

# Energy Distribution

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## 1 Configurations and Weights

For molecules in a system there is a probability that it will exist in a given energy state. In a system any individual molecule can exist in states with energies from  $\epsilon_0$  to  $\epsilon_1$  where  $\epsilon_0 \equiv 0$ . That lowest energy state the **zero-point energy** is the baseline for measuring other energies in the system, and must be accounted for to obtain actual energies of the system.

For a system of  $N$  molecule there will be  $N_0$  in the  $\epsilon_0$  state,  $N_1$  in the  $\epsilon_1$  state and so on where  $\sum N_n = N$ . The set of populations  $N_0, N_1 \dots$  in the form  $\{N_0, N_1 \dots\}$  is an **Instantaneous Configuration**. For a system of  $N$  particles there are  $N(N-1)$  configuration total and  $1/2N(N-1)$  distinguishable configurations. Systems display the behavior of the state they are most likely to exist in. The number of ways a general configuration can be achieved the **weight**  $\mathscr{W}$  is based on number of ways each particle entered its state. For example there are  $N_0!$  ways for  $N_0$  molecules to be selected. Over the entire configuration that the weight can be given by the following.

$$\mathscr{W} = \frac{N!}{N_0!N_1!\dots}$$

This is because there are  $N!$  ways to select  $N$  particles, and for each state with  $N_i$  particles there are  $N_i!$  ways for those particles to be selected. When summed over the entire system that gives the number of ways to end up with a given configuration.

$\mathscr{W}$  can be approximated with the natural log and using the observation that  $\ln x! \approx x \ln x - x$  giving the following.

$$\ln \mathscr{W} = N \ln N - \sum_i N_i \ln N_i$$

### 1.1 Most Probable Distribution

The system will most likely exist in the configuration with the largest weight resulting in its properties matching that of the system. Since weight is a function of  $N_i$  that weight can be found by optimizing  $\mathscr{W}(N_i)$ . This gives the following derivative

$$d \ln \mathscr{W} = \sum_i \left[ \frac{\partial \mathscr{W}}{\partial N_i} \right] dN_i = 0$$

Given the reality that states will not all share the same energy the configuration with the greatest weight must also satisfy the condition

$$\sum_i N_i \epsilon_i = E$$

meaning the total energy must remain constant as  $N_i$  changes. The number of molecules is also fixed meaning that adding a molecule to one state necessitates removing one from another.

Adding the constraints of constant total energy and number of particles to the weight differential

$$\sum_i \left[ \frac{\partial \mathscr{W}}{\partial N_i} + \alpha - \beta \epsilon_i \right] dN_i = 0$$

where  $\alpha$  and  $\beta$  are constants. By inserting the equation for  $\ln \mathscr{W}$  into this equation the result is

$$\frac{N_i}{N} = e^{\alpha - \beta \epsilon_i}$$

which is close to the Boltzmann Distribution. By then canceling out  $\alpha$  the Boltzmann Distribution is arrived at.

$$\frac{N_i}{N} = \frac{e^{-\beta \epsilon_i}}{\sum_i e^{-\beta \epsilon_i}}$$

Here  $\beta = 1/kT$  where  $k$  is Boltzmann's constant. The denominator of the equation is called a **partition coefficient**

## 1.2 Relative Population of States

Since the partition coefficient's cancel when the ratio is taken the relative population is given by

$$\frac{N_i}{N_j} = e^{-\beta(\epsilon_i - \epsilon_j)}$$

It explains that at absolute zero the entire system is in the ground state, and the exponential decrease in populations with the increasing energy gap of states.

## 2 Molecular Energies

A molecule's mean energy is relative to its ground state energy  $E$  divided by the total number of molecules  $N$ . The most probable state at a temperature  $T$  is given by the Boltzmann distribution. This can be simplified to give the mean energy in terms of a partition function.

$$\langle \epsilon \rangle = \frac{\epsilon}{e^{\beta \epsilon} + 1}$$

Molecules have three modes of motion contributing to their energy states: translational, vibrational, and rotational.

The translational component of energy is given by  $\langle \epsilon \rangle = 3/2kT$  where the 3 is from the number of dimension, and the  $1/2kT$  is from the partition function.

The rotational component is dependent of layout of the molecule.