

Part IA Mathematics Examples Paper 8 Solutions

(1)

$$1. \quad f(t) = \frac{1}{2}a_0 + \sum_{n=1}^{\infty} a_n \cos nt + \sum_{n=1}^{\infty} b_n \sin nt$$

$$\text{where } a_0 = \frac{1}{\pi} \int_0^{2\pi} f(t) dt = \frac{T}{\pi}$$

$$a_n = \frac{1}{\pi} \int_0^{2\pi} \cos nt dt = \frac{1}{\pi} \left[\frac{\sin nt}{n} \right]_0^{2\pi} = \frac{\sin nT}{n\pi}$$

$$b_n = \frac{1}{\pi} \int_0^{2\pi} \sin nt dt = \frac{1}{\pi} \left[-\frac{\cos nt}{n} \right]_0^{2\pi} = \frac{1 - \cos nT}{n\pi}$$

The simplest approximation for low-pass filtering would be to truncate the series at a certain cut-off frequency. In practice, a more sophisticated approach would be needed: the amplitudes a_n, b_n would be attenuated smoothly as some decreasing function of n . Also, in practice the phase of the harmonics might be affected by transmission. This will result in the a_n and b_n being mixed together: more easily treated mathematically via the complex Fourier series.

2. The function is even: $f(\theta) = f(-\theta)$.

Cosines are all even, sines are all odd, so the Fourier series will contain only cosines.

$$\text{So } f(\theta) = \frac{1}{2}a_0 + \sum_{n=1}^{\infty} a_n \cos n\theta$$

$$\text{where } a_0 = \frac{1}{\pi} \int_0^{2\pi} f(\theta) d\theta = \pi$$

$$a_n = \frac{1}{\pi} \int_0^{2\pi} f(\theta) \cos n\theta d\theta = \underbrace{\frac{2}{\pi} \int_0^{\pi} \theta \cos n\theta d\theta}_{\times 2 \text{ because } -\pi \rightarrow 0 \text{ gives same answer}}$$

2 cont. $\therefore a_n = \frac{2}{\pi} \left\{ \left[\theta \frac{\sin n\theta}{n} \right]_0^\pi - \int_0^\pi \frac{\sin n\theta}{n} d\theta \right\}$ (parts)

$$= \frac{2}{\pi} \left\{ 0 + \left[\frac{\cos n\theta}{n^2} \right]_0^\pi \right\} = \begin{cases} 0, & n \text{ even} \\ -\frac{4}{n^2\pi}, & n \text{ odd} \end{cases}$$

$$\text{So } f(\theta) = \frac{\pi}{2} - \sum_{\substack{n=1 \\ \text{odd}}}^{\infty} \frac{4}{n^2\pi} \cos n\theta$$

$$\text{So } \frac{df}{d\theta} = \sum_{\substack{n=1 \\ \text{odd}}}^{\infty} \frac{4}{n\pi} \sin n\theta, \text{ as obtained for square wave in lectures.}$$

3.

$f(\theta):$
 $\begin{array}{c} | \quad 2\delta(\theta) \quad | \quad 2\delta(\theta-2\pi) \\ \hline | -\pi \quad 0 \quad \pi \quad 2\pi \\ | \quad -2\delta(\theta+\pi) \quad | \quad -2\delta(\theta-\pi) \end{array}$

The function is even, so Fourier series needs only cosines. Convenient to choose a range for integration which doesn't end right on a delta-function — any range of 2π will do, because the function is periodic.

So write $f(\theta) = \frac{1}{2} a_0 + \sum_{n=1}^{\infty} a_n \cos n\theta$

where $a_0 = \frac{1}{\pi} \int_{-\pi/2}^{3\pi/2} f(\theta) d\theta = \frac{1}{\pi} (2 - 2) = 0$

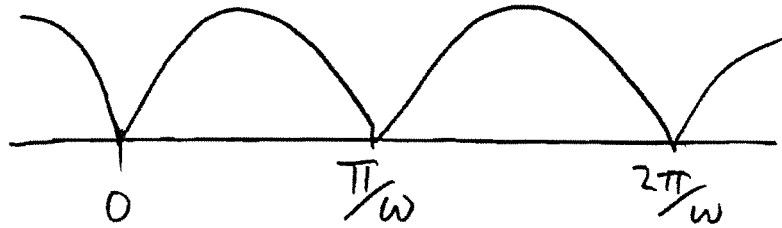
$$a_n = \frac{1}{\pi} \int_{-\pi/2}^{3\pi/2} f(\theta) \cos n\theta d\theta$$

$$= \frac{2}{\pi} \int_{-\pi/2}^{3\pi/2} (\delta(\theta) - \delta(\theta-\pi)) \cos n\theta d\theta$$

$$= \frac{2}{\pi} (1 - \cos n\pi) = \frac{2}{\pi} (1 - (-1)^n) = \begin{cases} 0 & n \text{ even} \\ \frac{4}{\pi} & n \text{ odd} \end{cases}$$

Integrating term by term recovers the square wave $\sum_{n=1, \text{ odd}}^{\infty} \frac{4}{n\pi} \sin n\theta$

4

 $f(t)$ 

$f(t)$ has period $\frac{T}{\omega}$, and is an even function. So it will have a Fourier series

$$f(t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos 2n\omega t$$

Either put $\theta = 2\omega t$, which changes the period to 2π , and use that Maths Data Book formula

or use the Electrical & Information Data Book directly, to deduce

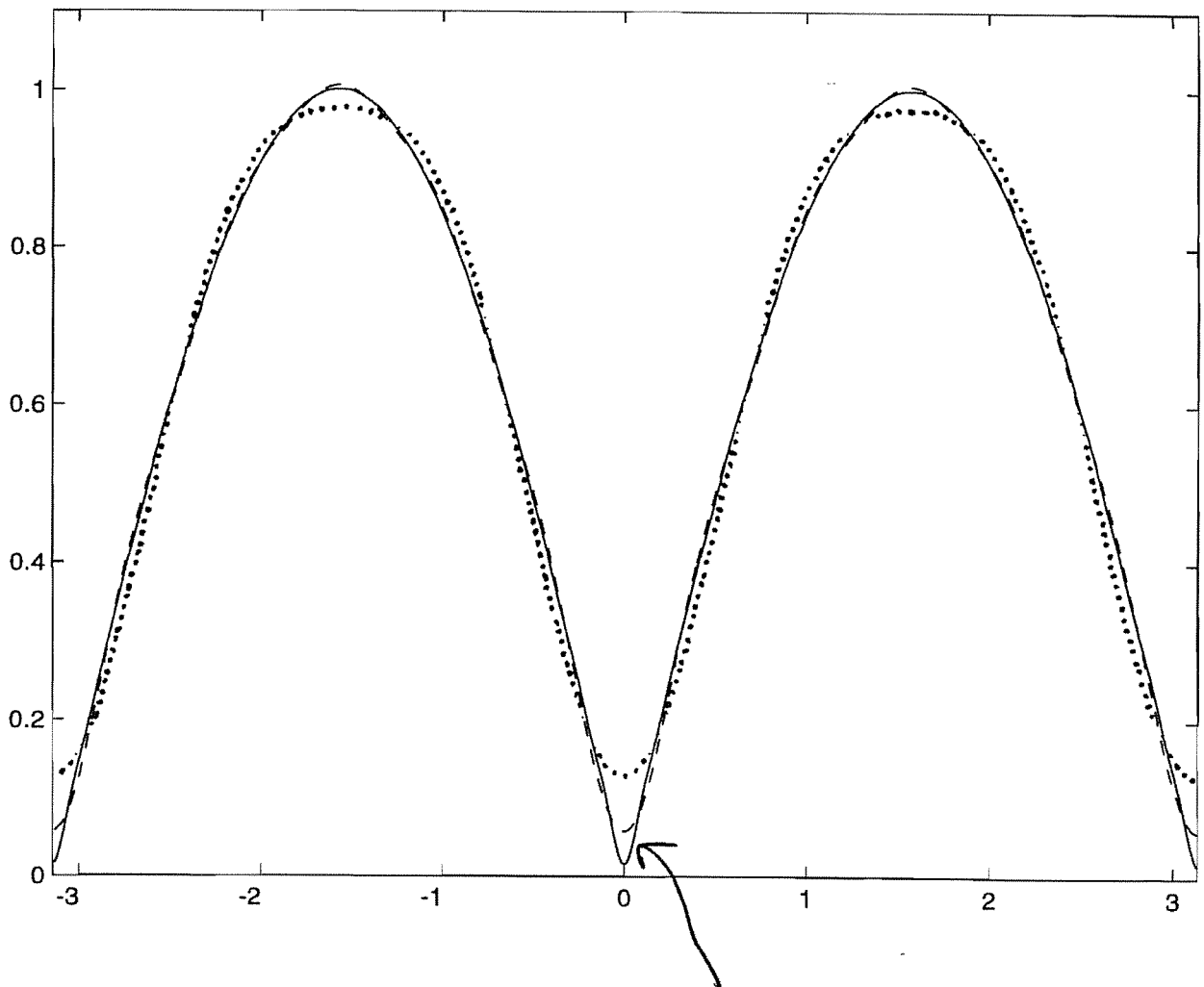
$$a_0 = \frac{2\omega}{\pi} \int_0^{\frac{\pi}{\omega}} \sin \omega t \, dt = \frac{2\pi}{\omega} \left[\frac{-\cos \omega t}{\omega} \right]_0^{\frac{\pi}{\omega}} = \frac{4}{\pi}$$

and

$$\begin{aligned} a_n &= \frac{2\omega}{\pi} \int_0^{\frac{\pi}{\omega}} \sin \omega t \cos 2n\omega t \, dt \\ &= \frac{\omega}{\pi} \int_0^{\frac{\pi}{\omega}} [\sin \omega(2n+1)t - \sin \omega(2n-1)t] \, dt \\ &= \frac{\omega}{\pi} \left[\frac{-\cos \omega(2n+1)t}{\omega(2n+1)} + \frac{\cos \omega(2n-1)t}{\omega(2n-1)} \right]_0^{\frac{\pi}{\omega}} \\ &= \frac{1}{\pi} \left[\frac{1 - \cos(2n+1)\pi}{2n+1} - \frac{1 - \cos(2n-1)\pi}{2n-1} \right] \\ &= \frac{2}{\pi} \left(\frac{1}{2n+1} - \frac{1}{2n-1} \right) = \frac{-4}{\pi(4n^2-1)} \end{aligned}$$

4 cont.

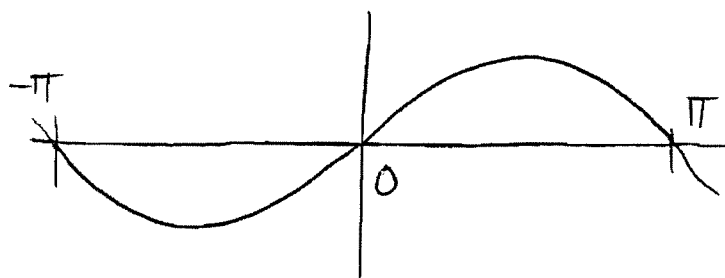
Computed waveforms:
... 2 terms (ie $n = 1, 2$, plus a_0 term)
-- 5 terms
— 20 terms



Notice that we need high terms to get this sharp corner.

(5)

$$5. \quad y = \begin{cases} x(\pi + x) & -\pi \leq x \leq 0 \\ x(\pi - x) & 0 \leq x \leq \pi \end{cases}$$



Only places where discontinuities could occur are $x = -\pi, 0, \pi$.

It is clear that y is continuous at these points. What about $\frac{dy}{dx}$? $\frac{dy}{dx} \Big|_{-\pi+} = -\pi = \frac{dy}{dx} \Big|_{\pi-}$

$$\frac{d^2y}{dx^2} ? \quad \frac{d^2y}{dx^2} = \begin{cases} 2 & -\pi \leq x \leq 0 \\ -2 & 0 \leq x \leq \pi \end{cases} \quad \frac{dy}{dx} \Big|_{0-} = \pi = \frac{dy}{dx} \Big|_{0+} \text{ so continuous}$$

so has discontinuities at $0, \pi$. So y, y' are continuous but y'' is not, so expect Fourier coefficients of the order of $1/n^3$ as $n \rightarrow \infty$.

y is an odd function, so need only sines:

$$\begin{aligned} y(x) &= \sum_{n=1}^{\infty} b_n \sin nx, \quad \pi b_n = \int_{-\pi}^{\pi} y(x) \sin nx \, dx \\ &= 2 \int_0^{\pi} y(x) \sin nx \, dx = 2 \int_0^{\pi} x(\pi - x) \sin nx \, dx \\ &= \left[-2x(\pi - x) \frac{\cos nx}{n} \right]_0^{\pi} + 2 \int_0^{\pi} (\pi - 2x) \frac{\cos nx}{n} \, dx \\ &= \frac{2}{n} \left[(\pi - 2x) \frac{\sin nx}{n^2} \right]_0^{\pi} - \frac{2}{n} \int_0^{\pi} (-2) \frac{\sin nx}{n} \, dx \end{aligned}$$

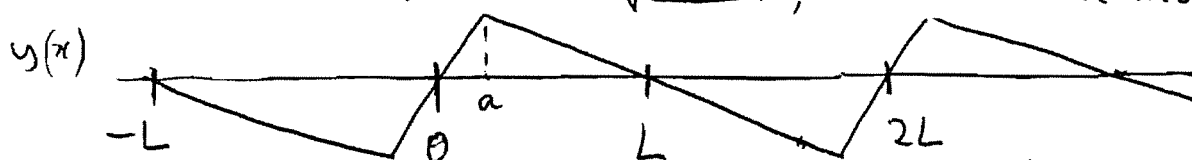
(6)

5 cont.

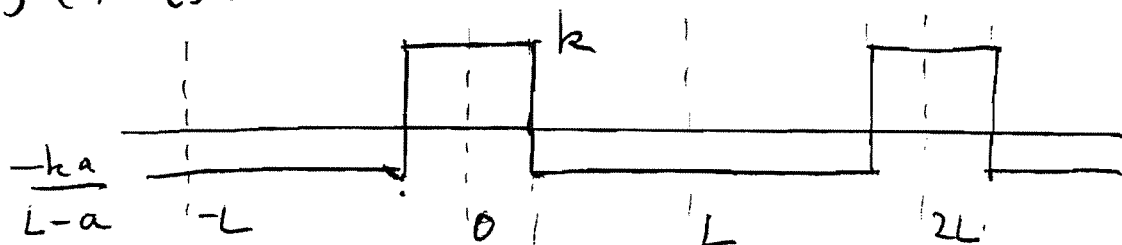
$$\begin{aligned}
 \therefore \pi b_n &= \frac{4}{n^2} \left[-\frac{\cos nx}{n} \right]_0^\pi \\
 &= \frac{4}{n^3} (1 - \cos n\pi) = \frac{4}{n^3} (1 - (-1)^n) \\
 &= \begin{cases} 8/n^3, & n \text{ odd} \\ 0, & n \text{ even} \end{cases}
 \end{aligned}$$

So the coefficients do indeed die away like n^{-3} , as expected.

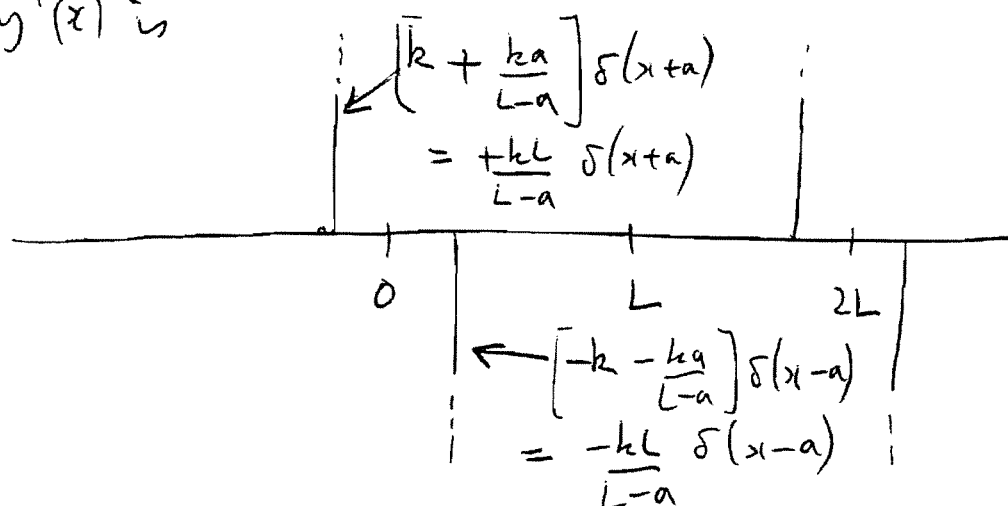
6. We have no choice how to extend the function here: the modes of the string are $\sin \frac{n\pi x}{L}$, so we use a Fourier sine series, so that the extended function is odd. Other extensions could represent the function, but wouldn't be modes.



To evaluate, we can either do it directly or use the trick from Q3:
 $y'(x)$ is:



So $y''(x)$ is



6 cont. $y''(x)$ is an odd function. Let its Fourier series be $\sum_{n=1}^{\infty} b'_n \sin \frac{n\pi x}{L}$ [NB. period is $2L$]

$$\text{Then } b'_n = \frac{2}{2L} \int_{-L}^L y''(x) \sin \frac{n\pi x}{L} dx = \frac{2}{L} \int_0^L y''(x) \sin \frac{n\pi x}{L} dx$$

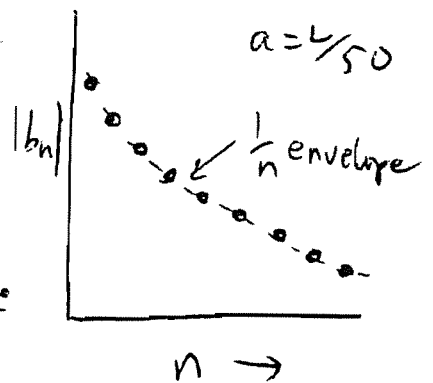
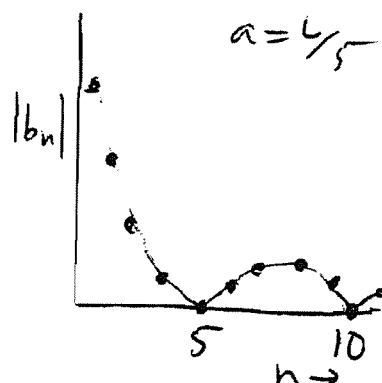
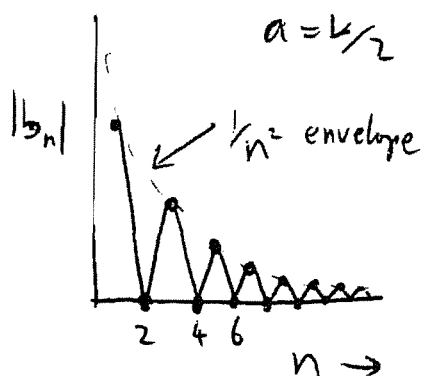
$$= \frac{2}{L} \int \left[\frac{-kL}{L-a} \delta(x-a) \right] \sin \frac{n\pi x}{L} dx$$

$$= \frac{-2k}{L-a} \sin \frac{n\pi a}{L}$$

Now integrate term by term to get the Fourier series for $y(x) = \sum_{n=1}^{\infty} b_n \sin \frac{n\pi x}{L}$

$$\text{where } b_n = -\frac{b'_n}{(n\pi/L)^2} = \frac{L^2}{n^2\pi^2} \cdot \frac{2k}{L-a} \sin \frac{n\pi a}{L}$$

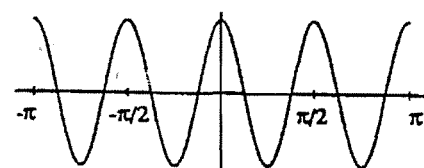
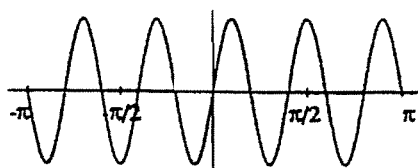
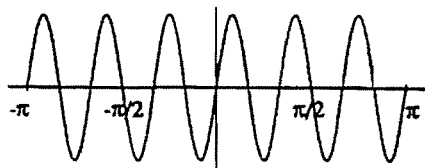
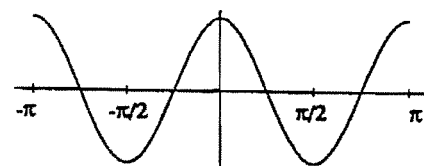
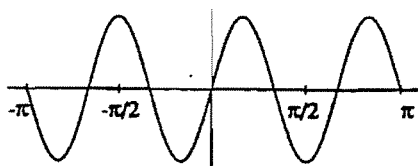
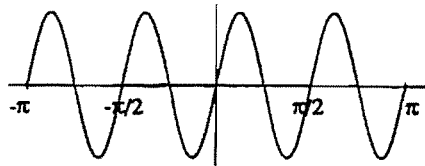
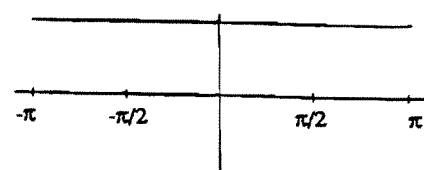
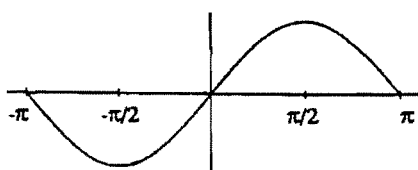
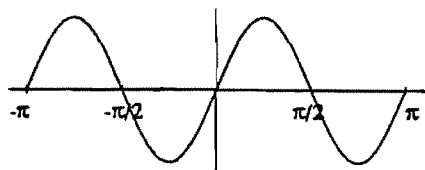
$\approx \frac{aL}{n\pi} \cdot \frac{2k}{L-a} \text{ when } \frac{n\pi a}{L} \ll 1$



$y(x)$ is continuous, but y' is not, so the Fourier coefficients die away like n^{-2} at large n . If the "corner" at the plucking point were rounded, $y'(x)$ would then be continuous, so $|b_n|$ would decay at least as fast as n^{-3} . So a rounded corner gives a less "bright" sound, because it has less high-frequency content. Contrast a guitar string plucked with the flesh of a finger, and one plucked with a narrow, hard plectrum (like a harpsichord string).

7.

The key to answering this question is to consider the symmetry of the terms in the three series.



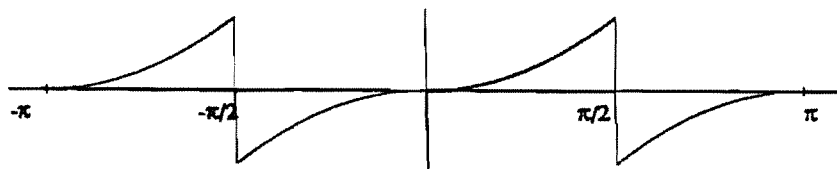
(i) Terms antisymmetric about $x=0$ & antisymmetric about $x=\pm\pi/2$

(ii) Terms antisymmetric about $x=0$, symmetric about $x=\pm\pi/2$

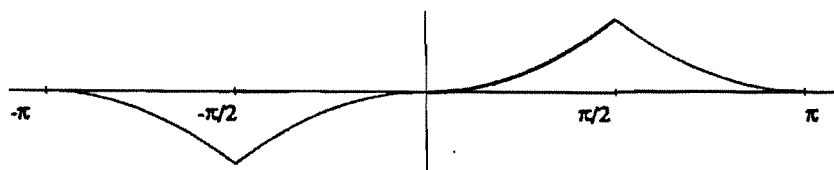
(iii) Terms symmetric about $x=0$, symmetric about $x=\pm\pi/2$.

Functions represented by these series must have the same symmetries.

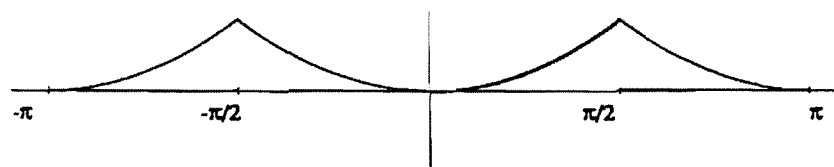
Case (i)



Case (ii)



Case (iii)



Cases (ii) & (iii) have discontinuities of slope, so $b_m, c_m = O(\frac{1}{m^2})$

Case (i) has discontinuities of value, so $a_m = O(\frac{1}{m})$ and much slower convergence.

8

$$y(t) = \begin{cases} e^{-\alpha t} & 0 < t < T/2 \\ 0 & T/2 < t < T \end{cases} = \sum_{n=-\infty}^{\infty} C_n e^{2\pi i n t / T}$$

(9)

$$\Rightarrow C_n = \frac{1}{T} \int_0^T y(t) e^{-2\pi i n t / T} dt = \frac{1}{T} \int_0^{T/2} e^{-(\alpha + 2\pi i n / T)t} dt$$

$$= \frac{1}{T} \left[\frac{e^{-(\alpha + 2\pi i n / T)t}}{-(\alpha + 2\pi i n / T)} \right]_0^{T/2}$$

$$= \frac{1 - e^{-\alpha T/2 - i n \pi}}{\alpha T + 2\pi i n}$$

If $\frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos \frac{2n\pi t}{T} + \sum_{n=1}^{\infty} b_n \sin \frac{2n\pi t}{T} = y(t) = \sum_{n=-\infty}^{\infty} C_n e^{2\pi i n t / T}$

then $\sum_{n=-\infty}^{\infty} C_n e^{2\pi i n t / T} = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \frac{e^{2\pi i n t / T} + e^{-2\pi i n t / T}}{2} + \sum_{n=1}^{\infty} b_n \frac{e^{2\pi i n t / T} - e^{-2\pi i n t / T}}{2i}$

$$\Rightarrow a_0 = 2C_0$$

and for $n \geq 1$ $\frac{a_n - ib_n}{2} = C_n \Rightarrow a_n = 2\operatorname{Re}(C_n); b_n = -2\operatorname{Im}(C_n)$

$$\therefore a_0 = \frac{2(1 - e^{-\alpha T/2})}{\alpha T}; a_n = 2\operatorname{Re} \frac{\alpha T - 2\pi i n}{\alpha^2 T^2 + 4n^2 \pi^2} (1 - \cos n\pi e^{-\alpha T/2})$$

ie. $a_n = \frac{2\alpha T (1 - e^{-\alpha T/2} \cos n\pi)}{\alpha^2 T^2 + 4n^2 \pi^2}$

$$\left[e^{-i n \pi} = \cos n\pi - i \sin n\pi \right]$$

and $b_n = \frac{+4n\pi (1 - e^{-\alpha T/2} \cos n\pi)}{\alpha^2 T^2 + 4n^2 \pi^2}$

JW/TPH