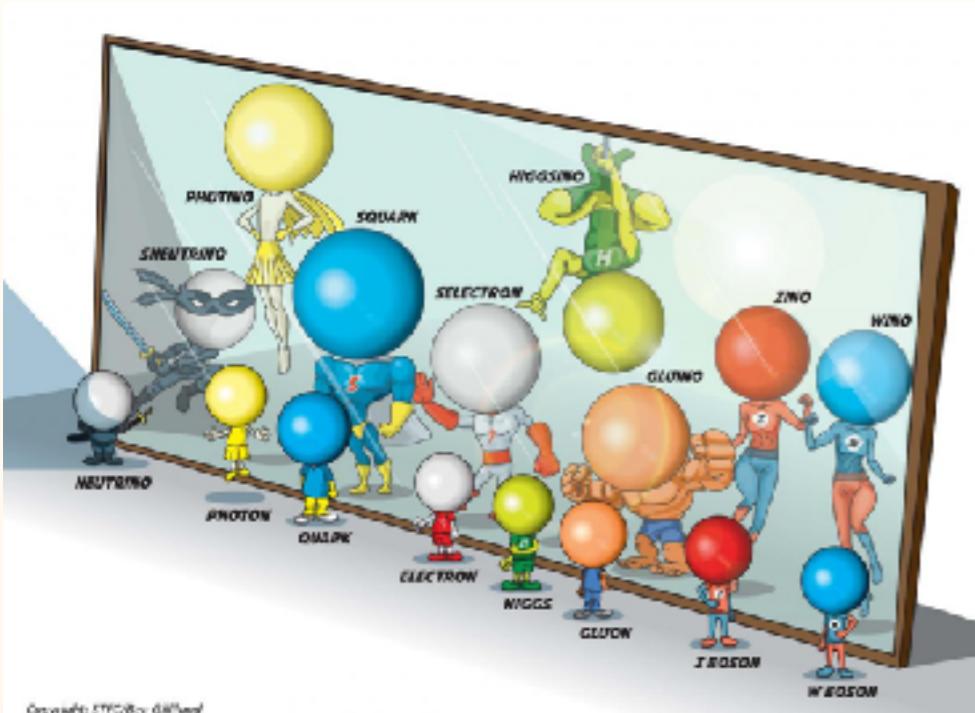
$$\Delta m_h^2 = \frac{3\Lambda^2}{8\pi^2 v^2} [4m_t^2 - 2m_W^2 - m_Z^2 - m_h^2] + \mathcal{O}\left(\log \frac{\Lambda}{v}\right)$$

or in other words: why is the **Higgs boson** so light?

SUPERSYMMETRY

(STILL) BEST MOTIVATED FRAMEWORK BEYOND SM

What people think about SUSY :



SUSY features:

- great simplification of the theory
- Coleman-Mandula theorem
- elegant solution to Hierarchy Problem
- needed by String Theory
- coupling unification
- DM candidate for free

What SUSY *really* is:

$$\begin{aligned}\{Q_\alpha, Q_\beta\} &= \{\bar{Q}_{\dot{\alpha}}, \bar{Q}_{\dot{\beta}}\} = 0 , \\ \{Q_\alpha, \bar{Q}_{\dot{\beta}}\} &= 2(\sigma^\mu)_{\alpha\dot{\beta}} P_\mu , \\ [Q_\alpha, P_\mu] &= [\bar{Q}_{\dot{\alpha}}, P_\mu] = 0 , \\ [Q_\alpha, M_{\mu\nu}] &= -\frac{1}{2}(\sigma_{\mu\nu})_\alpha{}^\beta Q_\beta , \\ [\bar{Q}_{\dot{\alpha}}, M_{\mu\nu}] &= -\frac{1}{2}(\bar{\sigma}_{\mu\nu})_{\dot{\alpha}}{}^{\dot{\beta}} \bar{Q}_{\dot{\beta}} .\end{aligned}$$

$$\mathcal{L}_{SUSY} = \int d^2\theta W(\{\Phi_i\}) + \frac{1}{16g^2} \int d^2\theta \text{Tr}(W_\alpha W^\alpha) + \int d^2\theta d^2\bar{\theta} \bar{\Phi} e^{2gV} \Phi + \text{h.c.}$$

SUSY bugs:

- hasn't been found yet...
(however this bug might turn out to be a feature...)

SUPERSYMMETRY AND THE DARK MATTER

Of course, our everyday world is **not** supersymmetric - SUSY has to be broken - if it happens roughly at energy scales in reach of the LHC

→ we call it "low scale" SUSY

Lepton & Baryon conservation can be ensured if one impose *R-parity*:

$$P_R = (-1)^{3(B-L)+2s}$$

i.e. superparticles : $P_R = -1$, particles $P_R = +1$.

A corollary: lightest superparticle (LSP) is automatically stable!

Moreover it can have properties required of dark matter particle and even better - one that interacts with the SM strongly enough to potentially give signals

↑ if "low scale" SUSY
8

WIMP

WEAKLY INTERACTING AND MASSIVE

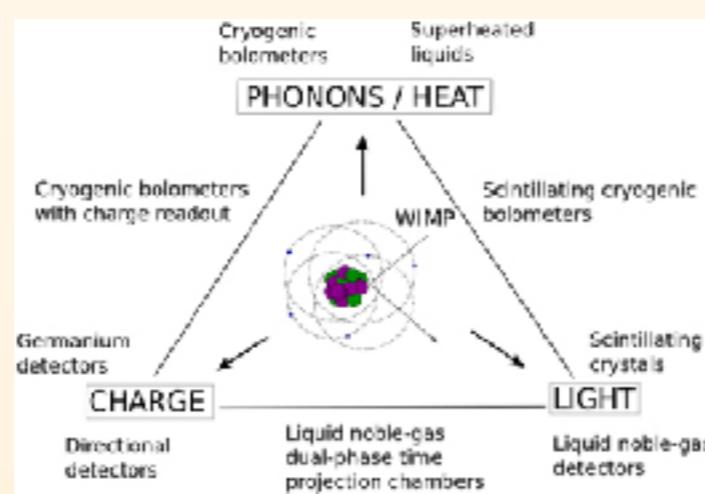
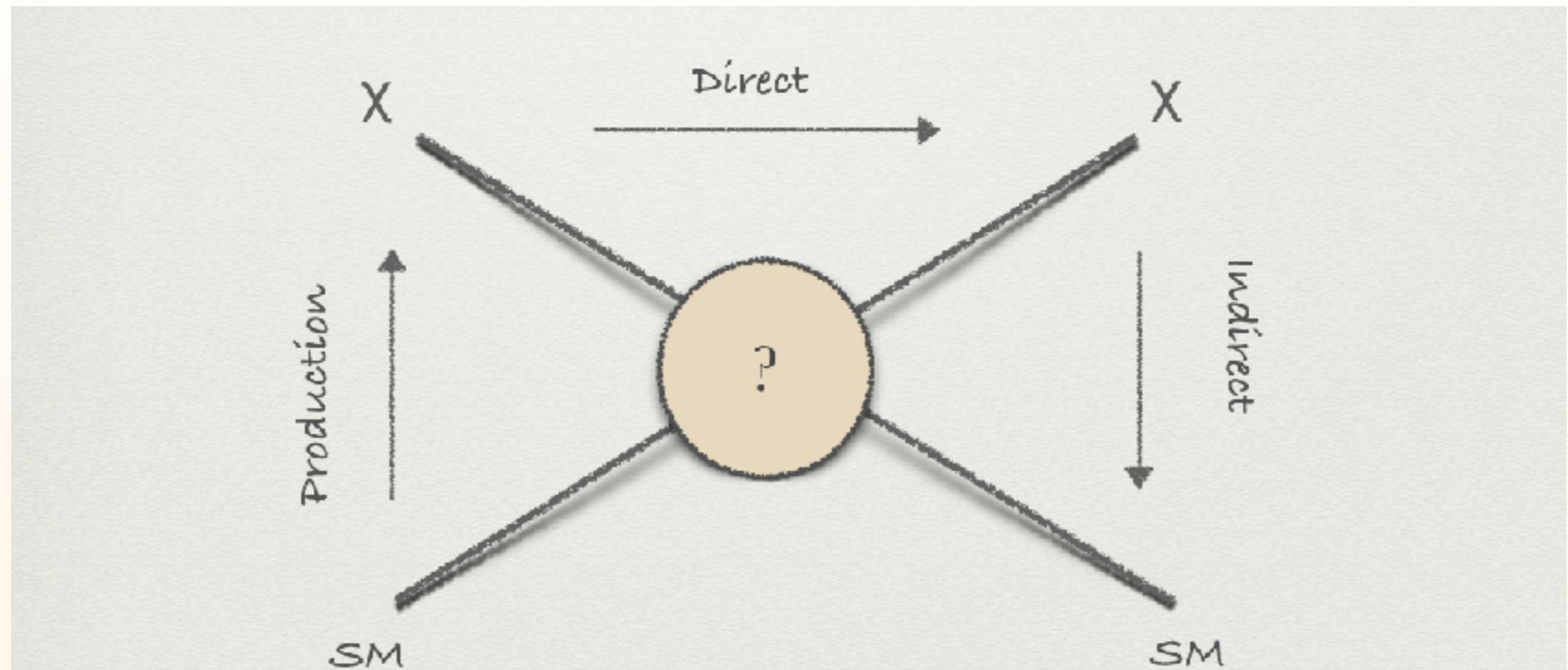
In a weak sense:

DM **cannot interact too strongly**, as we would see it and has to have a **mass** to contribute to observed gravitational potential (now and during the structure formation)

In a strong sense:

interacting through **SM weak interactions** and (therefore) not too light

WIMP DETECTION



THERMAL RELIC DENSITY AND THE „WIMP MIRACLE”

Dark matter could be created in many different ways...

...but every massive particle with not-too-weak interactions with the SM will be produced thermally, with relic abundance:

Lee, Weinberg '77; + others

$$\Omega_\chi h^2 \approx 0.1 \frac{3 \times 10^{-26} \text{cm}^3 \text{s}^{-1}}{\langle \sigma v \rangle}$$

This is dubbed the **WIMP miracle** because it coincidentally seem to point at the same energy scale as suggested by the Hierarchy Problem



MOTIVATION

THERMAL RELIC DENSITY

Theory:

I. Natural

Comes out **automatically** from the expansion of the Universe

Naturally leads to **cold DM**

II. Predictive

No dependence on **initial conditions**

Fixes coupling(s) \Rightarrow signal in DD, ID & LHC

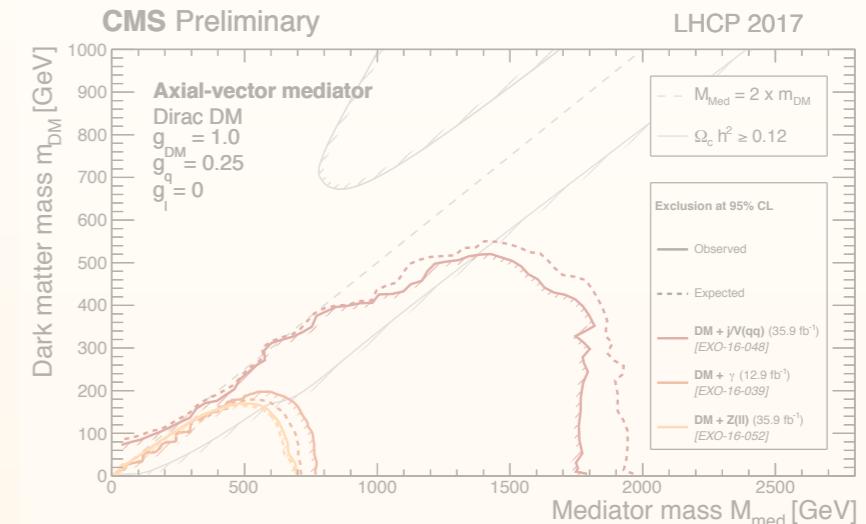
III. It is not optional

Overabundance constraint

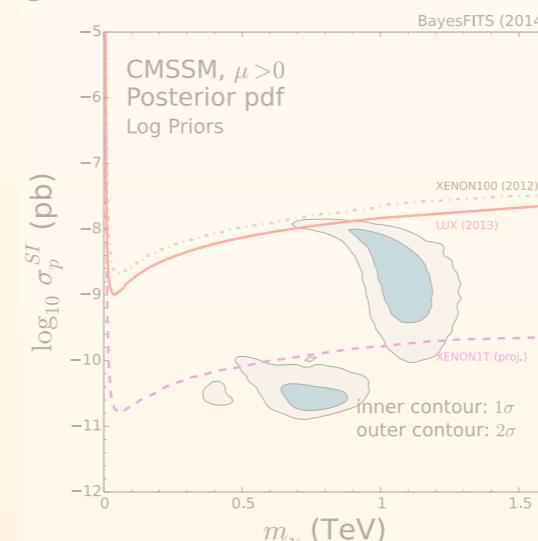
To avoid it one needs **quite significant deviations** from standard cosmology

Experiment:

...as a constraint:



...as a target:



...as a pin:

When a **dark matter signal** is (finally) found:
relic abundance can **pin-point** the
particle physics interpretation

"(...) besides the Higgs boson mass measurement and LHC direct bounds, the constraint showing **by far the strongest impact** on the parameter space of the MSSM is the **relic density**"

Roszkowski et al. '14

THERMAL RELIC DENSITY

STANDARD APPROACH

Boltzmann equation for $f_\chi(p)$:

$$E (\partial_t - H \vec{p} \cdot \nabla_{\vec{p}}) f_\chi = \mathcal{C}[f_\chi]$$

integrate over p
(i.e. take 0th moment)

$$\frac{dn_\chi}{dt} + 3Hn_\chi = -\langle \sigma_{\chi\bar{\chi} \rightarrow ij} \sigma_{\text{rel}} \rangle^{\text{eq}} (n_\chi n_{\bar{\chi}} - n_\chi^{\text{eq}} n_{\bar{\chi}}^{\text{eq}})$$

where the thermally averaged cross section:

$$\langle \sigma_{\chi\bar{\chi} \rightarrow ij} v_{\text{rel}} \rangle^{\text{eq}} = -\frac{h_\chi^2}{n_\chi^{\text{eq}} n_{\bar{\chi}}^{\text{eq}}} \int \frac{d^3 \vec{p}_\chi}{(2\pi)^3} \frac{d^3 \vec{p}_{\bar{\chi}}}{(2\pi)^3} \sigma_{\chi\bar{\chi} \rightarrow ij} v_{\text{rel}} f_\chi^{\text{eq}} f_{\bar{\chi}}^{\text{eq}}$$

Critical assumption:
kinetic equilibrium at chemical decoupling

$$f_\chi \sim a(\mu) f_\chi^{\text{eq}}$$

*assumptions for using Boltzmann eq:
classical limit, molecular chaos,...

...for derivation from thermal QFT
see e.g., 1409.3049

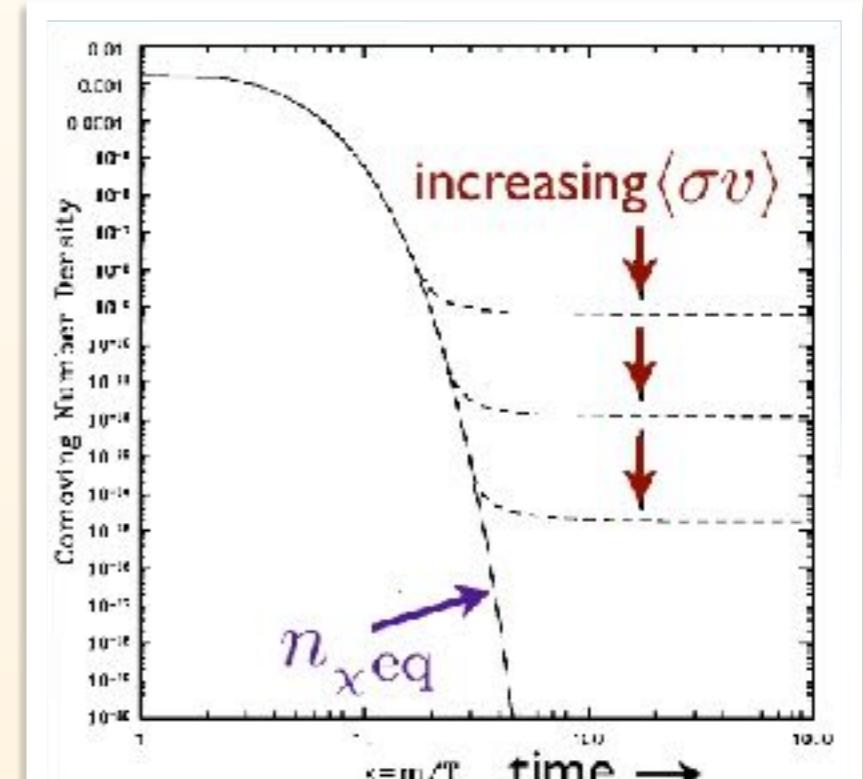
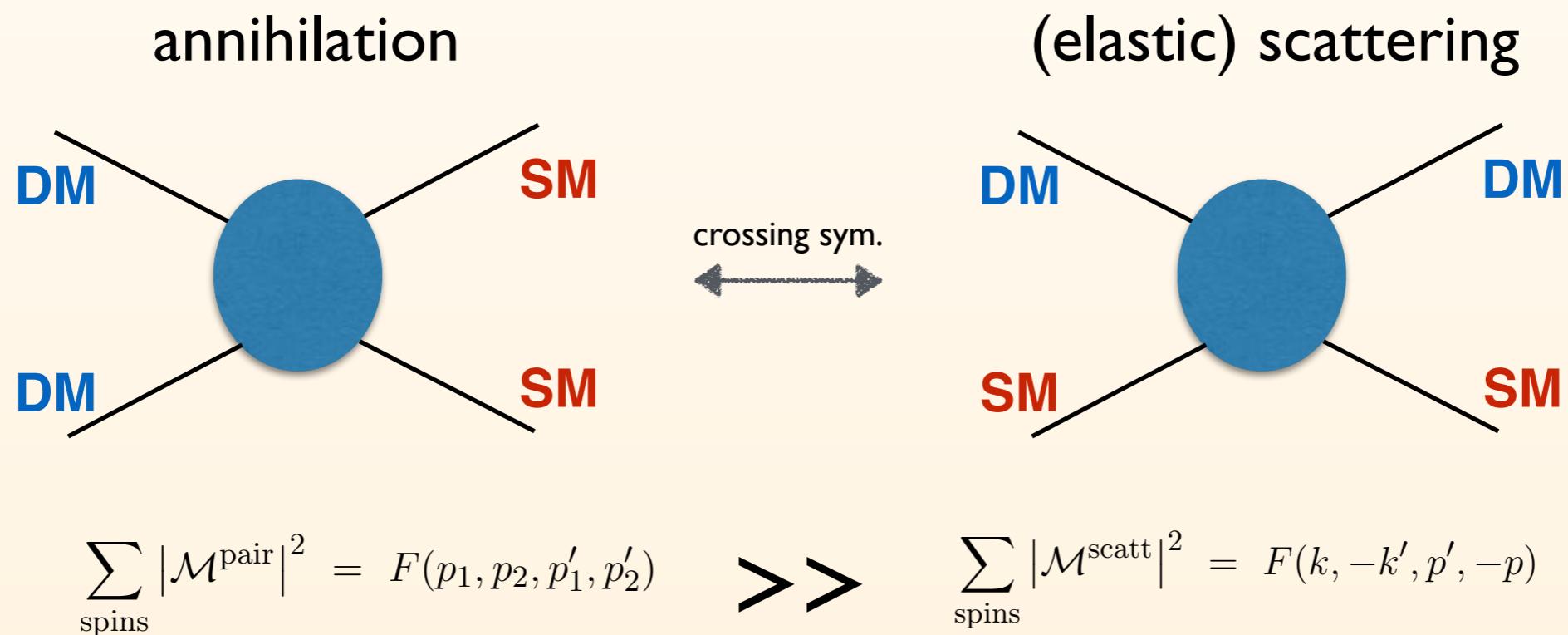


Fig.: Jungman, Kamionkowski & Griest, PR'96

INTERLUDE: WHAT IF KD VERY EARLY?

Recall: in *standard* thermal relic density calculation:

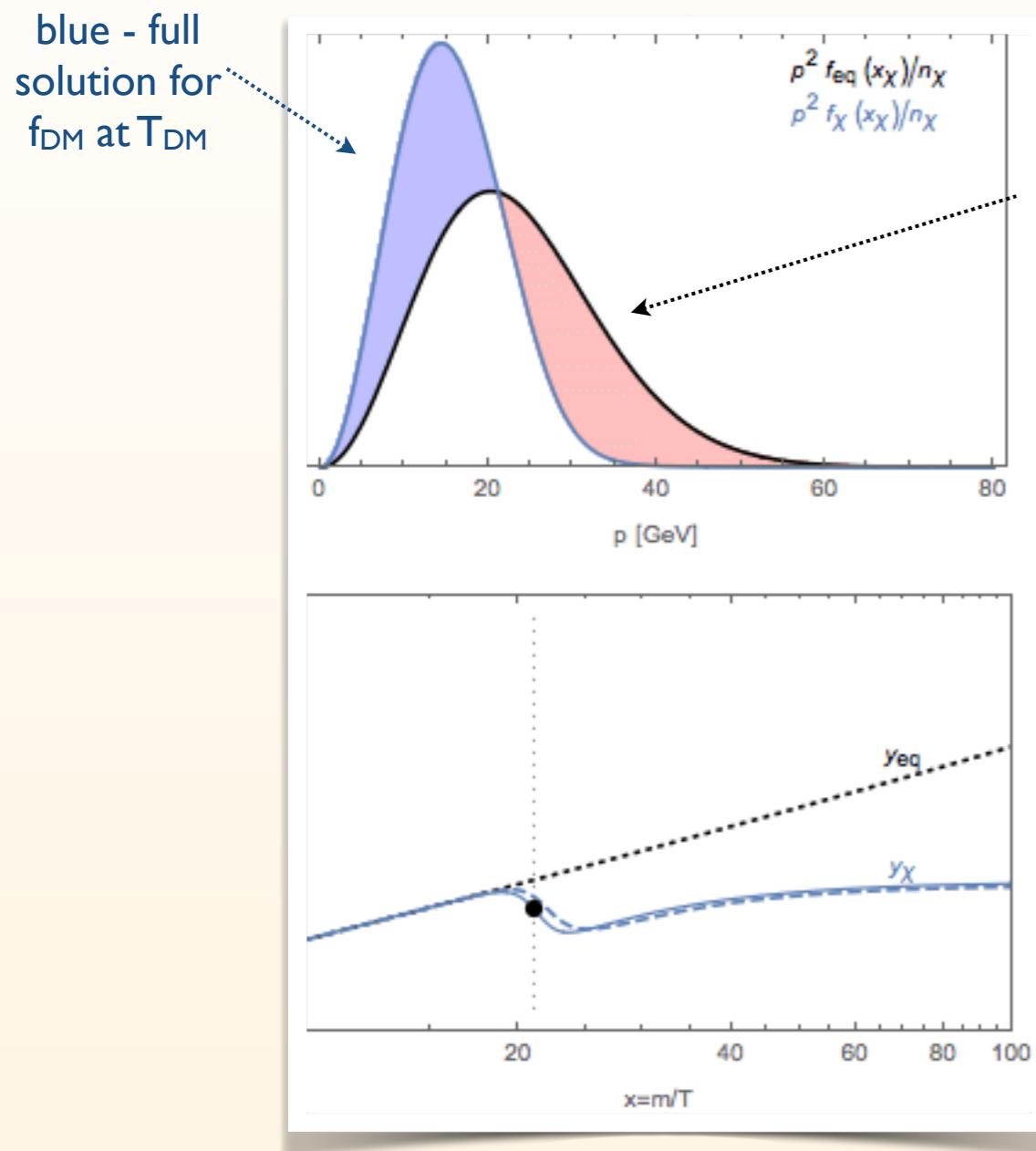
Critical assumption:
kinetic equilibrium at chemical decoupling



... then *standard* thermal relic density calculation **fails!**

FULL PHASE-SPACE EVOLUTION

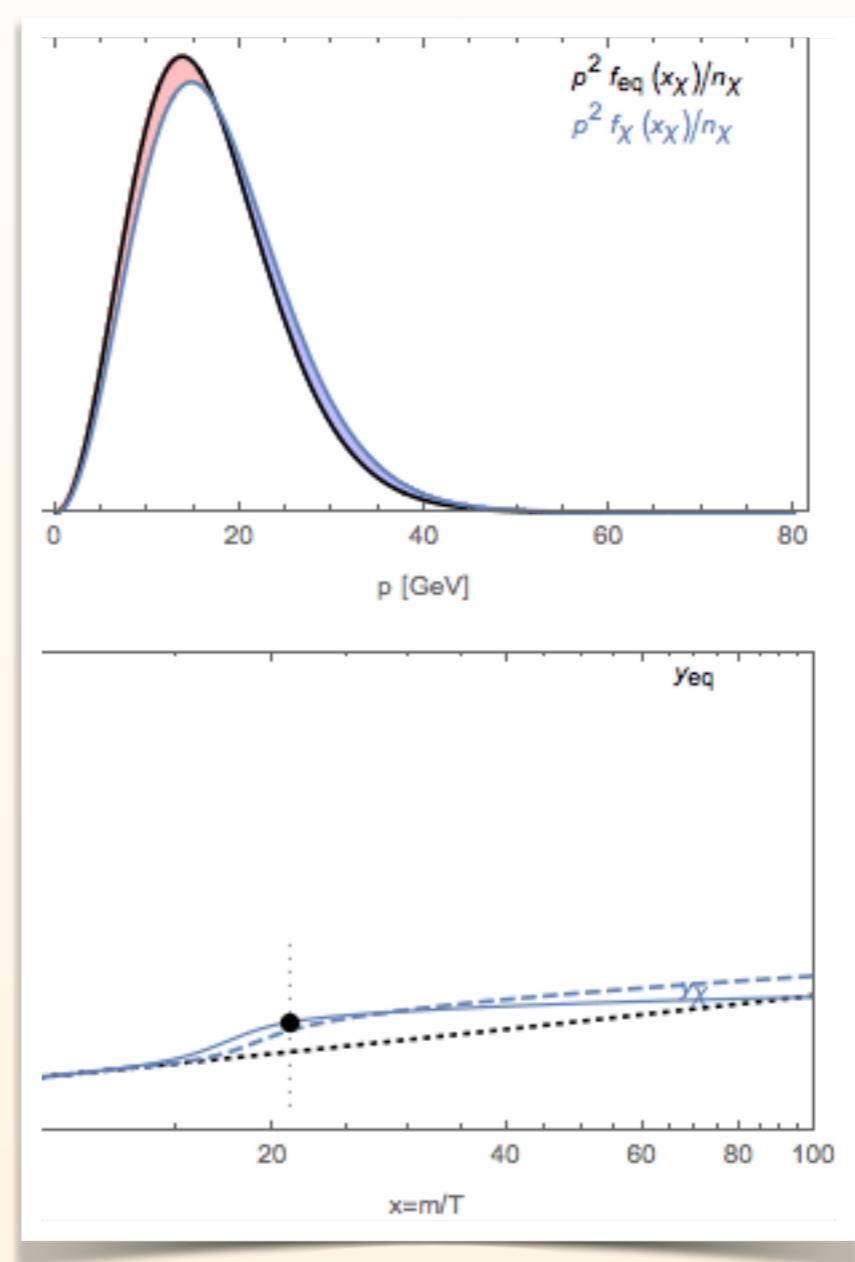
$m_{DM} = 58 \text{ GeV}$



significant deviation from equilibrium shape **already around freeze-out**

→ effect on relic density largest, both from different T and f_{DM}

$m_{DM} = 62.5 \text{ GeV}$



large deviations **at later times**, around freeze-out not far from eq. shape

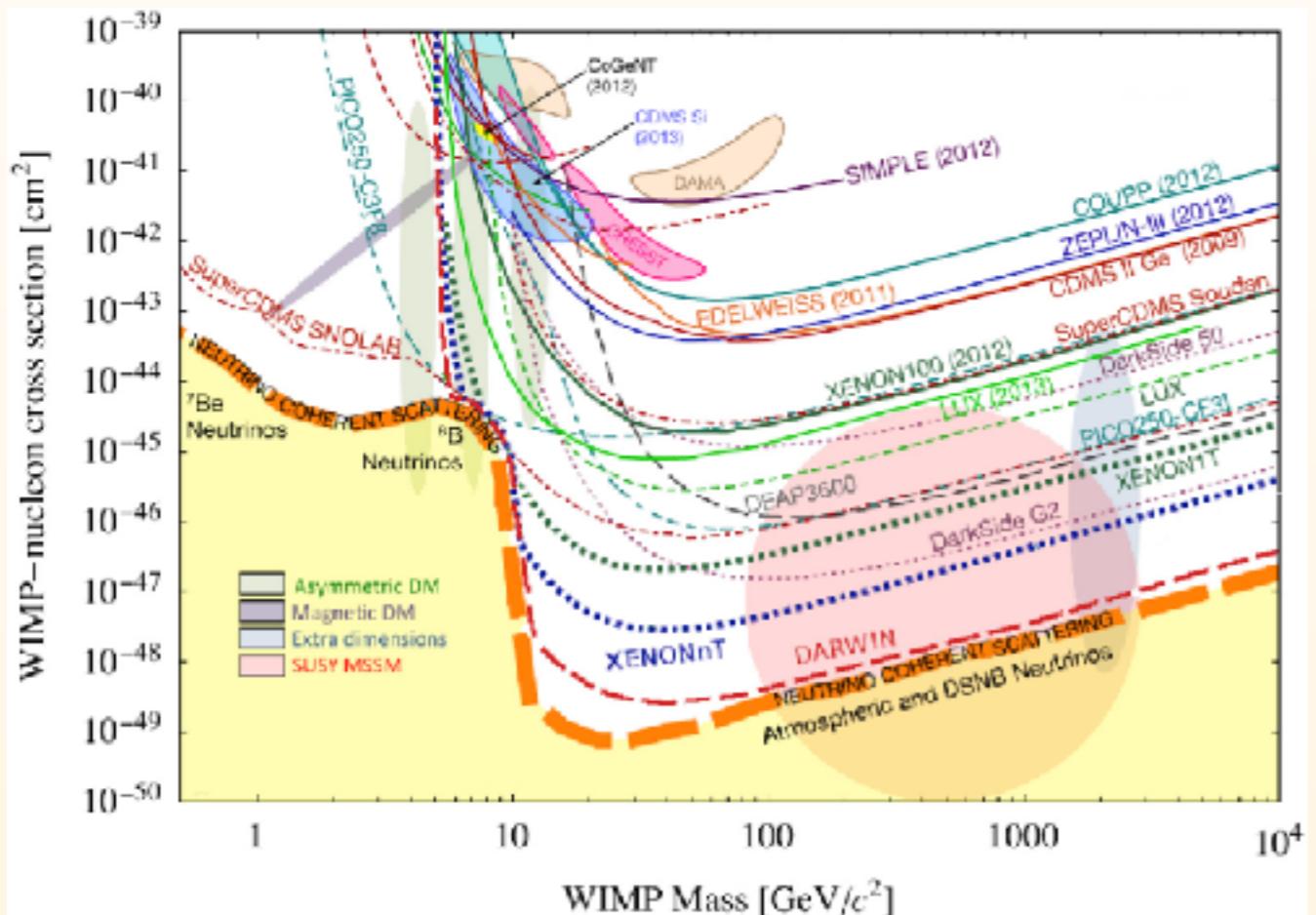
→ effect on relic density ~only from different T

CURRENT LIMITS AND DECLINE OF THE WIMP PARADIGM

”The great tragedy of science - the slaying of a beautiful hypothesis by an ugly fact”

Aldous Huxley

On both Direct Detection and LHC front no* signal of DM particle!



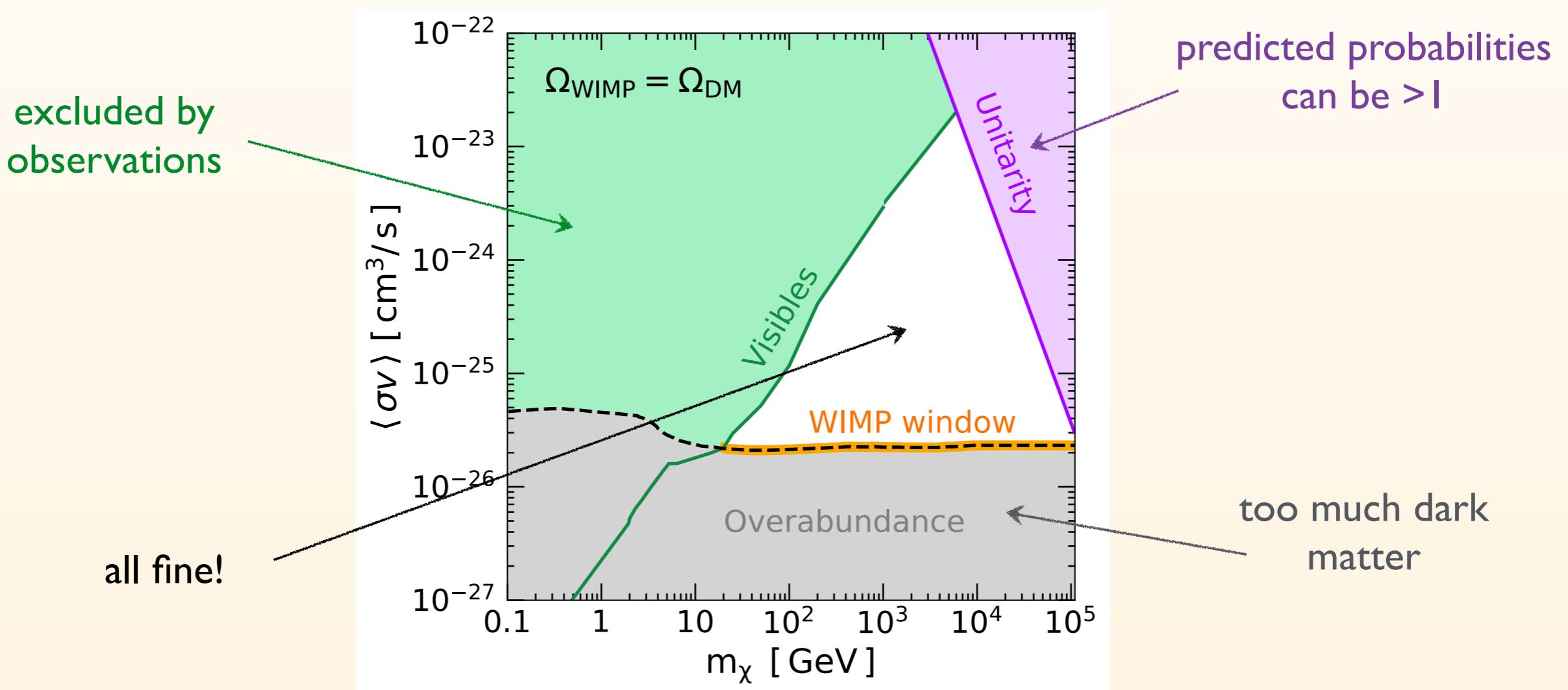
Common feeling: low scale SUSY in "dire restraints"

... BUT IN FACT WIMP NOT EVEN SLIGHTLY DEAD

Most of the (strongest) limits are based on **assumptions** motivated by theoretical prejudice (or convenience)



this can lead to a very broad-brush conclusions



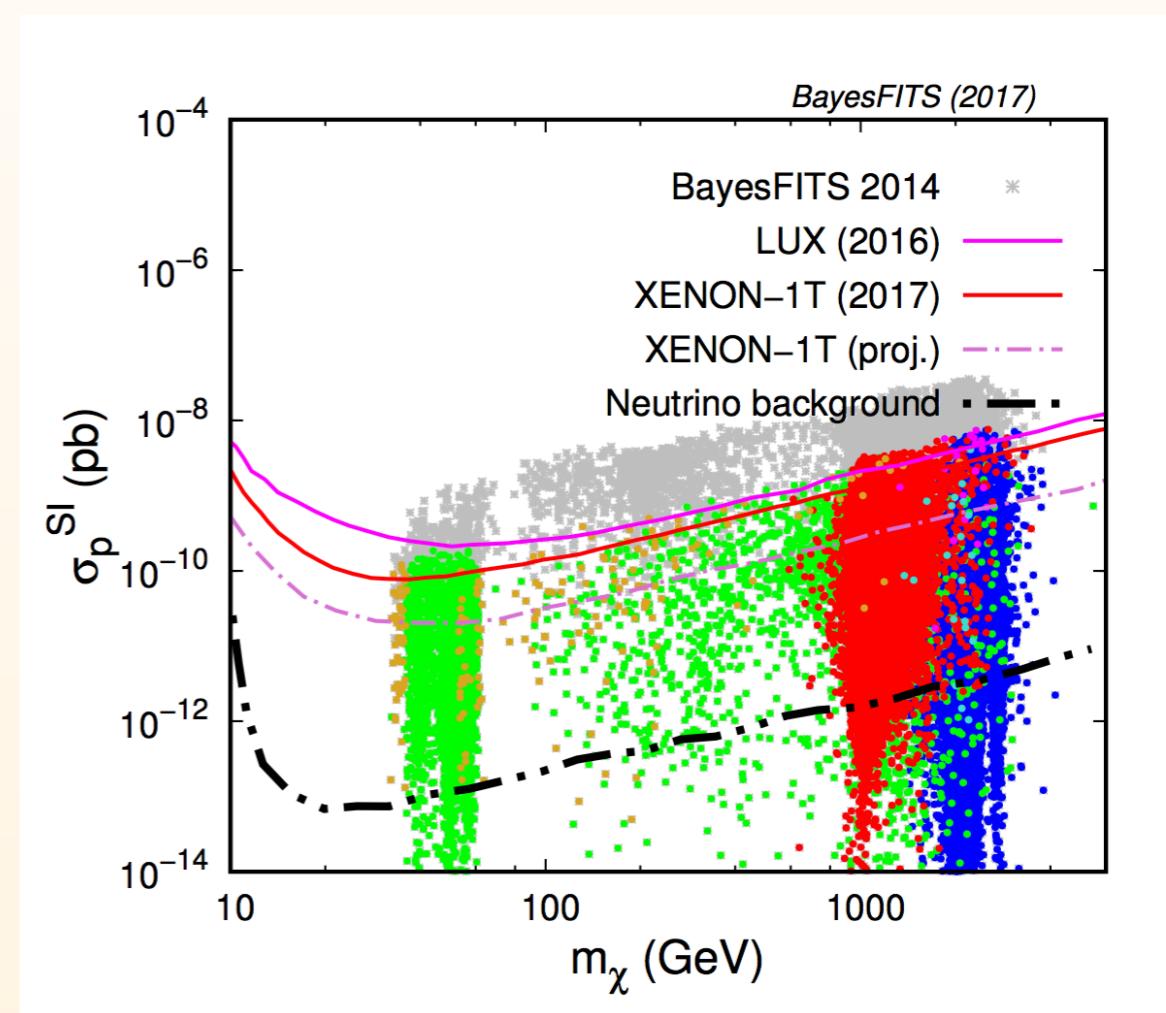
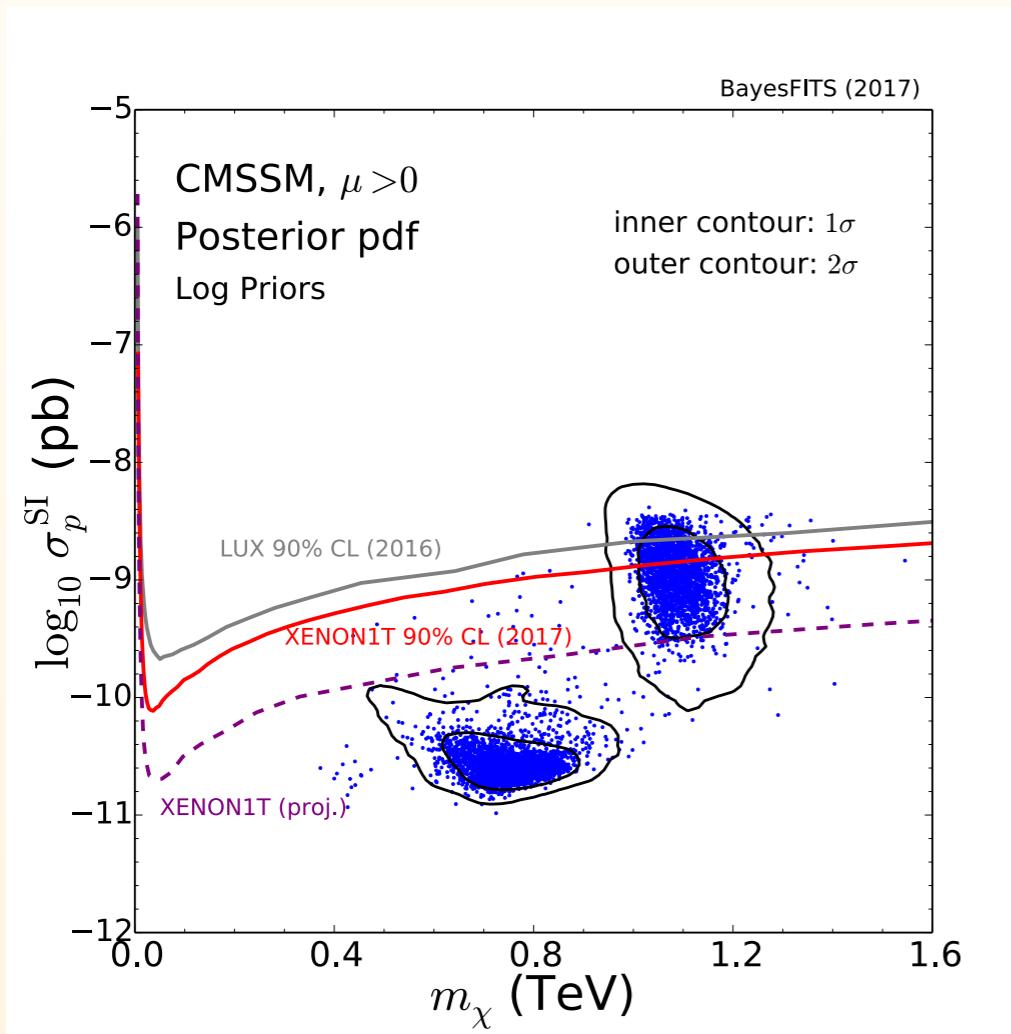
R. Leane et al; 1805.10305

SUSY WIMP

ALSO ACTUALLY QUITE OK

CMSSM points satisfying all the constraints and giving good DM candidate:

generalization to the full pMSSM:



CMSSM- constrained Minimal Supersymmetric Standard Model

pMSSM- phenomenological Minimal Supersymmetric Standard Model

CHAPTER #2

DM AT THE TEV SCALE

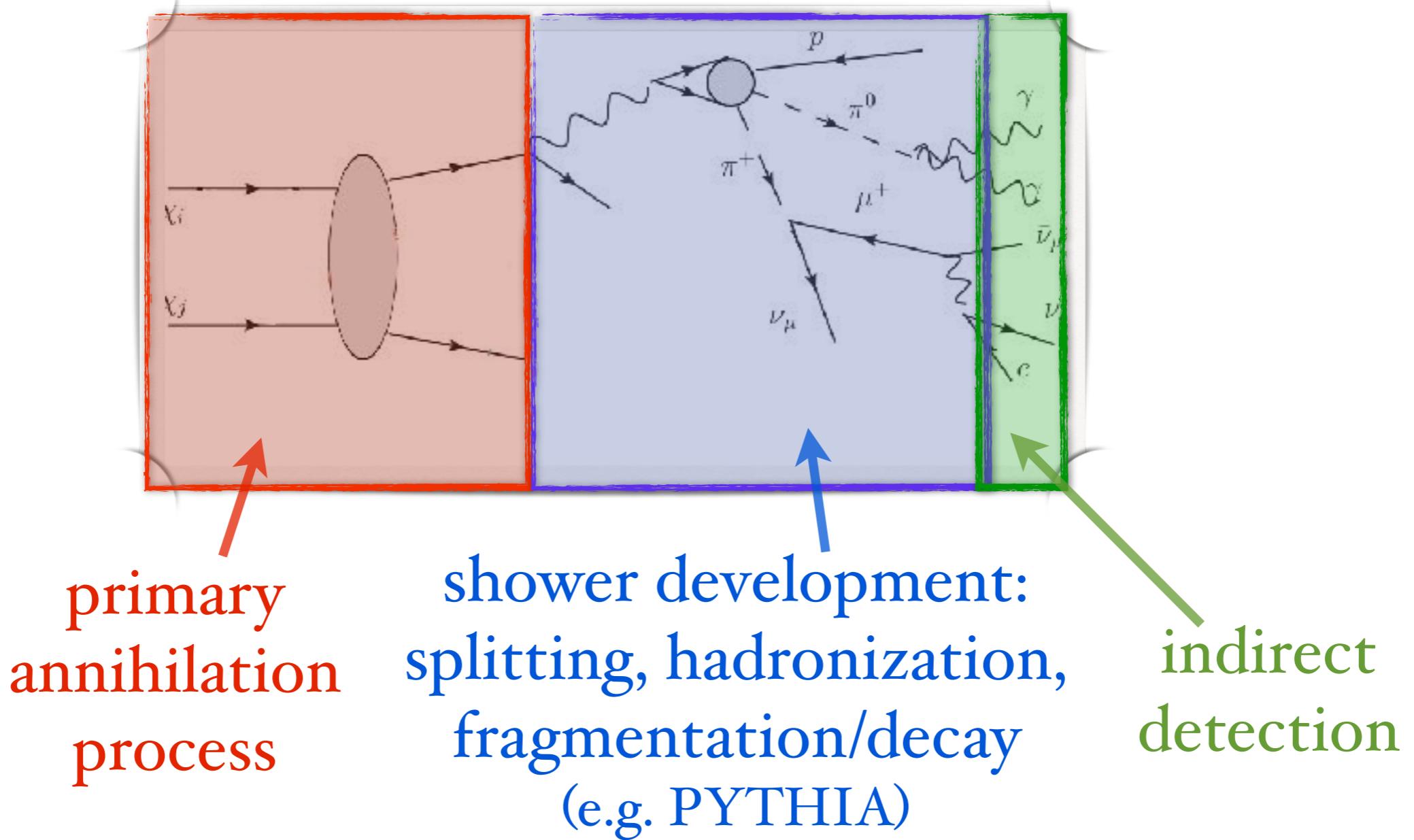
WHY NOT TO GO TO TEV...

- Little Hierarchy Problem: further away from the lamppost (LHC), fine tuning gets worse for simplest models (e.g. CMSSM)
 - Thermal abundance requires large couplings (unitarity bound) or specific mechanism
-

...AND WHY IT IS WORTH IT

- There is no reason in principle not to consider full thermal range up to unitarity limit (apart from naturalness mentioned above)
- We have already seen that even SUSY has regions in that regime and there are many more models on the market
- Fun: new phenomena and new challenges appear

INDIRECT DM DETECTION

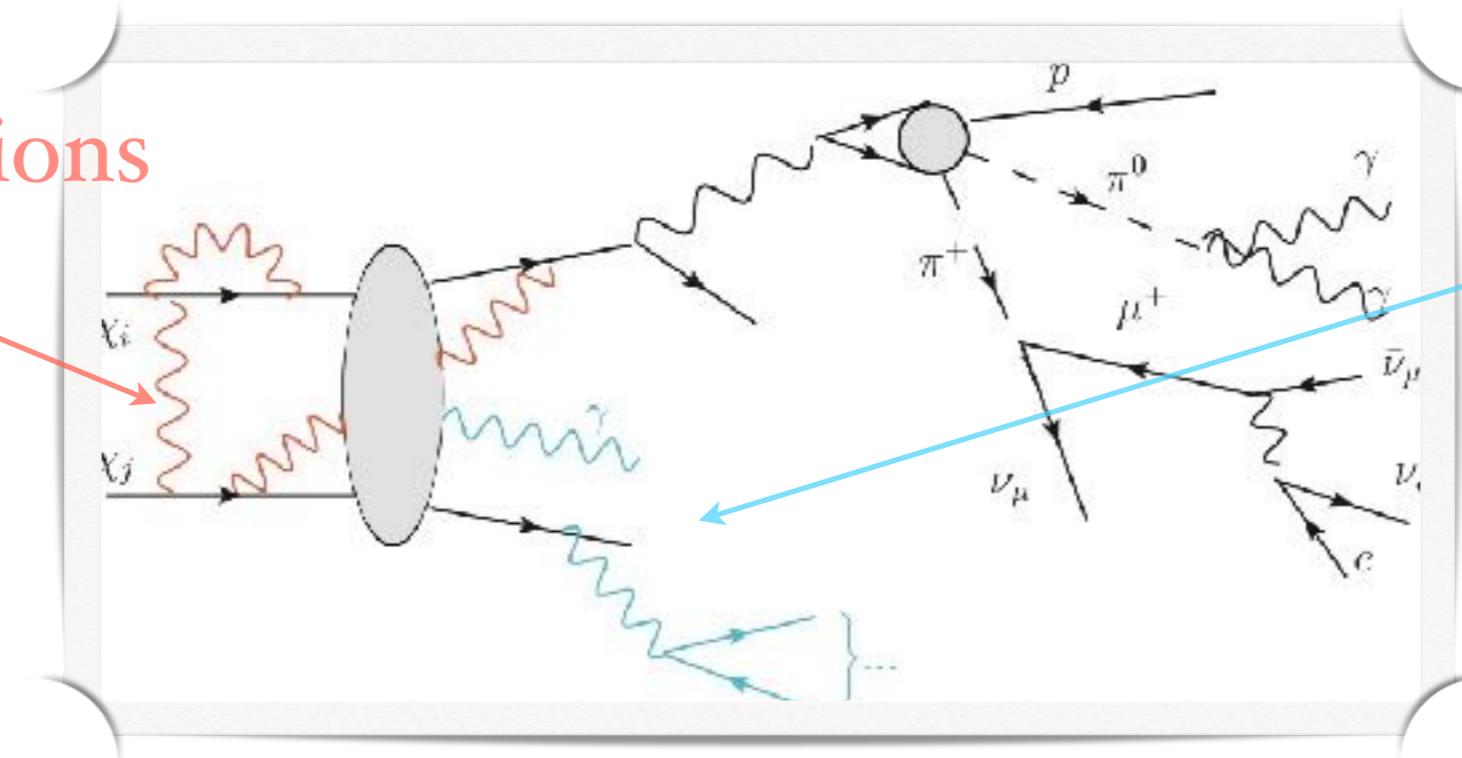


This Feynman diagram is an approximation of lowest order
in perturbation theory!

Actual process can contain many more interactions

EW CORRECTIONS

loop corrections



internal
bremsstrahlung

enhancement by large (Sudakov) logarithms:

$$\alpha_2 \log \frac{m^2}{m_W^2}$$

$$\alpha_2 \left(\log \frac{m^2}{m_W^2} \right)^2$$

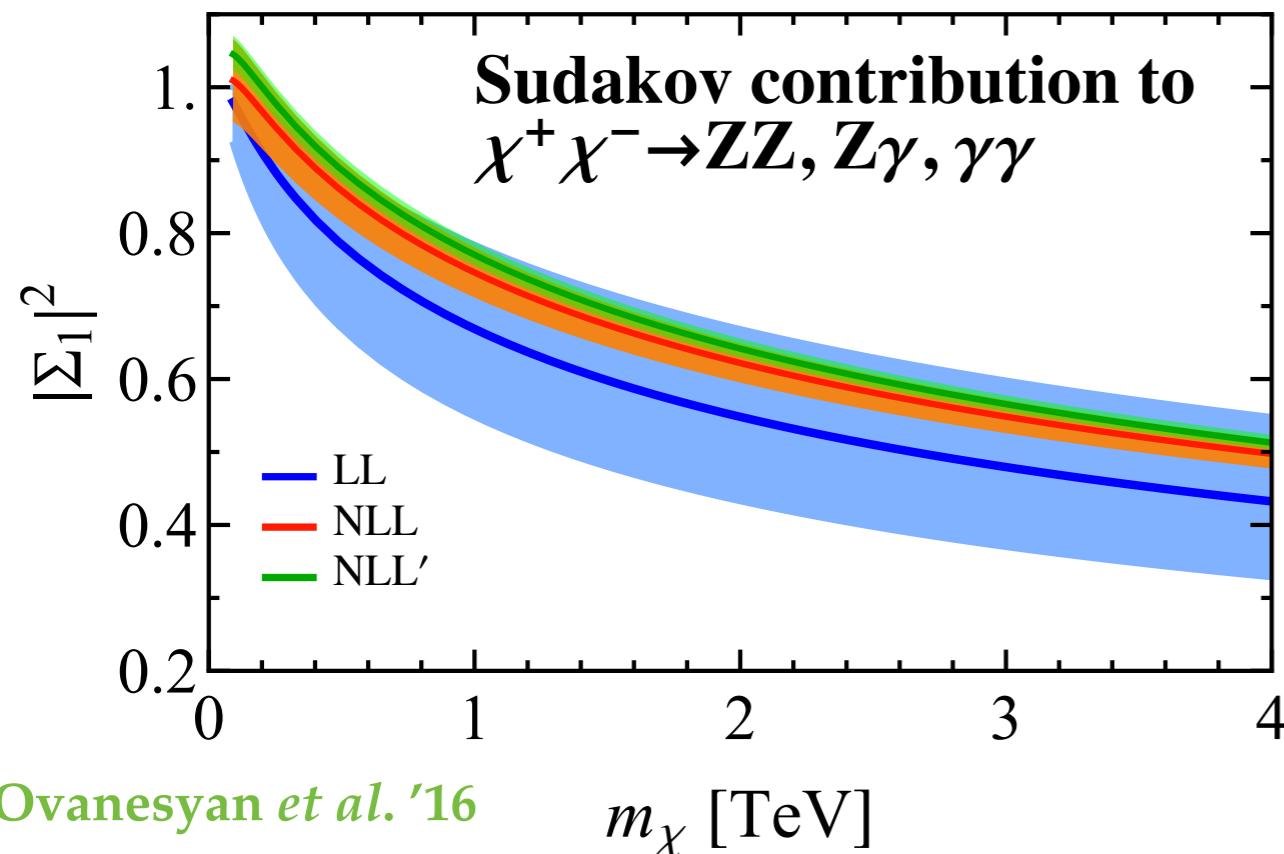
$$m = 1 \text{ TeV}, \alpha_2 \approx \frac{1}{30} \Rightarrow \approx 0.17$$

$$\approx 0.86$$

LL RESSUMATION

Using EFT techniques the contribution for **large logarithms** can be summed to all orders:

$$\ln \frac{C}{C^{\text{tree}}} \sim \sum_{k=1}^{\infty} \left[\underbrace{\alpha_2^k \ln^{k+1}}_{\text{LL}} + \underbrace{\alpha_2^k \ln^k}_{\text{NLL}} + \underbrace{\alpha_2^k \ln^{k-1}}_{\text{NNLL}} + \dots \right]$$



This is a relatively complicated computation, which does **not** have to be done if DM is lighter!

SOMMERFELD EFFECT

re-summation

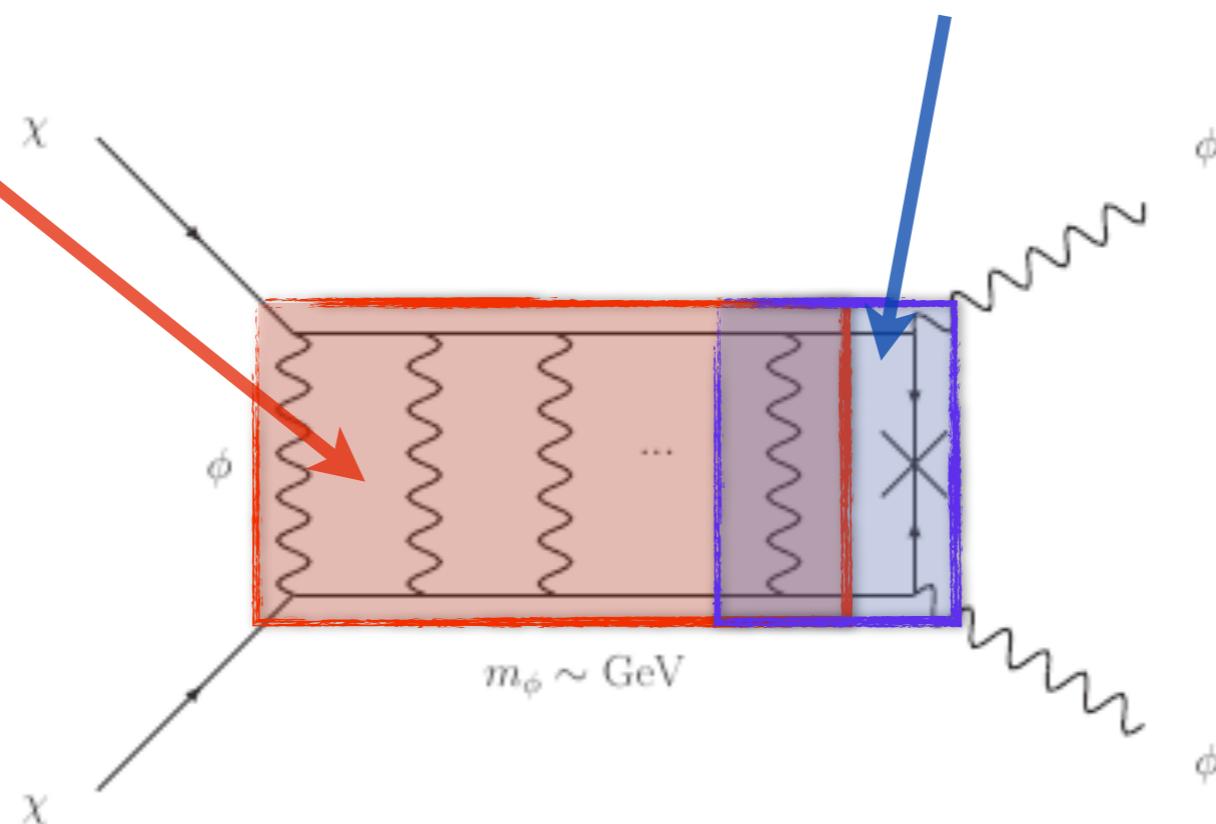
$$\frac{1}{m_\phi} \gtrsim \frac{1}{\alpha m_\chi}$$

force range
Bohr radius

$$m_\chi v^2 \lesssim \alpha^2 m_\chi$$

kinetic energy
Bohr energy

$$\text{one-loop} \propto \alpha \frac{m_\chi}{m_\phi}$$



$$\sigma_{\text{SE}} = S(v) \sigma_0$$

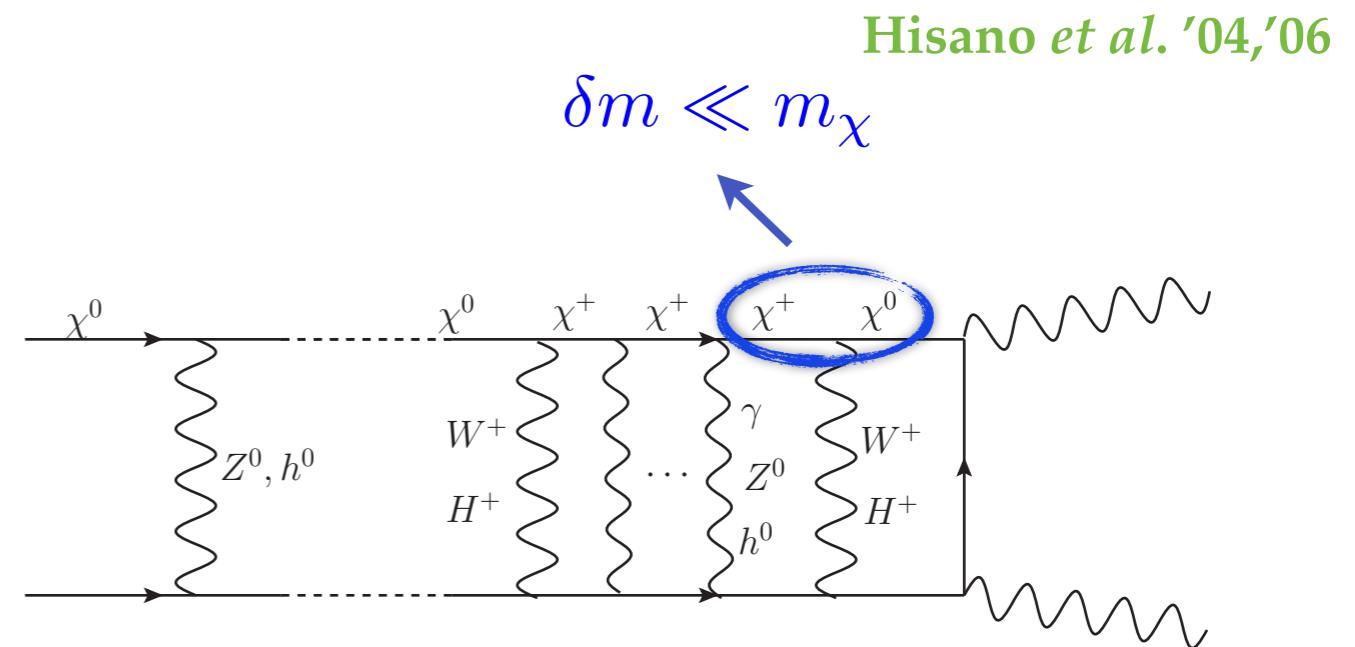
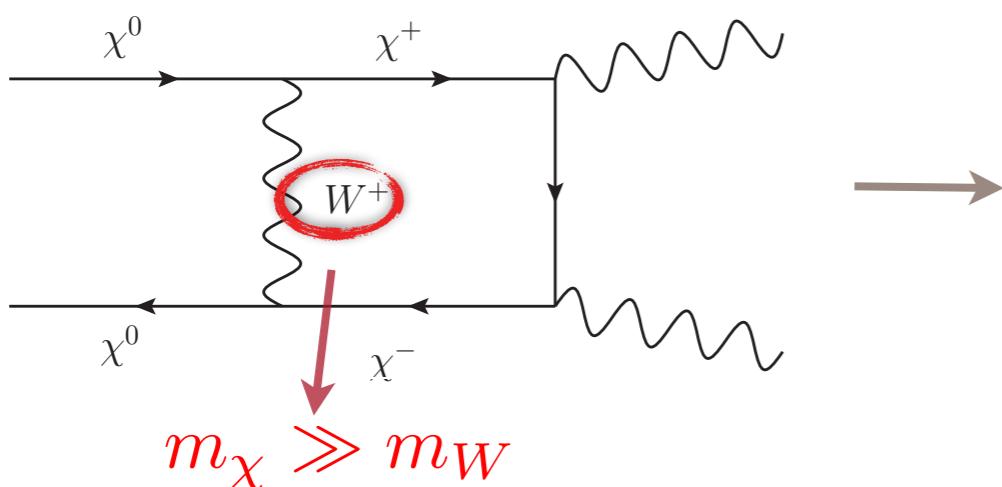
Arkani-Hamed *et al.* '09

→ in a special case of Coulomb force: $S(v) = \frac{\pi\alpha/v}{1 - e^{-\pi\alpha/v}} \approx \pi \frac{\alpha}{v}$

THE SOMMERFELD EFFECT FROM EW INTERACTIONS

force carriers in the MSSM:

~~χ~~ , W^\pm , Z^0 , h_1^0 , h_2^0 , H^\pm



at TeV scale \Rightarrow generically effect of $\mathcal{O}(1 - 100\%)$

on top of that **resonance** structure

↳ effect of $\mathcal{O}(\text{few})$
for the relic density

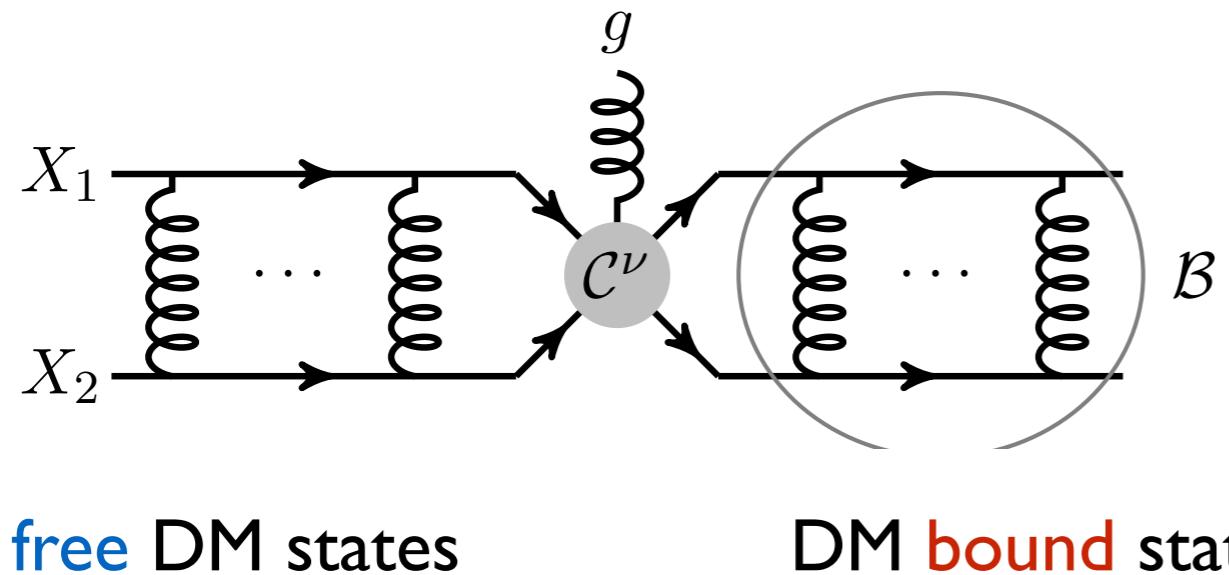
AH, R. Iengo, P. Ullio. '10

AH '11

AH *et al.* '17, M. Beneke *et al.*; '16 25

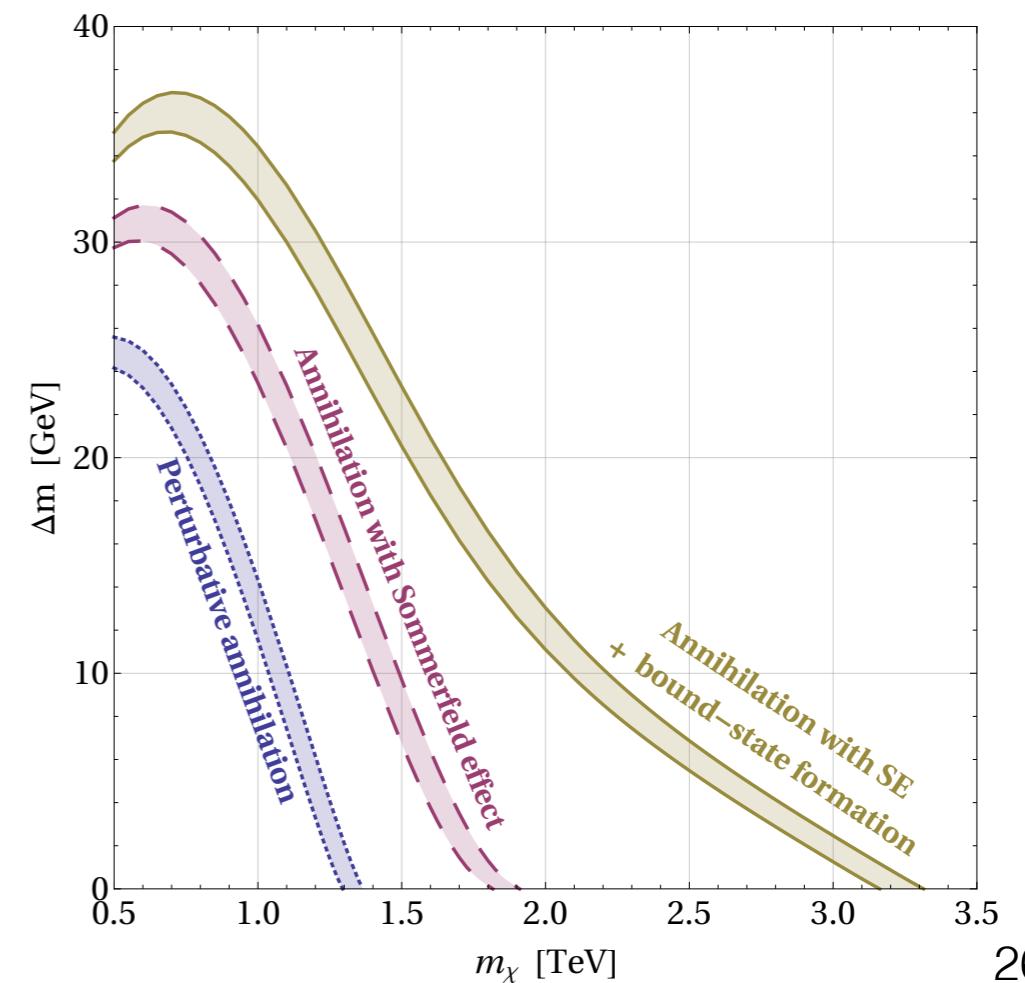
BOUND STATE FORMATION

see papers by K. Petraki et al. '14-18



can happen when long range force between DM states is present

significant modification of the annihilation rate - large effects on the DM models, especially in the TeV scale

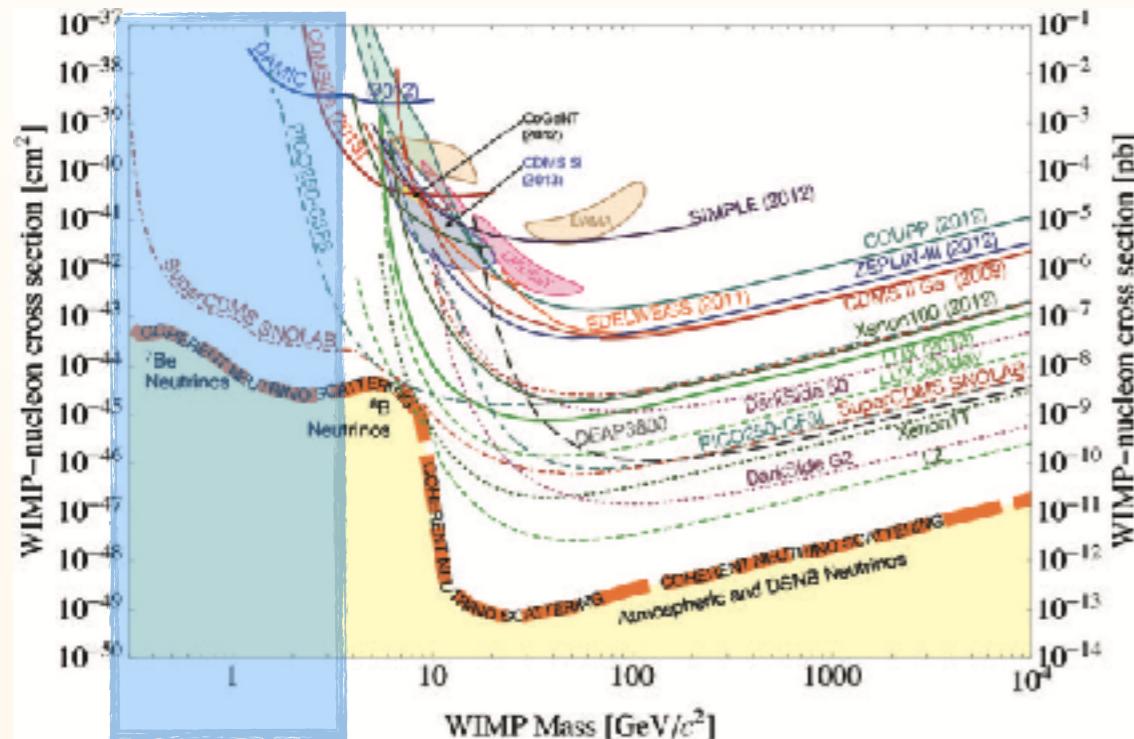


CHAPTER #3

LIGHT DM

MEV-GeV SCALE DM

Below the sensitivity threshold for most
Direct Detection experiments

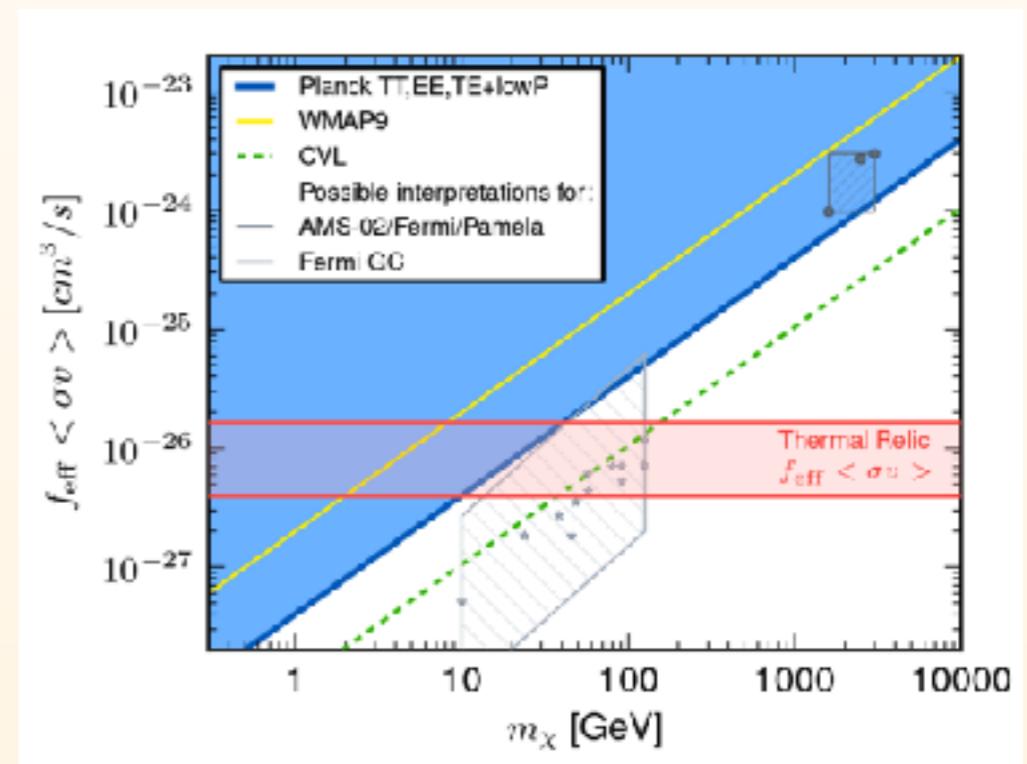


Are there any
upsides?



But needs **very weak** coupling to visible sector:

- I. light = easier to produce at colliders
- II. limits from **cosmology** are quite strong

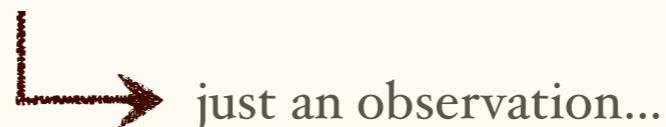


- III. thermal relic density $\sim (\text{mass}/\text{coupling})^2$

LIGHT DM MOTIVATION

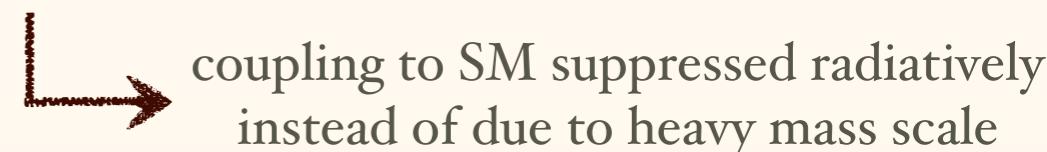
Few arguments in favor of MeV-GeV dark sector:

Mass scale of known visible matter



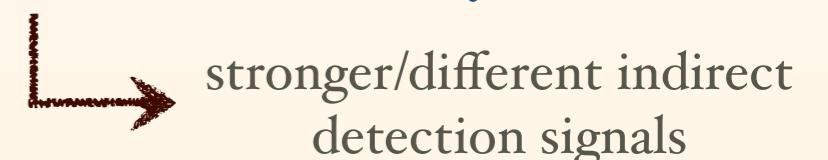
just an observation...

Natural region for hidden dark sectors coupled to SM at loop level



coupling to SM suppressed radiatively
instead of due to heavy mass scale

For fixed mass density - lighter DM means higher number density



stronger/different indirect
detection signals

DM self-interactions: $0.1 \text{ cm}^2 \text{ g}^{-1} < \langle \sigma_T \rangle_{30} / m_\chi < 10 \text{ cm}^2 \text{ g}^{-1}$



the lighter, the stronger

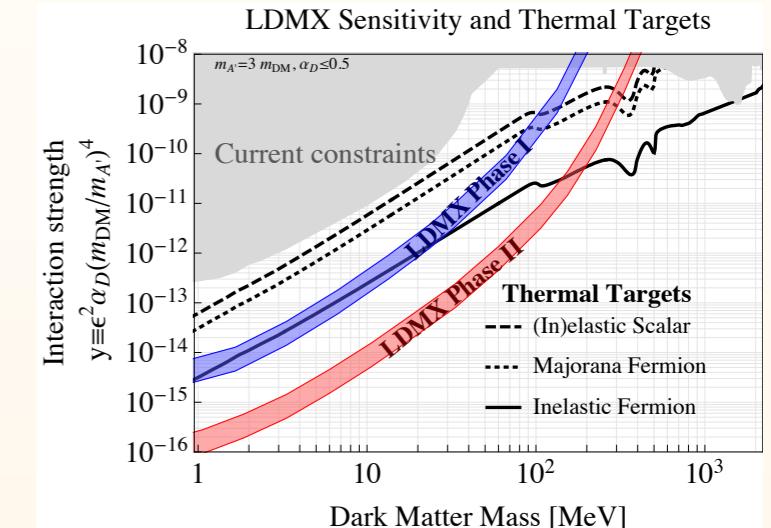
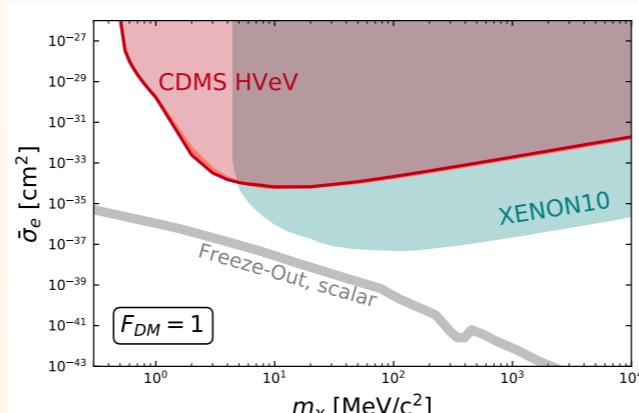
DETECTION OF LIGHT DM

List of strategies:

1. standard Indirect Detection - looking at data in lower energies

2. Fixed target accelerator experiments, e.g. LDMX →

3. Direct Detection - scattering on electrons

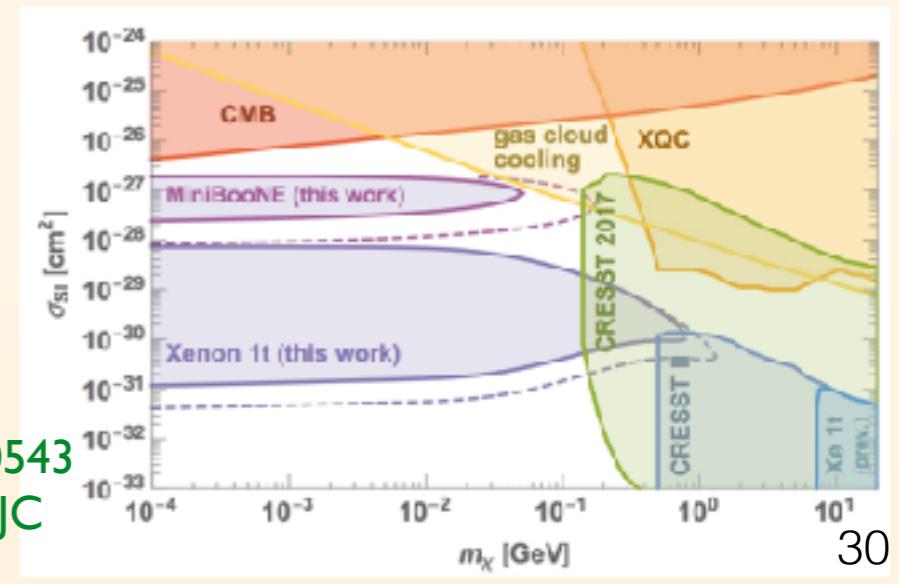


4. Scattering (elastic or inelastic) on CRs in present and Early Universe

see e.g. Capiello et al. 1810.07705

5. ...

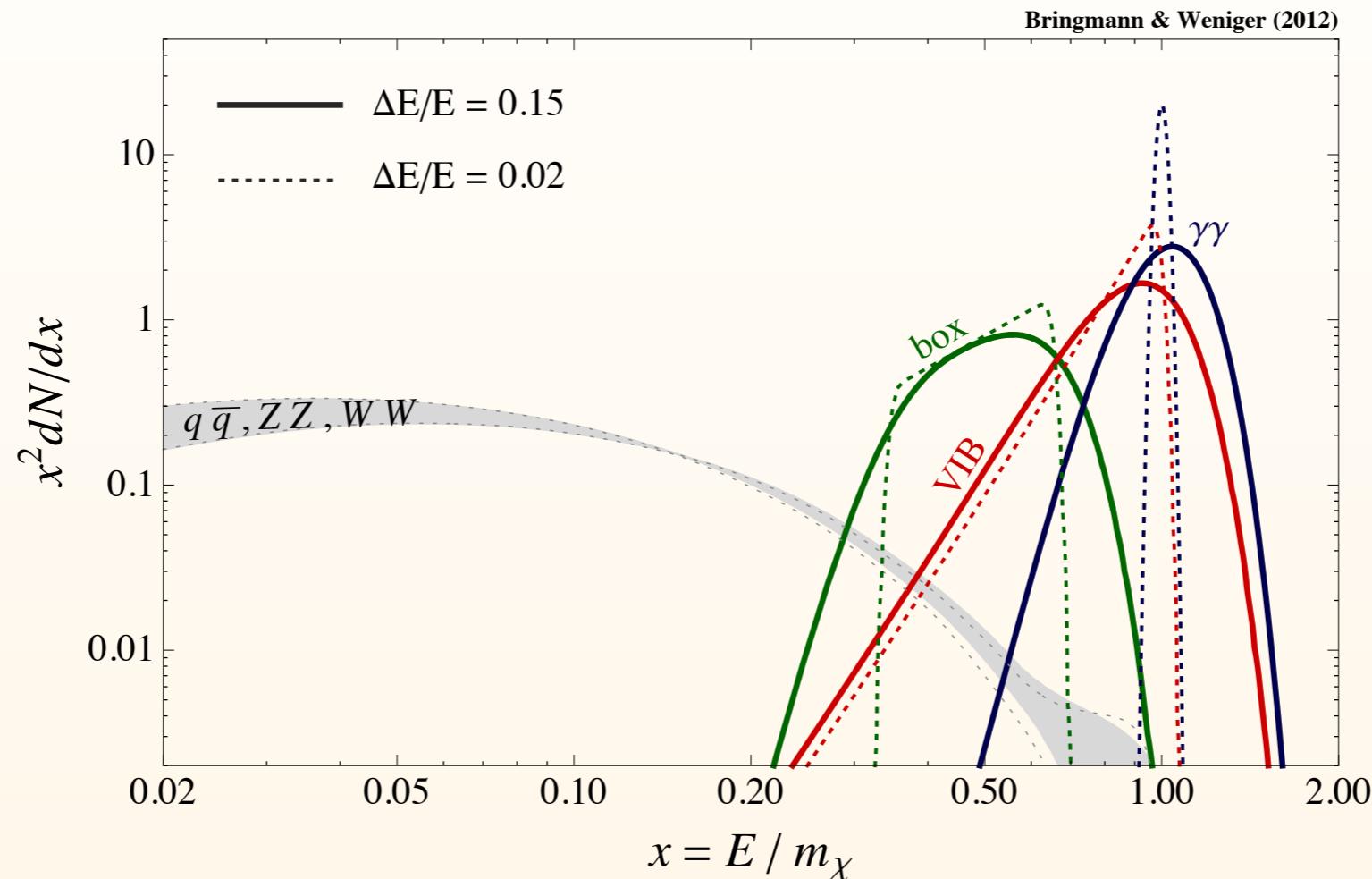
Some of the limits:



T. Bringmann, M. Pospelov; 1810.10543
- to see more join our Bi-weekly JC
today at 2.15pm!

SPECTRAL FEATURES

WHAT CAN WE LOOK FOR?



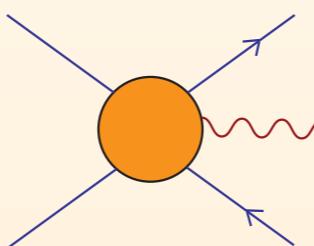
Gamma-ray lines

$$\chi\chi \rightarrow \gamma\gamma$$

generically loop-suppressed

Internal Bremsstrahlung

$$\chi\chi \rightarrow f\bar{f}\gamma$$



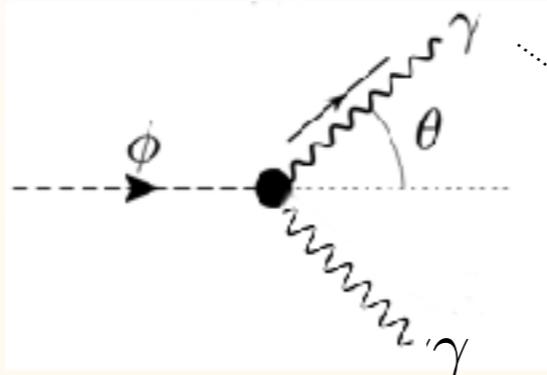
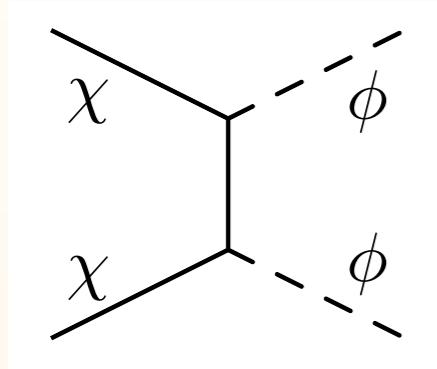
Box-shaped

$$\chi\chi \rightarrow \phi\phi \implies \phi \rightarrow \gamma X$$

tree-level;
cascade decay

GAMMA-RAY BOXES

Consider a process: $\chi\chi \rightarrow \phi\phi \Rightarrow \phi \rightarrow \gamma\gamma$



In the LAB frame:

$$E_\gamma = \frac{m_\phi^2}{2 m_{DM}} \left(1 - \cos \theta \sqrt{1 - \frac{m_\phi^2}{m_{DM}^2}} \right)^{-1}$$

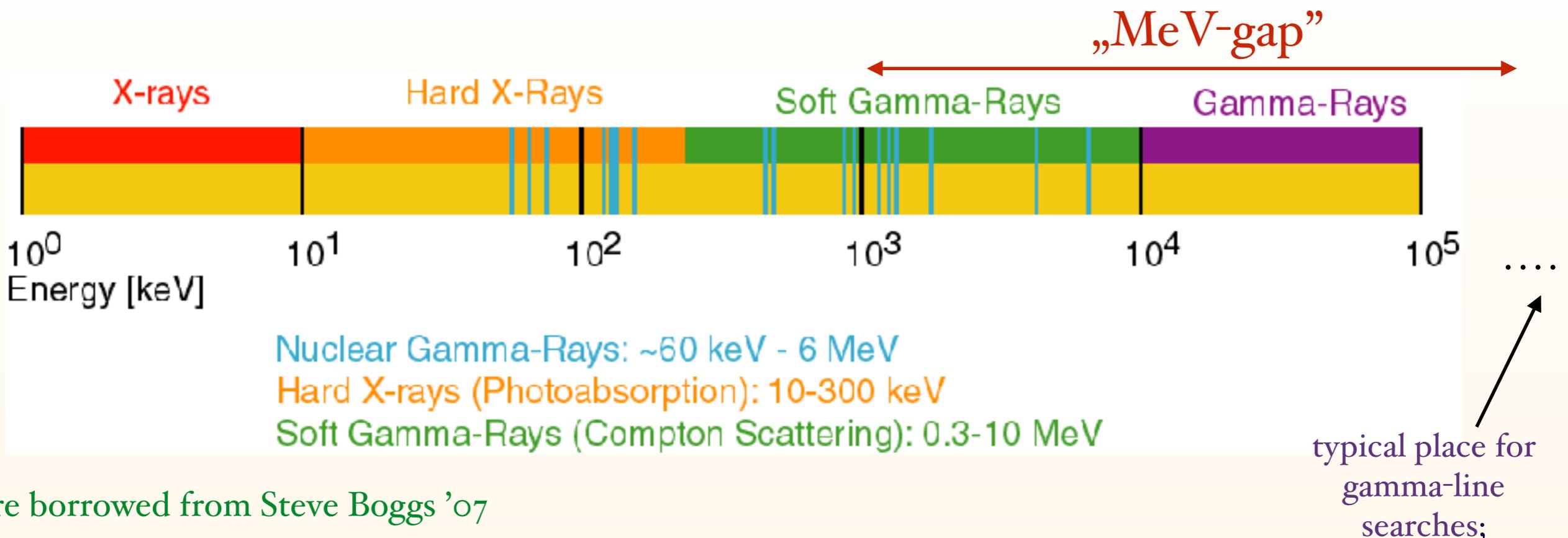
If ϕ produced at rest \rightarrow monochromatic line...

...if not, boosted to give a box shaped spectrum:

$$\frac{dN_\gamma}{dE} = \frac{2}{\Delta E} [\Theta(E - E_-) - \Theta(E - E_+)]$$

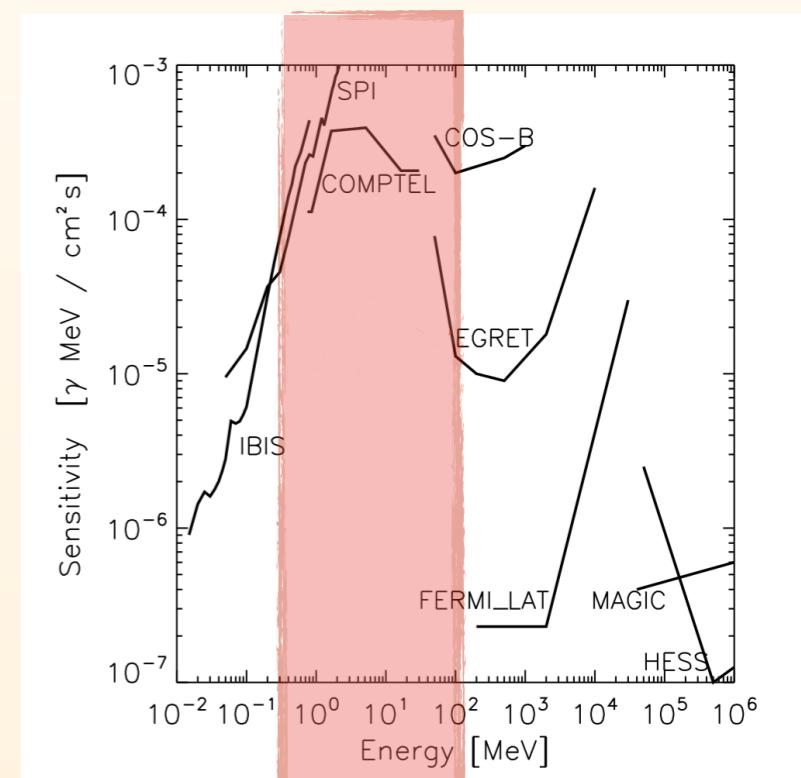
(For narrow boxes I may use the box and line terms interchangeably...)

KNOWN GAMMA-RAY LINES



ture borrowed from Steve Boggs '07

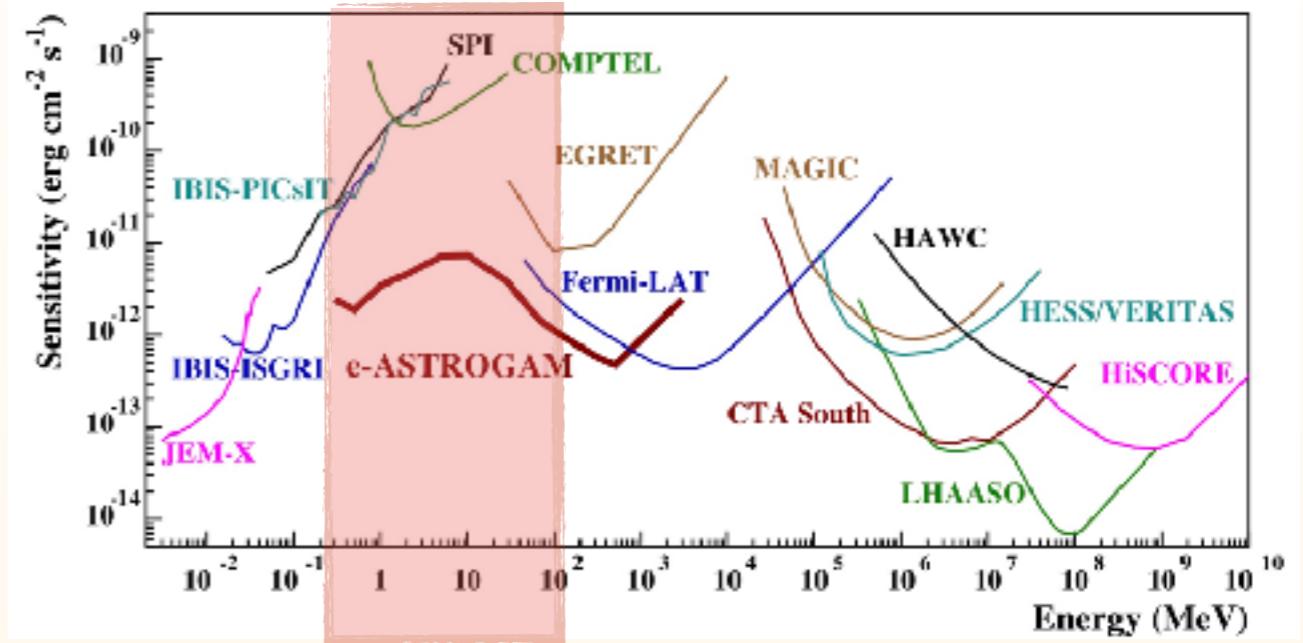
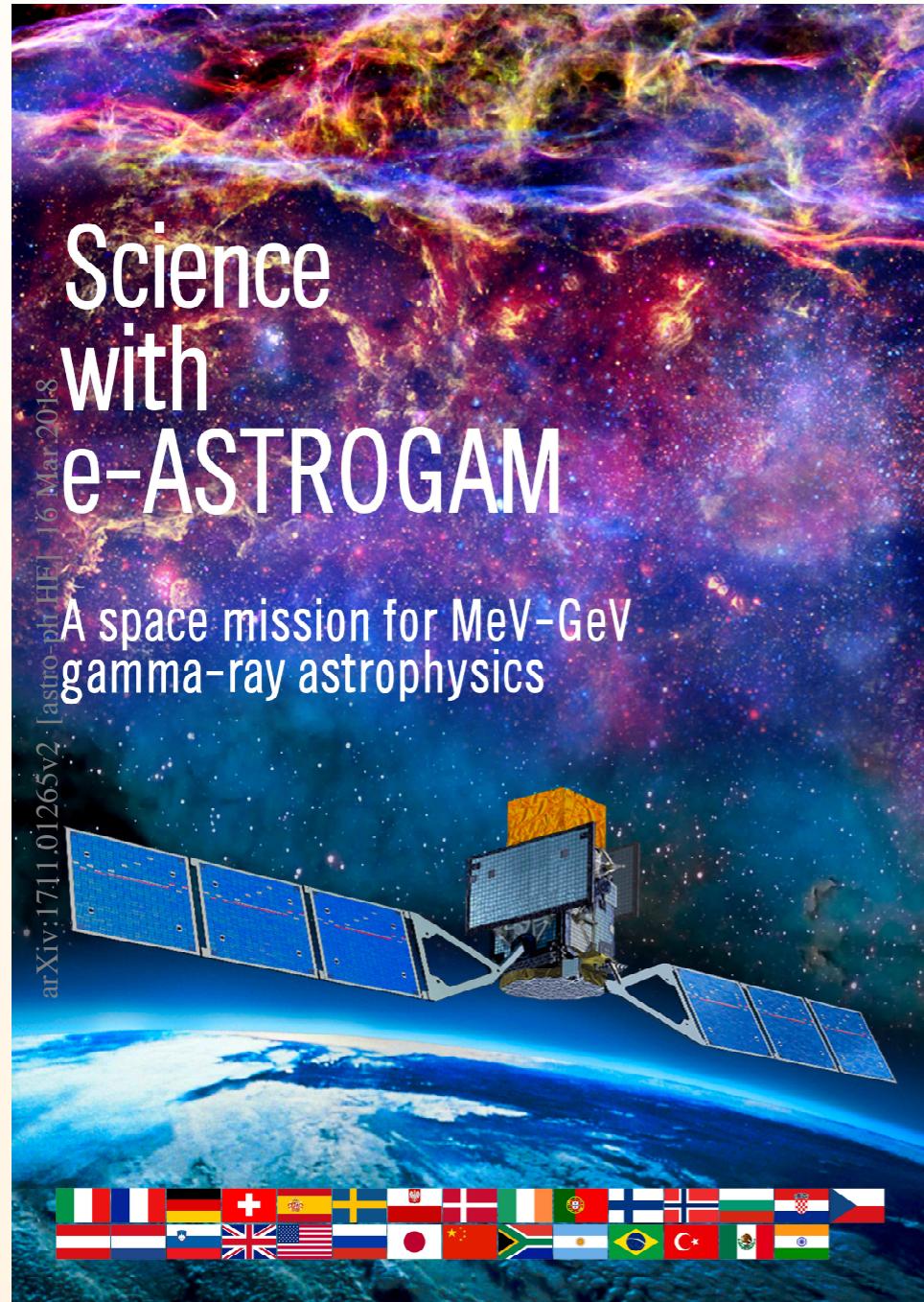
The MeV-gap contains a timely sweet spot for spectral features searches:
no background lines + scarce complementary data



„MeV-gap”

E-ASTROGAM <http://eastrogam.iaps.inaf.it>

submitted (w/o success) to M5 ESA Call



Main features:

- Broad energy coverage (0.3 MeV to 3 GeV);
- Large FoV (>2.5 sr), ideal to detect transient sources and hundreds of GRBs;
- Pioneering polarimetric capability for both steady and transient sources;
- Optimized source identification capability obtained by the **best angular resolution** (about 0.15 degrees at 1 GeV);
- <ms trigger and alert capability for GRBs and other transients;
- Combination of **Compton and pair-production** detection techniques

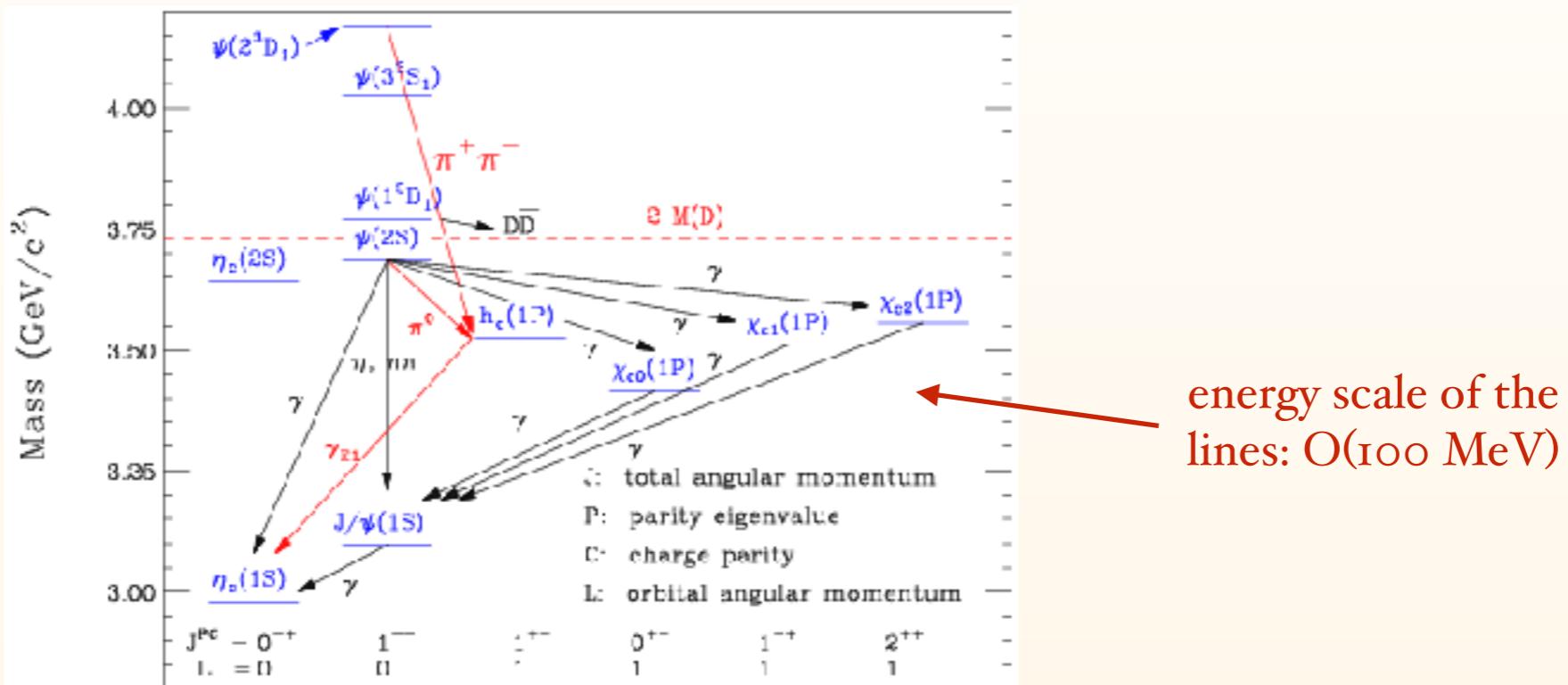
A. De Angelis (ed) et al., JHEAp 19 (2018) 1-106

@UiO: T. Bringmann, AH, A. Raklev, J. Van den Abeele

MESON SPECTROSCOPY

Transitions between **meson states** lead to monochromatic pions or photons:

E.g.:



B and D mesons are composed from **one light** and **one heavy quark**

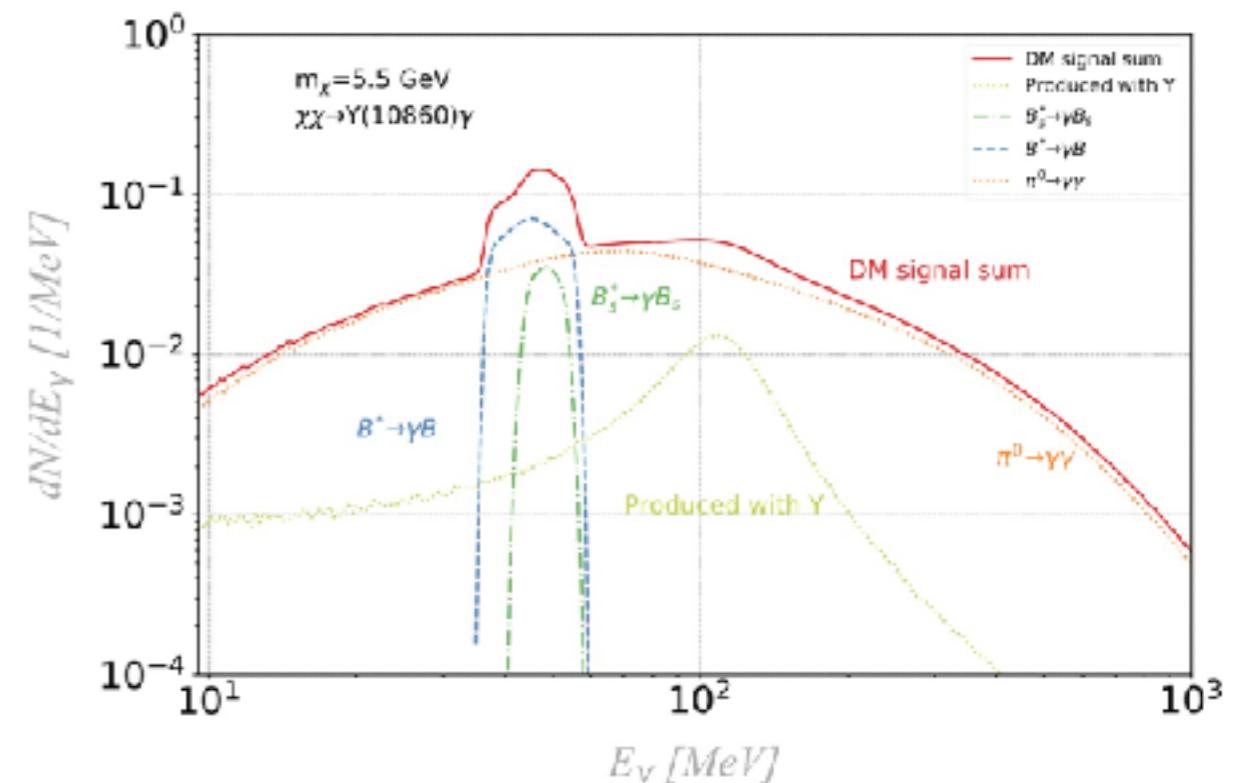
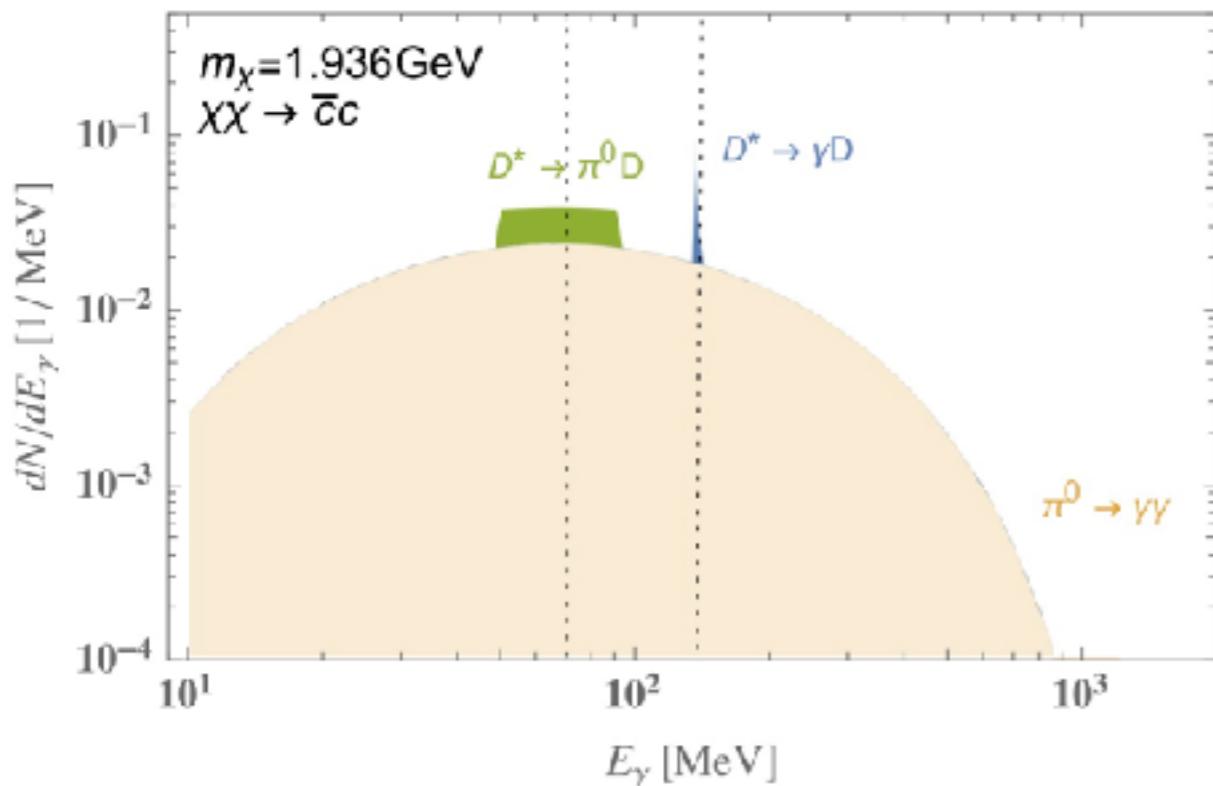
can be produced in annihilation to $b\bar{b}$ and $c\bar{c}$

do **not** show up in astrophysical background

MEV SPECTRAL FEATURES

T. Bringmann, A. Galea, AH and Ch. Weniger; Phys.Rev. D95 (2017)

Spectral boxes coming from **excited meson** decays:



...and from formation of **bound states** with accompanying photon emission

A. Raklev, I. Strümke, J. van den Abeele

A NEW HOPE?

(INSTEAD OF CONCLUSIONS)

A New Era in the Quest for Dark Matter

Gianfranco Bertone¹ and Tim M.P. Tait^{1,2}

ABSTRACT

There is a growing sense of ‘crisis’ in the dark matter community, due to the absence of evidence for the most popular candidates such as weakly interacting massive particles, axions, and sterile neutrinos, despite the enormous effort that has gone into searching for these particles. Here, we discuss what we have learned about the nature of dark matter from past experiments, and the implications for planned dark matter searches in the next decade. We argue that diversifying the experimental effort, incorporating astronomical surveys and gravitational wave observations, is our best hope to make progress on the dark matter problem.

Nature, volume 562, pages 51–56 (2018)

... but precision cosmology & astrophysics has a potential to provide the so-much needed observational input and show which way to follow

From HEP perspective it all may feel quite depressing...

(...) the new guiding principle should be “no stone left unturned”.



In any case a lot of physics still awaits to be understood in how the Dark Matter came to life and how it can reveal itself to us

P.S.

... SO, WHAT WILL IT BE THEN?

On one hand a **decent, robust and well motivated** theory (WIMP) on which you can still safely bet... though **perhaps bit rusty now**

On the other **new challenging ideas** and mechanisms, that might bring some new fresh air... but **first have to prove their worth**



vs.



(Though we'll probably wait bit longer than a month to see what DM theory will prevail)