



U.S. DEPARTMENT OF  
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Nuclear Computational Low-Energy Initiative

# Eigenvector Continuation Emulators for Chiral EFT

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The Ohio State University

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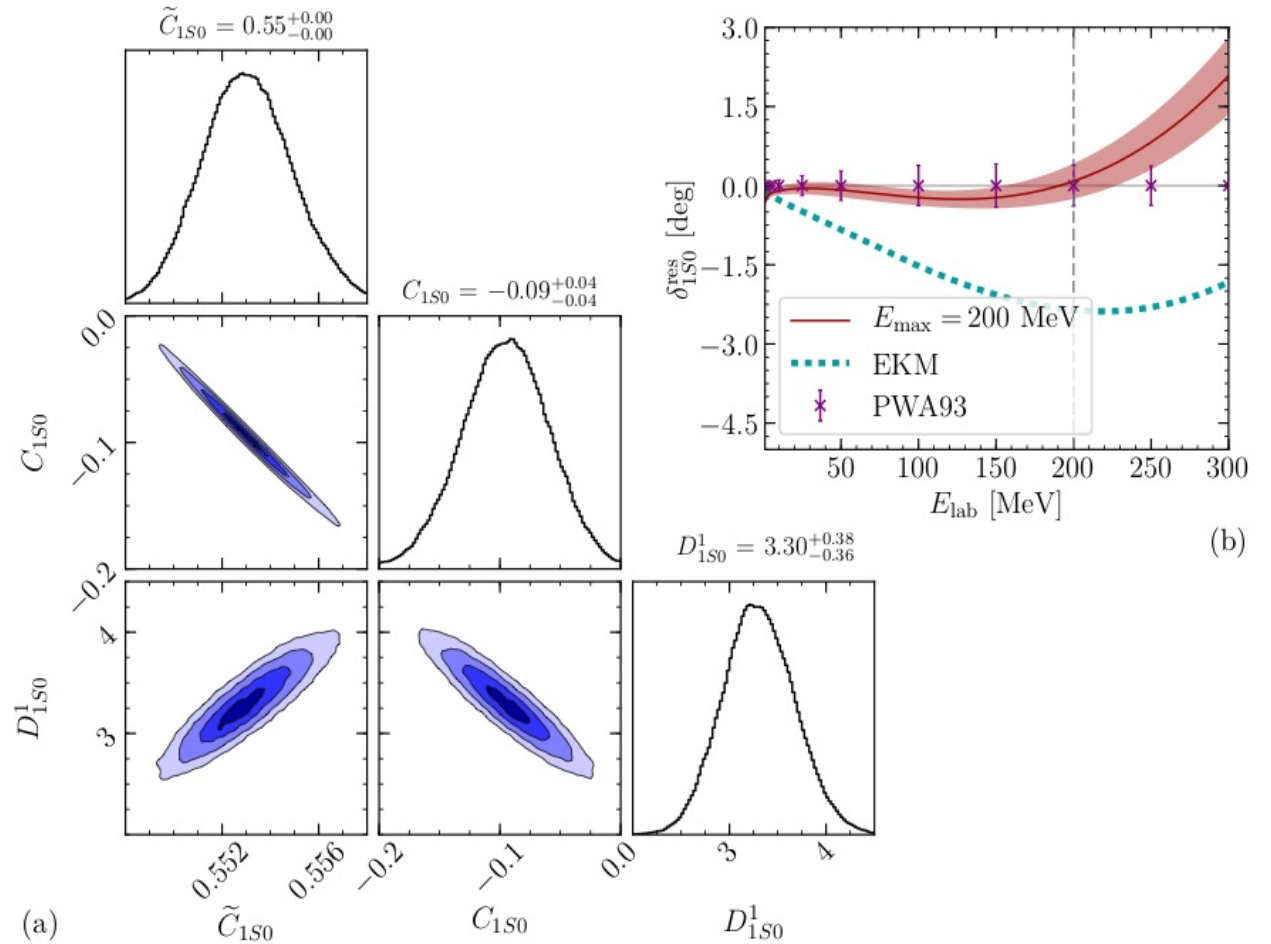
Collaborators: R.J. Furnstahl, J.A. Melendez, C. Drischler, Xilin Zhang

*J.A. Melendez, C. Drischler, ajg, R.J. Furnstahl, and Xilin Zhang,*

*arXiv: 2106.15608 , Phys. Lett. B 821, 136608 (2021)*

# Phase shifts for nucleon-nucleon (NN) scattering with UQ

- Full sampling for Bayesian UQ can be expensive using direct calculations
- Alternative: sample from a previously trained computer model
- Linear parameter dependence in  $\chi$ EFT allows for fast calculations



*Wesolowski et al., J. Phys. G 46, 045102 (2019)*

# Chiral EFT potentials for NN scattering

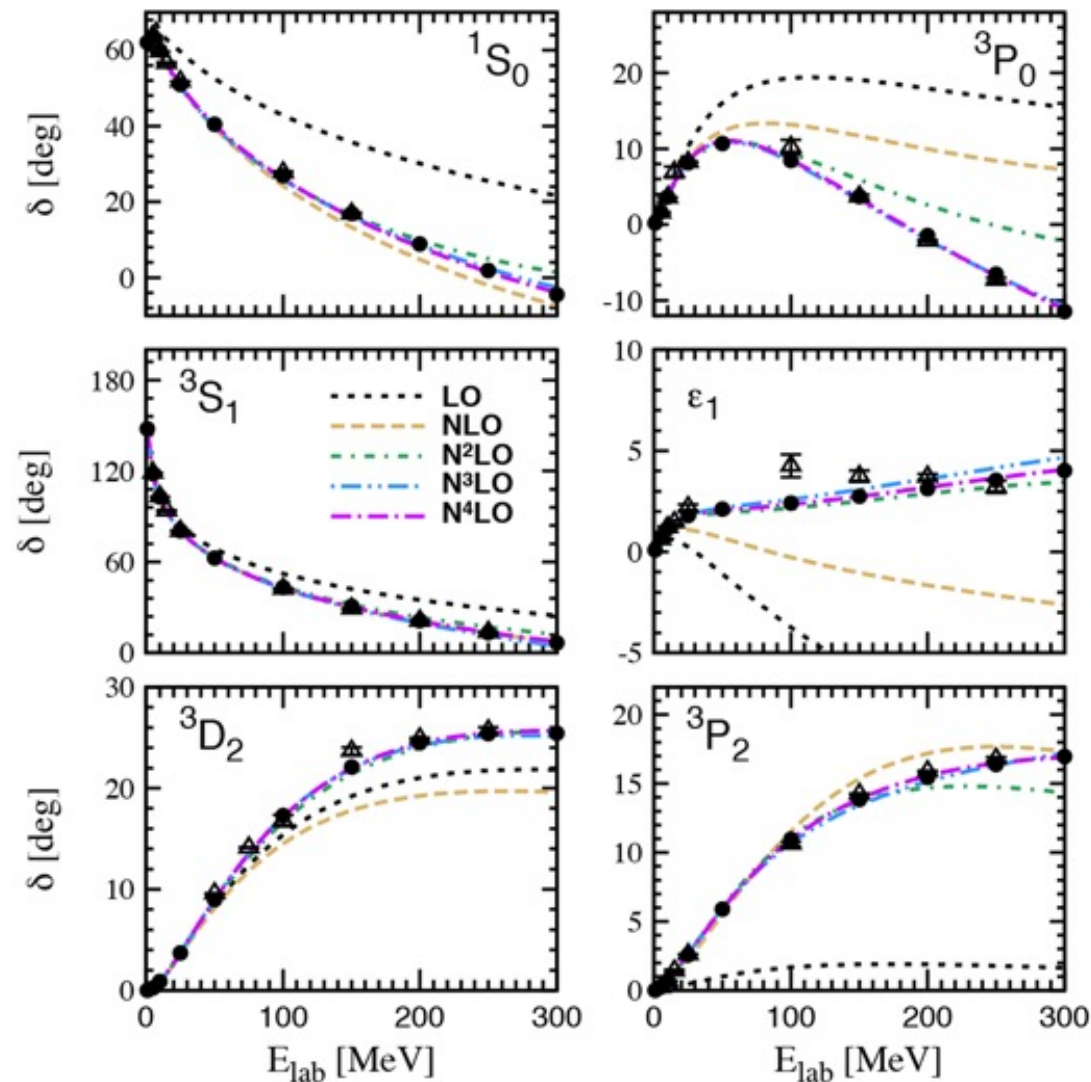
*P. Reinert et al., Eur. Phys. J B 54, 86 (2018)*

- Here: SMS semi-local momentum-space regularized potential
- Candidate for full Bayesian UQ using **emulators**
- Can take advantage of linearity between matrix elements and LECs:

$$V = \boxed{C_0}V^{(0)} + \boxed{C_2}V^{(2)} + \boxed{C_4}V^{(4)}$$

→ only calculate matrix elements once!

- Test emulators on neutron-proton **scattering** at cutoff  $\Lambda = 450 \text{ MeV}$



# Emulating the Lippmann-Schwinger (LS) equation

LS equation:

Sets of parameters:

K-matrix formulation:

$$K(\vec{a}) = V(\vec{a}) + V(\vec{a}) G_0(E_q) K(\vec{a}) \rightarrow \{\vec{a}_i\} \rightarrow K_\ell(E_q) = -\tan \delta_\ell(E_q)$$

$$E_q = q^2/2\mu$$

Newton variational principle (NVP):

$$\tilde{K}(\vec{\beta}) = \sum_{i=1}^{n_t} \beta_i K_i \rightarrow \mathcal{K}[\tilde{K}] = V + V G_0 \tilde{K} + \tilde{K} G_0 V - \tilde{K} G_0 \tilde{K} + \tilde{K} G_0 V G_0 \tilde{K}$$

$$\mathcal{K}[K_{\text{exact}} + \delta K] = K_{\text{exact}} + (\delta K)^2$$

Implementation:

$$\langle \phi' | \mathcal{K}(\vec{a}, \vec{\beta}) | \phi \rangle = \langle \phi' | V(\vec{a}) | \phi \rangle + \vec{\beta}^T \vec{m}(\vec{a}) - \frac{1}{2} \vec{\beta}^T M(\vec{a}) \vec{\beta}$$

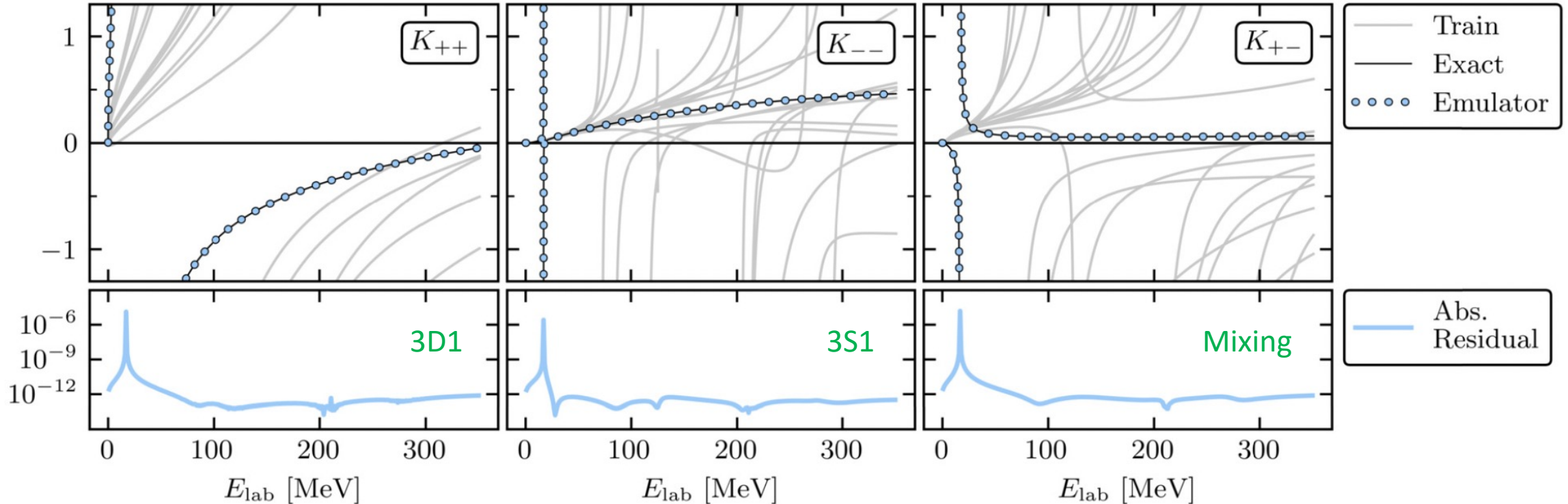
$$\frac{d\mathcal{K}}{d\vec{\beta}} = 0 \rightarrow \langle \phi' | \mathcal{K}(\vec{a}, \vec{\beta}) | \phi \rangle \approx \langle \phi' | V(\vec{a}) | \phi \rangle + \frac{1}{2} \vec{m}^T M^{-1}(\vec{a}) \vec{m}$$

*J. A. Melendez et al., Phys. Lett. B 821, 136608 (2021)*

# NVP emulation: SMS chiral potential

- Emulation of 3S1-3D1 coupled channel
- Basis size of 12 at  $N^4\text{LO}+$

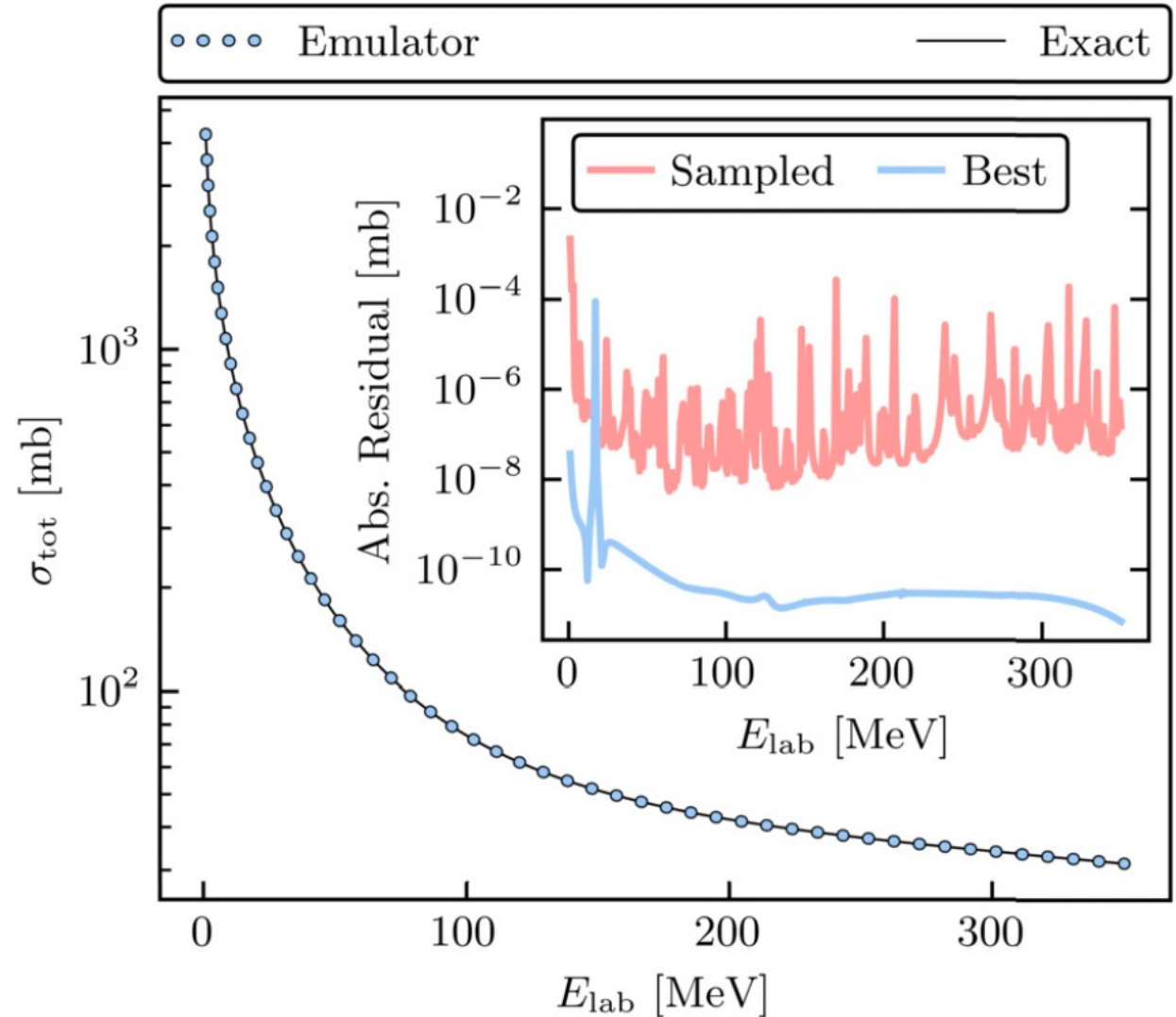
Dealing with  
anomalies/singularities:  
*C. Drischler et al.,  
arXiv: 2108.08269 (2021)*



*J. A. Melendez et al., Phys.  
Lett. B 821, 136608 (2021)*

# NVP emulation: SMS chiral potential

- Observable: Total cross section for partial waves up to  $j = 20$
- Sampling: randomly chose values in an interval of  $[-5, 5]$
- Errors: consistent for cutoffs  $\Lambda = 400 - 550 \text{ MeV}$
- **Negligible** compared to other uncertainties
- Timing: Speed up of  $> 300x$  compared to exact calculation
- $> 1000x$  if mesh size is doubled





# Eigenvector continuation (EC) for scattering

Hamiltonian:

Sets of parameters:

K-matrix formulation:

$$\hat{H}(\boldsymbol{\theta}) = \hat{T} + \hat{V}(\boldsymbol{\theta}) \quad \longrightarrow \quad \{(\boldsymbol{\theta})_i\} \quad \longrightarrow \quad \mathcal{K}_\ell(E) = \tan \delta_\ell(E)$$

Kohn variational principle (KVP):

S-wave:  $\ell = 0$   
 $p \equiv \sqrt{2\mu E}$

*R. J. Furnstahl et al.,  
 Phys. Lett. B 809,  
 135719 (2020)*

$$|\psi_{trial}\rangle \xrightarrow{r \rightarrow \infty} \frac{1}{p} \sin(pr) + \frac{[\mathcal{K}_0(E)]_{trial}}{p} \cos(pr)$$

$$\delta\beta[|\psi_{trial}\rangle] = \delta \left[ \frac{[\mathcal{K}_0(E)]_{trial}}{p} - \frac{2\mu}{\hbar^2} \langle \psi_{trial} | \hat{H}(\boldsymbol{\theta}) - E | \psi_{trial} \rangle \right] = 0 \quad \longrightarrow \quad \beta[|\psi_{exact}\rangle] = \frac{[\mathcal{K}_0(E)]_{exact}}{p}$$

EC implementation:

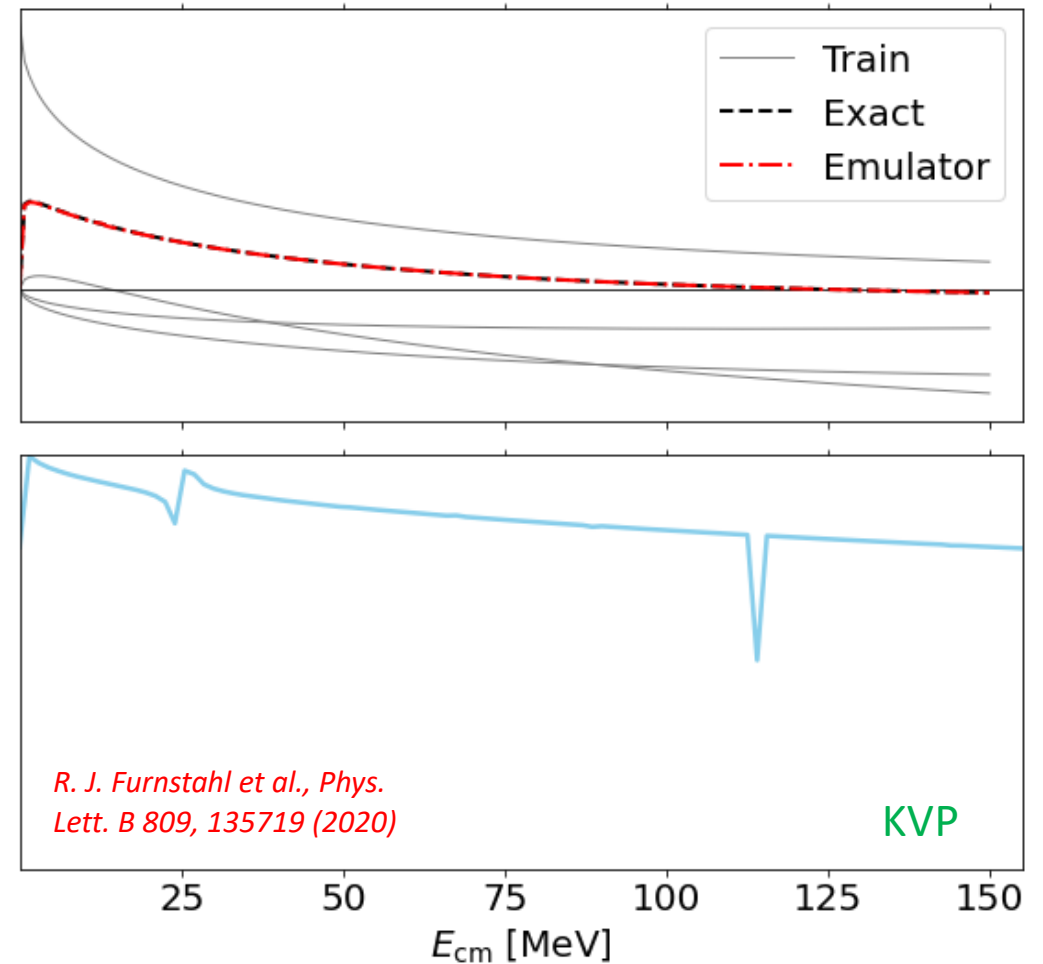
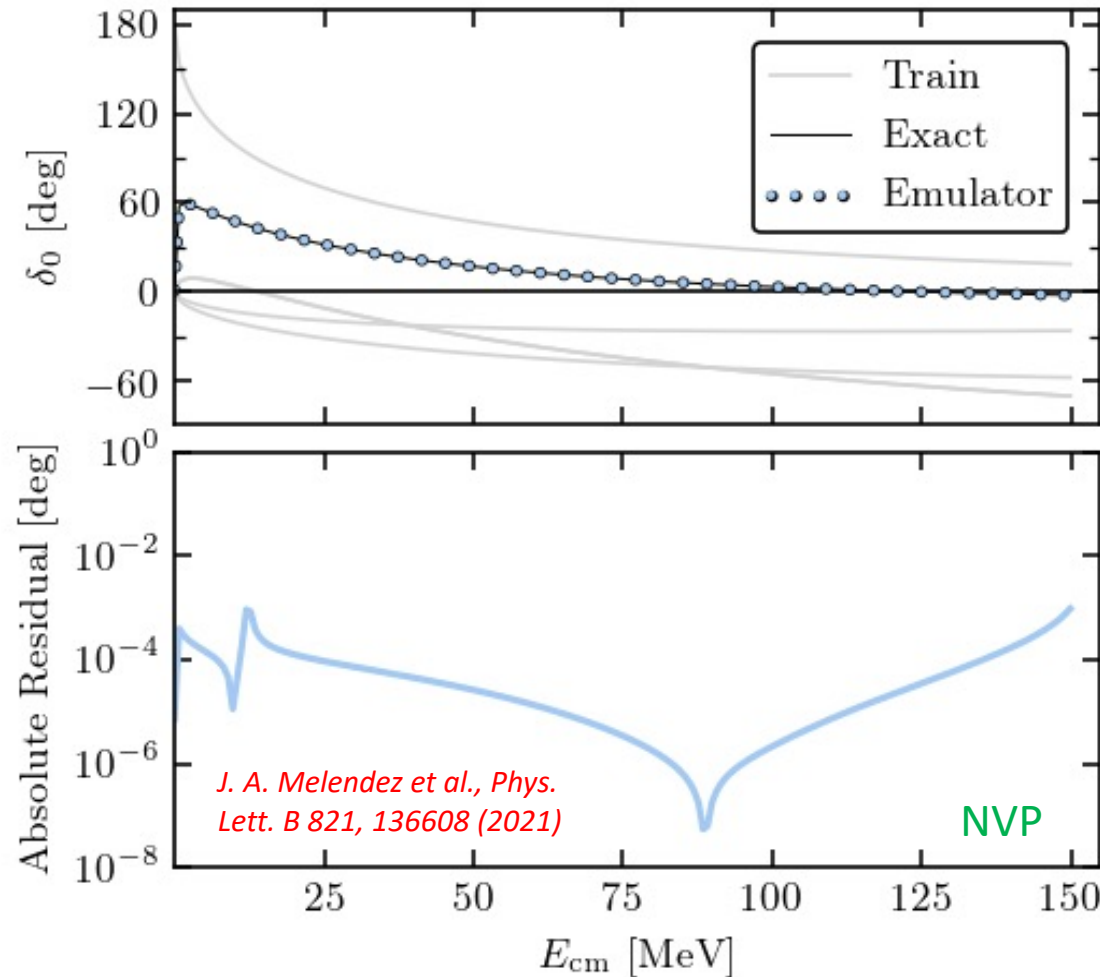
$$|\psi_{trial}\rangle = \sum_{i=1}^{N_b} c_i |\psi_E(\boldsymbol{\theta}_i)\rangle \quad \longrightarrow \quad \sum_j (\Delta U^T + \Delta U)_{ij} c_j = \sum_j \Delta \tilde{U}_{ij} c_j = \frac{\mathcal{K}_0^{(i)}(E)}{p} - \lambda$$

$$\Delta U_{ij} \equiv \frac{2\mu}{\hbar^2} \iint dk dp \psi_E(p; \theta_i) [V(k, p; \theta) - V_j(k, p)] \psi_E(k; \theta_j)$$

# Comparing both methods: Minnesota potential

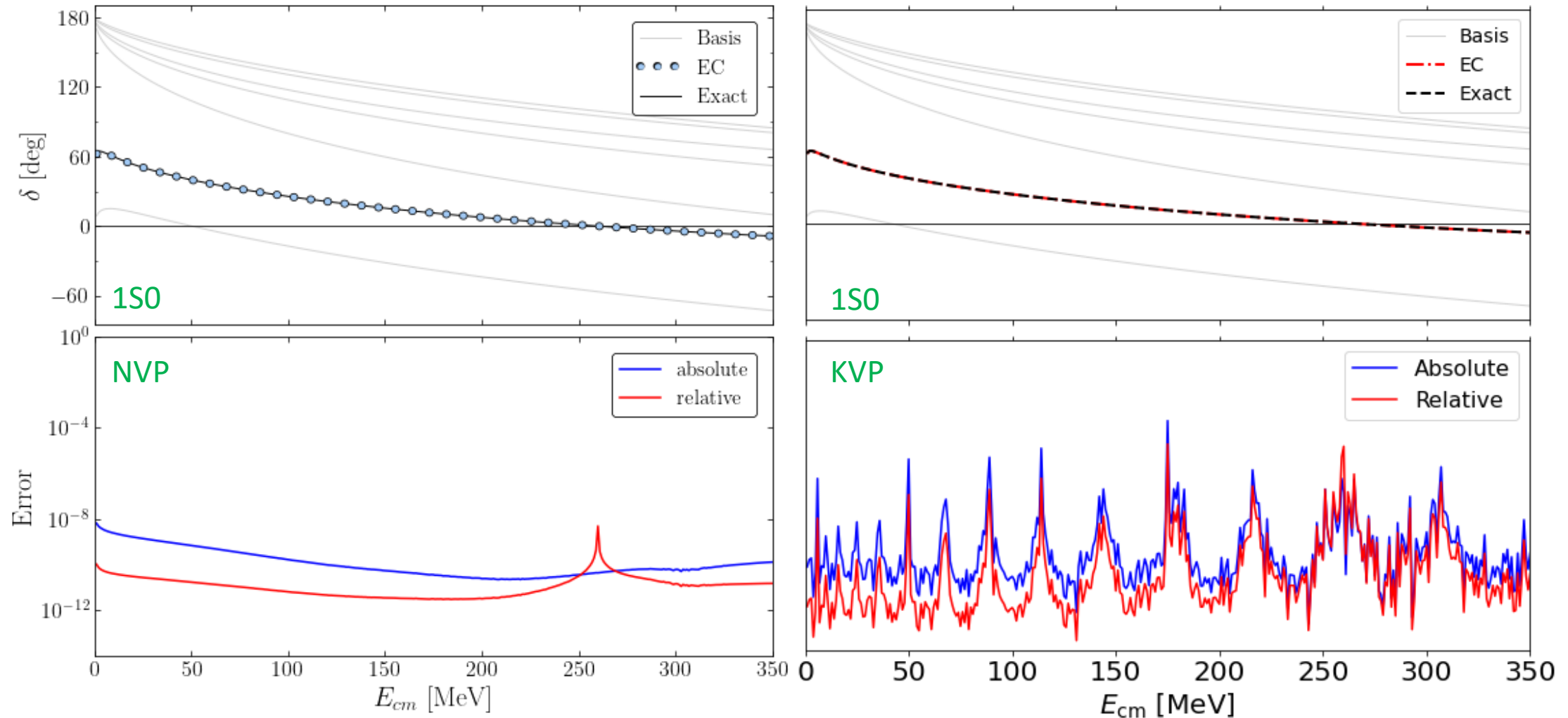
$$V(r) = V_{0R}e^{-\kappa_R r^2} + V_{0s}e^{-\kappa_s r^2}$$

$$V_{0R} = 200.0 \text{ MeV} \quad k_R = 1.487 \text{ fm}^{-2}$$
$$V_{0s} = -91.85 \text{ MeV} \quad k_s = 0.465 \text{ fm}^{-2}$$





# Comparing both methods: SMS chiral potential



# Summary

- NVP provides a general method of creating emulators for scattering
- Wave functions are not required for emulation of two-body scattering observables!
- Preliminary results show that NVP and KVP have similar accuracy

## Ongoing work

- Understand sensitivity of LECs
- Compare NVP and KVP emulation of observables
- Compare timing between NVP and KVP
- EC applications to three-body scattering (Xilin Zhang)
- What else can be emulated?

Thank you!