



U.S. DEPARTMENT OF  
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Nuclear Computational Low-Energy Initiative

# Eigenvector Continuation Emulators for NN Scattering

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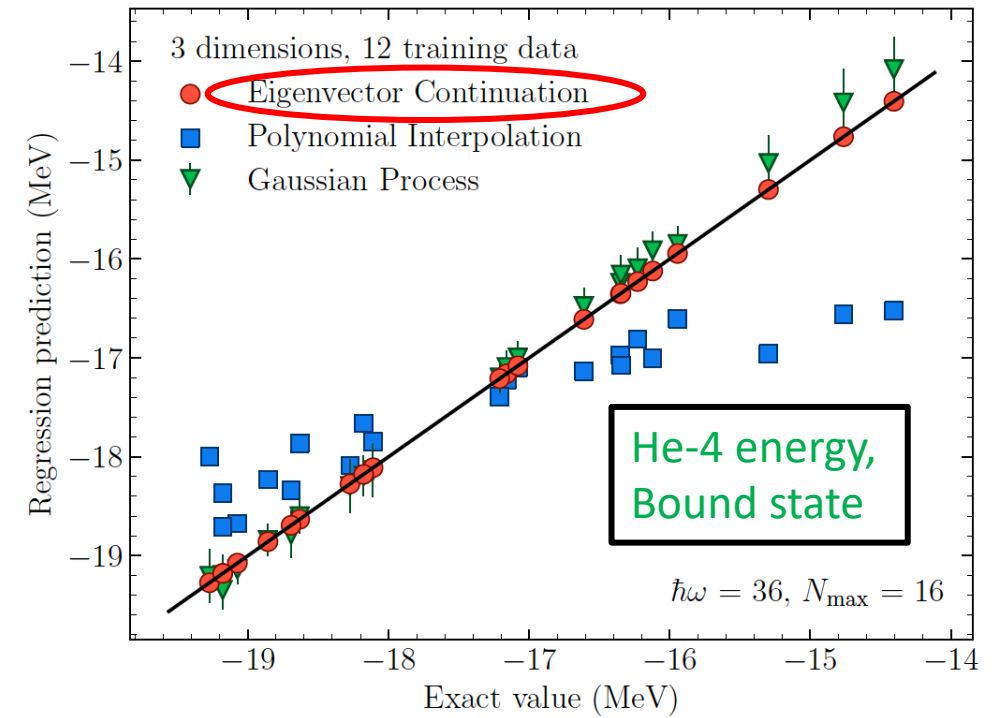
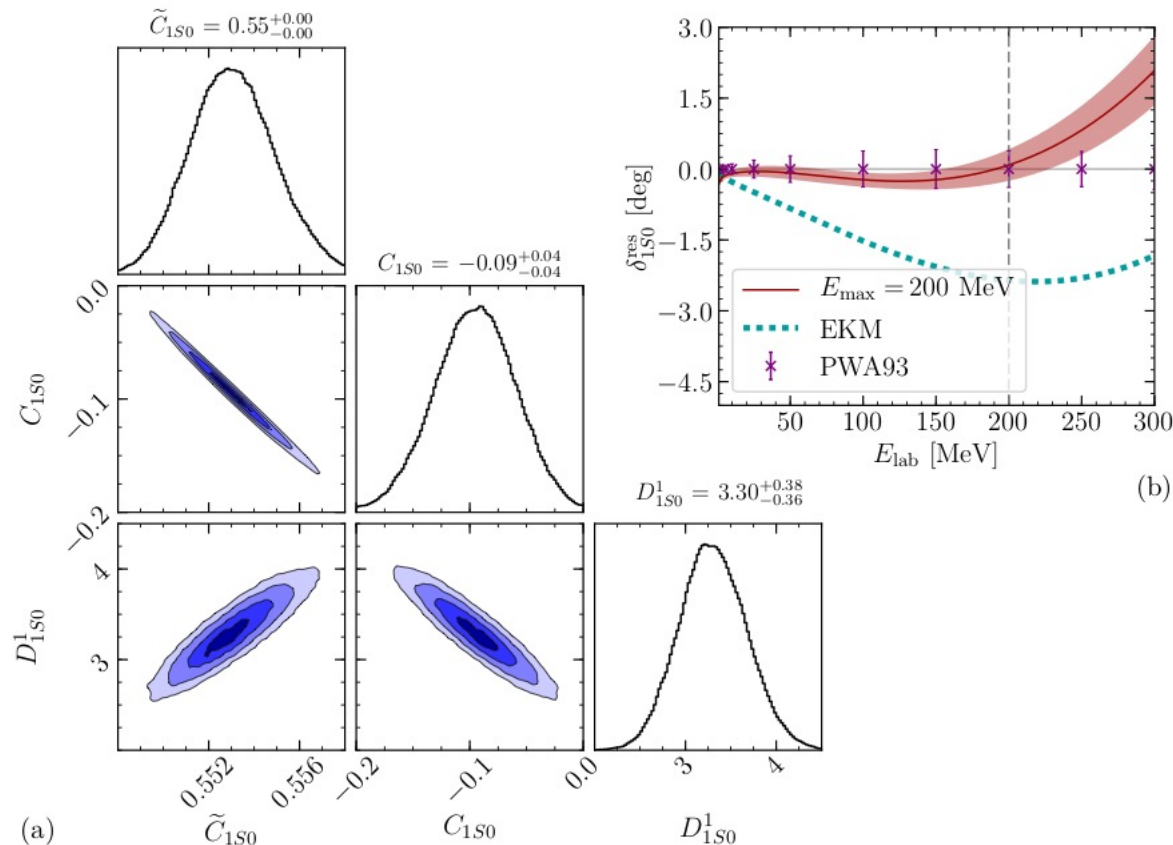
The Ohio State University

(Virtual) APS April Meeting 2021

Collaborators: R.J. Furnstahl, P.J. Millican, Xilin Zhang

# Phase shifts for NN scattering with UQ

- Full sampling for Bayesian UQ can be expensive using direct calculations
- Alternative: sample from a previously trained computer model
- Linear parameter dependence in  $\chi$ EFT allows for fast calculations



König et al, *Phys. Lett. B* 810, 135814 (2020)

Wesolowski et al., *J. Phys. G* 46, 045102 (2019)

# Chiral EFT potentials for NN scattering

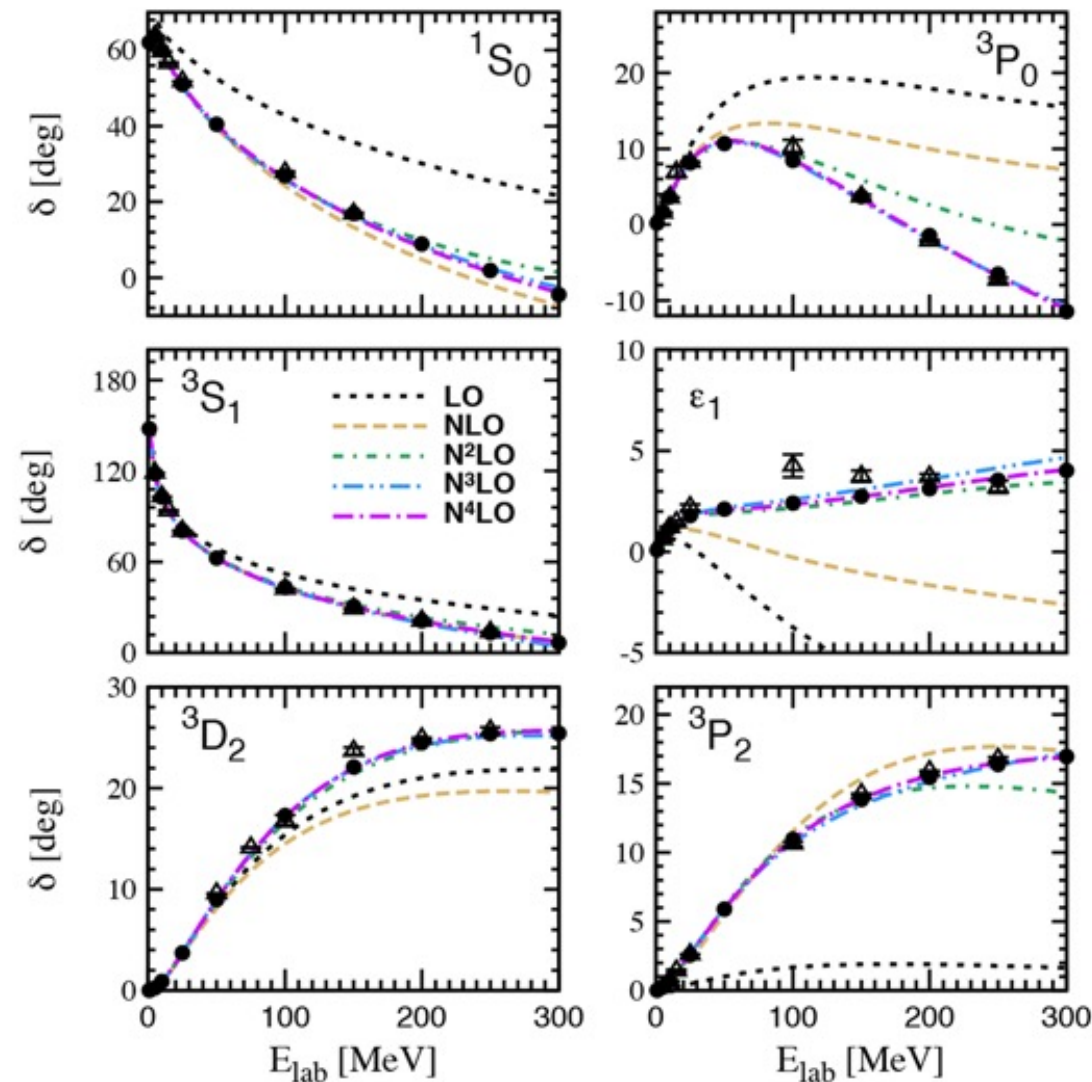
*P. Reinert et al., Eur. Phys. J B 54, 86 (2018)*

- Here: RKE semi-local momentum-space regularized potential
- Candidate for full Bayesian UQ using eigenvector continuation (EC)
- Can take advantage of linearity between matrix elements and LECs:

$$V = \boxed{C_0} V^{(0)} + \boxed{C_2} V^{(2)} + \boxed{C_4} V^{(4)}$$

→ only calculate matrix elements once!

- Test EC neutron-proton *scattering* for the  $^1S_0$  channel at cutoff  $\Lambda = 450 \text{ MeV}$



# Eigenvector continuation (EC) for scattering

Hamiltonian:

Sets of parameters:

K-matrix formulation:

$$\hat{H}(\boldsymbol{\theta}) = \hat{T} + \hat{V}(\boldsymbol{\theta}) \quad \longrightarrow \quad \{(\boldsymbol{\theta})_i\} \quad \longrightarrow \quad \mathcal{K}_\ell(E) = \tan \delta_\ell(E)$$

Kohn variational principle (KVP):

S-wave:  $\ell = 0$

$$p \equiv \sqrt{2\mu E}$$

$$|\psi_{trial}\rangle \xrightarrow{r \rightarrow \infty} \frac{1}{p} \sin(pr) + \frac{[\mathcal{K}_0(E)]_{trial}}{p} \cos(pr)$$

$$\delta\beta[|\psi_{trial}\rangle] = \delta \left[ \frac{[\mathcal{K}_0(E)]_{trial}}{p} - \frac{2\mu}{\hbar^2} \langle \psi_{trial} | \hat{H}(\boldsymbol{\theta}) - E | \psi_{trial} \rangle \right] = 0 \quad \longrightarrow \quad \beta[|\psi_{exact}\rangle] = \frac{[\mathcal{K}_0(E)]_{exact}}{p}$$

EC implementation:

$$|\psi_{trial}\rangle = \sum_{i=1}^{N_b} c_i |\psi_E(\boldsymbol{\theta}_i)\rangle \quad \longrightarrow \quad \sum_j (\Delta U^T + \Delta U)_{ij} c_j = \sum_j \Delta \tilde{U}_{ij} c_j = \frac{\mathcal{K}_0^{(i)}(E)}{p} - \lambda$$

$$\Delta \tilde{U}_{ij} = \frac{2\mu}{\hbar^2} \langle \psi_E(\boldsymbol{\theta}_i) | 2\hat{V}(\boldsymbol{\theta}) - \hat{V}(\boldsymbol{\theta}_i) - \hat{V}(\boldsymbol{\theta}_j) | \psi_E(\boldsymbol{\theta}_j) \rangle$$

*R. J. Furnstahl et al., Phys. Lett. B 809, 135719 (2020)*

- Simple **matrix inversion** + **cancellation** of Coulomb force!
- Stationary approximation to K-matrix (not an upper/lower bound)

# Momentum space implementation

- **EC matrix** can be calculated using **momentum space wave function**

$$\Delta \tilde{U}_{ij} = \frac{2\mu}{\hbar^2} \langle \psi_E(\boldsymbol{\theta}_i) | 2\hat{V}(\boldsymbol{\theta}) - \hat{V}(\boldsymbol{\theta}_i) - \hat{V}(\boldsymbol{\theta}_j) | \psi_E(\boldsymbol{\theta}_j) \rangle$$

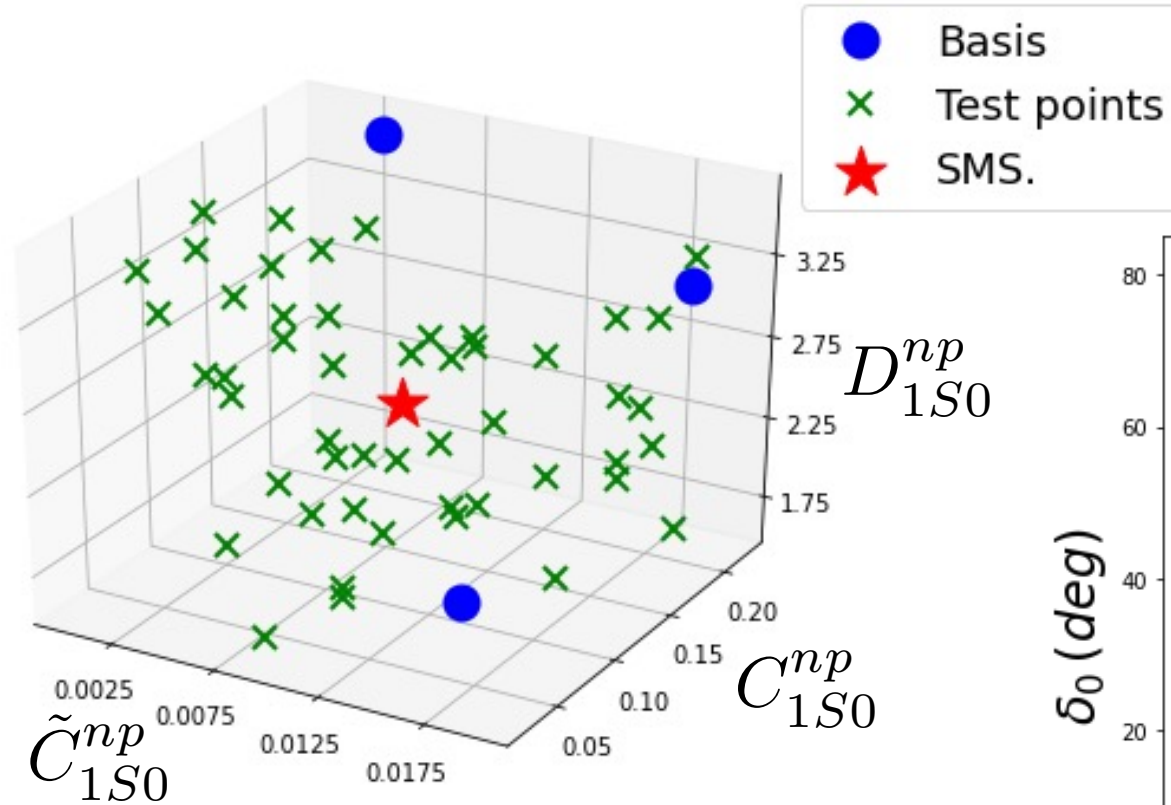
$$\psi_E(k; \theta) = \frac{1}{k^2} \delta(k - k_0) - \frac{2}{\pi} \mathbb{P} \frac{R(k, k_0; \theta)}{k^2 - k_0^2}$$

- Matrix elements are now calculated using

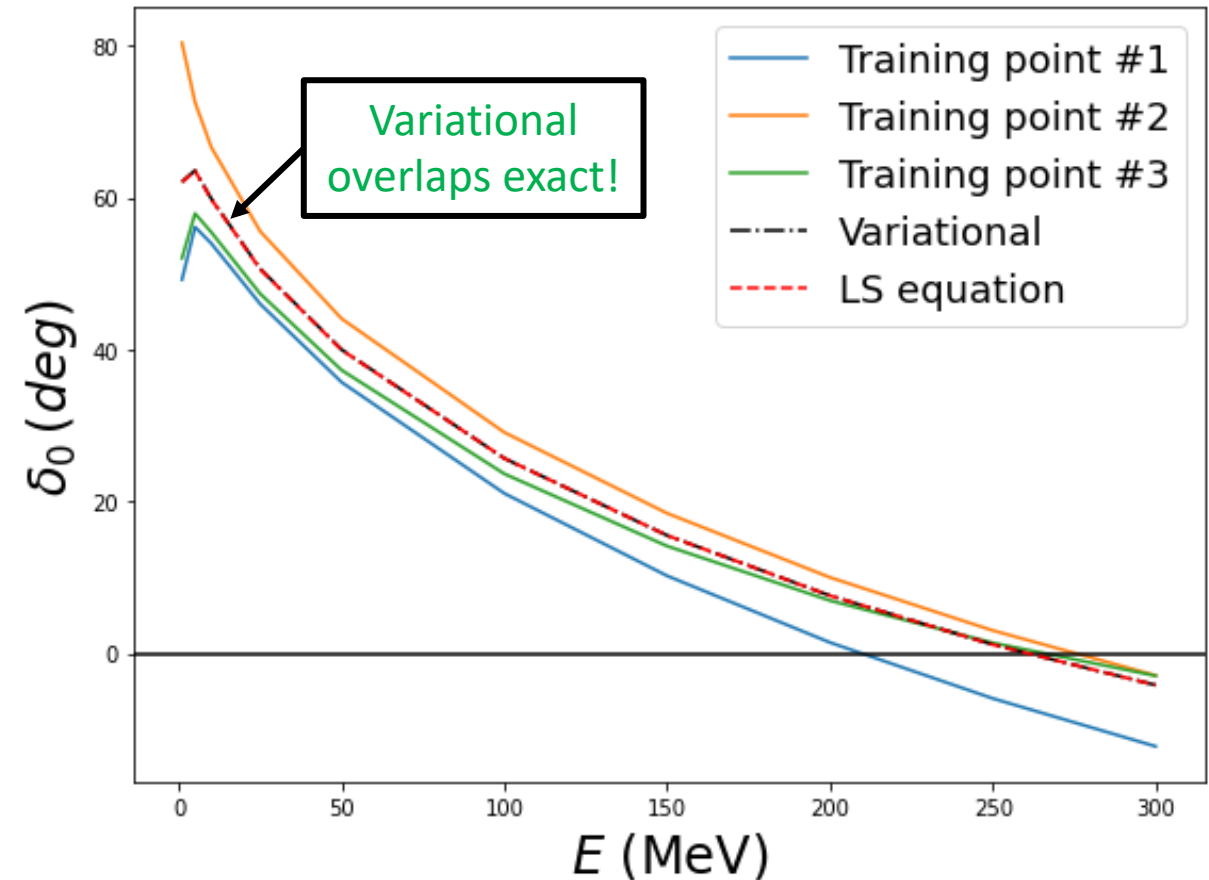
$$\Delta U_{ij} \equiv \frac{2\mu}{\hbar^2} \iint dk dp \psi_E(p; \theta_i) [V(k, p; \theta) - V_j(k, p)] \psi_E(k; \theta_j)$$

- The principal value integrals are evaluated with the same numerical method as conventionally used for the Lippmann-Schwinger equation

# SMS Chiral potential results



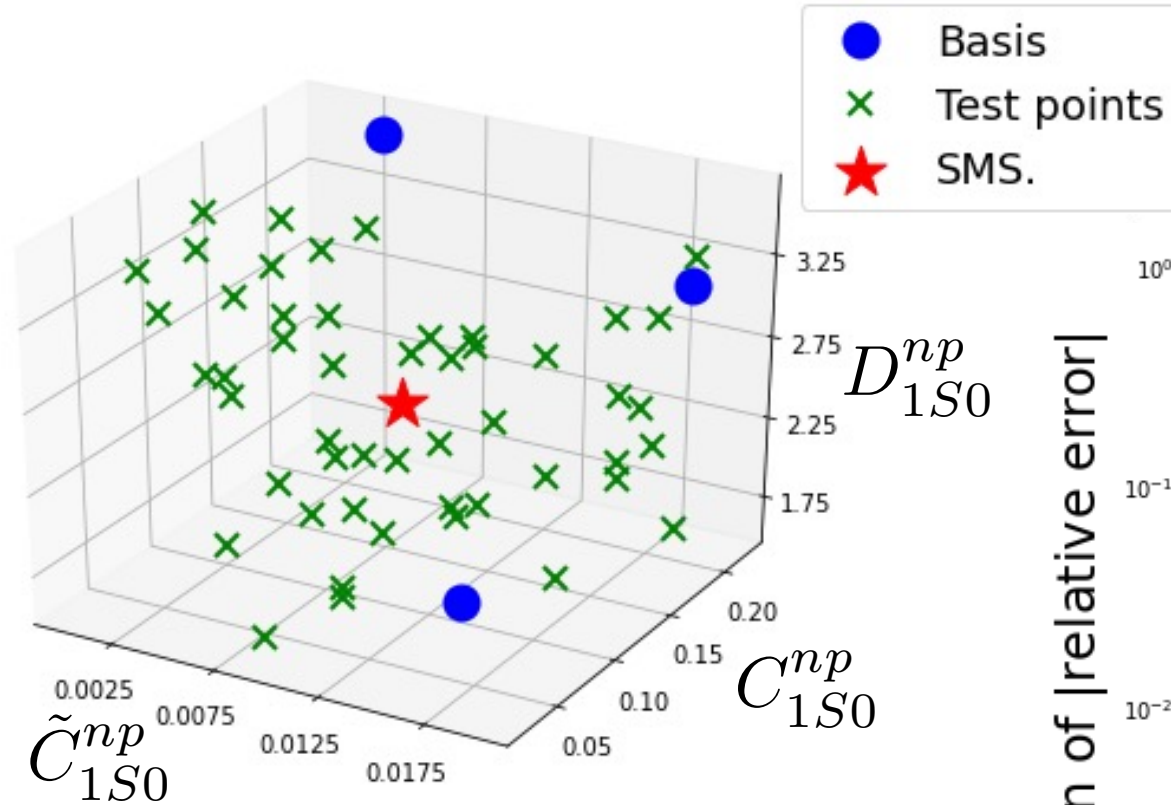
★ = (0.010, 0.121, 2.579)



- Sampled points in coupling constant space for the  $1S_0$  channel
- 50 test points with best parameters set as a star (Reinert, et al.)

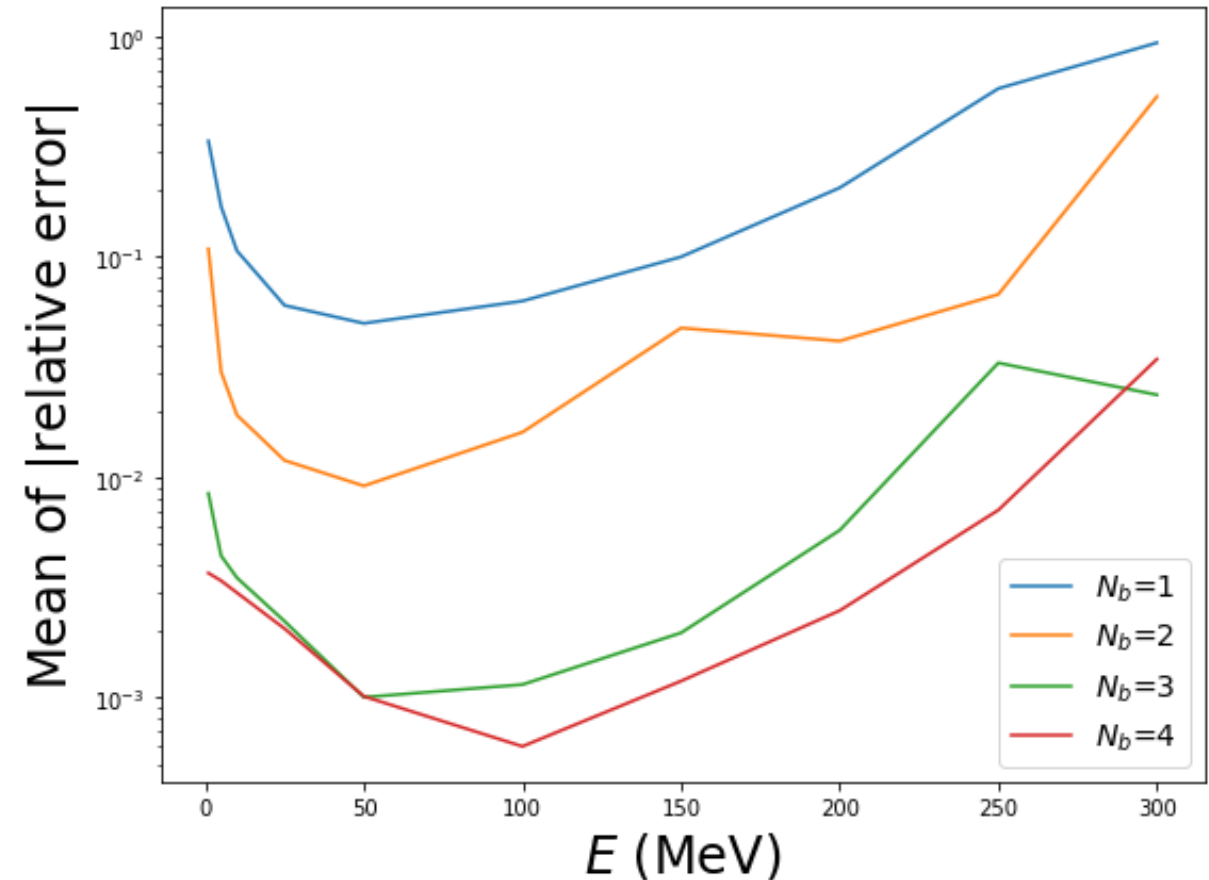


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# Numerics currently being studied

- The basis gets close to being linearly dependent with increased size causing the condition number to grow
- Leads to the matrix becoming **ill-conditioned**
  - Use Moore-Penrose pseudo-inverse with appropriate cutoff
  - Regularize using a “nugget”
- Accuracy currently limited
  - Should get better based on other calculations
- EC calculation for each additional parameter set is **1000x times faster** than calculating directly
  - Need to pre-calculate necessary information needed for EC such as wave functions and potentials



# Summary

- EC works for momentum space applications ([S13.00004: P.J. Millican](#))
- EC emulators can be used for sampling LECs

## Ongoing work

- Understand numerical issues and sensitivity of LECs
- Techniques for a more efficient EC calculation
- Full calculation for all partial waves
- Calculate all observables
- Test EC for different potentials in coordinate and momentum space
- EC applications to three-body scattering ([KP01.00027: Xilin Zhang](#))

Thank you!