TEPERATURA

mjera Ex neutreatenog gibanja čestica od kojih je sačinjeno tijelo (ili toplimoti seustar) → na odredknim temp čestice se dvugačuje gibaju M₁ V₂ M₃

Maxwellow jed po brzinama (milenoskopski)

makroskopske voličire Skala temperature i nulli zakon termodinarniki

Celsiusora skala: vocliste, lediste vode pri stema, flaten

Termodinamicka skala: tropa bothe vode ekstapolacija nule za p-+ ili 1 K = (+4v - to) /273, 16 -C

T[K] = t [°C] +273 IS

Toplina [J]

-nuredono pibanje - mi pretvorbi a men. rad postoje gubita (disipationi učma)
* može se prenositi ako su u kontaktu, a da se pritom re vrši rad

Siraye tyela sagrijavanjen

-> limamo sirayi . l= lo (1+daT)

- volumno sirreyje : V= No (1+80+)

de Tool 8 = dV ~ 3d

Molekulamo-kinetička teorija plinova

Idealni plin:

- medu mol sile - gernemanive - volumen mol zamimariv u odnosu na volumen pline

- sudani su savršeno elastični (ZOKO, ZOEL)

-spedenja brzina <u>pojedime</u> mod je <u>mula</u> (Brownovo gibanje) - mijerne medara je zemernanivo u odnoru na mijerne između sudara

- sye testice imaju stalnu i istu Marju

Tak idealnoz plima u MKT

$$P_{x}, P_{y}$$

$$P_{x}' = -P_{x}$$

$$P_{x}' = 2P_{x}; = P_{x} - P_{x}$$

$$P_{x}' = 2P_{x}; = 2P_{x};$$

$$\mp_i = \frac{2m(v_i x)^2}{2\alpha} = \frac{m(v_i x)^2}{2\alpha}$$

$$= \frac{2m(v_i x)^2}{2\alpha}$$

$$= \frac{m}{\alpha} \sum_{i=1}^{\infty} (v_i x)^2 = \sqrt{v_i^2}$$

$$= \sqrt{v_i^2}$$

$$= \sqrt{v_i^2}$$

$$= \sqrt{v_i^2}$$

 $F = \frac{M}{a} \sum_{n=1}^{\infty} (N \cdot x)^2 = N \sqrt{x^2}$ $= \sqrt{1 - x^2}$ $= \sqrt{1 - x^2}$ =

$$\frac{\sum (V_{ix})^{2} = \sum (V_{iy})^{2} = \sum (V_{i2})^{2}}{\sum (V_{ix})^{2} + \sum (V_{iy})^{2} + \sum (V_{iz})^{2} = 3\sum (V_{ix})^{2} = \sum V_{i}^{2} = NV^{2}}$$

$$P(+lae) = \frac{F}{S} = \frac{m}{(3)} \cdot \frac{1}{3} \cdot NV^{2} \cdot \frac{1}{(2)} = \frac{1}{3} \cdot \frac{N}{V} \cdot 2 \cdot \left(\frac{1}{2} \cdot mV^{2}\right) = P = \frac{2}{3} \cdot \frac{N}{V} \cdot E_{V} \cdot V$$

Jednadžba stanja => $PV = \frac{2}{3} N \bar{E} L$ plina: pV = N KT $= \frac{2}{3} N \bar{E} L$

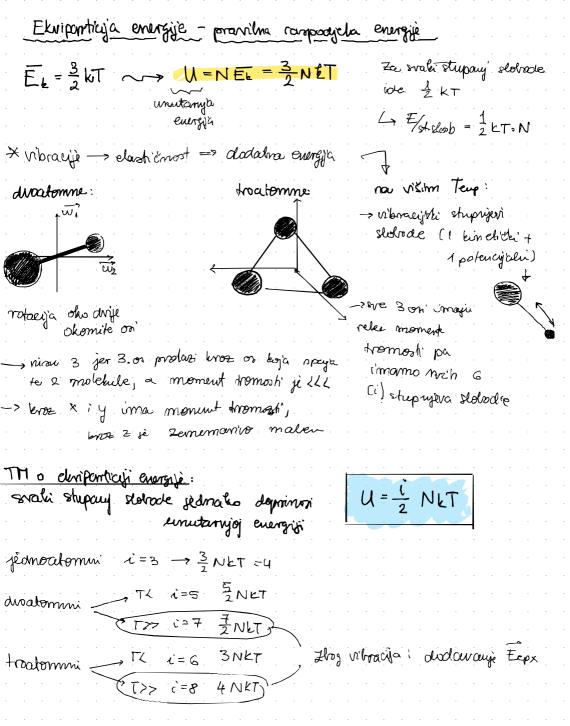
$$\frac{2}{3} \text{ MEz} = \text{MKT} \longrightarrow \frac{1}{2} \text{mV}^2 = \frac{3}{2} \text{ K}$$

$$\text{def appolulu}$$

$$\text{termodimonicle}$$

$$\text{temperature}$$

$$\frac{1}{2} \text{mV}^2 = \frac{3}{2} \text{ K}$$



Raspodjela čestka po bristnama

$$\vec{v} = \int_{-\infty}^{\infty} n(v)v dv \qquad \Rightarrow \vec{v} = \int_{-\infty}^{\infty} n_i v_i \\
\vec{v} = \int_{-\infty}^{\infty} n(v)v dv \qquad \Rightarrow \vec{v} = \int_{-\infty}^{\infty} n_i v_i \\
\vec{v} = \int_{-\infty}^{\infty} n(v)v dv \qquad \Rightarrow \vec{v} = \int_{-\infty}^{\infty} n_i v_i \\
\vec{v} = \int_{-\infty}^{\infty} n(v)v dv \qquad \Rightarrow \vec{v} = \int_{-\infty}^{\infty} n(v)v dv \qquad \Rightarrow \int_{-\infty}^{\infty} n(v)v dv dv \qquad \Rightarrow \int_{-\infty}^{\infty} n(v)v dv dv \qquad \Rightarrow \int_{-\infty}^{\infty} n(v)v dv dv \qquad \Rightarrow \int_{-\infty}^{\infty}$$

F= VSLT/11m

prosjètna kvad. enzuma) \ \vec{r}^2 = \ \frac{3kt}{m}

prosjecina Dizira

v= f(v)·vav =

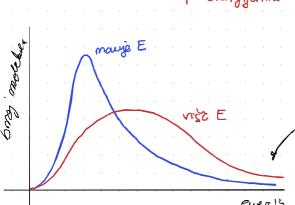
 $=\frac{\sum_{n=0}^{\infty} n_i v_i^n}{\sum_{n=0}^{\infty} n_i}$

ranpodicle otestica po en

I r f(r) dr = for r3 koust e - mr2 dr $= \sqrt{\frac{8kT}{\pi m}}$ $\int v^2 f(v) dv = \int_0^\infty v^4 koust = \frac{mv^2}{aut} dv$

Maxwell-Boltzmanova raspodyla - po energionna

 $E = \frac{1}{2}mv^{2}$ $\frac{dN}{dv} = \frac{dN}{dE} \cdot \frac{dE}{dv}$ $N_{E} = \frac{dN}{dE} = \frac{2N}{\sqrt{\pi L^{3}T^{3}}} \sqrt{E} \exp\left(-\frac{E}{kT}\right)$



. malo česh'ea misle E

· Steduja E 7 najvjengálnujó · posloje cestice ege imagiv vrlo visoke