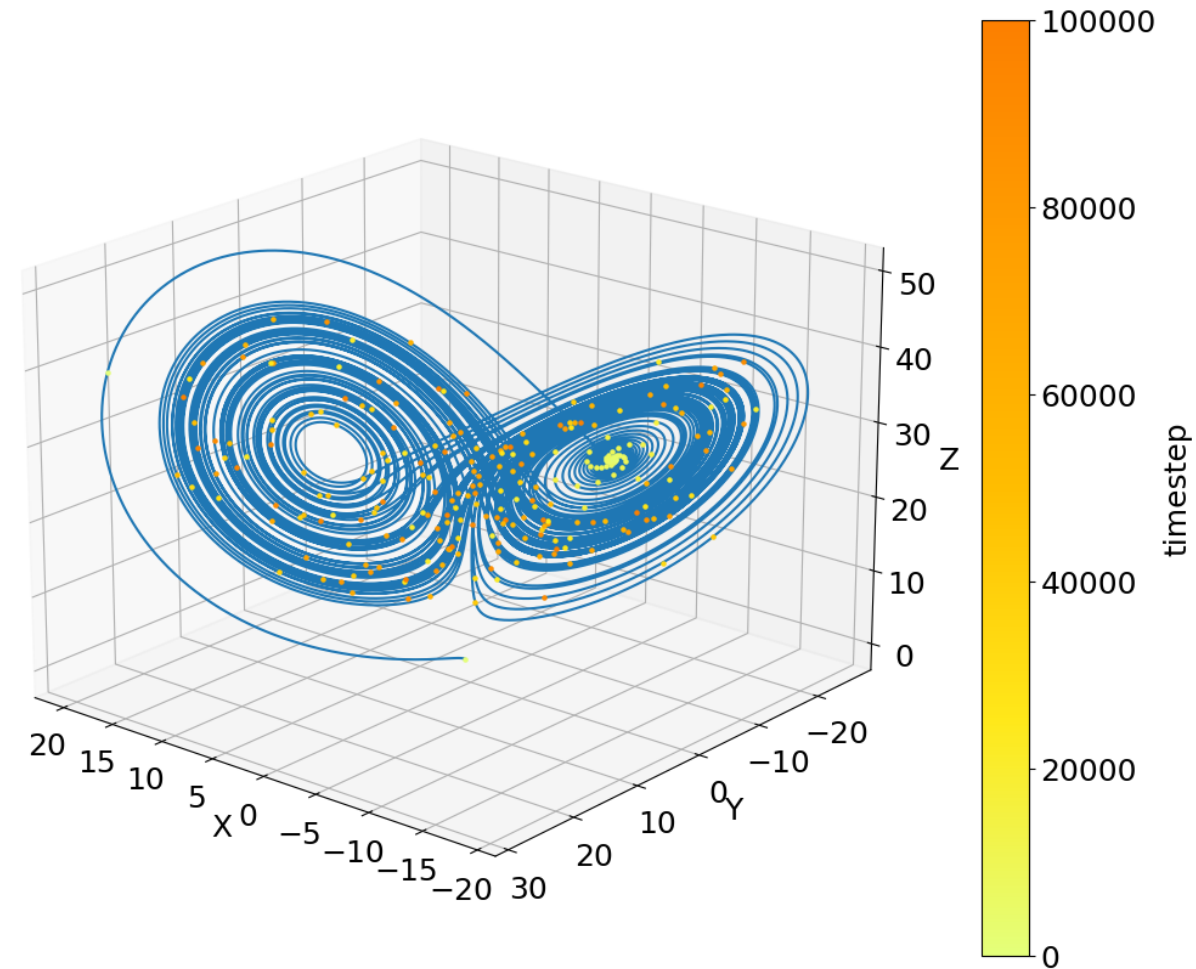


# Final Project

## The Lorenz Model and Numerical Methods

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# 1. Introduction

- Chaos is one of the important features of nonlinear systems
- Definition of chaos:  
*Chaos is aperiodic long-term behavior in a deterministic system that exhibits sensitive dependence on initial conditions* (Strogatz, 2015)
- Features of chaos: **sensitive to initial condition** / **strange attractor**
- **Sensitive to initial condition**
  - ✓ In a chaotic system, small differences in initial conditions would be amplified as the system develops
  - ✓ Whatever small the difference is, it would be magnified, and they have different states after sufficient time
- **Strange attractor**
  - ✓ The trajectory of a chaotic system has a strange attractor, different from the point attractor/limit cycle (see appendix slide)

# 1. Introduction

- Lorenz model, proposed by Lorenz (1963), is one of the most famous chaotic systems
- It is a simplified model of a convective atmosphere
- Use the Lorenz model as an example of a chaotic system
- Goals of this project/questions to be answered
  - Investigate the behavior of the Lorenz system and observe features of the chaotic system and its trajectory
  - Investigate the time development of the Lorenz system with different numerical schemes. Do they have completely different results or similar to each other?
  - How the system is sensitive to initial conditions and parameters?

## 2. Lorenz model

- The Lorenz model, proposed by Lorenz (1963)
- $x, y, z$ : 3-dimensional variables
- $\sigma, r, b$ : parameters: all positive
  - $\sigma$ : Prandtl number, kinematic viscosity/thermal conductivity
  - $r$ : Ratio of Rayleigh number/Reynolds number
- First, consider  $25 < r$  which is critical value
- Consider initial condition  $(x, y, z) = (0, 1, 0)$

$$\begin{cases} \frac{dx}{dt} = \sigma(y - x) \\ \frac{dy}{dt} = rx - y - xz \\ \frac{dz}{dt} = xy - bz \end{cases}$$

## 2. Lorenz model: features

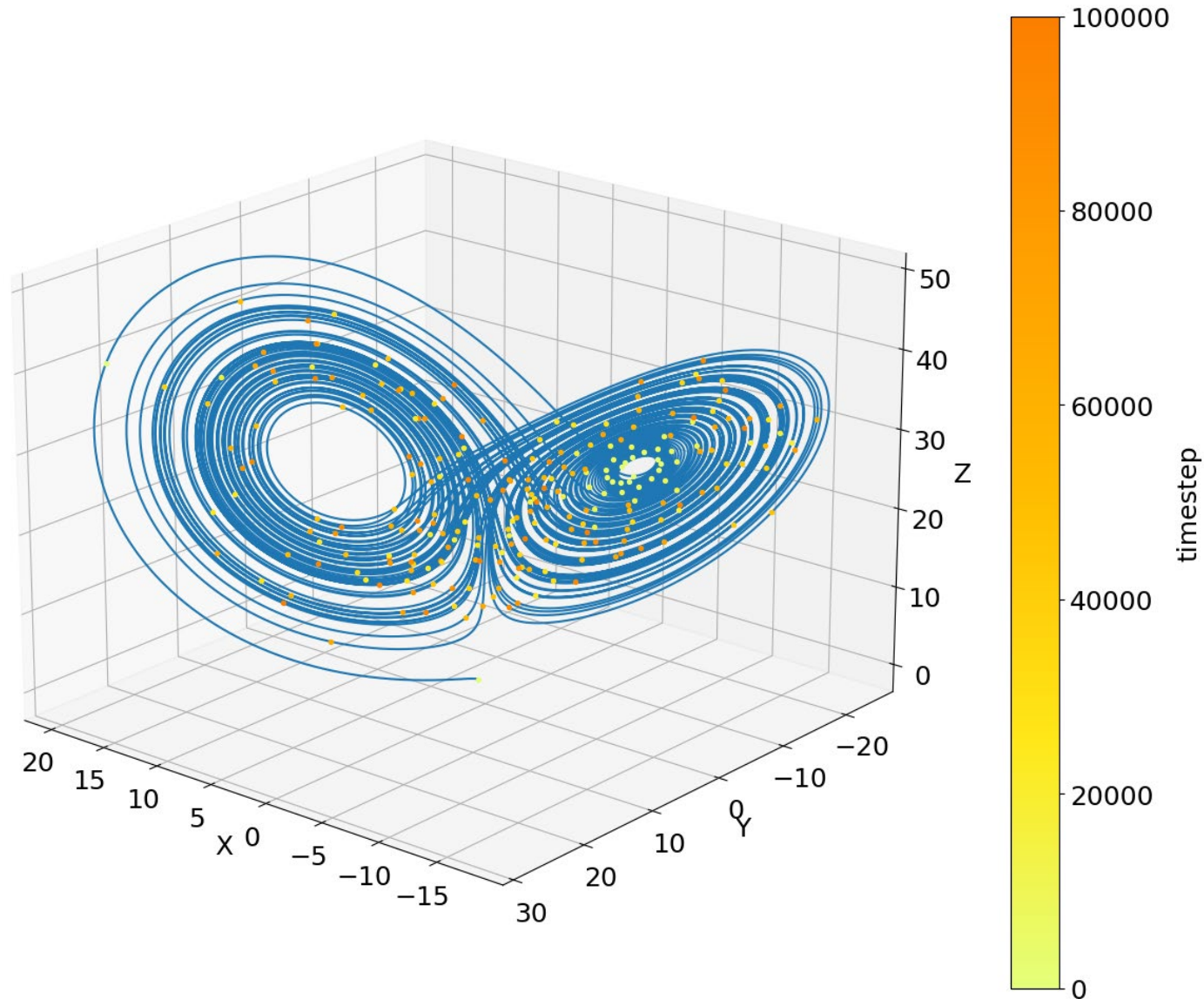
- Lorenz system is **dissipative**
  - Phase space volume decreases exponentially
  - Volume in phase space contract to 2D plane or 1D line or point
- Fixed points of the Lorenz system are surrounded by **unstable limit circles**
  - Trajectories fly away from them
  - Trajectories do not cross each other (never repeat the previous status)
  - Also, trajectories always stay in a finite region
- The above conditions look contradict each other:
  - Trajectories fly away from the limit circle but still stay in a finite region
  - Phase space volume decreases but infinitely transits to new status
- **Strange attractor satisfies all these conditions** (and that's why it's interesting!)

# 3. Numerical Methods

- **Three** methods were used
  - Predictor-Corrector
  - Runge-Kutta (RK4)
  - Backward Differentiation Formulas (BDF)
- For BDF, the Jacobi matrix with the initial condition was **weak stiff**
  - The ratio of eigenvalues
$$\frac{|\lambda_{max}|}{|\lambda_{min}|} \sim 8.58$$
- $\Delta t = 0.001$  and total integration time=100000 was use (See Appendix B)
- Since the Lorenz system is **dissipative**, we **cannot use Symplectic Integrators**

# 4. Results: Overview

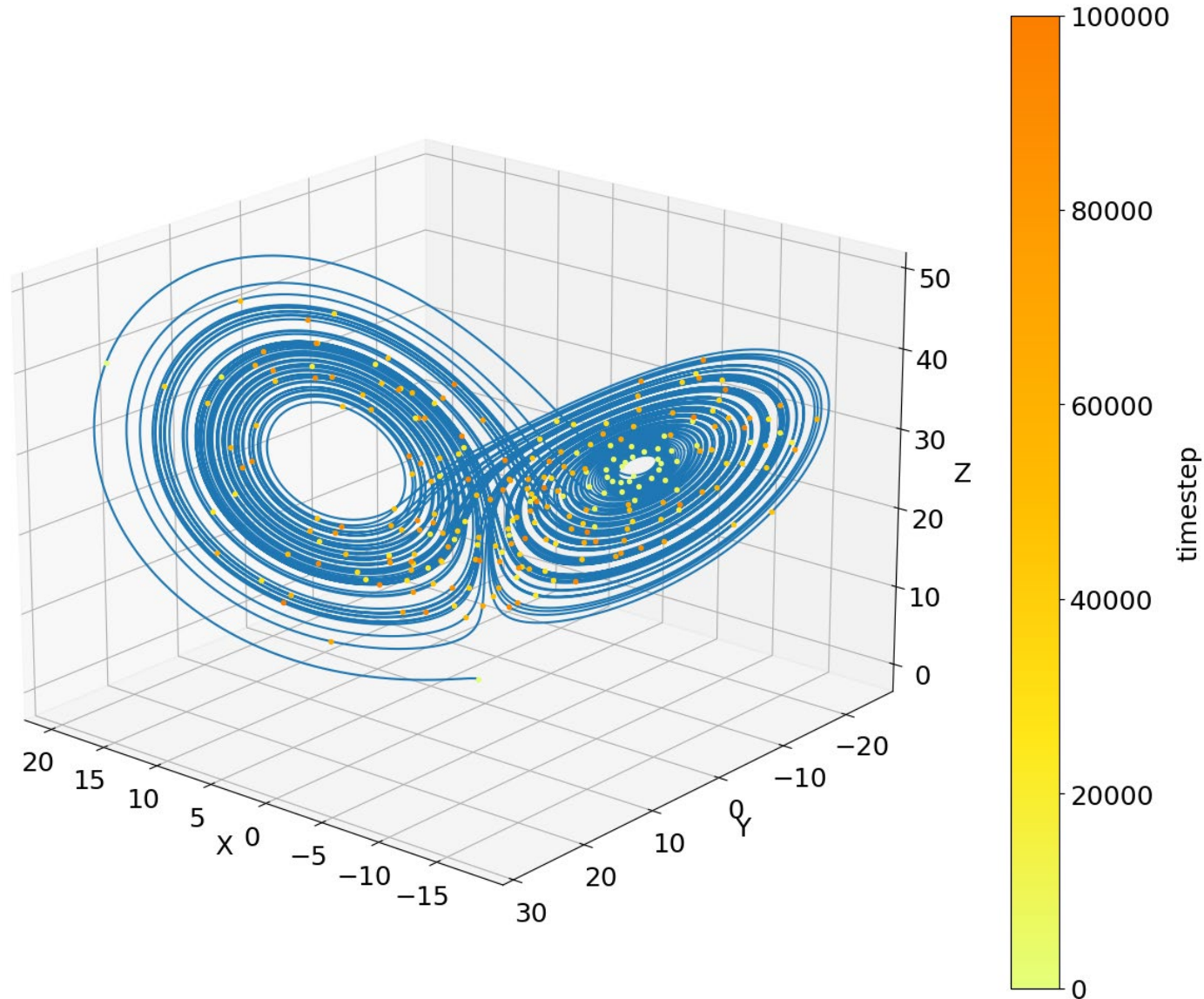
- Result from RK4
- $\Delta t=0.001$
- 2 flat disks exist in XYZ space
- At first, round in a small region  
(see the right **disk**)
- **Chaotic** behavior starts from  
around **16000 steps**
- Then, jump to another disk
- Trajectory contentiously jumps  
between the two disks after  
sufficient time





# 4. Results: Overview

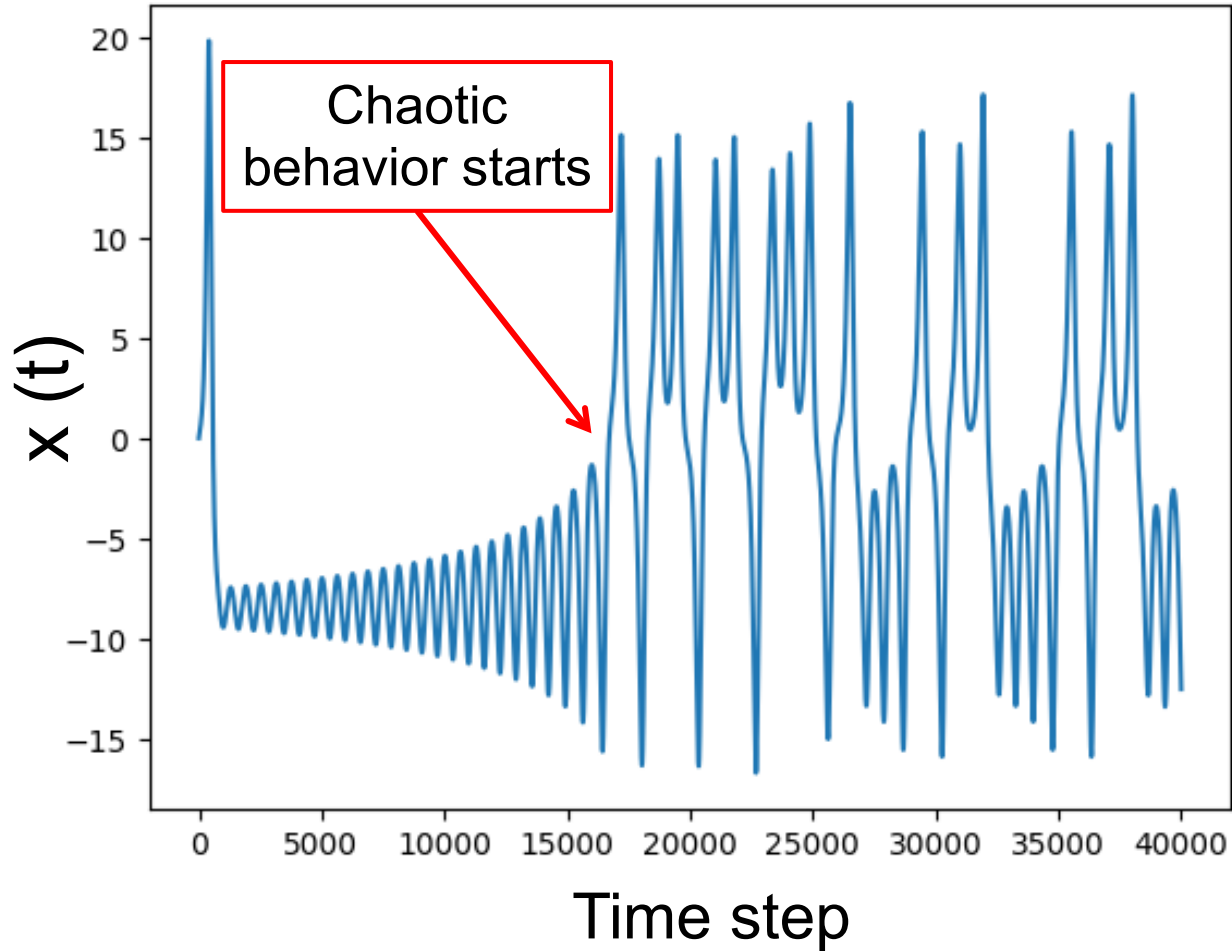
- This attractor satisfies the restriction we saw
  - Volume construction: trajectory is folded into 2-dimensional disks
  - Unstable limit circle: there are limit circles at the center of the disk, but they are unstable (apart from them)
  - Stay in finite volume: the starting point is the edge of the disk, and all trajectories stay inside of it



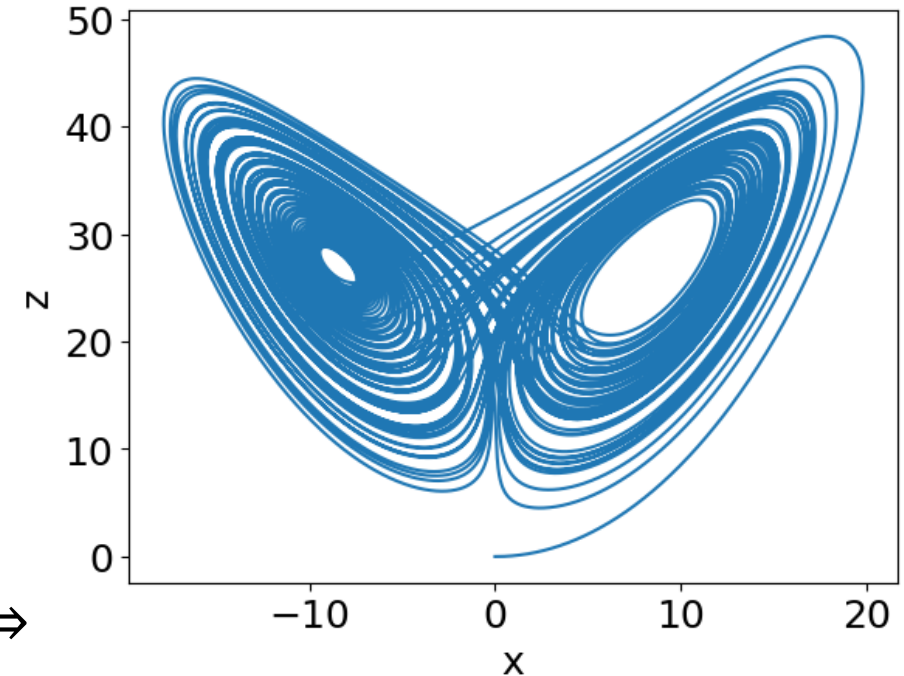
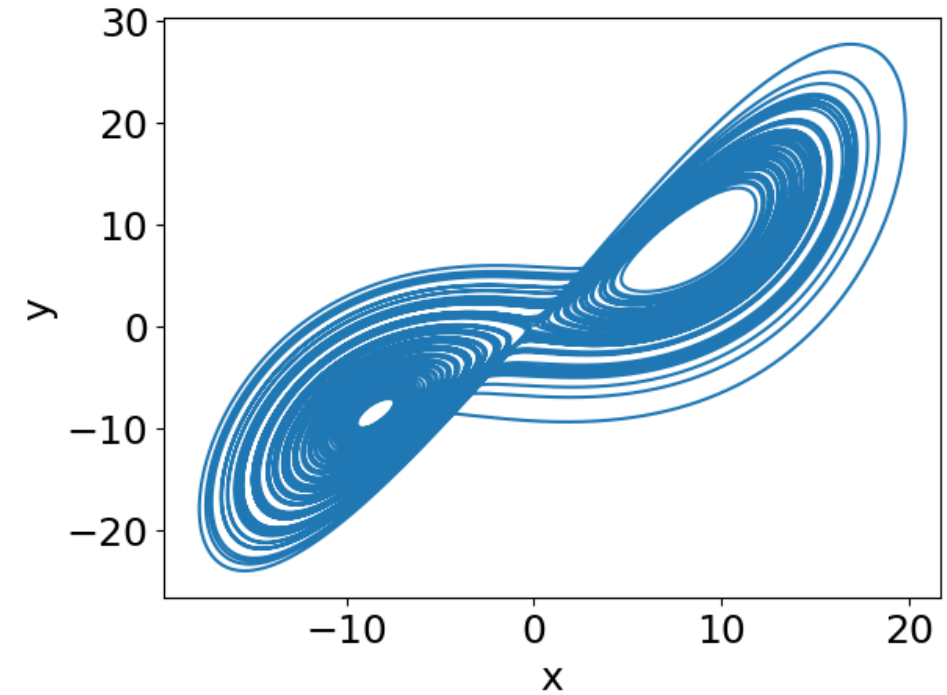


# 4. Results: Overview

Time development of strange attractor,  $x(t)$

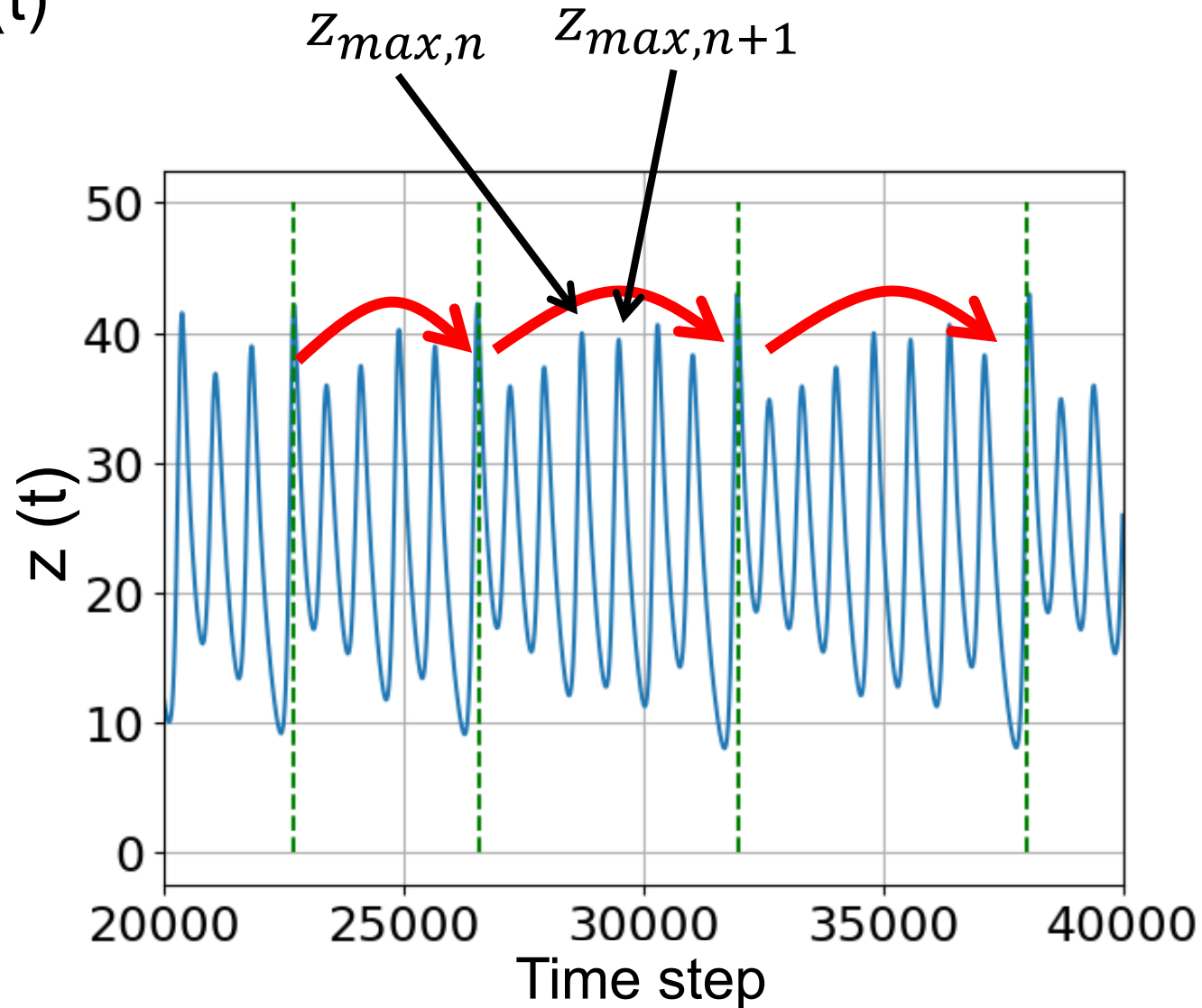


Projection of strange attractor to x-y, x-z plane  $\Rightarrow$



# 4. Results: Lorenz mapping

- There is a pattern on peaks of  $z(t)$
- Next to the largest peak is relatively small  
(see the red arrows)
  - Then increase
  - Decrease
  - Next largest peak
- Plot  $z_{max,n}$  versus  $z_{max,n+1}$ 
  - $z_{max,n}$ :  $n$ -the peak of  $z(t)$
  - $z_{max,n+1}$ : the next peak of  $z_{max,n}$
  - This plot is called **Lorenz mapping**

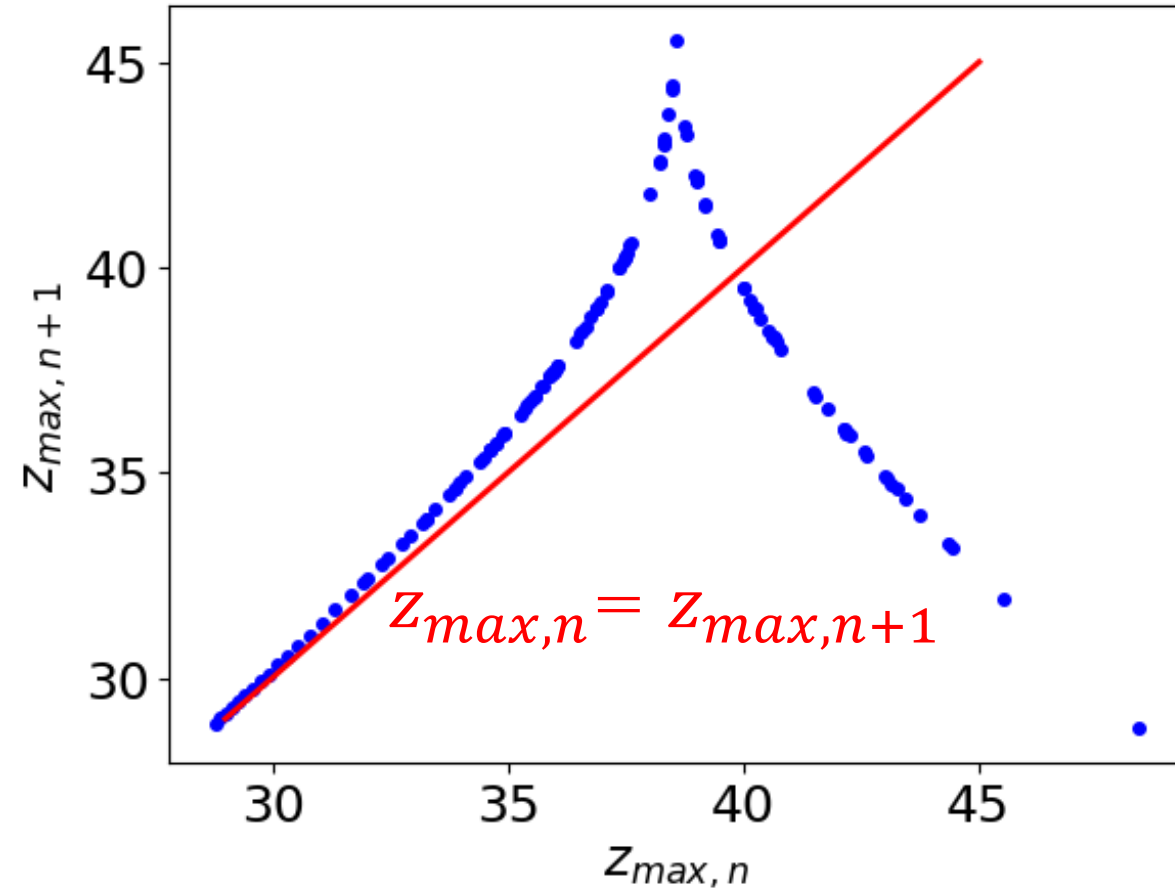


# 4. Results: Lorenz mapping

- Lorenz mapping:
  - Points fall into a line
  - Indicates there is a coherent structure
- At all points, the slope is always steeper than 1, i.e.

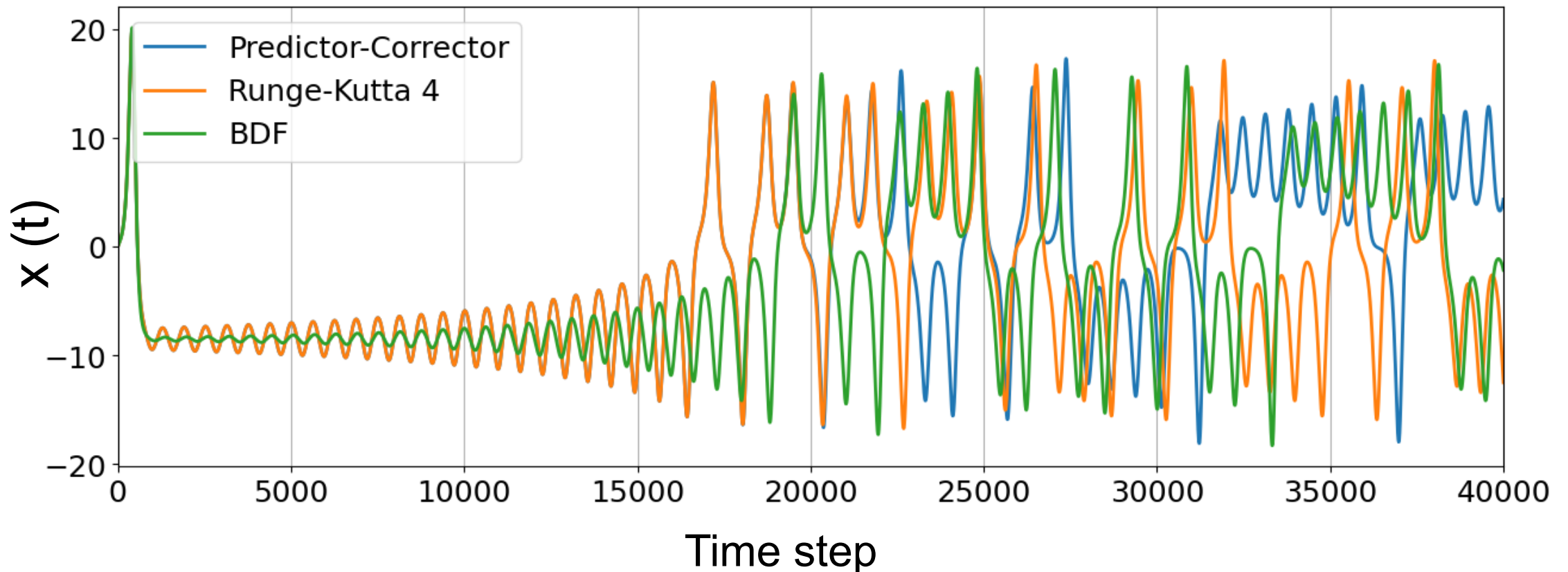
$$|f'(z_{max,n})| > 1$$

- This is used to prove that the fixed point of the system **unstable**
  - Small perturbation  $\eta$ , near fixed points, is amplitude
  - $\eta_{n+1} = |f'(z_{max,n})|\eta_n$   
 $\Rightarrow |\eta_{n+1}| > |\eta_n|$



# 4. Results: Comparison of numerical methods

- Starting point of chaotic behavior
  - Predictor-Corrector (PC) and RK4 start chaotic behavior earlier than BDF
  - BDF shows different output since the early steps
  - PC and RK4 shows the different outputs from around  $t \sim 22500$  timesteps

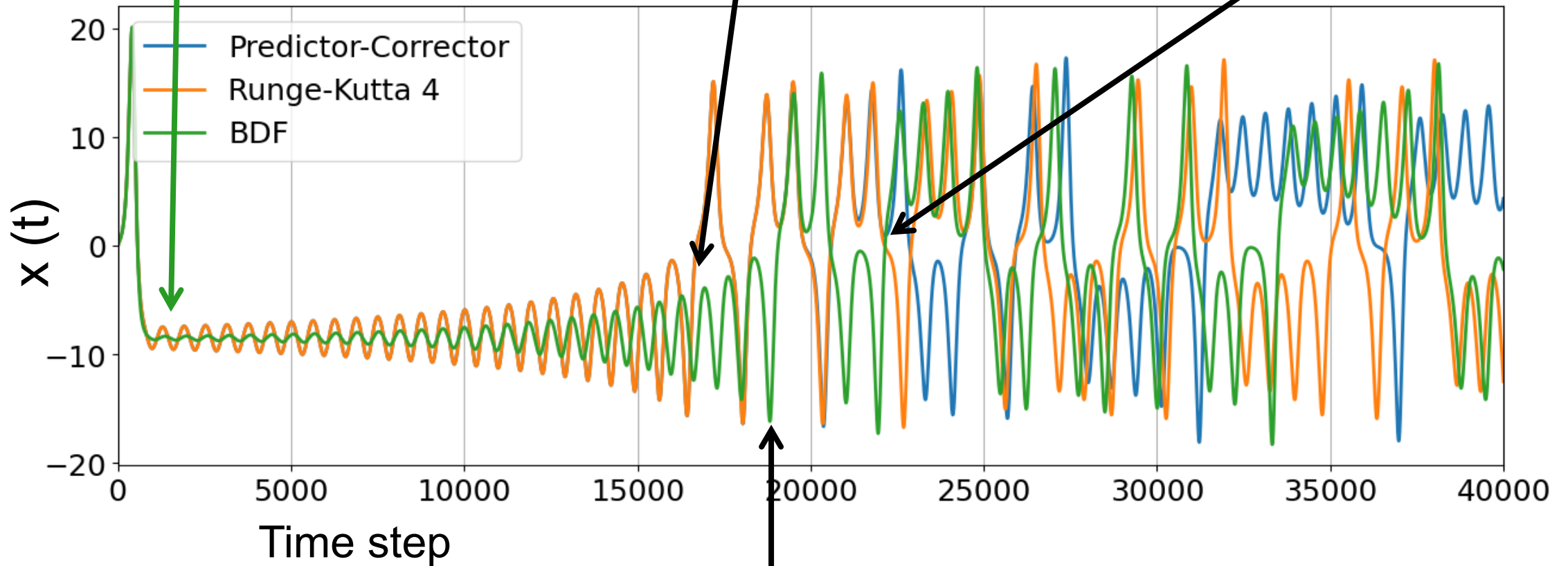


# 4. Results: Comparison of numerical methods

Different output: **BDF**

Chaotic behavior start:  
**PC, RK4**

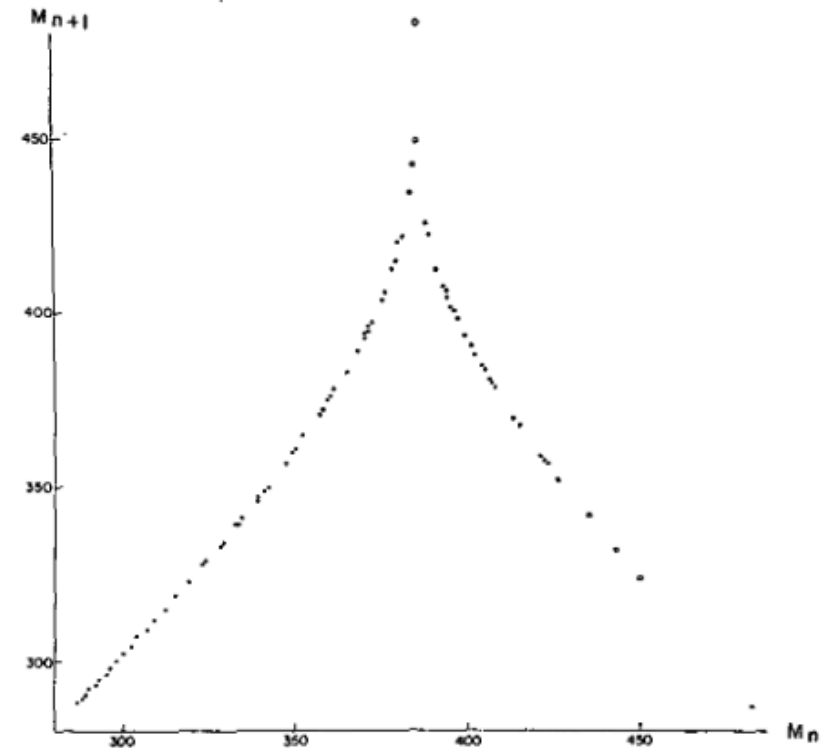
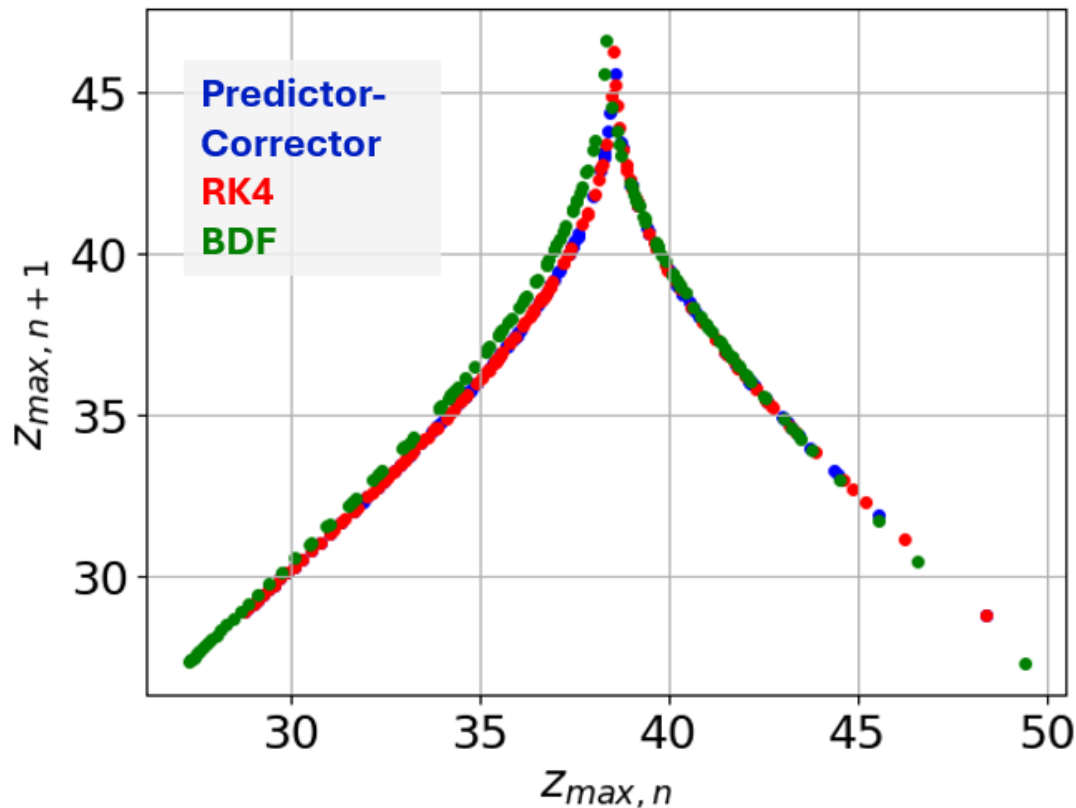
Different outputs:  
**PC, RK4**



Chaotic behavior starts: **BDF**

# 4. Results: Comparison of numerical methods

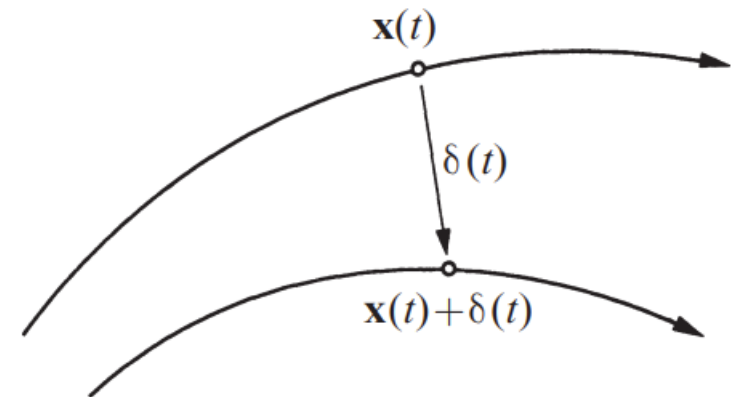
- Three methods shows almost the same Lorenz map
  - Predictor-Corrector (PC) and RK4 gave almost the same output
  - BDF method has the slightly smaller peak, but overall shape is identical to others
  - The shape is the same as Lorenz (1963)



Lorenz (1963)

# 4. Results: Sensitivity on initial conditions

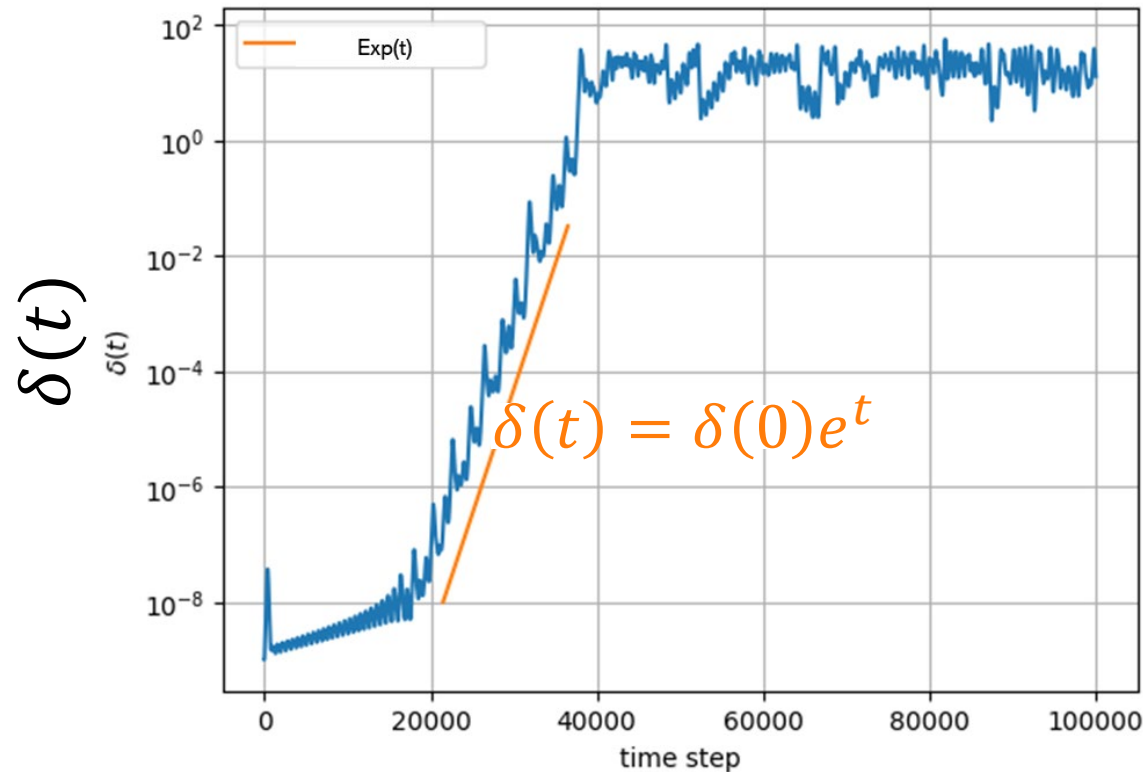
- The motion on the attractor exhibits sensitive dependence on initial condition
- How rapidly points apart from each other? Use **Liapunov exponent**
  - Consider two points  $x(t)$  and  $x(t) + \delta(t)$  initially close to each other
  - $\delta(t)$  is the distance of two point (see the figure)
  - $\delta(t)$  exponential increase, then we expect to have Liapunov exponent  $\lambda$  with
$$\delta(t) \sim \delta(0)e^{\lambda t}$$
  - $\delta(0)$  is the initial distance, say  $10^{-9}$
- Literature says,  $\lambda \sim 0.9$



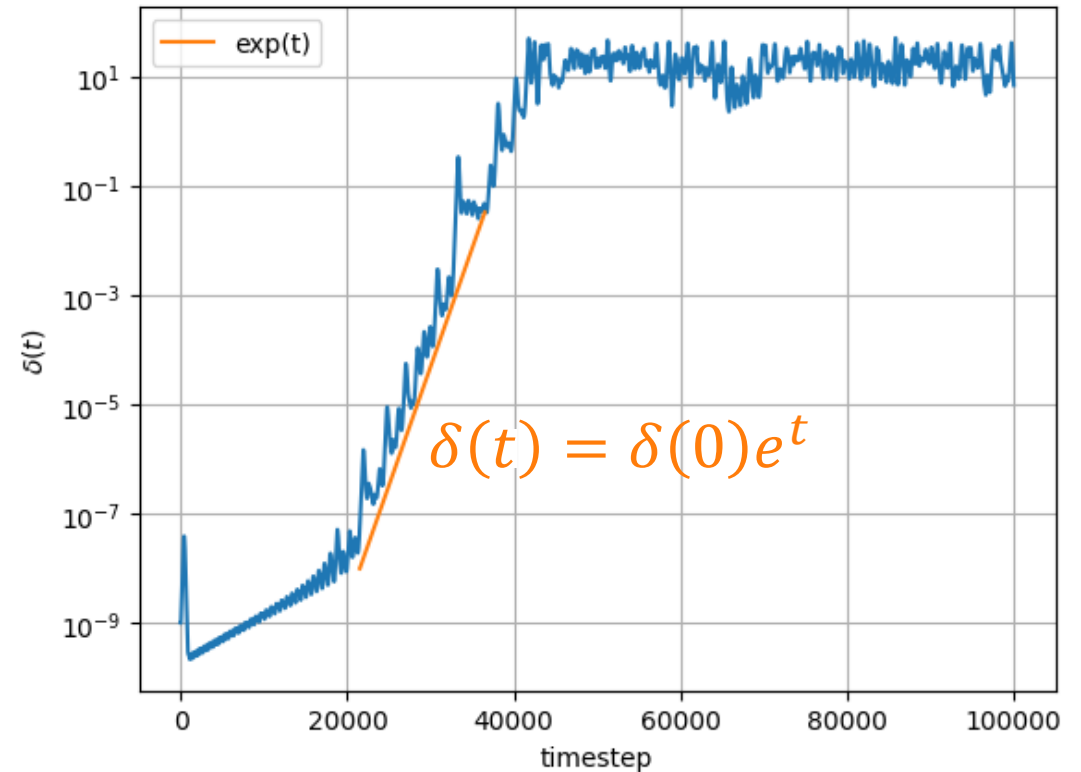


# 4. Results: Sensitivity on initial conditions

- Both RK4 and BDF method shows  $\lambda \sim 1$ , close to literatures
- $\delta(t)$  converge and constant, because the trajectory stay in a finite volume



RK4

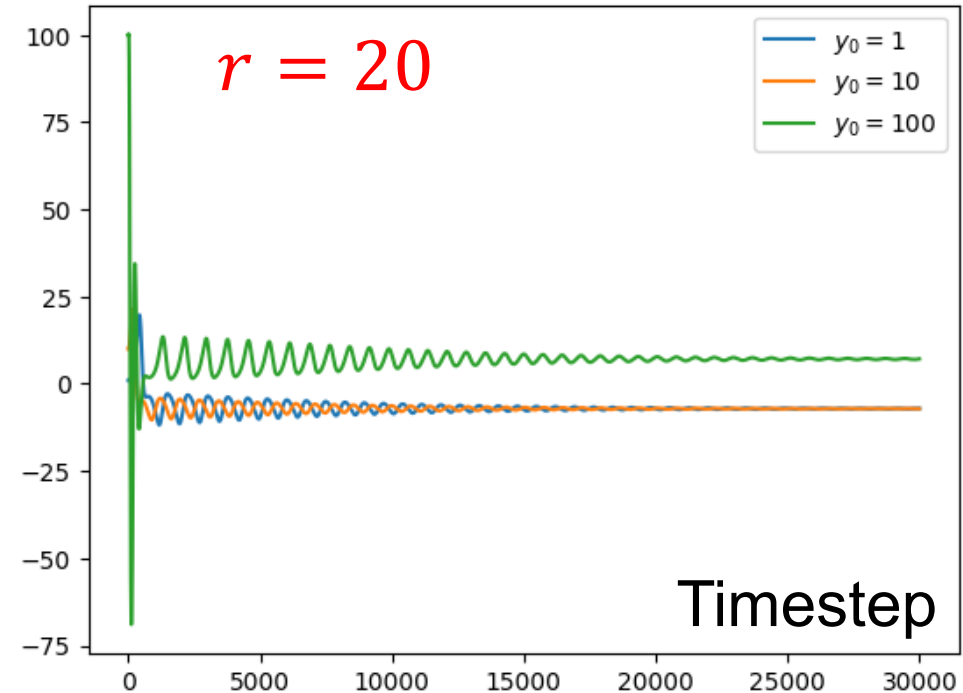
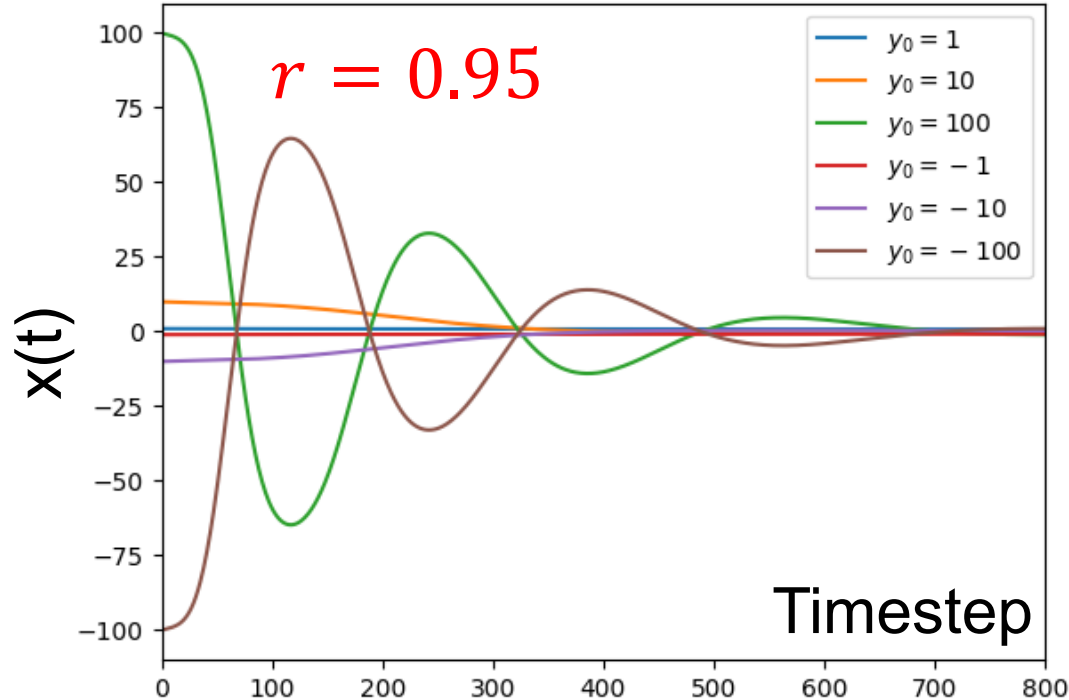


BDF

# 5. Parameter dependence

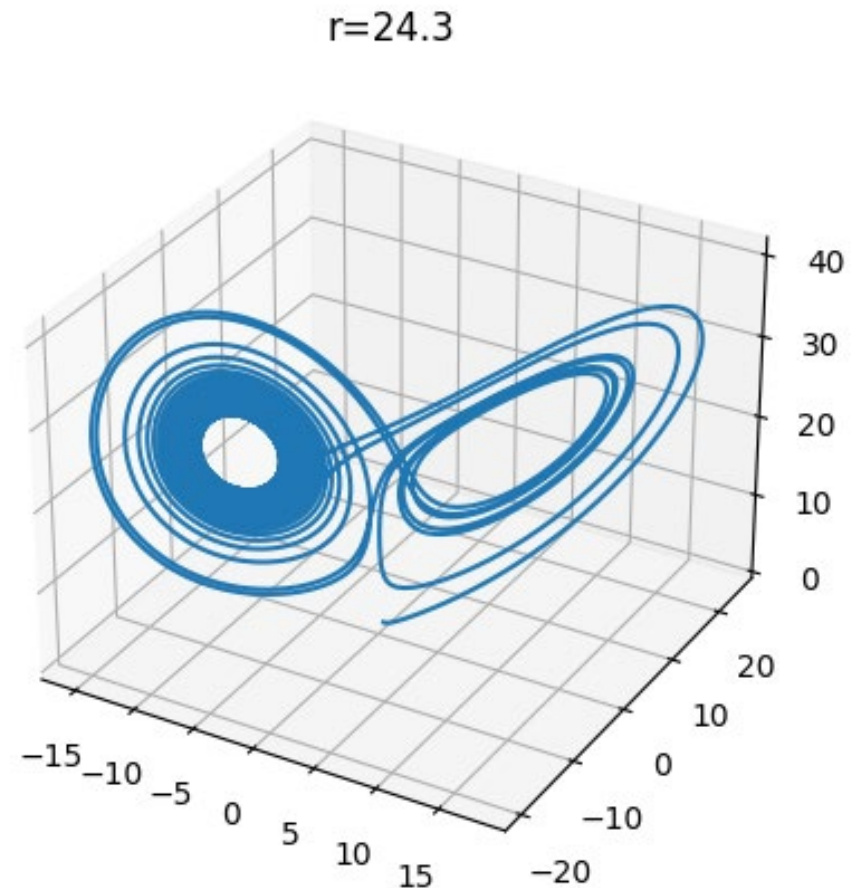
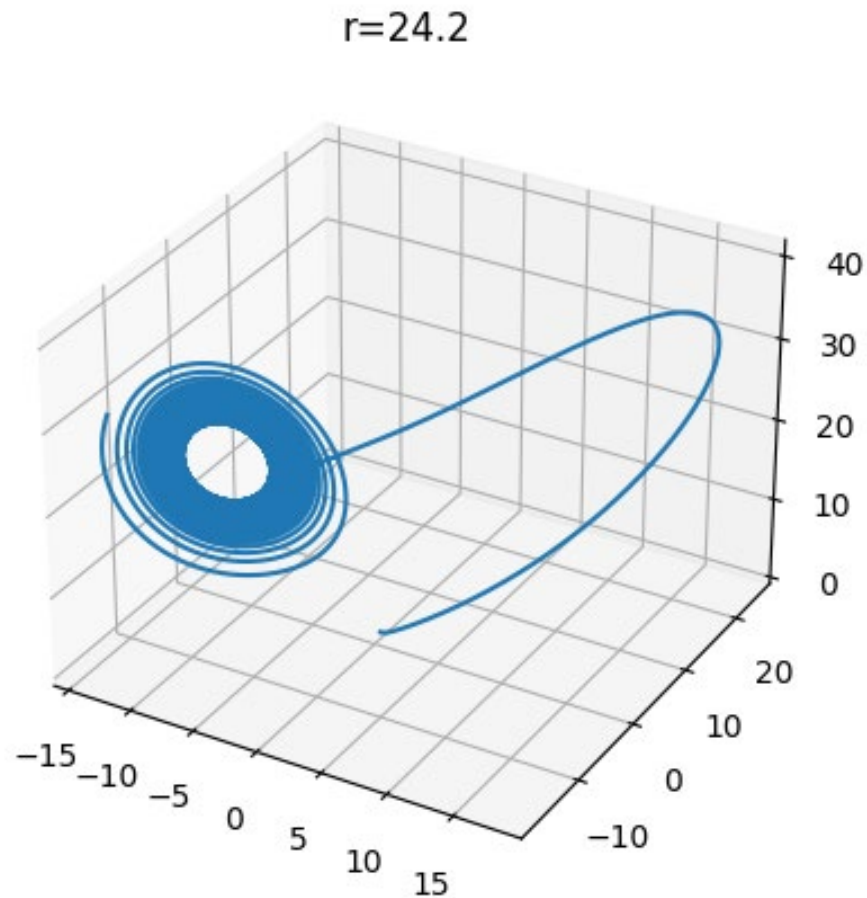
- Change the parameter  $r$
- When  $r < 1$ 
  - Trajectories fall into a point (point attractor)
- When  $1 < r < r_H$  ( $r_H$  critical value for behavior,  $\sim 25$ )
  - Two point attractors
  - Depends on the initial condition, trajectories fall into either two points

$$\begin{cases} \frac{dx}{dt} = \sigma(y - x) \\ \frac{dy}{dt} = rx - y - xz \\ \frac{dz}{dt} = xy - bz \end{cases}$$



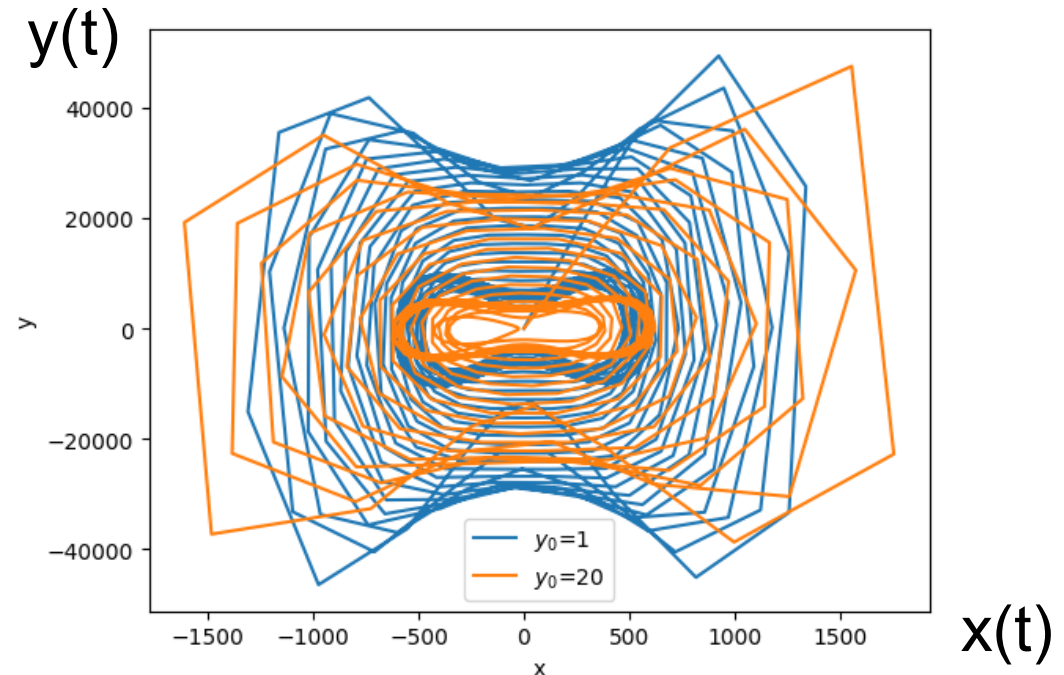
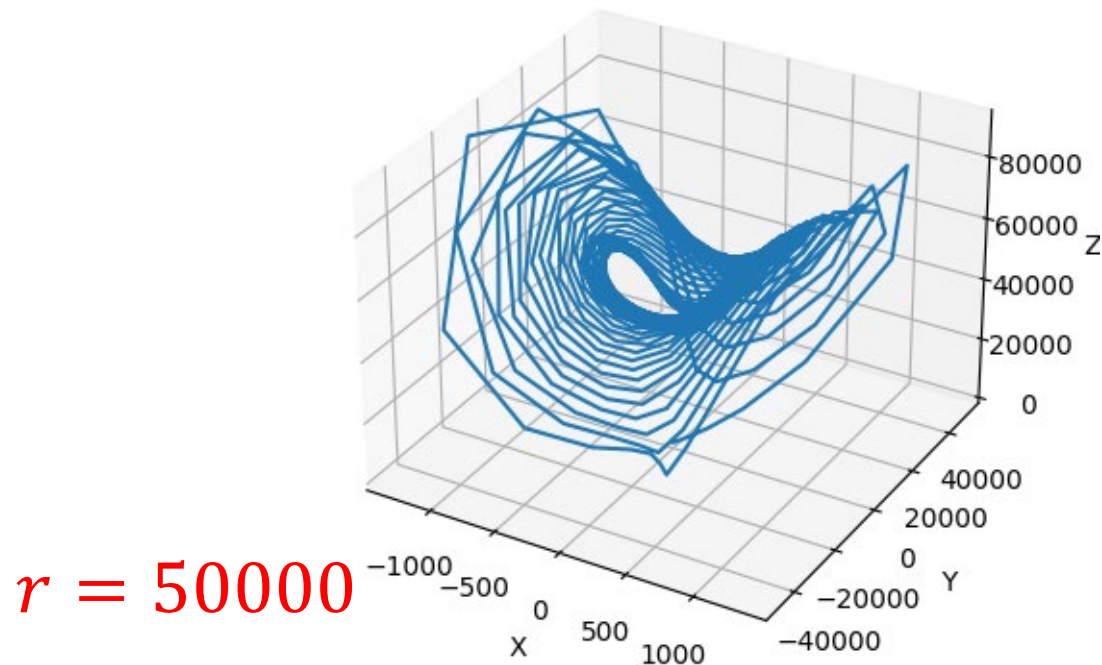
# 5. Parameter dependence

- When  $r \approx < r_H$ 
  - Transition from point attractor to strange attractor was observed around  $r \sim 24.2$



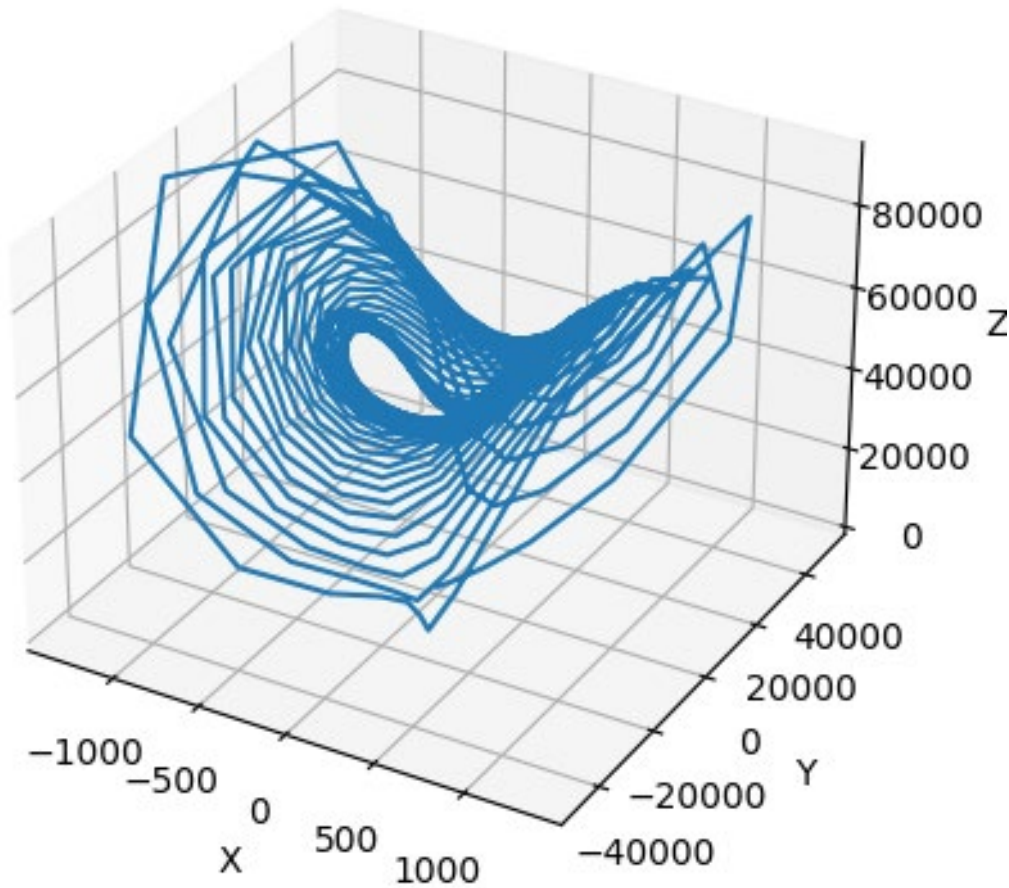
# 5. Parameter dependence

- When  $r_H \ll r$ 
  - Saddle shape trajectory was observed
  - Plotting  $x(t)$ - $y(t)$ , there is limit circle-like attractor
  - Trajectory falls to the circle after sufficient time
  - However, the circle shape depends on initial condition, thus not a limit circle with strictly definition (it should be independent from initial condition)

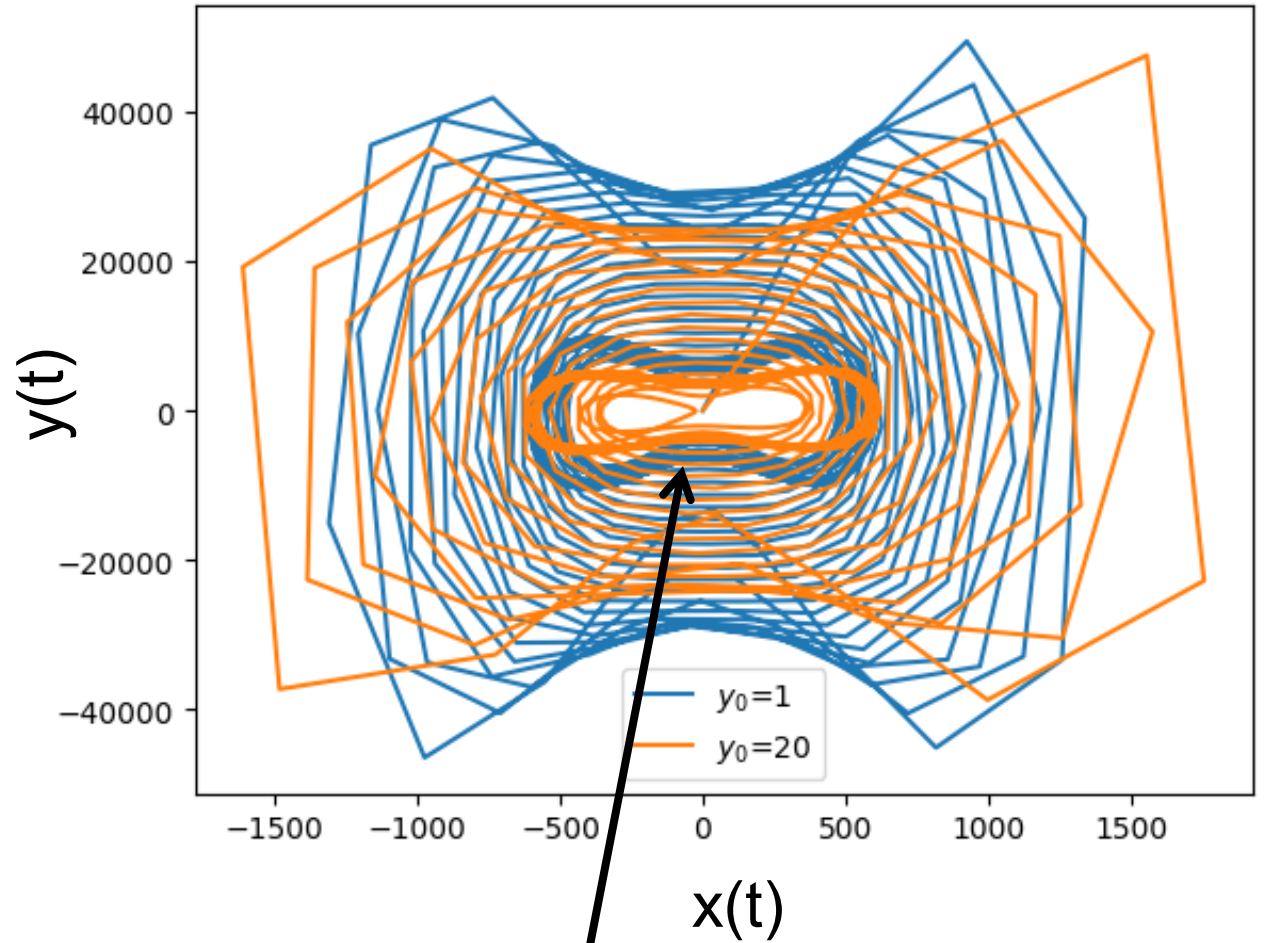




# 5. Parameter dependence



Saddle shape trajectory



Limit circle-like trajectory

# 6. Summary

- Lorenz model feature was investigated with three numerical methods
- Statistical quantities (Lorenz map, Liapunov exponent) were same as other research/ literature
- There were difference where the chaos starts, or location ( $x(t)$ ,  $y(t)$ ,  $z(t)$ ) after long time step
- With large  $r$ , saddle shape trajectory appeared
- There was limit circle-like trajectory, but the shape of the circle depends on the initial condition.

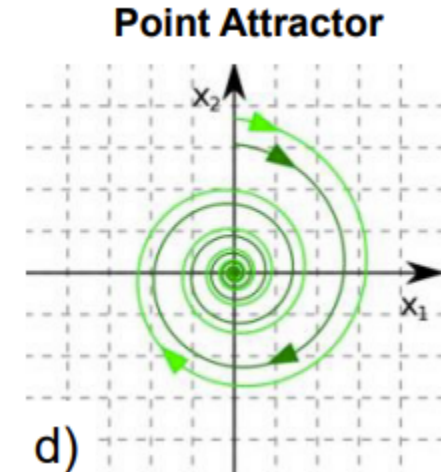
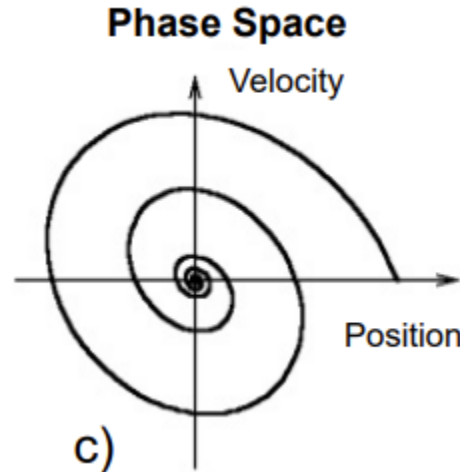
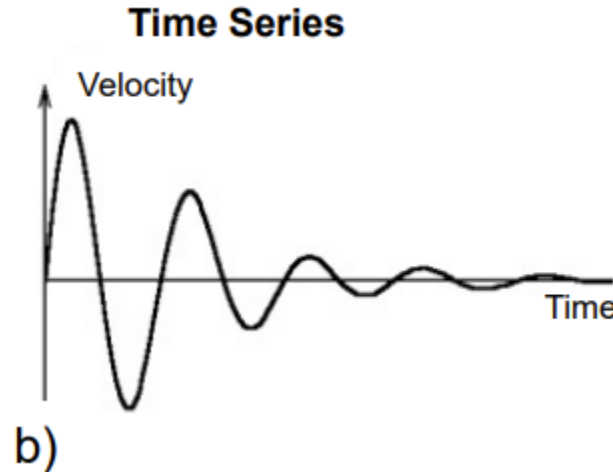
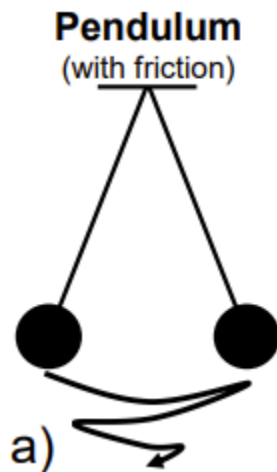
# Appendix A: What is attractor?

- Attractor :

➤ *an attractor is a set to which all neighboring trajectories converge. Stable fixed points and stable limit cycles are examples (Strogatz, 2015, p331)*

- Example 1: Pendulum with friction

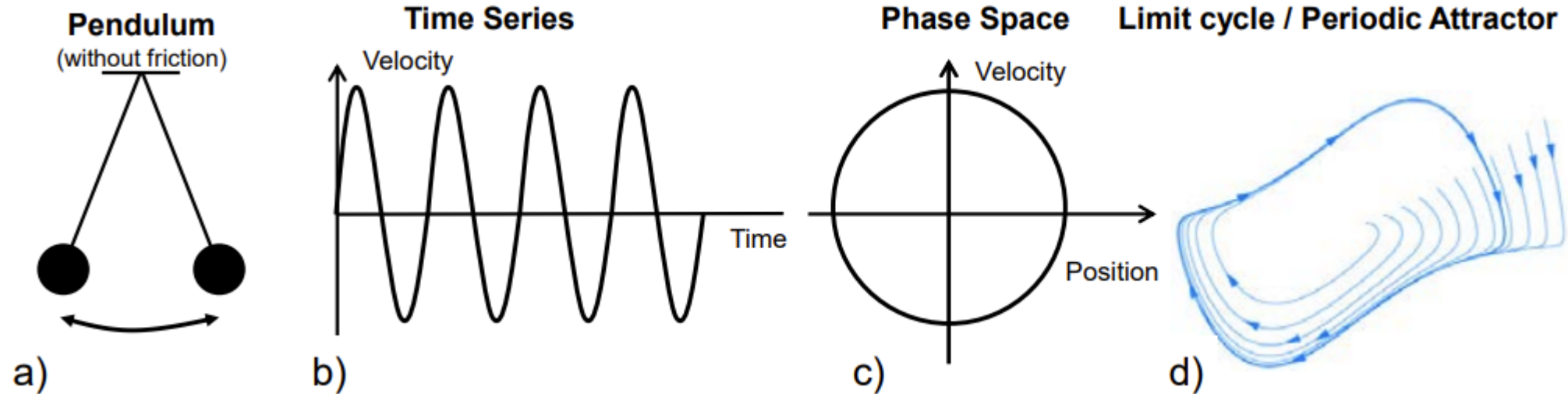
- The amplitude of the pendulum decreases over time and will stop at the end
- In phase space ( $\theta$ - $v$ ), Trajectory falls into a point = **Point Attractor**





# Appendix A: What is attractor?

- Example 2: Pendulum without friction
  - The amplitude of the pendulum is constant over time
  - In phase space ( $\theta$ - $v$ ), the trajectory has a circle shape and rounds on it  
= **Limit Circle**

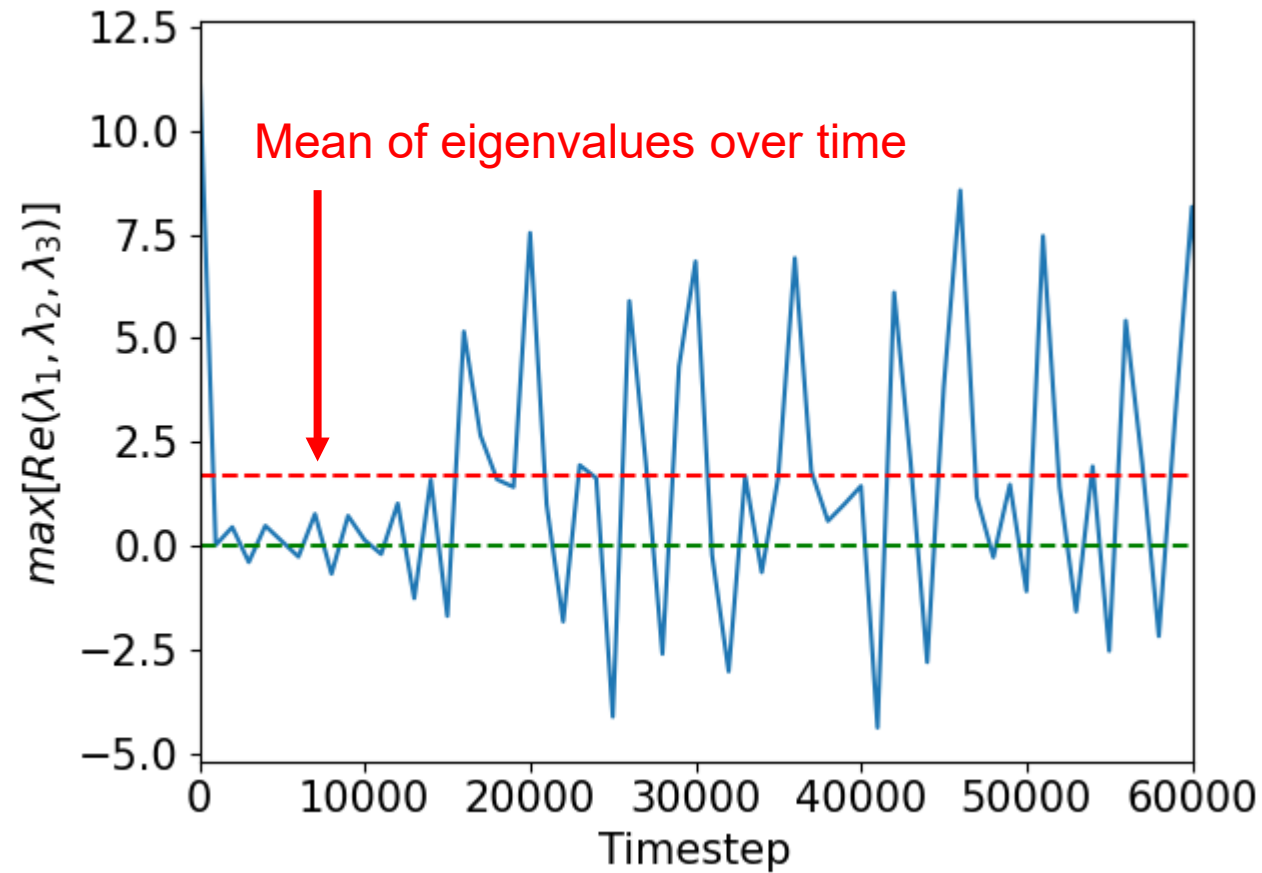


Hasse and Bekker, 2016

The strange attractor of chaos is neither point (point attractor) nor periodic (limit circle), but still stays in a finite region and moves inside it forever

# Appendix B: Numerical stability

- The eigenvalues of the initial condition are,  $8/3, -11 \pm \sqrt{1201}$
- Linear stability analyses: apply  $\Omega = -11 + \sqrt{1201}$   
 $\mathbf{u}(t) = \mathbf{u}(0)e^{\Omega t} \rightarrow \infty$
- Then, numerically unstable?
- The Jacobian (and its eigenvalues) changes over time, thus  $\Delta t * 100000 = 100$  might not too big to explode
- The maximums eigenvalue (real part) changes over time positive to negative, not immediately go to infinity (see left figure)



# Reference

- Robert G.L. Pryor, Jim E.H. Bright, Applying Chaos Theory to Careers: Attraction and attractors, Journal of Vocational Behavior, Volume 71, Issue 3, 2007, Pages 375-400
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- Strogatz, Steven. Nonlinear Dynamics and Chaos : with Applications to Physics, Biology, Chemistry, and Engineering. Second edition. Boulder, CO: Westview Press, a member of the Perseus Books Group, 2015.
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