

MA 2103 Assignment 02 1

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1. Problem 4 Pg. 632

The temperature Distribution $u(x,t)$ is given by $u(x,t) = X(x)T(t)$ which satisfies the Heat/Diffusion Equation gives as

$$\nabla^2 u(x,t) = \frac{1}{\alpha^2} \frac{\partial u}{\partial t}$$

The possible solutions for the above PDE is

$$u(x,t) = \begin{cases} e^{-k^2 \alpha^2 t} \sin kx, \\ e^{-k^2 \alpha^2 t} \cos kx \end{cases} \quad (1)$$

The given boundary conditions at $t=0$ are given as

$$\begin{cases} 0, & x \in (0,5) \\ 20, & x \in (5,10) \end{cases} \quad (2)$$

It is given that the surfaces are kept at 0. i.e $u(0,t) = u(10,t) = 0$

At $x = 0$, $u = 0$, thus we can discard all the $\cos kx$ solutions. At $x = 10$, $u = 0$, thus $\sin 10k = 0$

$$\sin 10k = 0$$

$$\text{we get, } 10k = n\pi$$

$$k = \frac{n\pi}{10}$$

Now our general solution is

$$u(x,t) = \sum_{n=1}^{\infty} a_n \sin \frac{n\pi x}{10} e^{-[(n\pi\alpha)/10]^2 t} \quad (3)$$

We will consider the initial condition ($t = 0$) to find the Fourier coefficient

$$u(x, t = 0) = u_0$$

$$u_0 = \sum_{n=1}^{\infty} a_n \sin \frac{n\pi x}{10} = \begin{cases} 0, & x \in (0,5) \\ 20, & x \in (5,10) \end{cases} \quad (4)$$

We will use the Fourier trick to find the coefficient

$$\int_0^5 (0) \sin \frac{n\pi x}{10} dx + \int_5^{10} 20 \sin \frac{n\pi x}{10} dx = \int_0^{10} \sum_{n=1}^{\infty} a_n \sin \frac{n\pi x}{10} \sin \frac{m\pi x}{10} dx$$

$$\frac{-200}{n\pi} \cos \frac{n\pi x}{10} \Big|_5^{10} = a_n \frac{10}{2} = 5a_n$$

$$a_n = \frac{40}{n\pi} \left[\cos n\pi - \cos n\frac{\pi}{2} \right]$$

$$a_n = \begin{cases} 0, & n = 4k \\ \frac{40}{n\pi}, & n = 2k + 1 \\ -\frac{80}{n\pi}, & n = 4k + 2 \end{cases}$$

Now our final general solution is given as follows:

$$u(x, t) = \frac{40}{\pi} \sum_{n \text{ is odd}}^{\infty} \frac{1}{n} \sin \frac{n\pi x}{10} e^{-[(n\pi\alpha)/10]^2 t} - \frac{80}{\pi} \sum_{n=2k+2}^{\infty} \frac{1}{n} \sin \frac{n\pi x}{10} e^{-[(n\pi\alpha)/10]^2 t}$$

where, $k \in \mathbb{N}$

2. Problem 6 Pg. 632

The given Boundary conditions is given as

$$u(0, 0) = 20 \qquad u(L, 0) = 150 \qquad (5)$$

$$u(0, t) = 20 \qquad u(L, t) = 50 \qquad (6)$$

The temperature distribution of the bar of length L is given by Eq(3.15) given as follows

$$u(x, t) = \frac{200}{\pi} \sum_{n=1}^{\infty} \frac{(-1)^{n-1}}{n} e^{-[(n\pi\alpha)/L]^2 t} \sin n\pi x \qquad (7)$$

This equation satisfies the Laplace's condition $\nabla u = 0$ Hence we get $u = ax + b$. We will find the coefficients a and b at initial steady state distribution.

Thus at $x = 0$, $u_0 = 20$, we get $b = 20$. Now at $x = L$, $u_0 = 150$, we get $\frac{130}{L}$.

$$u_0 = \frac{130}{L}x + 20 \qquad (8)$$

The final steady state distribution is given by u_f . Thus at $x = 0$, $u_f = 20$, we get $b = 20$. Now at $x = L$, $u_f = 50$, we get $\frac{130}{L}$.

$$u_f = \frac{30}{L}x + 20 \qquad (9)$$

The temperature distribution for both ends kept at 0 is given by

$$u(x, t) = \sum_{n=1}^{\infty} a_n e^{-[(n\pi\alpha)/L]^2 t} \sin n\pi x \qquad (10)$$

Now if we add a linear question to a Fourier series it leaves the Fourier series unchanged. Thus,

$$u(x, t) = u_f + \sum_{n=1}^{\infty} a_n e^{-\left(\frac{n\pi\alpha}{L}\right)^2 t} \sin n\pi x \qquad (11)$$

At $t = 0$,

$$\begin{aligned} u(x, 0) &= u_f + \sum_{n=1}^{\infty} a_n e^{-\left(\frac{n\pi\alpha}{L}\right)^2 t} \sin n\pi x \\ \frac{130}{L}x + 20 &= \frac{30}{L}x + 20 + \sum_{n=1}^{\infty} a_n \sin n\pi x \\ \sum_{n=1}^{\infty} a_n \sin n\pi x &= \frac{100}{L}x \end{aligned}$$

Now we will find the Fourier coefficient a_n .

$$\begin{aligned} a_n &= \frac{2}{L} \int_0^L \frac{100x}{L} \sin \frac{n\pi x}{L} dx \\ a_n &= \frac{200}{L^2} \left[\frac{-Lx}{n\pi} \cos \frac{n\pi x}{L} + \frac{L^2}{n^2\pi^2} \sin \frac{n\pi x}{L} \right] \\ a_n &= \frac{200}{n\pi} (-1)^{n-1} \end{aligned}$$

Thus we get the final solution as follows

$$u(x, t) = u_f + \frac{200}{\pi} \sum_{n=1}^{\infty} (-1)^{n-1} e^{-(\frac{n\pi\alpha}{L})^2 t} \sin n\pi x \quad (12)$$

3. Problem 7 Pg. 633

Both the ends of the bar of length L are insulated thus by Neumann condition.

$$\left. \frac{\partial u}{\partial x} \right|_{x=0} = \left. \frac{\partial u}{\partial x} \right|_{x=L} = 0$$

It is given that Initially ($t = 0$) the temperature distribution is given as by

$$u(x, 0) = x \quad (13)$$

Our required heat equation is given as $u(x, t) = F(x)T(t)$ which satisfies the heat/ diffusion equation. The solutions are given as follows

$$u(x, t) = \begin{cases} e^{-k^2\alpha^2 t} \sin kx, \\ e^{-k^2\alpha^2 t} \cos kx \end{cases} \quad (14)$$

As $\frac{\partial u}{\partial x} = 0$ at $x = 0$. Thus we can discard the $\sin kx$ solution. As $\frac{\partial u}{\partial x} = 0$ at $x = L$ we get $\frac{\partial \cos kx}{\partial x} = \sin(kL) = 0$, thus we get $k = \frac{n\pi}{L}$

The solution satisfies the diffusion equation given below now putting the value $u(x, t) = F(x)T(t)$

$$\begin{aligned} \nabla^2 u(x, t) &= \frac{1}{\alpha^2} \frac{\partial u}{\partial t} \\ \frac{1}{F(x)} \frac{\partial^2 u}{\partial x^2} &= \frac{1}{\alpha^2 T(x)} \frac{\partial u}{\partial t} = 0 \\ \frac{\partial^2 F(x)}{\partial x^2} &= -k^2 & \frac{\partial T(t)}{\partial t} &= 0 \\ \text{For } k &= 0 \\ \frac{\partial^2 F(x)}{\partial x^2} &= 0 & \frac{\partial u}{\partial t} &= 0 \end{aligned}$$

Lets solve the left PDE we get the solution

$F(x) = ax + b$ Now as per Neumann condition

$$\left. \frac{\partial F(x)}{\partial x} \right|_{x=0} = 0 \quad \left. \frac{\partial T(t)}{\partial t} \right|_{x=0} = 0$$

Thus we get $a = 0$ and $T(t) = d = \text{constant}$. Thus we get $u(x, t)|_{k=0} = b * d = \text{constant}$. Thus the value for $k=0$ of $u(x, t)$ is a constant which we will be added later in the heat distribution

For $k > 0$, we got before $k = \frac{n\pi}{L}, n \in \mathbb{N}$. To include the constant for $k = 0$ in the general solution is given by

$$u(x, t) = \sum_{n=0}^{\infty} b_n e^{[(n\pi\alpha)/L]^2 t} \cos \frac{n\pi x}{L} \quad (15)$$

For initial condition $t = 0, u(x, 0) = u_0 = x$

$$u_0 = \sum_{n=0}^{\infty} b_n \cos \frac{n\pi x}{L}$$

$$u_0 = b_0 + \sum_{n=1}^{\infty} a_n \cos \frac{n\pi x}{L}$$

Lets find b_0 and b_n as follows,

$$b_0 = \frac{1}{L} \int_0^L u(x, 0) dx$$

$$b_0 = \frac{1}{L} \int_0^L x dx$$

$$b_0 = \frac{L}{2}$$

$$b_n = \frac{2}{L} \int_0^L x \cos \frac{n\pi x}{L} dx$$

$$b_n = \frac{2}{L} \frac{L}{n\pi} \left[x \sin \frac{n\pi x}{L} \Big|_0^L - \int_0^L \cos \frac{n\pi x}{L} dx \right]$$

$$b_n = \frac{2L}{n^2\pi^2} [\cos n\pi - 1]$$

$$b_n = \frac{2L}{n^2\pi^2} [(-1)^n - 1]$$

If n is odd $b_n = \frac{-4L}{n^2\pi^2}$

If n is even $b_n = 0$

Thus the general solution is

$$u(x, t) = \frac{L}{2} - \frac{4L}{\pi^2} \sum_{\text{odd } n} b_n e^{-[(n\pi\alpha)/L]^2 t} \cos \frac{n\pi x}{L} \quad n \in \mathbb{N}$$

4. Problem 9 Pg. 633

The bar is insulated at $x = 0$, hence we can apply Neumann condition $\frac{\partial u}{\partial x} \Big|_{x=0} = 0$. The heat distribution is given by $u(x, t) = F(x)T(t)$ which satisfies the heat/diffusion equation. The solution of the distribution is given as

$$u(x, t) = \begin{cases} e^{-k^2\alpha^2 t} \sin kx, \\ e^{-k^2\alpha^2 t} \cos kx \end{cases} \quad (16)$$

Now applying the Neumann condition $\frac{\partial u}{\partial x} \Big|_{x=0} = 0$, we can see that the $\sin(kx)$ terms can be eliminated. At $x = 2$, $u = 100$.

$$u = 100 = e^{-k^2 \alpha^2 t} \cos 2k$$

Thus we will have to add 100 to $u(x, t)$.

$$u(x, t) = 100 + e^{-k^2 \alpha^2 t} \cos kx$$

Now, $\cos 2k = 0$, i.e $k = \frac{2n-1}{2} \frac{\pi}{2} = \frac{2n-1}{4} \pi$ The general solution is given as follows.

$$u(x, t) = 100 + \sum_{n=1}^{\infty} b_n e^{-[(2n-1)\pi\alpha/4L]^2 t} \cos \frac{2n-1}{4} \frac{\pi x}{L}$$

Initially ($t=0$) $u(x, 0) = 0$

$$-100 = \sum_{n=1}^{\infty} b_n \cos \frac{2n-1}{4} \frac{\pi x}{L}$$

Lets find the fourier coefficient b_n .

$$\begin{aligned} b_n &= \frac{2}{2}(-100) \int_0^2 \cos \frac{2n-1}{4} \pi x dx \\ b_n &= \frac{-400}{(2n-1)\pi} \sin \frac{2n-1}{4} \pi x \Big|_0^2 \\ b_n &= \frac{-400(-1)^{n+1}}{(2n-1)\pi} \end{aligned}$$

For convenience we will fit our summation from $n = 0$ to $n = \infty$. This converts our $k = \frac{2n-1}{4}$ to $k = \frac{2n+1}{4}$ and similarly $b_n = \frac{-400(-1)^n}{(2n+1)\pi}$. Thus the final solution we obtain is.

$$u(x, t) = 100 - \frac{400}{\pi} \sum_{n=0}^{\infty} \frac{(-1)^n}{2n+1} e^{-[(2n+1)\pi\alpha/4L]^2 t} \cos \frac{2n+1}{4} \frac{\pi x}{L}$$

5. Problem 12 Pg. 633

Given : $\Psi(x, 0) = \sin^2 \pi x$ on the interval $(0, 1)$.

The Schrödinger's equation is given by

$$-\frac{\hbar^2}{2m\Psi} \nabla^2 \Psi + V = i\hbar \frac{1}{T} \frac{\partial T}{\partial t} \quad (17)$$

For a particle in a box problem we take $V(x) = 0$ for $x \in (0, 1)$ Thus we get

$$-\frac{\hbar^2}{2m\Psi} \nabla^2 \Psi = i\hbar \frac{1}{T} \frac{\partial T}{\partial t} \quad (18)$$

By solving the Schrödinger's equation we get the following solution

$$\Psi = \begin{cases} e^{-[iEt]/\hbar} \sin kx \\ e^{-[iEt]/\hbar} \cos kx \end{cases} \quad (19)$$

where $k = \frac{\sqrt{2mE}}{\hbar}$

At $x = 0$, $\Psi = 0$ so we can discard all the $\cos(kx)$ solutions.

At $x = L$, $\Psi = 0$ $\sin(k) = 0$, hence we get $k = n\pi = \frac{\sqrt{2mE}}{\hbar}$, which gives us $E_n = \frac{\hbar^2 n^2 \pi^2}{2m}$. Now our general solution is ,

$$\Psi(x, t) = \sum_{n=1}^{\infty} a_n \sin(n\pi x) e^{-[iE_n t]/\hbar}$$

At $t = 0$,

$$\Psi(x, t = 0) = \sum_{n=1}^{\infty} a_n \sin(n\pi x) = \sin^2(\pi x)$$

We will now find the Fourier coefficients as follows,

$$\begin{aligned} a_n &= 2 \int_0^1 \sin^2(n\pi x) \sin(n\pi x) dx \\ a_n &= 2 \int_0^1 \frac{[1 - \cos(2n\pi x)]}{2} \sin(n\pi x) dx \\ a_n &= \int_0^1 \sin(n\pi x) dx - \int_0^1 \cos(2n\pi x) \sin(n\pi x) dx \\ a_n &= \frac{1}{n\pi} [1 - \cos n\pi] + \frac{1}{2} \left[\frac{\cos(n+2)\pi}{(n+2)\pi} - \frac{1}{n+2} + \frac{\cos(n-2)\pi}{(n-2)\pi} - \frac{1}{n-2} \right] \end{aligned}$$

$$a_n = \begin{cases} 0, & \text{for even } n \\ \frac{8}{n\pi(4-n^2)}, & \text{for odd } n \end{cases} \quad (20)$$

Thus the general solution is

$$u(x, t) = \frac{8}{\pi} \sum_{\text{Odd } n}^{\infty} \frac{\sin(n\pi x)}{n(4-n^2)} e^{-[iE_n t]/\hbar} \quad \text{where } E_n = \frac{n^2 \pi^2 \hbar^2}{2m}$$