

COULOMBOV ZAKON I EL. POLJE

Coulombov zakon - sila kojom djeluju dva naboja

$$\vec{F} = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r^2} \hat{r}$$

} sila kojom međudjeluju
DVA naboja

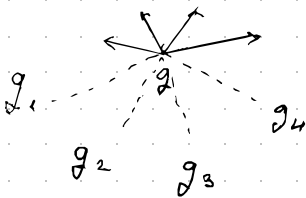
Puno veći od onog G iz gravitacije

ϵ_0 - permitivnost
vakuma

$$q_{\min} = e = 1,602 \times 10^{-19} \text{ C} \rightarrow [C] = [\text{sek Amp}]$$

! naboj je očuvan UVISEK! $Q_{\text{uk}}(\text{prije}) = Q_{\text{uk}}(\text{poslije})$

$$\vec{F} = \sum_{i=1}^n \vec{F}_i$$



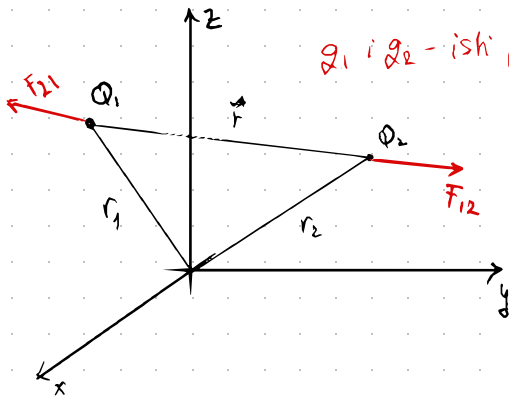
- Coulombova sila je odbojna
kada naboji imaju isti predznak

-- privlačna -- suprotni (+ i -)

q_1 i q_2 - isti predznak

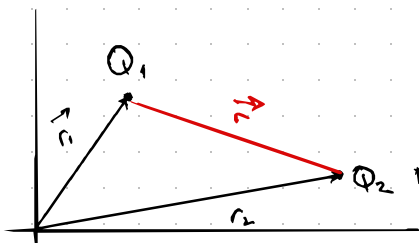
- neovisno jesu li naboji
čestica istoj ili suprotnos
predznaka, uvijek vrijedi:

$$F_{12} = -F_{21} = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{r_{12}^2} \hat{r}_{12}$$



ELEKTRIČNO POLJE

- prostor u kojem postoje zamjenjive el. silnice iz izvora Q_1



$$E_1 = \frac{F}{Q_2}$$

*za el. polja vrijedi načelo superpozicije

$$E_i = \frac{1}{4\pi\epsilon_0} \frac{Q_i}{r^2} \hat{r} \Rightarrow \vec{E}_i = \sum_{j \neq i} \vec{E}_j = \vec{E}_{i1} + \vec{E}_{i2} + \dots$$

$$\vec{F}_i = \sum_{j \neq i} \vec{F}_j = \vec{F}_{i1} + \vec{F}_{i2}$$

Coulombova sila koja odvaja dvije

$$F_{12} = \frac{1}{4\pi\epsilon_0} \frac{Q_1 \cdot Q_2}{|r_2 - r_1|^2} \cdot \frac{(r_2 - r_1)}{|\text{ostane smjer}|} \Rightarrow F_{12} = \frac{1}{4\pi\epsilon_0} \frac{Q_1 Q_2}{|r_2 - r_1|^2} \cdot \hat{r}$$

Ako silu F_{12} napišemo u obliku el. djelovanja Lorentzove sile: $F_{12} = q_2 E_1(r_2)$

$$\rightarrow \frac{1}{4\pi\epsilon_0} \frac{Q_1 Q_2}{|r_2 - r_1|^2} \hat{r} = q_2 E_1$$

$$\left| E_1 = \frac{1}{4\pi\epsilon_0} \frac{Q_1}{|r_2 - r_1|^2} \hat{r} \right.$$

Ako je raspodjela el. naboja u prostoru opisana volumnom, površinskom ili linijskom gustoćom naboja

\Rightarrow naboj $q \rightarrow dq'$ (nalazi se u okolini točke r')

\Rightarrow sumu zamjenjujemo integralom

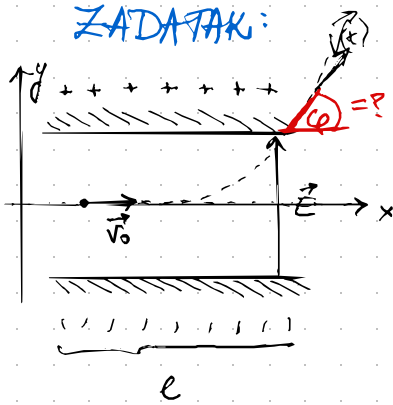
$$\underline{E[r] = \frac{1}{4\pi\epsilon_0} \int \frac{dq'}{|r - r'|^3} (r - r')}$$

linearna gustoća $\Rightarrow \frac{dq'}{dl} = \lambda \Rightarrow dq' = \overset{\text{linijska gustoća}}{\lambda[r']} \cdot dl$

površinska gustoća $\Rightarrow \frac{dq'}{dS} = \sigma \Rightarrow dq' = \overset{\text{površinska}}{\sigma[r']} dS$

volumna gustoća $\Rightarrow \frac{dq'}{dV} = \rho \Rightarrow dq' = \overset{\text{volumna}}{\rho[r']} dV$

ZADATAK:



$$\vec{v}_0 = v_0 \hat{x}$$

$$\vec{E} = E \hat{y}$$

$$F = m \cdot a = qE$$

$$m a_x = 0$$

$$\frac{dv_x}{dt} = 0$$

$$\int_{v_{x0}}^{v_x(t)} dv_x = 0$$

$$v_x - v_0 = 0$$

$$\underline{v_x = v_0 \text{ const.}}$$

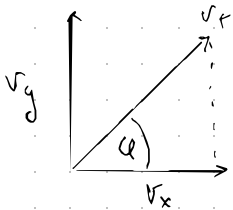
$$\frac{dx}{dt} = v_0$$

$$x(t) = v_0 t$$

$$\tan \phi = \frac{v_y}{v_x}$$

$$\phi = \arctan \left(\frac{qE}{mv_0} \cdot \frac{1}{v_x} \right) = \arctan \left(\frac{qE}{mv_0^2} \right)$$

$$v_0 = \frac{l}{t} \rightarrow t = \frac{l}{v_0}$$



$$y$$

$$m a_y = qE$$

$$\frac{dv_y}{dt} = \frac{qE}{m} \quad \int dt$$

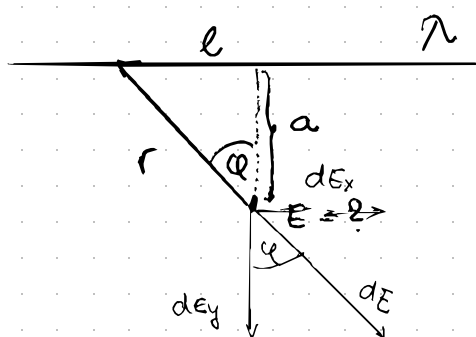
$$\int_0^{v_y(t)} dv_y = \frac{qE}{m} \int_0^t dt$$

$$\underline{v_y(t) = \frac{qE}{m} \cdot t}$$

↑ neodređeni
konstanta
odredjuje l (put)

$$v_y(t) = \frac{qE}{m} \cdot \frac{l}{v_0}$$

ZADATAK 3:



$$\vec{E} = \frac{1}{4\pi\epsilon_0} \int \frac{dq}{r^2} \hat{r}$$

$$\frac{dq}{dl} = \lambda \rightarrow dq = \lambda \cdot dl$$

$$E = \frac{1}{4\pi\epsilon_0} \int \frac{\lambda dl}{r^2} \hat{r}$$

$$dE = \frac{1}{4\pi\epsilon_0} \frac{\lambda dl}{r^2} \hat{r}$$

$$dE_x = 0$$

$$dE_y = (dE) \cos \varphi$$

$$E_y = \frac{1}{4\pi\epsilon_0} \int \frac{\lambda dl}{r^2} \cos \varphi$$

r želimo izraziti preko $\cos \varphi$ (a)

jer nam je to konstanta,

a minus bi bio sa l koji je nedefiniran

$$\cos \varphi = \frac{a}{r} \rightarrow r = \frac{a}{\cos \varphi}$$

$$r^2 = \frac{a^2}{\cos^2 \varphi}$$

$$\frac{l}{a} = \frac{\varphi}{a} \rightarrow l = \frac{l}{a} \cdot a \rightarrow \left(\frac{l}{a}\right)' = \frac{1}{\cos^2 \varphi}$$

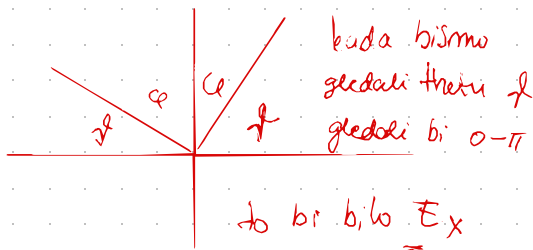
$$E_y = \frac{1}{4\pi\epsilon_0} \int \frac{\cos \varphi}{a^2} \cdot \lambda \cdot \frac{a}{\cos^2 \varphi} d\varphi$$

$$E_y = \frac{1}{4\pi\epsilon_0} \int \frac{\lambda \cos \varphi}{a} d\varphi$$

$$E_x = \frac{1}{4\pi\epsilon_0} \frac{\lambda}{a} \int_0^\pi \cos \varphi d\varphi$$

$$E_y = \frac{1}{4\pi\epsilon_0} \frac{\lambda}{a} \int_{-\pi/2}^{\pi/2} \cos \varphi d\varphi$$

koje kutove za granice uvesti?

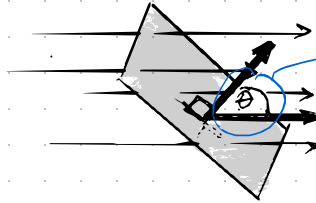
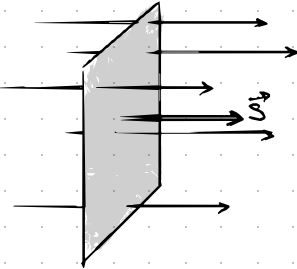


TOK električnog polja = Φ

$$\Phi_E = \oint_S \vec{E} \cdot d\vec{S}$$

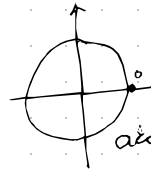
zatvorena
ploha

- promatramo površinu S kroz
č. polje \rightarrow kroz nju prolazi više ili
manje silnica



kut između silnice el. polja
i smjera normale

skalarni umnožak $|\vec{E}| |\vec{S}| \cos \theta$



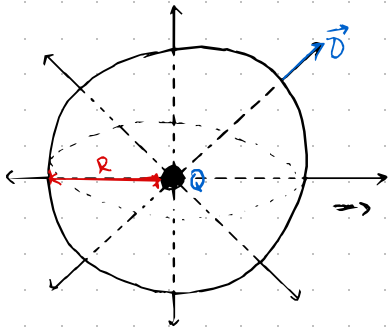
isto je kut neki θ

vektor \vec{S} ima iznos vektora površine
• smjer vektora normale

$$\Phi_E = \oint_S E dS \cos \theta$$

**TOK EL. POLJA TOČKASTOG
NABOJA Q (zatvorena ploha)**

\rightarrow zamišljena ploha S je
kugla (sfera)



silnice okomito probijaju kuglinu plohu

• iznos vektora D jednak je u svim točkama
kugline plohe

\rightarrow tok el. polja jednak umnošku iznosa vektora \vec{D} i ukupne
površine kugle: $\Phi = 4R^2 \pi D = 4R^2 \pi \epsilon E$

u sredini naboj Q ,
polje na površini

$$E = \frac{Q}{4\pi \epsilon R^2}$$

TOK električnog polja kroz neku
kuglinu plohu proporcionalan je
naboj u središtu te kugle
(i neovisan je o R)

\downarrow
Gauss

$$\rightarrow D = \epsilon E = \frac{Q}{4\pi R^2}$$

$$\rightarrow \Phi = D \cdot S = \frac{Q}{4\pi R^2} \cdot 4R^2 \pi = Q$$

JAKOST EL. POLJA:

$$\oint \vec{E} \cdot d\vec{S} = 0$$

Gaussov zakon

→ d. toč po zatvorenoj plosni ~ ukupnom naboju koji smo obuhvatili tom plosnom

$$\oint_S \vec{E} d\vec{S} = \frac{Q_{unutra}}{\epsilon_0}$$

→ Fizikalna posljedica toga što elektrostatičke nle. opada s kvadratom udaljenosti $F \sim \frac{1}{r^2}$

Gaussov teorem - teorem o divergenciji

$$\oint \vec{E} d\vec{S} = \int \nabla \cdot \vec{E} dV$$

- zatvorena plosna S obuhvata neki volumen V

volumna gustoća: $\rho_g = \rho[r] dV$

$$Q = \int \rho dV$$

$$\oint \vec{E} d\vec{S} = \frac{Q}{\epsilon} \Rightarrow \frac{1}{\epsilon} \int \rho dV = \int \nabla \cdot \vec{E} dV \Rightarrow$$

PRVA MAXWELLOVA

$$\frac{\rho}{\epsilon_0} = \nabla \cdot \vec{E}$$

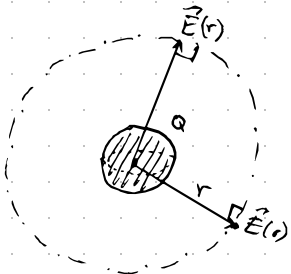
JEDNADŽBA!

→ Gaussov zakon vrijedi bez obzira na oblik

pošto točkasti naboj

EL. polje naboje čestice

$$\oint_S \vec{E} d\vec{S} = \frac{Q}{\epsilon} = \Phi_e$$



$$dV = dx dy dz$$

$$dS = dx dy$$

$$V = \frac{4}{3} \pi r^3$$

$$dS = r^2 \sin \varphi d\varphi d\psi$$

$$S = r^2 \int_0^\pi \sin \varphi d\varphi \int_0^{2\pi} d\psi$$

ploha koja obuhvata naboj Q

• $E(r)$ je \perp na sferu zbog prostorne simetrije

- polje ima isključivo komponentu koja je okomita na g jer tangencijalna ne može biti izražena (Riči sustav ne sadrži ništa jedno svojstvo pomoću kojeg bismo odredili smjer)

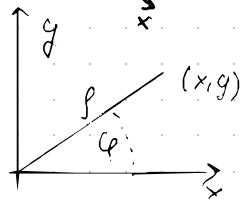
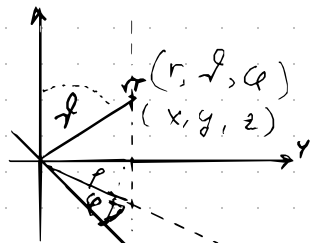
$$\oint_S \vec{E} d\vec{S} = \frac{Q}{\epsilon_0} = \Phi_e$$

ubacimo u kalkulator broj.

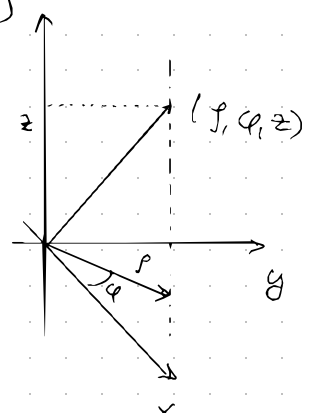
$$ES = \frac{Q}{\epsilon_0} \Rightarrow E r^2 \int_0^\pi \sin \varphi d\varphi \int_0^{2\pi} d\psi$$

$$\Phi_e = 4\pi r^2 E = \frac{Q}{\epsilon_0}$$

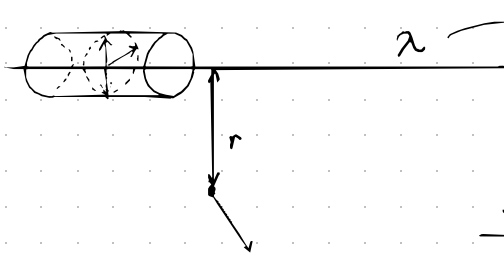
2 SFERNI KOORD.



CILINDRIČN:



POLE jednoducho nabitene žice



linearna hustota $\lambda = \frac{Q}{l}$
 $\hookrightarrow Q = \lambda l$

$$E = \frac{\lambda}{2\pi\epsilon_0 r} \quad (\text{dobro postupak})$$

- riešavame pomocou cylindrických plocha

$$d\vec{S} = dx dy \hat{z}$$

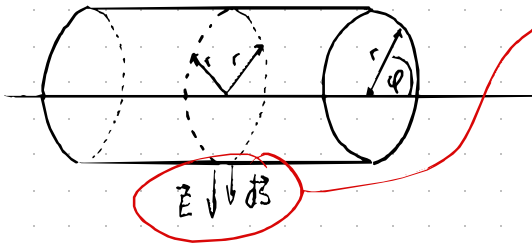
$$dS = r^2 \sin \theta d\theta d\phi$$

$$dS = r d\phi dz$$

$$dV = dx dy dz$$

$$dV = r^2 dr \sin \theta d\theta d\phi$$

$$dV = r dr d\phi dz$$



vekt. umnoženie
 kút je 0, $\cos = 1$

$$\oint \vec{E} d\vec{S} = \int E dS \cos(\vec{E}, d\vec{S})$$

$$\oint \vec{E} d\vec{S} = \oint E dS = \oint \frac{E(r)}{\text{konst}} r d\phi dz$$

$$E(r) r \int_0^{2\pi} d\phi \int_{-l}^l dz = \underline{\underline{E(r) r 2\pi 2l}}$$

Gaussov zákon: $\oint \vec{E} d\vec{S} = \frac{Q}{\epsilon_0}$

$$E(r) r 2\pi 2l = \frac{Q}{\epsilon_0}$$

$$Q = \lambda l, \quad l = 2l \rightarrow Q = 2l \lambda$$

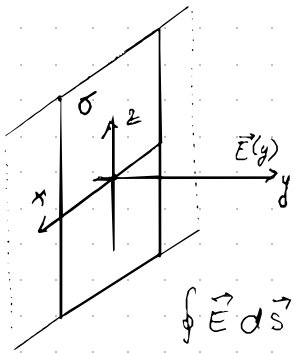
$$E(r) r 2\pi 2l = \frac{\lambda 2l}{\epsilon_0}$$

$$E = \frac{\lambda}{2\pi r \epsilon_0}$$

$$\underline{\underline{E = \frac{\lambda}{2\pi r \epsilon_0}}}$$

* pole umier i izvan
 nabijene kugle \Rightarrow ADI TORNE

POLJE jednakno nabojeve plohe



- Polje se ne mijenja po x i z , samo po y jer je normalan po y približavamo ili udaljavamo od plohe

$$\sigma = \frac{Q}{S} - \text{plošna gustoća naboja}$$

→ okomita je na ravninu

$$\sigma = \frac{dq}{dS} \quad dS = dx \cdot dz \text{ (površina)}$$

$$S = \int_{-x}^x dx \int_{-z}^z dz$$

$$S = 2x \cdot 2z$$

$$\oint E(y) dx dz = \oint E(y) dx dz = \oint 2E(y) dx dz = 2E(y) \int_{-x}^x dx \int_{-z}^z dz$$

$$\Rightarrow 2E(y) \cdot 2x \cdot 2z = \frac{Q}{\epsilon_0} \rightarrow \sigma \cdot S = \frac{\sigma S}{\epsilon_0} \cdot 2x \cdot 2z$$

otkud taj 2?
zbog zatvorenog S?

$$2E(y) = \frac{\sigma}{\epsilon_0} \rightarrow \boxed{E_y = \frac{\sigma}{2\epsilon_0}}$$

MAXWELLOVA JEDNADŽBA

► POLJE unutar i izvan jednolično nabojeve kugle

$$\rho = \frac{dQ}{dV} \rightarrow \text{gustoća naboja}$$

$$\hookrightarrow Q = \int \rho dV$$

- svaka zatvorena ploha obuhvata neki volumen

$$\oint E = \frac{Q}{\epsilon_0} \quad \oint \vec{E} d\vec{S} = \frac{Q}{\epsilon_0} = \frac{\int \rho dV}{\epsilon_0} = \frac{1}{\epsilon_0} \int \rho dV$$

↓
odakle slijedi Gaussov TM (TM o divergenciji)

ta ploha obuhvata neki volumen (zatvoreni S)

Gaussov TM:

$$\rightarrow \oint \vec{F} \cdot d\vec{S} = \int \vec{\nabla} \cdot \vec{F} dV$$

$$\Rightarrow \oint \vec{E} d\vec{S} = \frac{Q}{\epsilon_0} \quad \oint \vec{E} d\vec{S} = \int \vec{\nabla} \cdot \vec{E} dV$$

$$\Rightarrow \frac{1}{\epsilon_0} \int \rho dV = \int \vec{\nabla} \cdot \vec{E} dV$$

$$\leadsto \int \frac{\rho}{\epsilon_0} dV = \int \vec{\nabla} \cdot \vec{E} dV$$

$$\Rightarrow \boxed{\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}}$$

Ako imamo zatvorenu plohu i ako tražimo S možemo ga naći pomoću divergencije

analogno će vrijediti za mag polje, ali o tome kasnije
 $\nabla \cdot \vec{B} = 0$

1. MAXWELLOVA jednačica

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

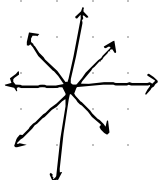
$$\vec{\nabla} = \frac{d}{dx} \hat{x} + \frac{d}{dy} \hat{y} + \frac{d}{dz} \hat{z}$$

$\vec{\nabla} \cdot U$ - gradient

$\vec{\nabla} \cdot \vec{E}$ - divergencija

$\vec{\nabla} \times \vec{E}$ - rotacija

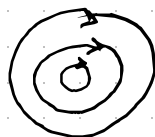
divergencija $\neq 0$
($\vec{\nabla} \cdot \vec{E} \neq 0$)



ili



divergencija = 0
($\vec{\nabla} \cdot \vec{E} = 0$)



Neko Marko računanje
gradienta i ostalih majija:

$$U(r) = \int \frac{\partial U}{\partial r} d\vec{r}' \rightarrow \text{integracija gradienta koju znamo otprije}$$

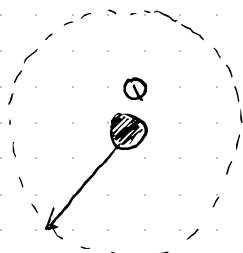
$$\hookrightarrow \text{isto možemo zapisati: } \int U dr = \int \frac{\partial U}{\partial r} dr'$$

$$\int E dr' = \int (\vec{\nabla} \times \vec{E}) d^2r \quad \text{vekt}$$

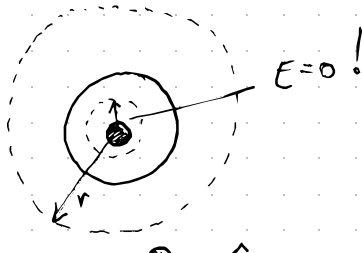
$$\rightarrow \int \vec{E} d^3r = \int (\vec{\nabla} \cdot \vec{E}) d^3r \quad \text{skal} \quad \text{HUH?}$$

Unutar kugle

* na ljusci



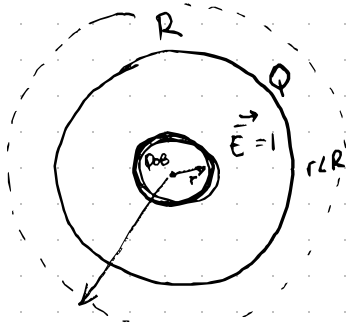
$$\vec{E} = \frac{Q}{4\pi\epsilon_0 r^2} \hat{r}$$



$$E(r) = \frac{Q}{4\pi\epsilon_0 r^2} \hat{r}$$

▶ HOMOGENA KUGLA:

$$\rho = \frac{Q}{V} \rightarrow \rho = \frac{dq}{dv} \rightarrow \int (\rho dV) = Q$$



$$\vec{E} = \frac{Q}{4\pi\epsilon_0 r^2} \hat{r}$$

KONZERVATIVNOST COULOMBOVE SILE & EL. POTENCIJAL

Prva Maxwellova jednačina: $\nabla \times \vec{E} = \frac{1}{c^2} \frac{\partial \vec{B}}{\partial t}$

Coulombova sila je konzervativna jer vrijedi: $\nabla \times \vec{F} = 0$

↳ to vrijedi samo u elektrostatici, ne i u magnetizmu *

↳ rotacija je 0

→ sledi: postoji skalarno polje Φ za koje vrijedi $F = -\nabla \Phi$ karmije smo nazvali potencijalom energije

analogno vrijedi: $\nabla \times \vec{E} = 0$ (d. polje je konzervativno)

⇒ -1- postoji negativni gradijent $E = -\nabla V$ * potencijal (skal. polje $V(x, y, z)$)

$$F = -\nabla \Phi = -\frac{\partial \Phi}{\partial l} \quad \text{po dužini vodica}$$

$$E = -\nabla V = -\frac{\partial V}{\partial l} \quad E = \frac{F}{q}$$

$$\Rightarrow V = -\int \vec{E} d\vec{l}$$

$$\Phi = -\int \vec{F} d\vec{l} \quad \text{elektr. potencijal}$$

$$\Rightarrow V = -\int \frac{F}{q} d\vec{l} = \frac{1}{q} \Phi \Rightarrow \Phi = qV \rightarrow \boxed{U = qV}$$

- el. potencijal po jednostavnoj analogiji s potencijalom energijom možemo prikazati kao:

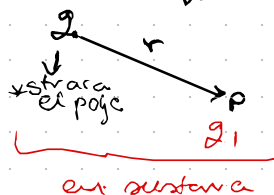
$$V = -\int \vec{E} d\vec{r} \rightarrow d\vec{r} \cdot \hat{r} \quad * E = \frac{1}{4\pi\epsilon_0} \cdot \frac{q}{r^2} \cdot \hat{r} \quad \text{jed. vektor smjera}$$

$$V = -\int \frac{q}{4\pi\epsilon_0 r^2} \cdot \hat{r} \cdot d\vec{r} \quad \text{LIN. ALG. } \hat{r} \cdot \hat{r} = 1$$

$$V = -\int \frac{q}{4\pi\epsilon_0 r^2} dr \rightarrow \boxed{V = \frac{q}{4\pi\epsilon_0 r} + C} \quad \text{* isto kao potencijalna energija}$$

- potencijal koji stvara neki q u udaljenosti r

$$\hookrightarrow U = q_1 V$$



- dovodimo q_1 iz beskonačnosti u točku potencijala V

⇒ možemo izračunati energiju sustava

$$U = q_1 V = \frac{q_1 q}{4\pi\epsilon_0 r}$$

- vezu el. polja i

el. potencijala:

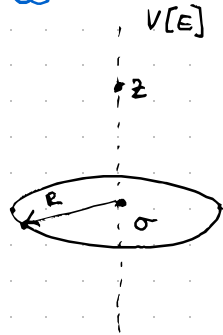
$$\vec{E} = -\nabla V \quad (\text{gradijent potencijala})$$

Zad. El. potencijal na bilo kojoj točki sred. osi homogeno nabijenog diska je dan izrazom: $V(z) = \frac{\sigma}{2\epsilon_0} (\sqrt{z^2 + R^2} - z)$.

Izvedite izraz za el. polje:

$$V(z) = \frac{\sigma}{2\epsilon_0} (\sqrt{z^2 + R^2} - z)$$

$$\sigma = \frac{Q}{R\pi} \quad (\text{plošna gustoba})$$



$$\vec{E} = \left(\frac{\partial V}{\partial x} \hat{x} + \frac{\partial V}{\partial y} \hat{y} + \frac{\partial V}{\partial z} \hat{z} \right)$$

$$\vec{E} = - \frac{\partial V}{\partial z} \hat{z}$$

$$E = - \frac{\sigma}{2\epsilon_0} \left(\frac{1}{z} - \frac{1}{\sqrt{z^2 + R^2}} \cdot 2z - 1 \right) = - \frac{\sigma}{2\epsilon_0} \left(\frac{2z}{z^2 + R^2} - 1 \right)$$

Ako je $\Phi_E = \oint \vec{E} d\vec{s}$ (Gaussov zakon)

$$\Phi_E = \oint \vec{E} d\vec{s} = \frac{Q_{\text{obuhvaćeni}}}{\epsilon_0} \quad (\Phi \sim \int \rho dV)$$

$$\rightarrow \oint \vec{E} d\vec{s} = \int \vec{\nabla} \cdot \vec{E} dV$$

nešto???