

DARK MATTER PRODUCTION OUT OF KINETIC EQUILIBRIUM: LATEST DEVELOPMENTS

Andrzej Hryczuk



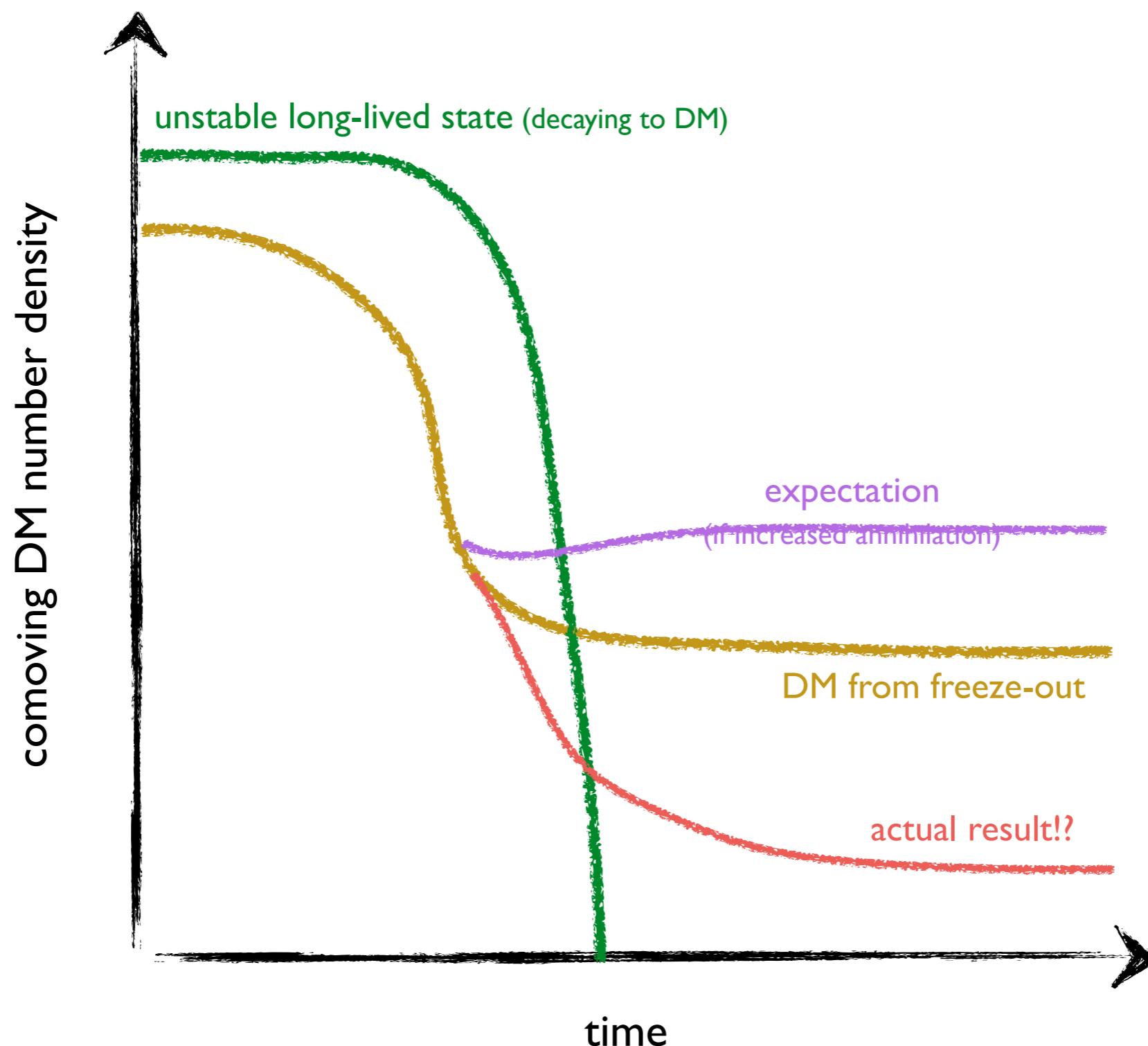
Based on:

T. Binder, T. Bringmann, M. Gustafsson & A.H. [1706.07433, 2103.01944](#)

A.H. & M. Laletin [2204.07078, 2104.05684](#)

work in progress with **S. Chatterjee**

IN CASE YOU'RE NOT INTERESTED IN WHAT FOLLOWS...



TO SEE WHY AND LEARN MORE STAY TUNED :)

THERMAL RELIC DENSITY

STANDARD SCENARIO

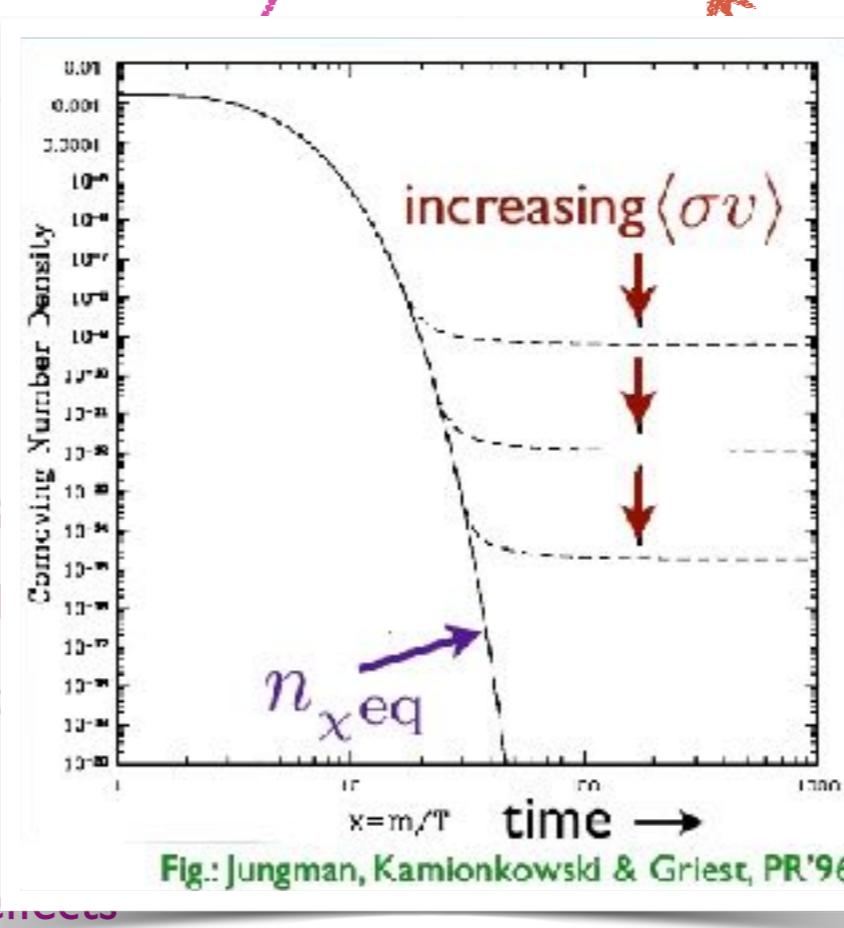
modified expansion rate

general
multi-
component
dark sector

$$\frac{dn_\chi}{dt} + 3Hn_\chi = \frac{dn_\chi}{dt} + 3Hn_\chi$$

modified cross
Sommer
Bound St
NLO
finite T effects

e.g., relentless DM, D'Eramo et al. '17, ...



numerical codes e.g.,
DarkSUSY, micrOMEGAs,
MadDM, SuperISORElic, ...

$$n_{\bar{\chi}} - n_\chi^{\text{eq}} n_{\bar{\chi}}^{\text{eq}}$$

$$n_{\bar{\chi}} - n_\chi^{\text{eq}} n_{\bar{\chi}}^{\text{eq}}$$

breakdown of necessary
assumptions leading to
different form of the
equation, e.g. violation of
kinetic equilibrium

where the thermally averaged cross section:

$$\langle \sigma_{\chi\bar{\chi} \rightarrow ij} v_{\text{rel}} \rangle^{\text{eq}} = -\frac{h_\chi^2}{n_\chi^{\text{eq}} n_{\bar{\chi}}^{\text{eq}}} \int \frac{d^3 \vec{p}_\chi}{(2\pi)^3} \frac{d^3 \vec{p}_{\bar{\chi}}}{(2\pi)^3} \sigma_{\chi\bar{\chi} \rightarrow ij} v_{\text{rel}} f_\chi^{\text{eq}} f_{\bar{\chi}}^{\text{eq}}$$

THERMAL RELIC DENSITY

STANDARD APPROACH

Boltzmann equation for $f_\chi(p)$:

$$E (\partial_t - H \vec{p} \cdot \nabla_{\vec{p}}) f_\chi = \mathcal{C}[f_\chi]$$

 integrate over p
 (i.e. take 0th moment)

$$\frac{dn_\chi}{dt} + 3Hn_\chi = -\langle \sigma_{\chi\bar{\chi} \rightarrow ij} \sigma_{\text{rel}} \rangle^{\text{eq}} (n_\chi n_{\bar{\chi}} - n_\chi^{\text{eq}} n_{\bar{\chi}}^{\text{eq}})$$

where the thermally averaged cross section:

$$\langle \sigma_{\chi\bar{\chi} \rightarrow ij} v_{\text{rel}} \rangle^{\text{eq}} = -\frac{h_\chi^2}{n_\chi^{\text{eq}} n_{\bar{\chi}}^{\text{eq}}} \int \frac{d^3 \vec{p}_\chi}{(2\pi)^3} \frac{d^3 \vec{p}_{\bar{\chi}}}{(2\pi)^3} \sigma_{\chi\bar{\chi} \rightarrow ij} v_{\text{rel}} f_\chi^{\text{eq}} f_{\bar{\chi}}^{\text{eq}}$$

Critical assumption:
 kinetic equilibrium at chemical decoupling

$$f_\chi \sim a(T) f_\chi^{\text{eq}}$$

*assumptions for using Boltzmann eq:
 classical limit, molecular chaos,...

...for derivation from thermal QFT
 see e.g., 1409.3049

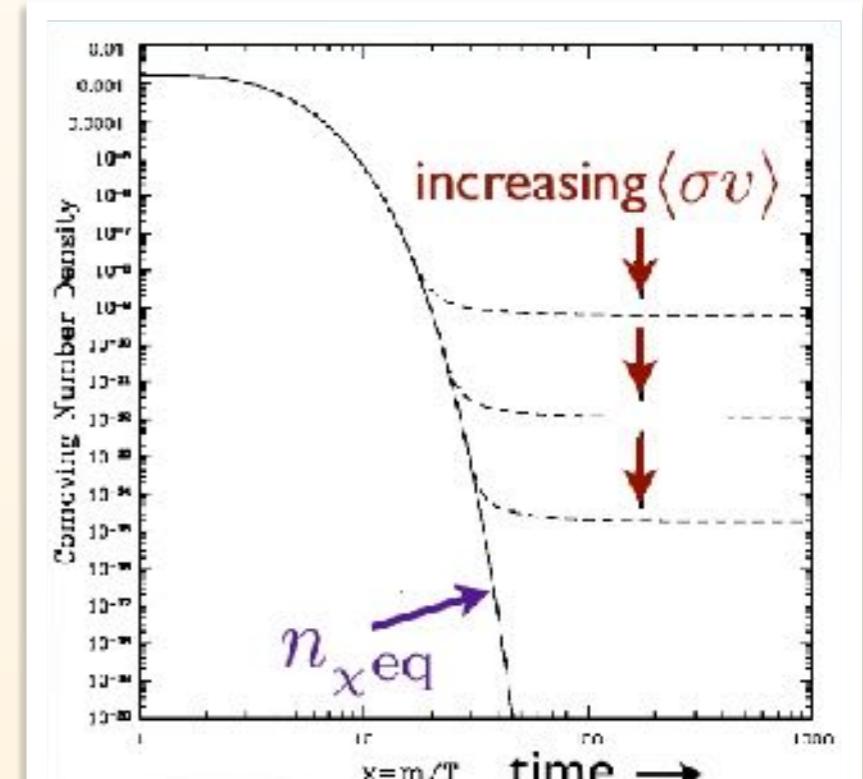
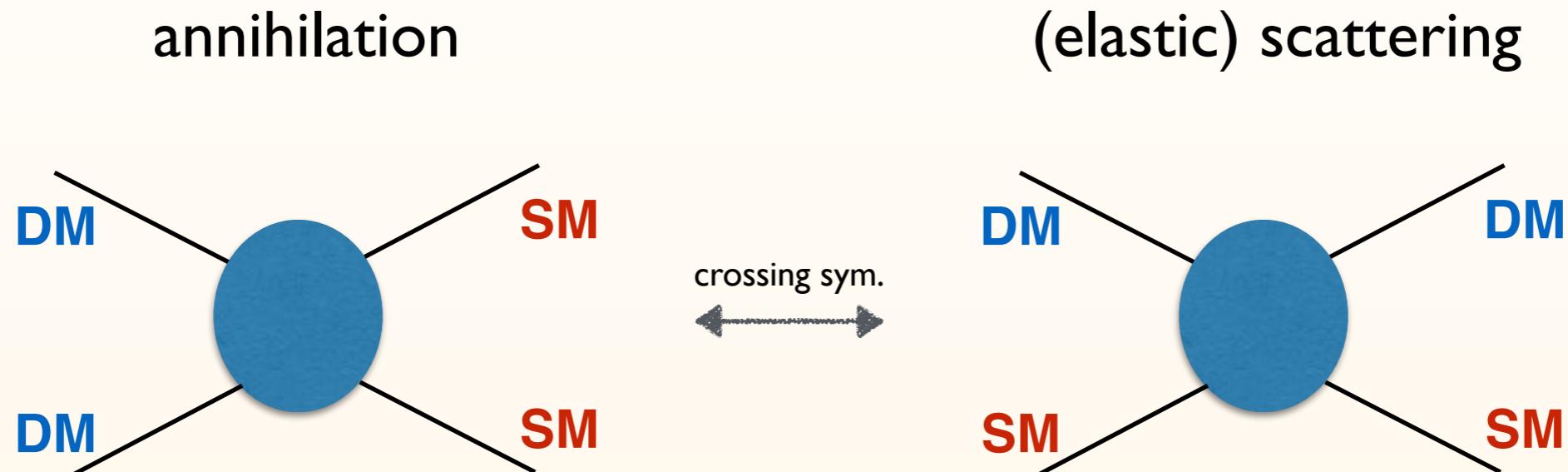


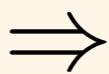
Fig.: Jungman, Kamionkowski & Griest, PR'96

FREEZE-OUT VS. DECOUPLING



$$\sum_{\text{spins}} |\mathcal{M}^{\text{pair}}|^2 = F(p_1, p_2, p'_1, p'_2) \quad \sim \quad \sum_{\text{spins}} |\mathcal{M}^{\text{scatt}}|^2 = F(k, -k', p', -p)$$

Boltzmann suppression of **DM** vs. **SM**



⇒ scatterings typically more frequent

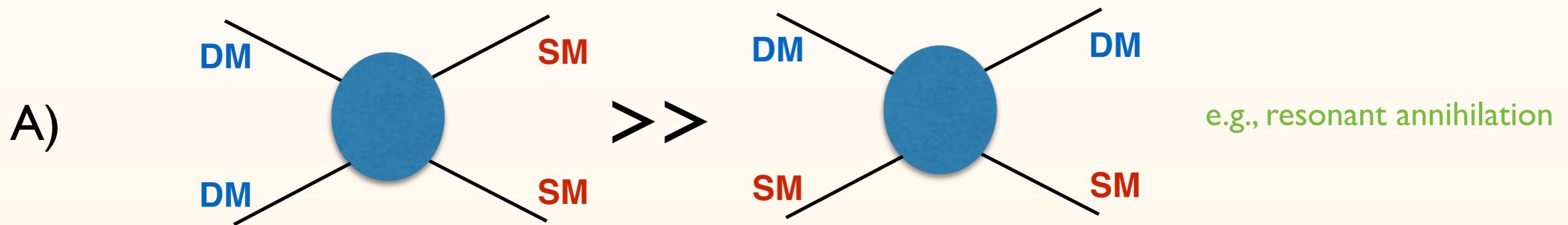
dark matter frozen-out but typically
still kinetically coupled to the plasma

Schmid, Schwarz, Widern '99; Green, Hofmann, Schwarz '05

EARLY KINETIC DECOUPLING?

A **necessary** and **sufficient** condition: scatterings weaker than annihilation
i.e. rates around freeze-out: $H \sim \Gamma_{\text{ann}} \gtrsim \Gamma_{\text{el}}$

Possibilities:



- B) Boltzmann suppression of **SM** as strong as for **DM**
e.g., below threshold annihilation (forbidden-like DM)

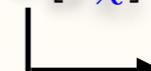
- C) Scatterings and annihilation have different structure
e.g., semi-annihilation, 3 to 2 models,...

- D) Multi-component dark sectors
e.g., additional sources of DM from late decays, ...

HOW TO GO BEYOND KINETIC EQUILIBRIUM?

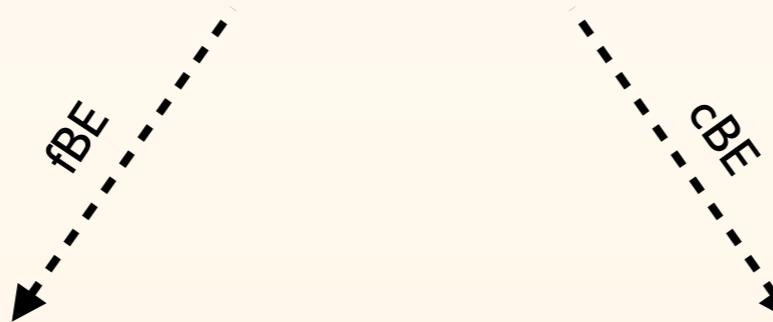
All information is in the full BE:
both about chemical ("normalization") and
kinetic ("shape") equilibrium/decoupling

$$E (\partial_t - H \vec{p} \cdot \nabla_{\vec{p}}) f_\chi = \mathcal{C}[f_\chi]$$



contains both **scatterings** and
annihilations

Two possible approaches:



solve numerically
for full $f_\chi(p)$

have insight on the distribution
no constraining assumptions

numerically challenging
often an overkill

consider system of equations
for moments of $f_\chi(p)$

partially analytic/much easier numerically
manifestly captures all of the relevant physics

finite range of validity
no insight on the distribution

0-th moment: n_χ
2-nd moment: T_χ
...

NEW TOOL!

GOING BEYOND THE STANDARD APPROACH

- Home
- Downloads
- Contact



Dark matter Relic Abundance beyond Kinetic Equilibrium

Authors: Tobias Binder, Torsten Bringmann, Michael Gustafsson and Andrzej Hryczuk

DRAKE is a numerical precision tool for predicting the dark matter relic abundance also in situations where the standard assumption of kinetic equilibrium during the freeze-out process may not be satisfied. The code comes with a set of three dedicated Boltzmann equation solvers that implement, respectively, the traditionally adopted equation for the dark matter number density, fluid-like equations that couple the evolution of number density and velocity dispersion, and a full numerical evolution of the phase-space distribution. The code is written in Wolfram Language and includes a Mathematica notebook example program, a template script for terminal usage with the free Wolfram Engine, as well as several concrete example models.
DRAKE is a free software licensed under GPLv3.

If you use DRAKE for your scientific publications, please cite

- DRAKE: Dark matter Relic Abundance beyond Kinetic Equilibrium,
Tobias Binder, Torsten Bringmann, Michael Gustafsson and Andrzej Hryczuk, [[arXiv:2103.01944](#)]

Currently, an user guide can be found in the Appendix A of this reference.
Please cite also quoted other works applying for specific cases.

v1.0 « [Click here to download DRAKE](#)

(March 3, 2021)

<https://drake.hepforge.org>

Applications:

DM relic density for
any (user defined) model*

Interplay between chemical and
kinetic decoupling

Prediction for the DM
phase space distribution

Late kinetic decoupling
and impact on cosmology

see e.g., [l202.5456](#)

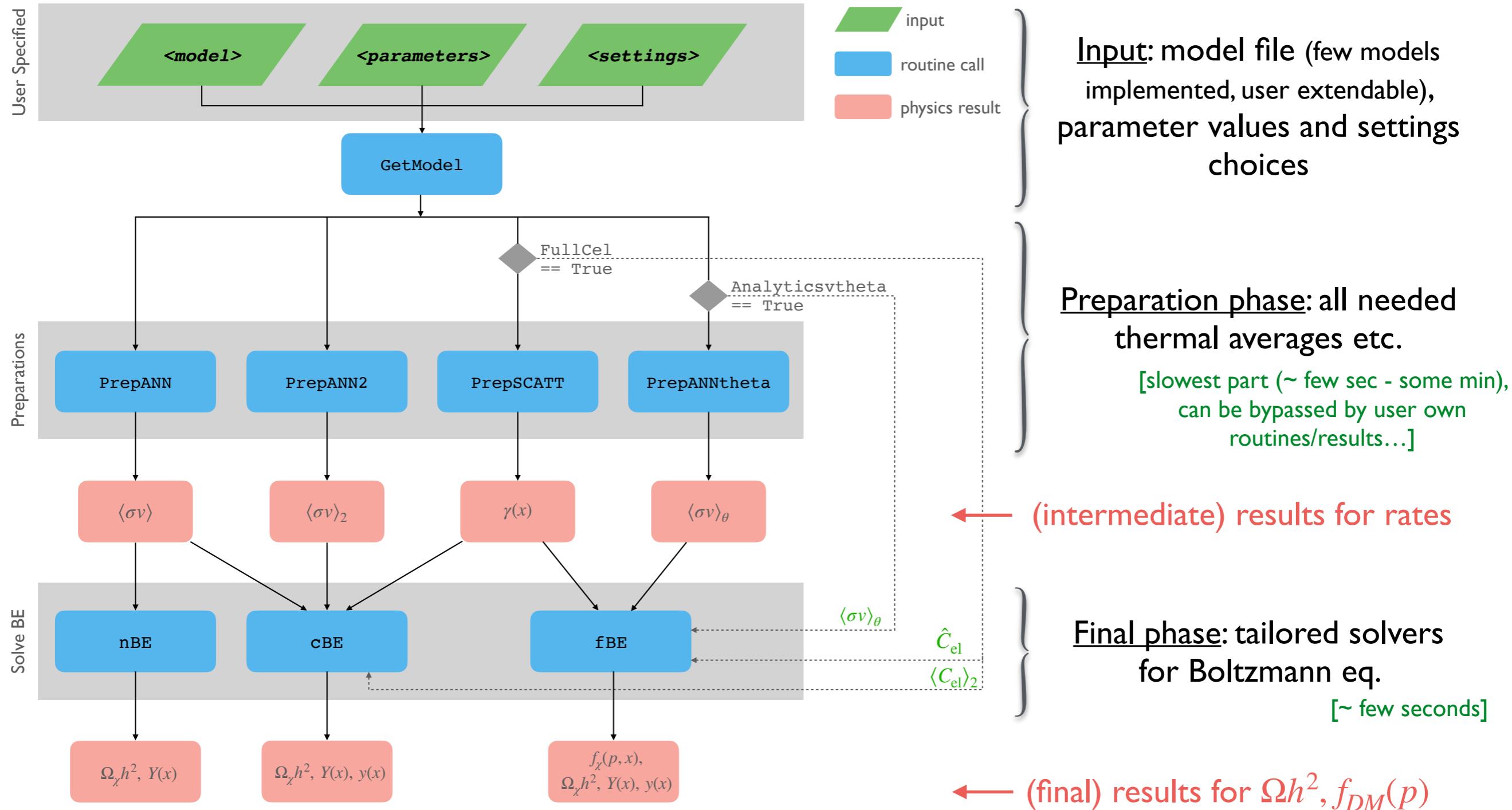
...

(only) prerequisite:
Wolfram Language (or Mathematica)

*at the moment for a single DM species and w/o
co-annihilations... but stay tuned for extensions!

FEW WORDS ABOUT THE CODE

written in *Wolfram Language*, lightweight, modular and simple to use both via script and front end usage



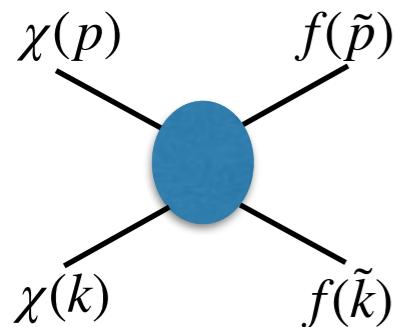
COLLISION TERM

$$E (\partial_t - H \vec{p} \cdot \nabla_{\vec{p}}) f_{\chi} = \mathcal{C}[f_{\chi}]$$



contains both scatterings and annihilations

Annihilation:



$$C_{\text{ann}} \sim \int d\tilde{\Pi} |\mathcal{M}|_{\chi\chi \rightarrow f\bar{f}}^2 \left(\underline{f_f^{\text{eq}}(\tilde{p})f_f^{\text{eq}}(\tilde{k})} - \underline{f_{\chi}(p)f_{\chi}(k)} \right)$$

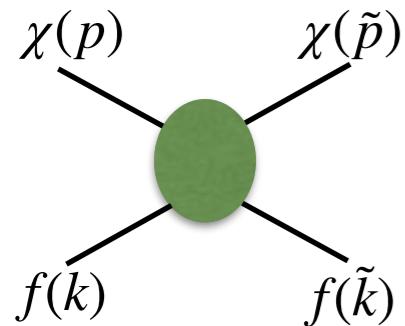
easy easy

easy: no unknown f_{χ} under integral
⇒ 1D integration

medium: no unknown f_{χ} under integral
⇒ 2-3D integration

hard: unknown f_{χ} under integral
⇒ 2-4D integration

El. scattering (on SM particles):



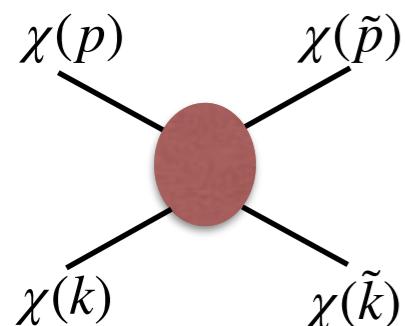
$$C_{\text{el}} \sim \int d\tilde{\Pi} |\mathcal{M}|_{\chi f \rightarrow \chi f}^2 \left(\underline{f_{\chi}(\tilde{p})f_f^{\text{eq}}(\tilde{k})(1 \pm f_f^{\text{eq}}(k))} - \underline{f_{\chi}(p)f_f^{\text{eq}}(k)(1 \pm f_f^{\text{eq}}(\tilde{k}))} \right)$$

hard medium

hard

medium

El. self-scattering (DM on DM):



$$C_{\text{self}} \sim \int d\tilde{\Pi} |\mathcal{M}|_{\chi\chi \rightarrow \chi\chi}^2 \left(\underline{f_{\chi}(\tilde{p})f_{\chi}(\tilde{k})} - \underline{f_{\chi}(p)f_{\chi}(k)} \right)$$

hard

An approximate method needed!

$$d\tilde{\Pi} = d\Pi_{\tilde{p}} d\Pi_k d\Pi_{\tilde{k}} \delta^{(4)}(\tilde{p} + p - \tilde{k} - k)$$

APPROACHES

I) Expand in „small momentum transfer”

Bringmann, Hofmann '06

$$\delta^{(3)}(\tilde{\mathbf{p}} + \tilde{\mathbf{k}} - \mathbf{p} - \mathbf{k}) \approx \sum_n \frac{1}{n!} (\mathbf{q} \nabla_{\tilde{\mathbf{p}}})^n \delta^{(3)}(\tilde{\mathbf{p}} - \mathbf{p})$$

$$M_{\text{DM}} \gg |\vec{q}| \sim T \gg m_{\text{SM}}$$

typical momentum transfer

Kasahara '09; Binder, Covi, Kamada, Murayama, Takahashi, Yoshida '16

$$f_3 \simeq f_1 + \tilde{\mathbf{q}}_i \frac{\partial f_1}{\partial \mathbf{p}_{1i}} + \frac{1}{2} \tilde{\mathbf{q}}_i \tilde{\mathbf{q}}_j \frac{\partial^2 f_1}{\partial \mathbf{p}_{1i} \partial \mathbf{p}_{1j}}$$

A.H. & S. Chatterjee, work in progress...
(on different expansion schemes)

\Rightarrow all lead to Fokker-Planck type eq.

II) Replace the backward term with a simpler one (i.e. a relaxation-like approximation)

Ala-Mattinen, Kainulainen '19

Ala-Mattinen, Heikinheimo, Kainulainen, Tuominen '22

$$\hat{C}_{\text{E},m}(p_1, t) \rightarrow -\delta f(p_1, t) \Gamma_{\text{E}}^m(p_1, t)$$

$$= (g_m(t) f_{\text{eq}}(p_1, t) - f(p_1, t)) \Gamma_{\text{E}}^m(p_1, t)$$

\Rightarrow simpler, but generally incorrect

III) Langevin simulations

Kim, Laine '23

$$(\hat{p}^i)' = -\hat{\eta} \hat{p}^i + \hat{f}^i, \quad \langle \hat{f}^i(x_1) \hat{f}^j(x_2) \rangle = \hat{\zeta} \delta^{ij} \delta(x_1 - x_2)$$

stochastic term, taking care of detailed balance

\Rightarrow perhaps promising...

IV) Fully numerical implementation

A.H. & M. Laletin [2204.07078](#) (focus on DM self-scatterings)

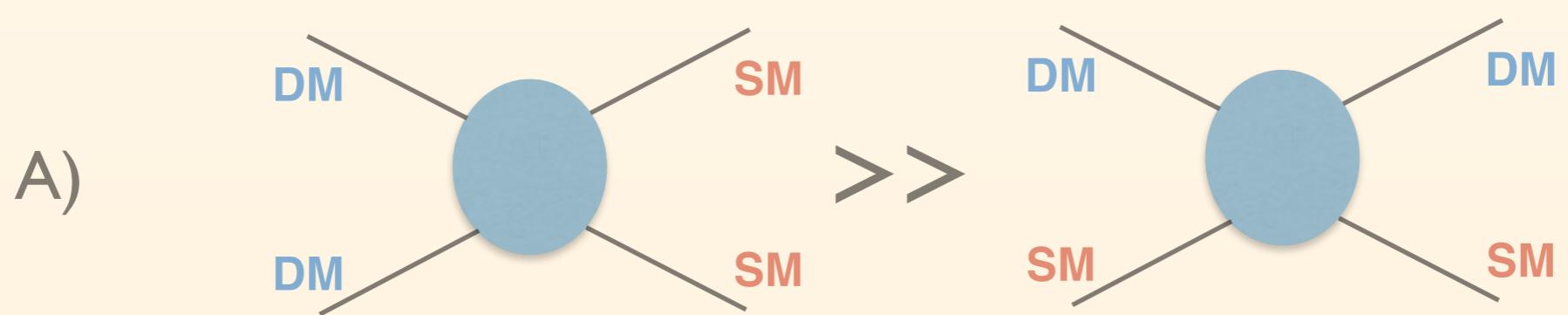
Ala-Mattinen, Heikinheimo, Kainulainen, Tuominen '22

Du, Huang, Li, Li, Yu '21

Aboubrahim, Klasen, Wiggering '23

\Rightarrow doable, but (very) CPU expensive

EXAMPLE A: SCALAR SINGLET DM



EXAMPLE A

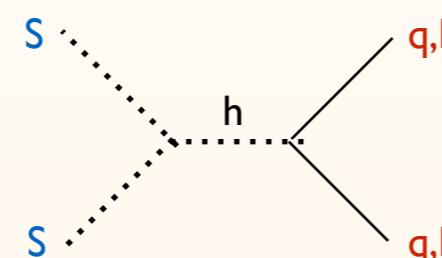
SCALAR SINGLET DM

To the SM Lagrangian add one singlet scalar field S with interactions with the Higgs:

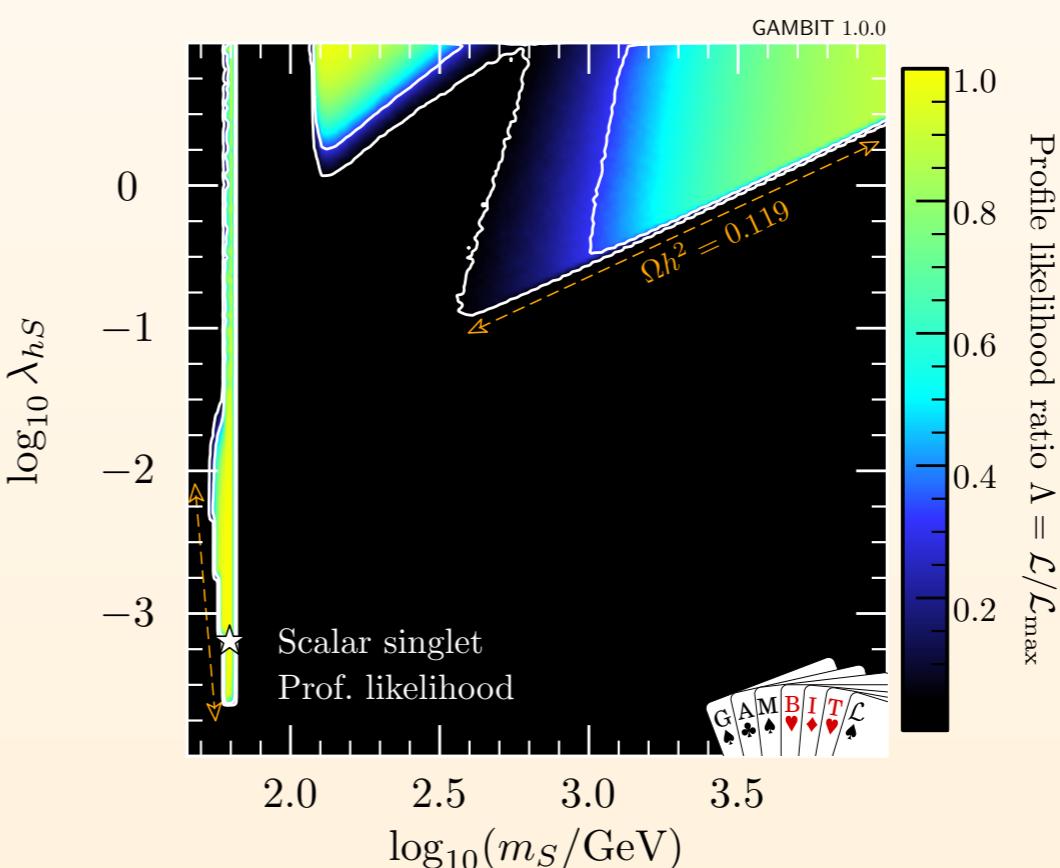
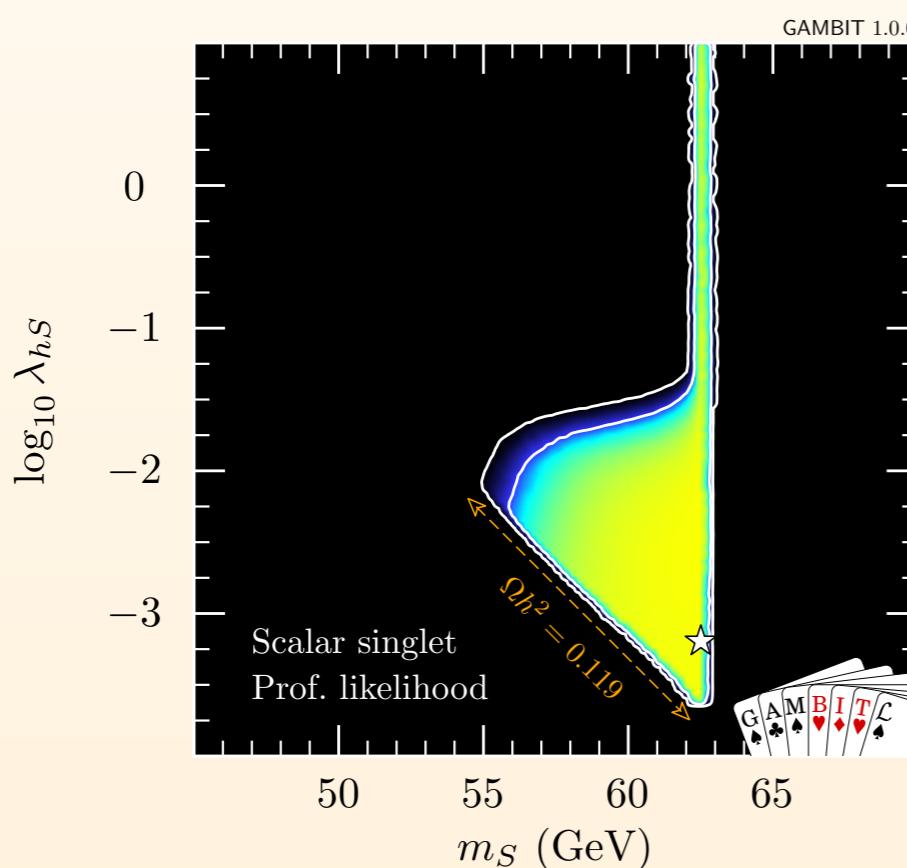
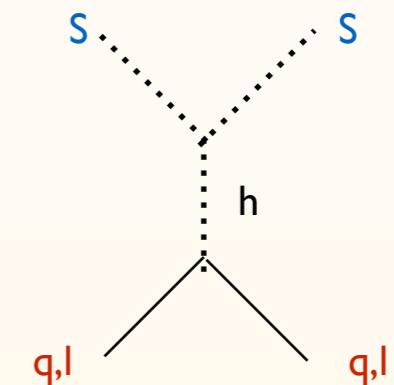
$$\mathcal{L}_S = \frac{1}{2}\partial_\mu S\partial^\mu S - \frac{1}{2}\mu_S^2 S^2 - \frac{1}{2}\lambda_s S^2 |H|^2$$

$$m_s = \sqrt{\mu_S^2 + \frac{1}{2}\lambda_s v_0^2}$$

**Annihilation
processes:
resonant**



**El. scattering
processes:
non-resonant**

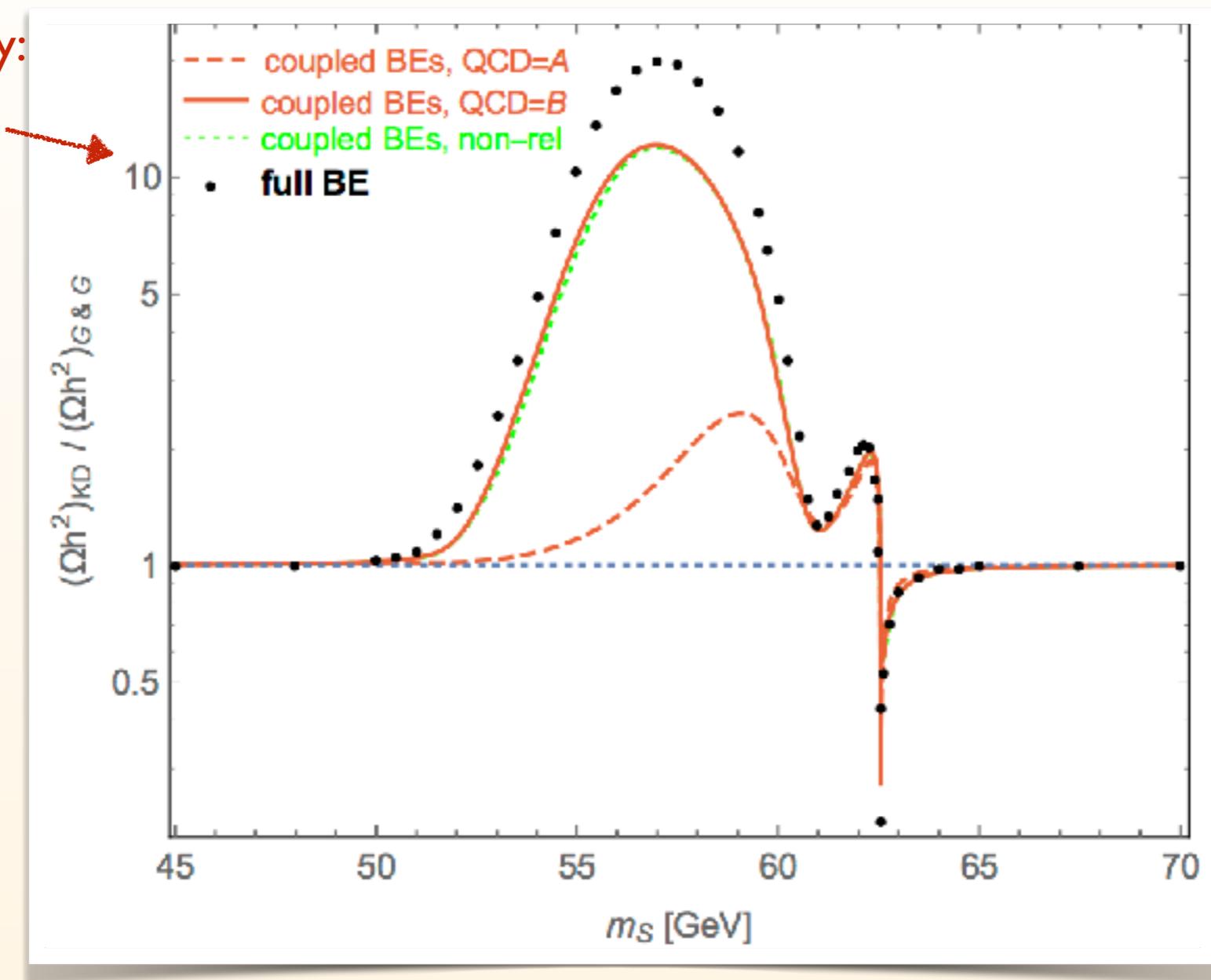


**GAMBIT collaboration
| 705.0793 |**

RESULTS

EFFECT ON THE Ωh^2

effect on relic density:
up to $O(\sim 10)$



[... Freeze-out at few GeV → what is the abundance of heavy quarks in QCD plasma?

two scenarios:

QCD = A - all quarks are free and present in the plasma down to $T_c = 154$ MeV
 QCD = B - only light quarks contribute to scattering and only down to $4T_c$...]

EXAMPLE D: WHEN ADDITIONAL INFLUX OF DM ARRIVES

D) Multi-component dark sectors

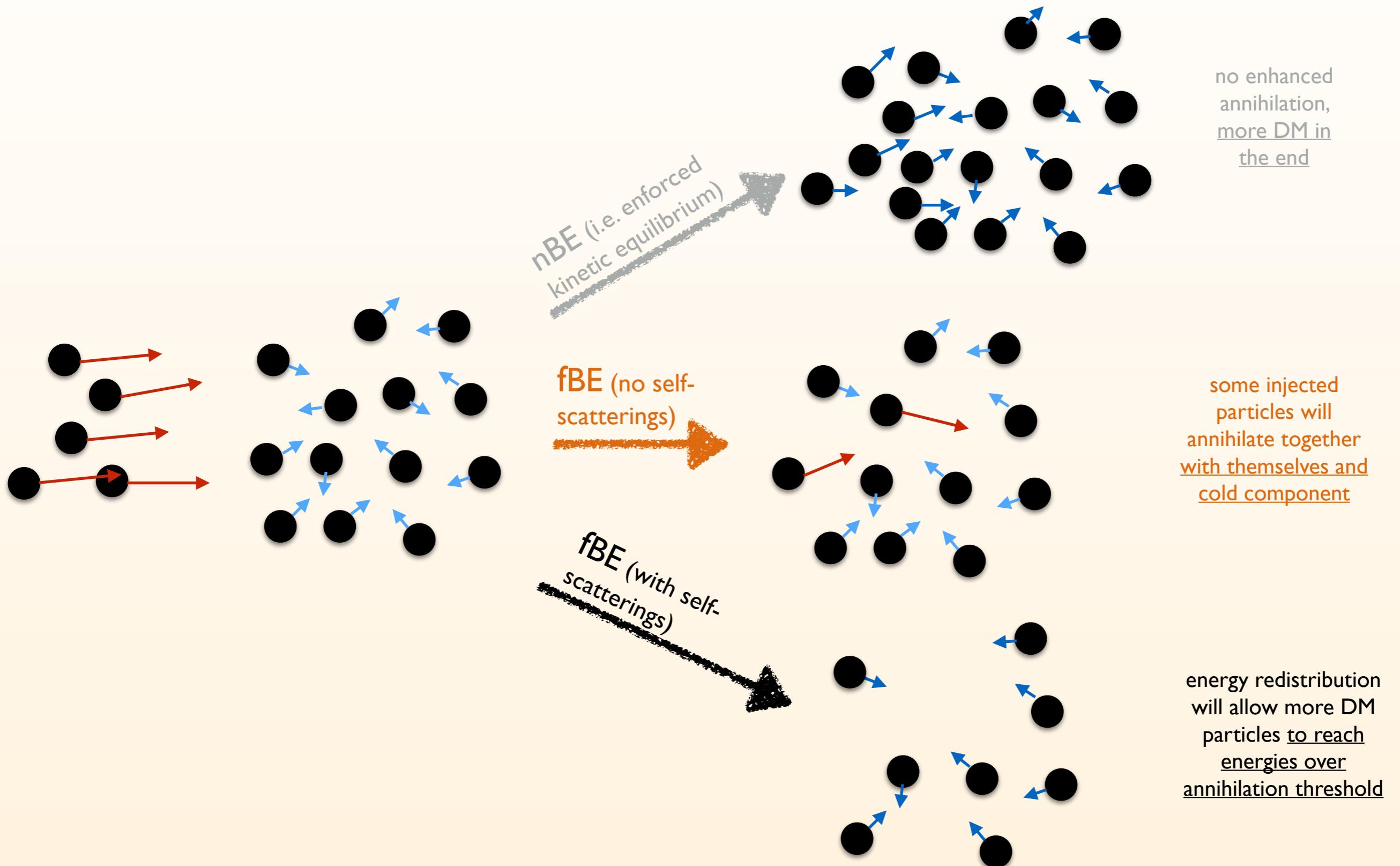
Sudden injection of more DM particles **distorts** $f_\chi(p)$
(e.g. from a decay or annihilation of other states)

- this can **modify the annihilation rate** (if still active)
- how does the **thermalization** due to elastic scatterings happen?

I) DM produced via:

- 1st component from thermal freeze-out
- 2nd component from a decay $\phi \rightarrow \bar{\chi}\chi$

2) DM annihilation has a **threshold**
e.g. $\chi\bar{\chi} \rightarrow f\bar{f}$ with $m_\chi \lesssim m_f$



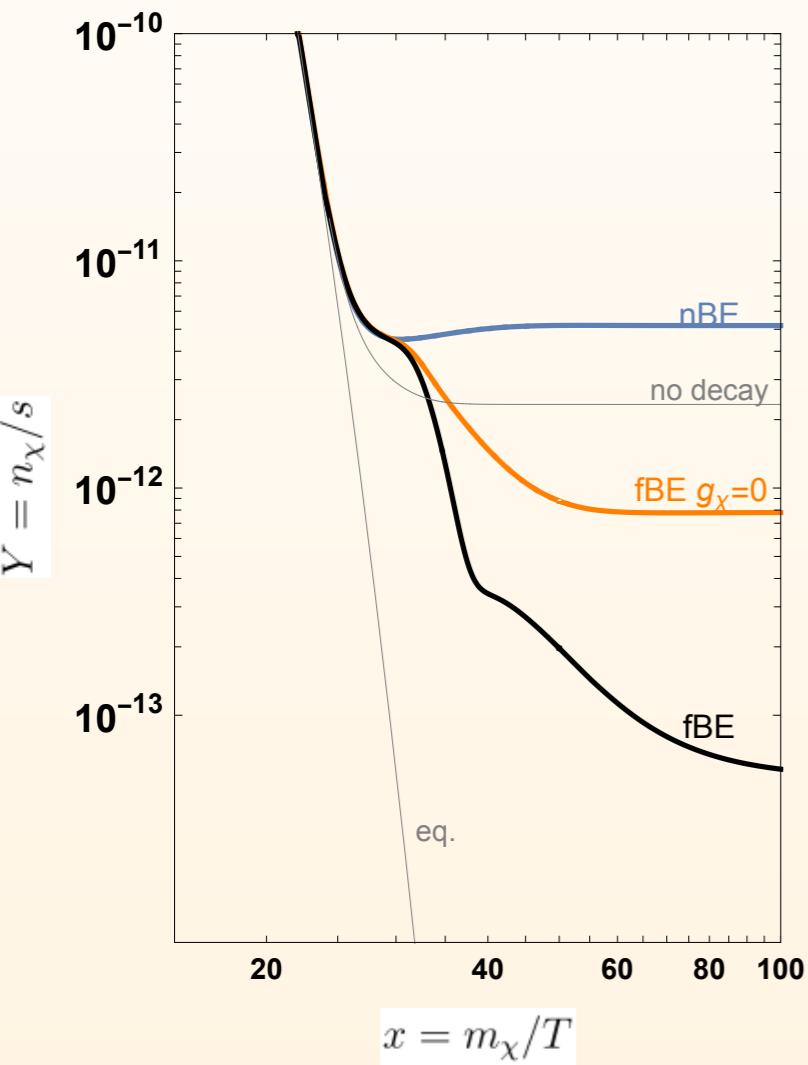
EXAMPLE EVOLUTION

I) DM produced via:

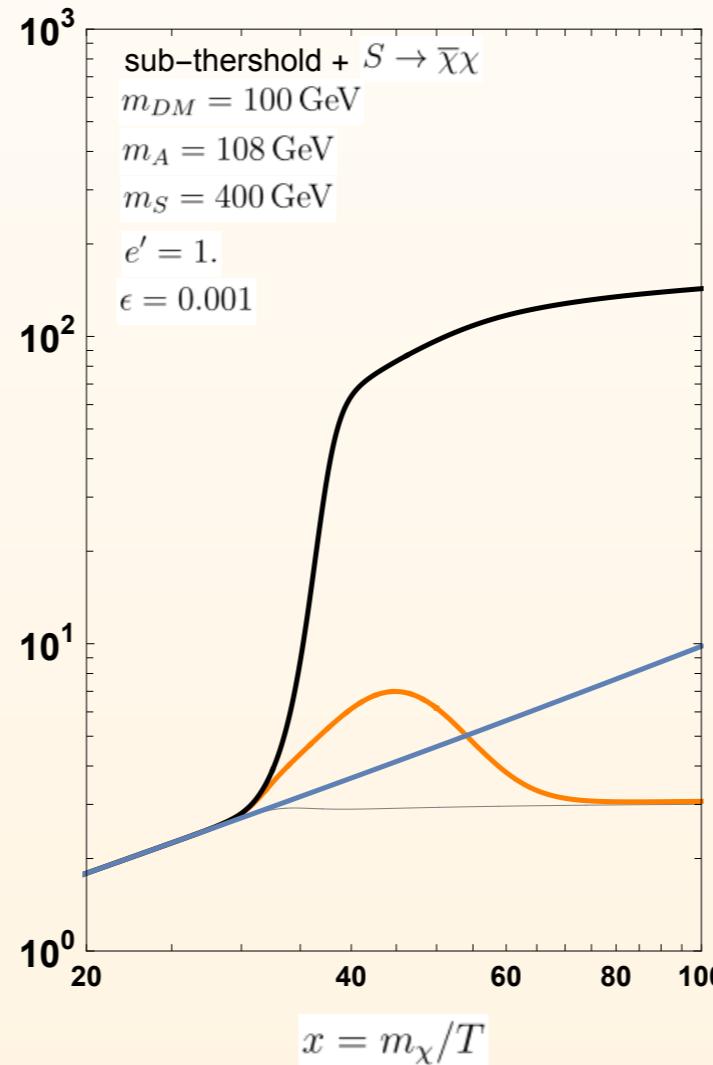
- 1st component from thermal freeze-out
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2) DM annihilation has a threshold
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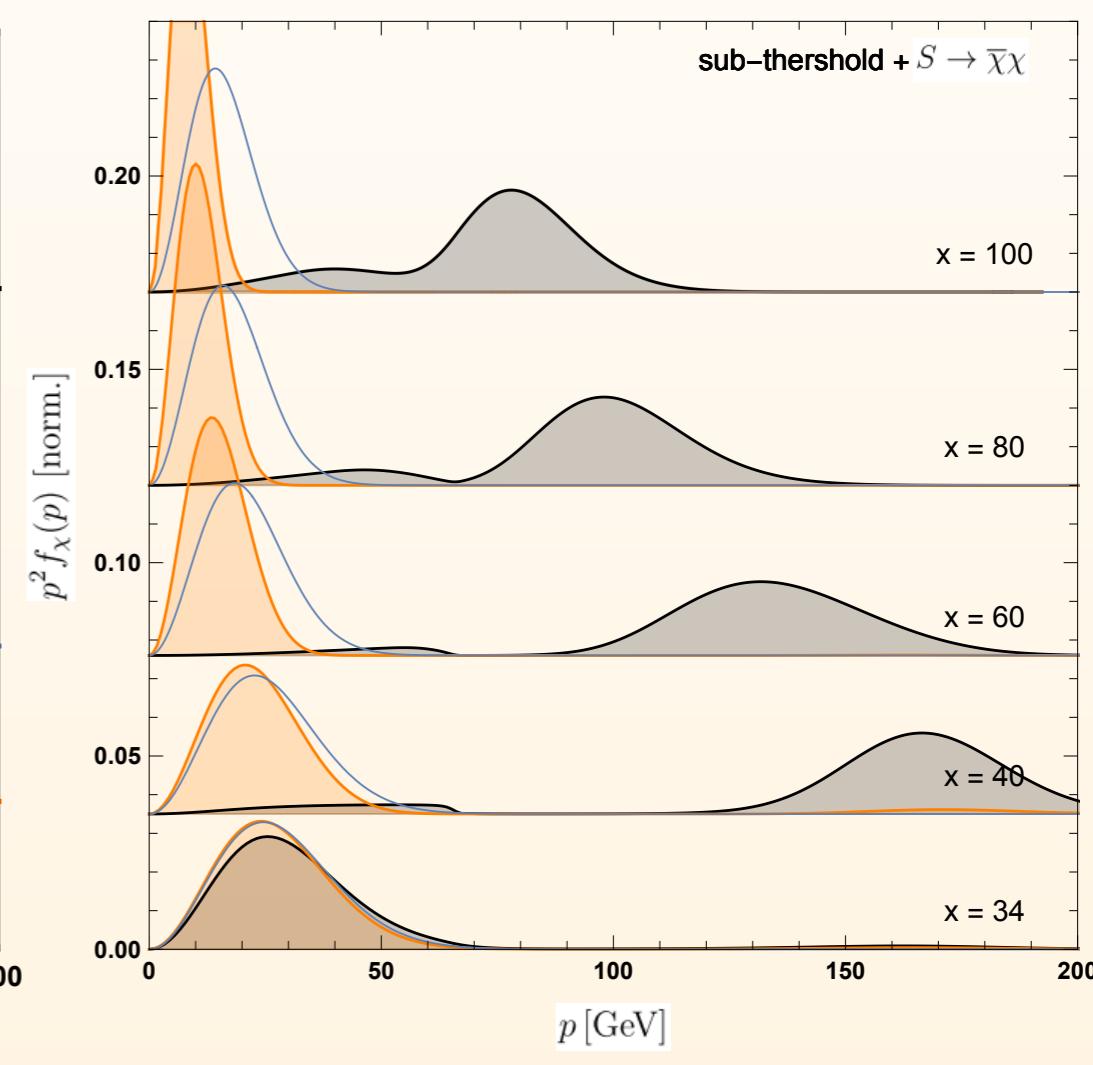
$Y \sim$ number density



$y \sim$ temperature



$p^2 f(p) \sim$ momentum distribution



SUMMARY

1. In recent years a **significant progress** in refining the relic density calculations (not yet fully implemented in public codes!)
2. **Kinetic equilibrium** is a necessary (often implicit) assumption for standard relic density calculations in all the numerical tools...
...while it is not always warranted!
3. Introduced coupled **system of Boltzmann eqs.** for **0th and 2nd moments** (cBE) allows for much more accurate treatment while the **full phase space Boltzmann equation** (fBE) can be also successfully solved for higher precision and/or to obtain result for $f_{\text{DM}}(p)$
(we also introduced **DRAKE**  a new tool to extend the current capabilities to the regimes **beyond kinetic equilibrium**)
4. **Multi-component sectors**, when studied at the fBE level, can reveal quite unexpected behavior